

**Assessment of gas hydrate saturations in the Makran
accretionary prism using the offset dependence of seismic
amplitudes**

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ABSTRACT

We estimate the saturations of gas hydrate and free gas based on measurements of seismic reflection amplitude versus offset (AVO) for a bottom simulating reflector (BSR) coupled with rock physics modeling. We apply the approach to data from a seismic line in the Makran accretionary prism in the Arabian Sea. The results reveal lateral variations of gas hydrate and free gas saturations between 4-29% and 1-7.5%, respectively, depending on the rock physics model used to relate seismic velocity to saturation. The approach is simple and easy to implement.

19 INTRODUCTION

20 Gas hydrates are crystalline solids of methane and water that are formed at high
21 pressure and low temperature (Kvenvolden, 1998; Sloan, 1998). They have attracted attention
22 because of their widespread occurrence in permafrost and outer continental margins; their
23 potential as future major energy resource; and their role in climate change and geo-hazards
24 (Dillon et al., 1991; Paul et al., 1991; Paull and Dillon, 2001; Milkov et al., 2004; Taylor and
25 Kwan, 2004; Makogon et al., 2007). The identification and quantification of gas hydrate are
26 essential to evaluate their resource potential and assess their associated environmental hazard.
27 Gas hydrate can be identified on multi-channel marine seismic reflection data by mapping an
28 anomalous bottom-simulating reflector (BSR), which marks the base of the high-pressure and
29 relatively low-temperature zone in which hydrates are stable (Andreassen et al., 1990,
30 Kvenvolden et al., 1993). Seismic reflections from BSRs exhibit a wide range of AVO
31 characteristics that depend upon the saturation and distribution of hydrate above and free gas
32 below the BSR (Hyndman and Spence, 1992; Andreassen et al., 1997; Ecker et al., 1998;
33 Carcione and Tinivella, 2000; Yuan et al., 1999). Hydrate may form part of the pore fill or
34 part of the rock matrix, and where present gas may be distributed uniformly or in patches
35 throughout the pore space (Dvorkin et al., 1999; Helgerud et al., 1999). Uniform and patchy
36 gas distributions yield different AVO behaviors because of velocity dispersion in partially
37 gas-saturated sediments (White, 1975; Lee, 2004). We assume a uniform distribution of free
38 gas below the BSR, although the method can be adopted for patchy saturation. Overpressure
39 in the trapped gas also affects the AVO responses, but it is hard to isolate its effect from the
40 effect of gas on P-wave reflection data because both overpressure and free gas can cause low
41 P-wave velocity (Tinivella, 2002). In the hydrocarbon industry, AVO attributes have gained
42 considerable popularity for predicting the lithology and reservoir characteristics (Castagna
43 and Backus, 1993; Castagna and Smith, 1994; Castagna et al., 1998). Several authors

(Hyndman and Spence, 1992; Andreassen et al., 1997; Ecker et al., 1998; Carcione and Tinivella, 2000; Yuan et al., 1999; Chen et al., 2007) have estimated hydrate and/or gas saturations from BSR's using conventional AVO modeling and various rock physics models. Here we present a simple approach that compares the AVO intercept (A) and gradient (B) estimated from the BSR with those from two gas-hydrate models to quantify the saturations of gas-hydrate and free gas across a BSR. For a theoretical calculation of A and B values, we used the Biot-Gassman approach of Lee (2002), and the effective medium theory (EMT) of Helgerud et al. (1999), in both of which hydrate are considered as part of the sediment frame. We demonstrate through a field example in the Makran accretionary prism (Minshull et al., 1992; Sain et al., 2000) that a plot of A versus B can be used to detect free gas below the BSR and quantify the amount of hydrate and gas along a seismic reflection profile.

THEORY

For a plane wave incident at an interface between two semi-infinite isotropic homogeneous elastic half-spaces, and for angles of incidence (θ) up to $\sim 30^\circ$, the P-wave reflection coefficient $R(\theta)$ can be expressed (Shuey, 1985) as

$$R(\theta) \approx A + B \sin^2 \theta, \quad \dots\dots\dots (1)$$

where,

$$A = \frac{1}{2} \left(\frac{\Delta V_P}{V_P} + \frac{\Delta \rho}{\rho} \right) \text{ and } B = A A_0 + \frac{\Delta \sigma}{(1 - \sigma)^2}, \quad (2)$$

$$A_0 = B_0 - 2(1 + B_0) \frac{1 - 2\sigma}{1 - \sigma}, \quad B_0 = \frac{\Delta V_P / V_P}{\Delta V_P / V_P + \Delta \rho / \rho}, \quad V_P = (V_{P2} + V_{P1}) / 2, \quad V_S = (V_{S2} + V_{S1}) / 2,$$

$$\rho = (\rho_2 + \rho_1) / 2, \quad \Delta V_P = (V_{P2} - V_{P1}), \quad \Delta V_S = (V_{S2} - V_{S1}), \quad \Delta \rho = (\rho_2 - \rho_1), \quad \Delta \sigma = (\sigma_2 - \sigma_1),$$

$$\sigma = (\sigma_2 + \sigma_1) / 2 \text{ and } \sigma_i = (0.5V_{Pi}^2 - V_{Si}^2) / (V_{Pi}^2 - V_{Si}^2), \quad i=1, 2.$$

66 V_p , V_s and ρ are the P-wave velocity, S-wave velocity and density respectively. The
 67 subscripts 1 and 2 represent the parameters of overlying and underlying layers, respectively.
 68 Equations 1-2 are valid for large negative A for a negative impedance contrast across an
 69 interface (Castagna et al., 1998; Shuey, 1985), and to about $A = 0.2$ for a positive impedance
 70 contrast (Shuey, 1985).

