# Aerodynamics \&Flight Mechanics Research Group 

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## Preamble

This document examines the application of an unsteady aerodynamic model to a sinusoidal variation in pitch angle. It is based on the concept of the Wagner or Kussner lift variations. Both of these models produce approximations consisting of a combination of unity and two exponential decays. To simplify the introductory analysis, only one exponential decay term is retained.

## Nomenclature

| Variable | Definition |
| :---: | :--- |
| $U$ | Forward Speed |
| $\rho$ | Air Density |
| $\lambda$ | Exponential Decay Factor |
| $\Delta \alpha$ | Incremental Pitch Angle Change |
| $\phi$ | Wagner Function |
| $\tau$ | Integrating time variable |
| $\Delta L$ | Incremental Lift Force |
| $\Delta C_{L}$ | Incremental Lift Coefficient |
| $\Delta C_{\text {Las }}$ | Incremental Lift Coefficient (Quasi-Steady) |
| $\Delta C_{\text {Lus }}$ | Incremental Lift Coefficient (Unsteady) |
| $d$ | Lift Decrement |
| $c$ | Wing Chord |
| $b$ | Wing Semi-Chord |
| $\alpha, \beta$ | Dummy Parameters |

## Basic Analysis

The Wagner lift variation is given by:

$$
\begin{equation*}
\Delta L=2 \pi \rho U^{2} b \cdot \Delta \alpha \cdot \phi(s) \tag{1.}
\end{equation*}
$$

The Wagner function, $\phi$, provides the unsteady lift variation using the reduced frequency as the independent variable.

The reduced frequency is defined by:

$$
\begin{equation*}
s=\frac{U t}{b} \tag{2.}
\end{equation*}
$$

Here, the term $b$ is the semi chord and the reduced frequency represents the wing movement in terms of semi-chord - or alternatively, the increase in the streamwise extension of the wake.

Equation (1) can be converted to a lift coefficient thus:

$$
\begin{equation*}
\Delta C_{L}=\frac{\Delta L}{\frac{1}{2} \rho U^{2} c}=\frac{2 \pi \rho U^{2} b \alpha \phi(s)}{\frac{1}{2} \rho U^{2} c}=2 \pi \alpha \phi(s) \tag{3.}
\end{equation*}
$$

Whence the ratio of the Unsteady lift coefficient change to that of the Quasi Steady result is given by:

$$
\begin{equation*}
\frac{\Delta C_{L U S}}{\Delta C_{L Q S}}=\phi(s) \tag{4.}
\end{equation*}
$$

The Wagner function can be considered to consist of the quasi-steady response of unity and a decrement function, $d$ :

$$
\begin{equation*}
\phi(s)=1-d(s) \tag{5.}
\end{equation*}
$$



Figure 1

We use a simplified version of the conventional Wagner function, namely an single exponential decay term:

$$
\begin{equation*}
d(s)=A e^{-\lambda s} \tag{6.}
\end{equation*}
$$

The shape of this function is shown in Figure 1.
The response, $R(t)$, to a forcing function $f(t)$, with an indicial response $g(t)$ is given by the Duhammel Superposition Integral:

$$
\begin{equation*}
R(t)=f(0) \cdot g(t)+\int_{0}^{t} \frac{d f(\tau)}{d \tau} \cdot g(t-\tau) d \tau \tag{7.}
\end{equation*}
$$

In our case, the function $g$ is equal to $\phi$, hence using (5) \& (6), (7) becomes:


$$
\begin{gather*}
R=f(0)\left[1-d\left(\frac{U t}{b}\right)\right] \\
+\int_{0}^{t} f^{\prime}(\tau) \cdot\left[1-d\left(\frac{U(t-\tau)}{b}\right)\right] d \tau \tag{8.}
\end{gather*}
$$

Which can be simplified by integrating by parts:

$$
\begin{gather*}
R=f(t)-f(0) d\left(\frac{U t}{b}\right) \\
-\int_{0}^{t} f^{\prime}(\tau) \cdot d\left(\frac{U(t-\tau)}{b}\right) d \tau \tag{9.}
\end{gather*}
$$

In our case, the decrement function is given by:

$$
\begin{equation*}
d(t)=A e^{-\frac{\lambda U}{b} t} \tag{10.}
\end{equation*}
$$

Whence, (9) becomes:

$$
\begin{align*}
& R=f(t)-f(0) A e^{-\frac{\lambda U}{b} t} \\
& -\int_{0}^{t} f^{\prime}(\tau) \cdot A e^{-\frac{\lambda U}{b}(t-\tau)} d \tau \tag{11.}
\end{align*}
$$

We therefore have the lift coefficient ratio as:

$$
\begin{align*}
& \frac{C_{L}}{C_{L Q S}}=f(t)-f(0) A e^{-\frac{\lambda U}{b} t} \\
& -A e^{-\frac{\lambda U t}{b}} \int_{0}^{t} f^{\prime}(\tau