71 To interpret seismic data in terms of hydrate content, one needs to establish a relation
 72 between hydrate saturation in sediments and their velocity. Lee et al. (1996) developed a
 73 weighted equation (WE) model, in which hydrate velocities are weighted sum of Wyllie's
 74 time average equation (Wyllie et al., 1958) and Wood's equation (Wood, 1941). However,
 75 the WE model is not based on physical principles and requires empirical calibration. The
 76 velocities predicted from the cementation theory proposed by Dvorkin and Nur (1996) are
 77 much higher than those normally observed in nature (Ecker et al., 1998). Helgerud et al.
 78 (1999) introduced an EMT that considers the hydrate as part of the rock frame and
 79 successfully applied their approach to the P-wave velocity data from Site 995 of ODP Leg
 80 164 in the Blake Ridge area. Lee (2002) proposed a method based on Biot–Gassmann theory
 81 (BGTL) to relate the elastic properties of sediments to those of the matrix and the pore fluid.
 82 Several researchers (Chand et al., 2004; Dai et al., 2004) compared various models and found
 83 modest variations in the pattern of velocity with hydrate saturation if the same mineralogy
 84 and values of elastic constants are assumed. Here, we use the EMT and BGTL models to
 85 calculate the A and B parameters for various gas-hydrate and free gas saturations.

86 In the absence of direct sampling, we assume quartz and clay as mineralogical
 87 constituents for the sediment matrix in the Makran (Sain et al., 2000). For a seafloor porosity
 88 (ϕ_0) of 60% (Fowler et al., 1985) and a compaction constant λ of 1.17 (Minshull and White,
 89 1989), the porosity at the depth of the BSR (~510 m below the seafloor) is calculated as
 90 ~39% using the Athy's law ($\phi(z) = \phi_0 e^{(-z/\lambda)}$). Other parameters used in our calculation taken

from Lee (2002) are shown in Table 1. The gas density and bulk modulus of methane are calculated to be 0.2 g/cc and 67 MPa respectively at $\sim 21^\circ \text{C}$ and 26.34 MPa pressure (Ghosh and Sain, 2008) at BSR depth using Batzle and Wang relation (1992). We use a background P-wave velocity (in absence of gas or hydrate) of 2000 m/s obtained from the normal move out (NMO) interval velocity (described in the next section) for both rock physics models at the BSR depth. This velocity and porosity corresponds to a composition of 10% clay and 90% quartz using the EMT model. Assuming the same composition of clay and quartz, the value of consolidation constant 'm' (Lee, 2004) is adjusted to 2.13 in the BGTL model to obtain background velocity of 2000 m/s at the BSR depth. We used the Biot coefficient $\beta = -68.7421/(1 + e^{(\phi+0.40635)/0.09425}) + 0.98469$ for unconsolidated marine sediment (Lee, 2002, 2004) in the BGTL model. The P- and S-wave velocities decrease with increasing β at a given porosity. A reasonable value of β can be calculated by least-squares fitting of the observed data for a wide range of porosities (Lee, 2002).

We first compute the A and B attributes using equation 2 by varying the hydrate saturation from 0 to 80% at intervals of 5% for fixed gas saturations of 0, 1, 2, 5 and 10% below the BSR. Above $\sim 10\%$ gas saturation, the P-wave velocity remains almost unchanged. The resulting variations of A as a function of B for $10 \pm 5\%$ clay content at 39% porosity and $39 \pm 2\%$ porosity at 10% clay content are shown in Figure 1 and Figure 2 respectively to demonstrate the effects of clay content and porosity. The porosity is more sensitive than the mineralogical composition to the change in A and B values. In A-B space, a hydrate/brine BSR falls in the 2nd quadrant (negative A and positive B) and is treated as a reference model. The A and B values clearly deviate from this reference trend in presence of gas below the BSR.

116 FIELD DATA

117 We compute the AVO attributes from seismic data (Figure 3) that show a clear BSR
118 in the Makran accretionary prism (Minshull et al., 1992; Sain et al., 2000). The data have 24
119 traces in each CDP with a maximum offset of 2559 m and dominant frequency of ~35 Hz. A
120 representative super-gather of CDPs 4378 and 4379 is displayed in Figure 4. The BSR
121 reflection coefficients are calculated from the amplitudes of the seafloor reflection, the BSR
122 and the first multiple of the seafloor reflection using the approach of Warner (1990). The data
123 processing sequence included a band pass filter (4-8-50-60 Hz), a spherical divergence
124 correction ($1/(\text{time} \times \text{velocity}^2)$), minimum phase spiking deconvolution, an NMO correction
125 and trace equalization. The trace to trace or very short wavelength variations are assumed to
126 be noise and need to be eliminated. Trace equalization compensates the amplitude variations
127 due to abnormal shot strength and receiver coupling etc. But it should be applied very
128 cautiously within a long-gate time window (e.g. seafloor to first multiple of seafloor) at the
129 end of all processing steps (Ostrander, 1984; Yu, 1985). Hydrophone directivity is corrected
130 using the function of Sheriff and Geldert (1995). No source directivity correction is applied
131 as the source was of limited spatial extent (Sain et al., 2000). Root-mean-square amplitudes
132 were picked for a time window of 40-50 ms around the BSR, between CDPs 4340 to 4500,
133 where the BSR is strong and relatively flat. The time window was chosen to yield a reliable
134 result despite the change in BSR waveform at CDP 4450. Velocity analysis has been carried
135 out at a 10 CDP interval (225 m). We used a vertically smoothed (20 ms) and spatially
136 averaged NMO interval velocities (Figure 5) to calculate angles of incidence at the BSR
137 using ray theory. Offsets of up to 1610 m have incidence angles of less than 30° and were
138 used in the analysis; results differed little if instead a 25° threshold was used. Reflector
139 depths were obtained by converting the velocity-time function into a velocity-depth function
140 with Dix's formula. A straight line was fit to the averaged interval velocities (Figure 5) which

gives a reference velocity of 2 km/s at the BSR depth. This velocity was used to calibrate the rock physics models. The A and B values and their associated uncertainties (Figure 6a) were derived by weighted least-squares fitting for each supergather of two consecutive CDPs across the section. A running average of A and B over CDP location was superimposed on the theoretical nomogram (Figure 1-2). The running average removes noise from estimated A and B values and potentially reveals lateral changes in hydrate and gas content. Figure 6b illustrates the estimated hydrate and gas saturations with their uncertainties due to uncertainties in A, B and porosity.

DISCUSSION

The A and B attributes are a measure of the normal incidence reflection coefficient (P-wave impedance contrast) and change in Poisson's ratio (a function of both P- and S-wave velocities), respectively. Hence, the A-B crossplot coupled with the rock physics modeling is an important tool for estimation of gas hydrate and free gas respectively. The BGTL model shows almost linear trends with hydrate and gas saturations, whereas the EMT model is more nonlinear (Figure 1-2). At low hydrate saturations ($< 10\%$), the V_p/V_s ratio increases with hydrate saturation because the EMT model predicts little variation in V_s at low hydrate saturation. This complex behavior suggests that it will be difficult to estimate hydrate and gas saturations with the EMT approach when hydrate content is low. The EMT theory predicts higher S-wave velocity than the BGTL theory for any proportion of clay (or quartz) at a specified porosity (for porosity $> 30\%$) and vice versa (Lee 2002; Ojha and Sain, 2008). Thus the predicted background V_p/V_s ratio for the EMT model is always less than that for the BGTL model, and the EMT model always predicts higher hydrate saturation than that of the BGTL model. The BGTL model is more realistic in case of quantification from both V_p and V_s together. The uncertainties in hydrate saturation due to uncertainties in A are small for

166 both the EMT and BGTL models (Figure 6). The maximum uncertainty in hydrate saturation
167 due to uncertainties in B is $\pm 1.5\%$ and $\pm 1.0\%$ using, EMT and BGTL models respectively.
168 The corresponding uncertainties in gas saturation for both models are less than $\pm 0.12\%$. The
169 largest uncertainties come from the EMT-based estimate due to error in porosity. Uncertainty
170 in estimation of gas saturation may increase with errors in density and bulk modulus of gas
171 which are dependent on pressure and temperature.

172 The presence of gas-rich/gas-poor dipping strata, which produce tuning related highs
173 and lows on amplitudes, and velocity variation at few CDPs in the gas zone below the BSR
174 (Figure 5), complicate the analysis. By using a supergather of two consecutive CDPs and
175 running average of A and B values over a CDP profile, the AB crossplot method gives an
176 overall idea of hydrate and gas in the region. AVO modeling of a synthetic seismogram
177 (Figure 7) shows that the crossplot technique works well where the gas layer is greater than
178 10 m thick. As the thickness of the gas layer decreases, amplitude increases with offset due to
179 tuning between the top and bottom of the gas layer. In an area where the gas layer thickness
180 is unknown, the crossplot technique can be used to extrapolate interpretation away from a
181 small number of selected locations where more time-consuming waveform inversion or
182 modeling techniques are applied, so long as the gas layer is greater than 10 m thick. Tuning
183 effects can lead to localized peaks and troughs in A and B values. By using running means of
184 A and B values over several CDPs, we are able largely to remove these artifacts. The full
185 waveform inversion analysis of two CDP supergatherers from the same dataset shows that the
186 gas layer is indeed greater than 10 m thick (Sain et al., 2000). Using high-frequency (up to
187 650 Hz) deep-towed multichannel seismic data, Chapman et al. (2002) delineated a thin, 4-8
188 m negative velocity gradient layer at the BSR on the Vancouver Island margin. However, at
189 lower frequencies, as used here, such thin layers may not be resolved and the BSR commonly
190 appears as a single interface overlying a thick low-velocity zone (e.g., Hyndman and Spence,

1992; Andreassen et al., 1997; Muller et al., 2007; Ecker et al., 1998; Yuan et al., 1999; Carcione and Tinivella, 2000).

Gas hydrates may be distributed in concentrations that are both laterally and vertically variable, controlled by heterogeneities in lithology, permeability and methane transport (Ruppel and Kinoshita, 2000; Trehu et al., 2004; Riedel et al., 2006; Collett et al., 2008; Malinverno et al., 2008). An identifiable reflector is always needed to apply an AVO technique like the A-B crossplot, and in case of gas-hydrate exploration, the reflector is the BSR. The crossplot technique can be used as an effective tool on a regional scale to assess the gas hydrate and free gas saturation at the BSR.

CONCLUSIONS

The A-B crossplot (Figure 1-2) is a nomogram for various rock physics models from which we can (i) understand whether the BSR is underlain by gas and (ii) estimate directly hydrate and gas content. The A and B values derived from the BSR in the Makran accretionary prism deviate clearly from the reference trend and thus indicate the presence of gas below the BSR. The EMT model indicates a hydrate saturation of 14-29% and a gas saturation of 1-6%, while the BGTL model indicates 4-15% hydrate and 1-7.5% gas, respectively, which varies laterally (Figure 6b). The AVO attribute analysis is simple and easy to implement across a wide area to estimate hydrate and gas saturations at a BSR.

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338

Figure captions

Figure 1. Theoretical crossplots between intercept (A) and gradient (B) computed using varying saturation of gas hydrates with several fixed saturation of free gas based on (a) BGTL and (b) EMT models respectively, prepared for $10\pm 5\%$ clay content and 39% porosity. Running average of A and B values (plus signs) estimated for various CDP locations are superimposed.

Figure 2. Theoretical crossplots between intercept (A) and gradient (B) computed using varying saturation of gas hydrates with several fixed saturation of free gas based on (a) BGTL and (b) EMT models respectively, prepared for $39\pm 2\%$ porosity (ϕ) and 10% clay content. Running average of A and B values (plus signs) estimated for various CDP locations are superimposed.

Figure 3. Seismic stacked section along a south-north seismic line in the Makran accretionary prism. Inset shows the study area. The BSR shows crosscutting with dipping strata, opposite polarity with respect to the seafloor reflection and mimicking the seafloor topography.

Figure 4. A representative NMO corrected super gather of CDP 4378 and 4379 showing reflections from the seafloor, BSR and the first multiple of seafloor reflection. Amplitudes of the multiple and seafloor are used to calculate reflection coefficients of the BSR.

Figure 5. Thin lines show smooth NMO interval velocity functions at 10 CDP intervals between CDPs 4340 and 4500. Thick solid line shows spatially averaged interval velocity with depth. Dashed line is a straight line fitted to the averaged interval velocity (solid line).

Figure 6. (a) A and B values (with error bars) against the CDPs, estimated based on weighted least square fit and then running average for less than 30° angle (1610 m offset) and (b) hydrate and gas saturations against CDPs, estimated by EMT and BGTL models respectively

362 at $39\pm 2\%$ porosity and 10% clay content. Error bars are due to errors in A (small bars) and B
363 (large bars) at 39% porosity.

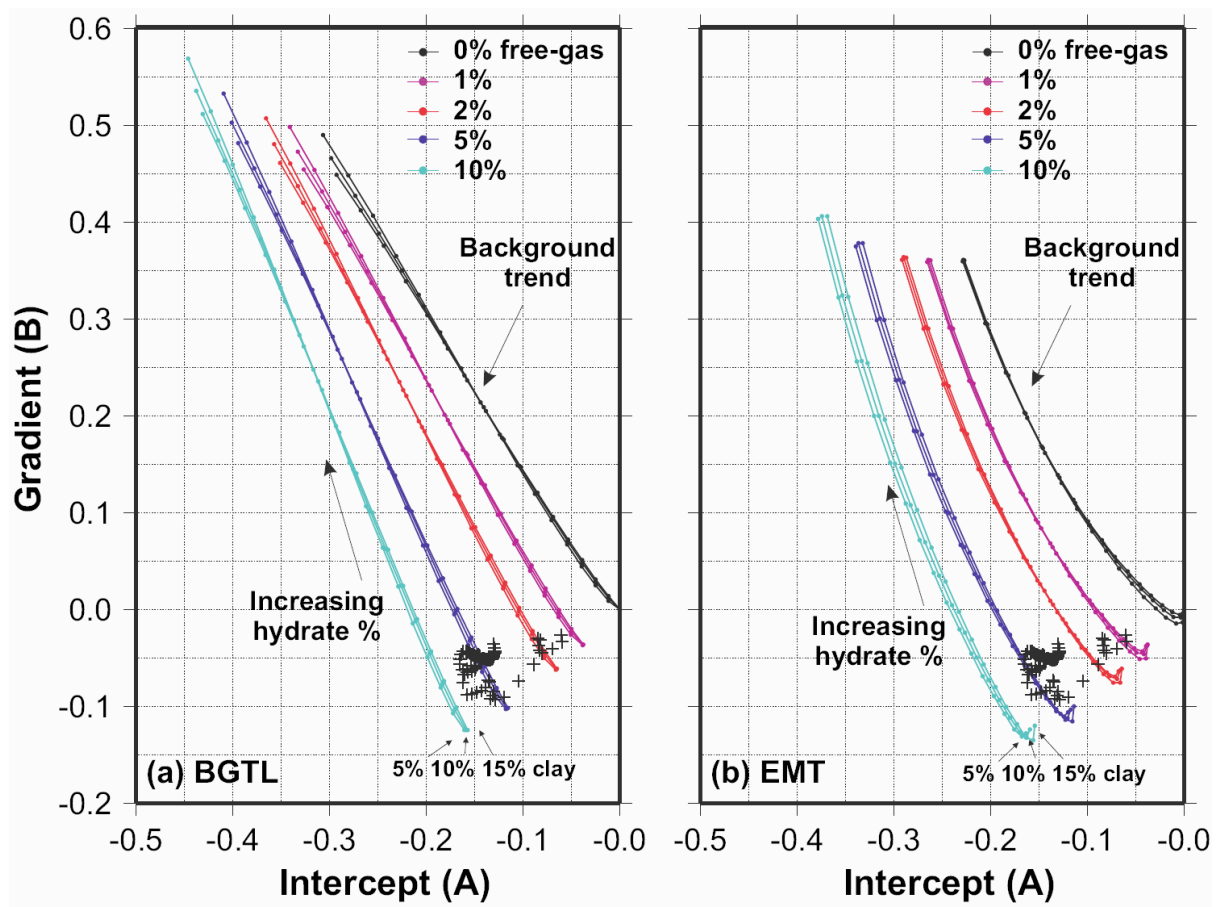
364 **Figure 7.** Modeled amplitudes of the BSR normalized with respect to the seafloor for several
365 thicknesses of hydrate and free gas layers (in sets). Velocity-depth functions are computed
366 using BGTL theory for 20% hydrate and 2% gas with Q-factor 250 and 150 (Singh and
367 Minshull, 1994) respectively. The synthetic seismograms are generated using the reflectivity
368 method (Fuchs and Mueller, 1971) with the frequency band of 2-10-80-94 Hz. Thick gas
369 layers result in negligible tuning effects (a and b). However, considerable effects are
370 observed when underlying gas layer is less than 10 m thick (c).

371

Table 1. Parameters used in theoretical calculations (Lee, 2002)

Component	K (GPa)	G (GPa)	ρ (g/cc)
Quartz	36	45	2.65
Clay	20.9	6.85	2.58
Hydrate	6.41	2.54	0.91
Methane*	0.067	-	0.20
Water	2.25	-	1.00
Critical Porosity 36%, No of grain/contact 9, N=1			
* Ghosh and Sain (2008)			

382



383

384 Figure 1

385

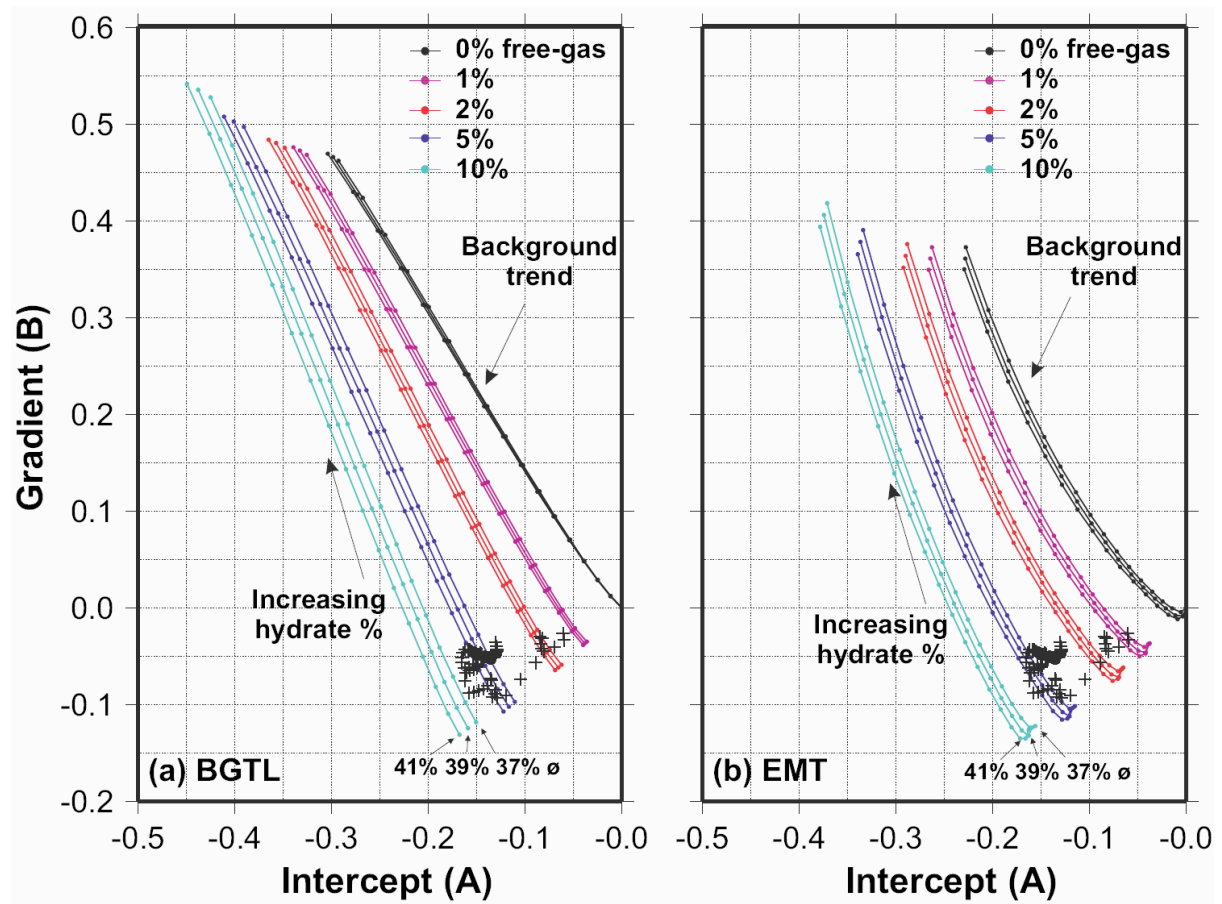


Figure 2

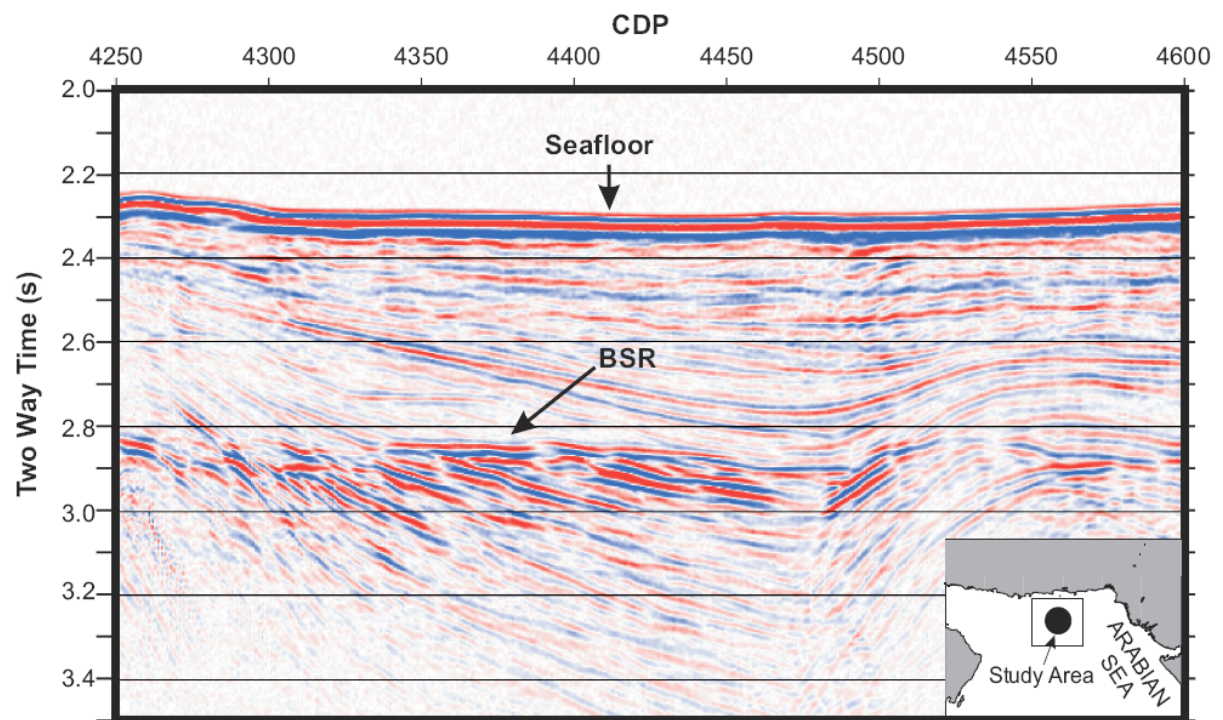


Figure 3

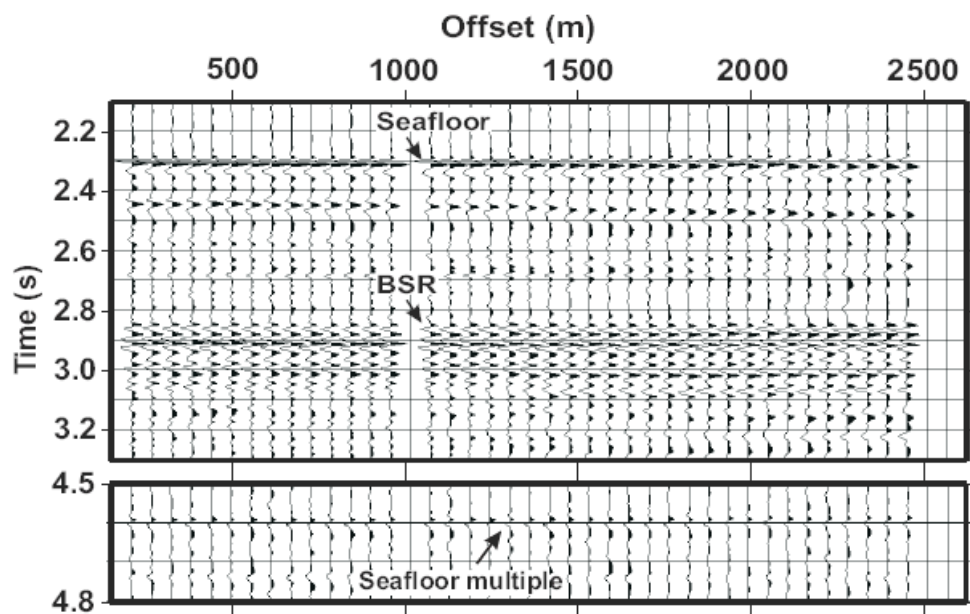


Figure 4

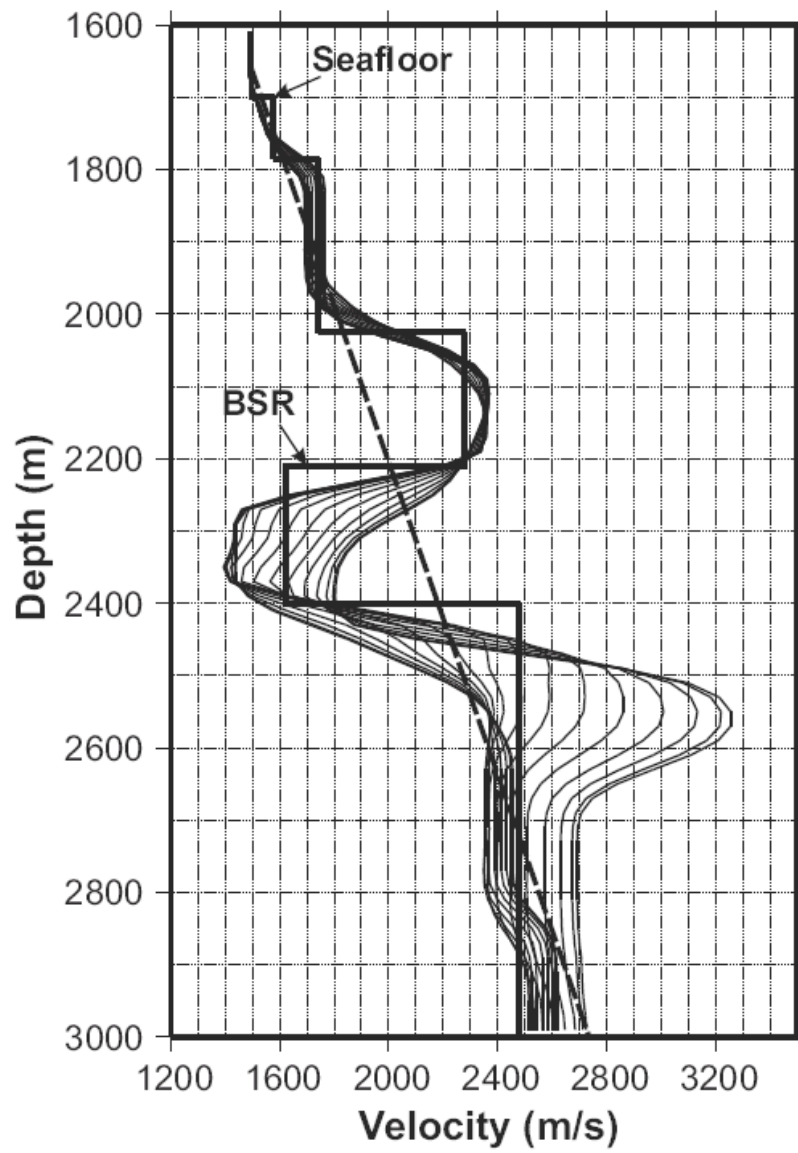


Figure 5

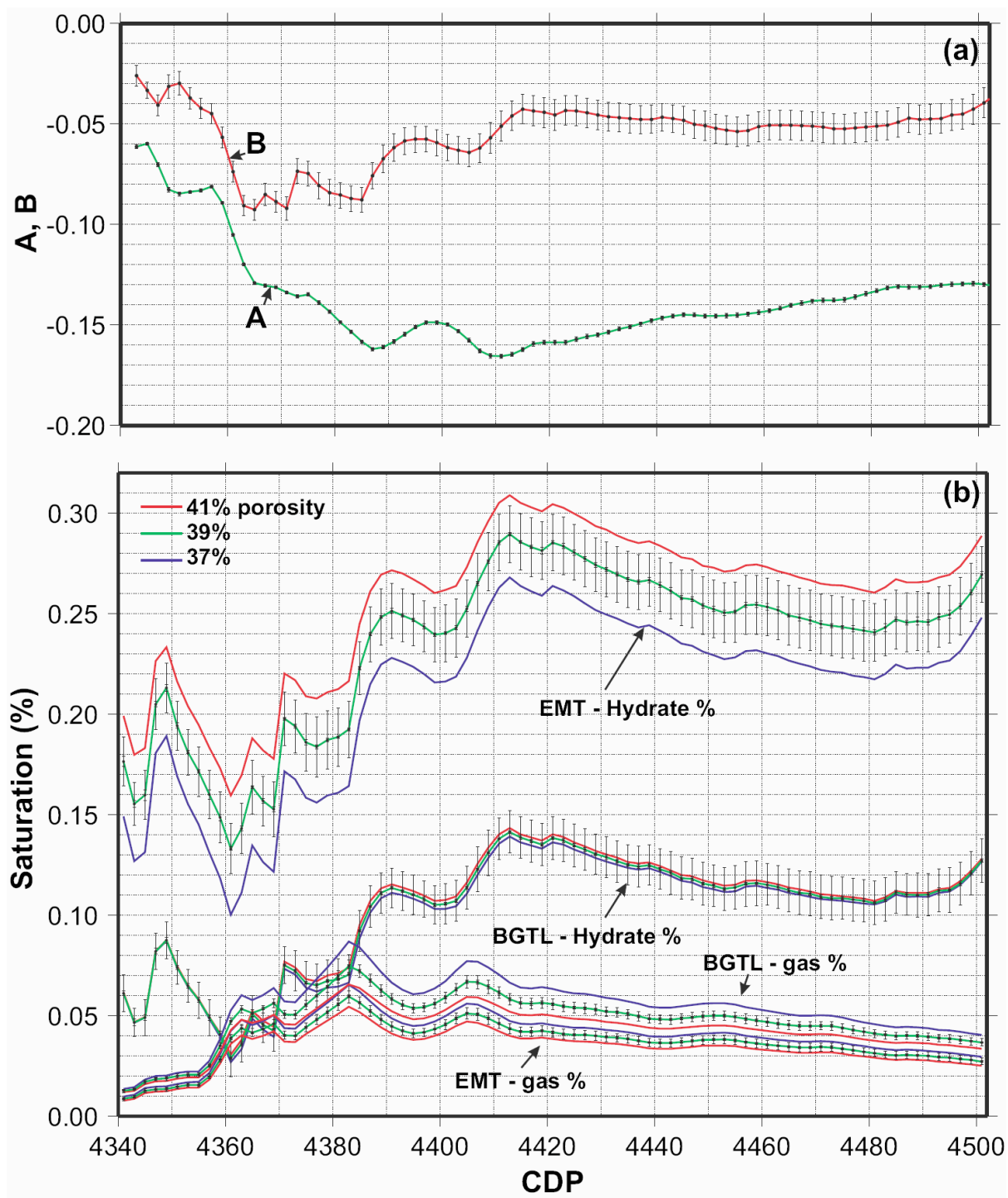


Figure 6

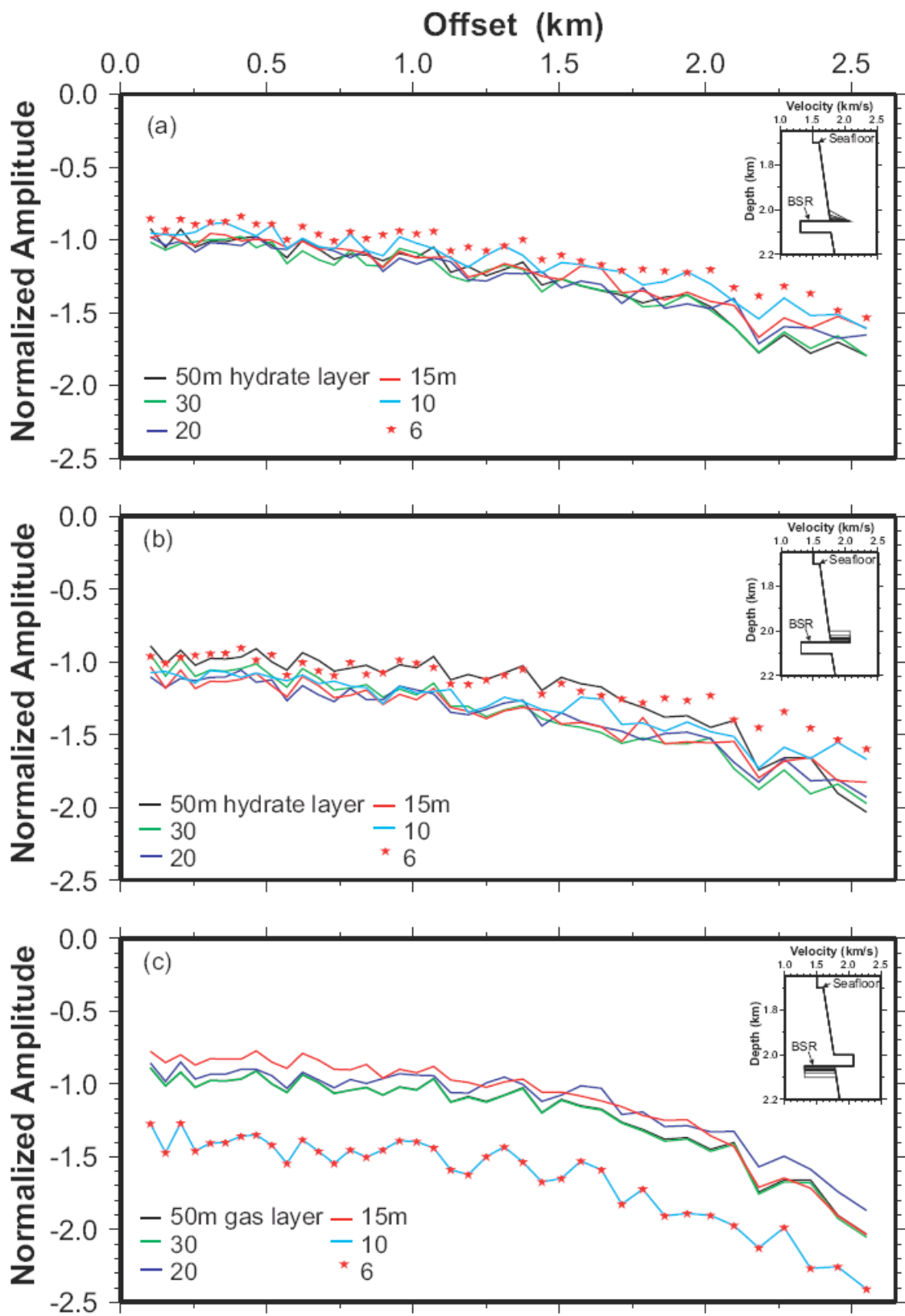


Figure 7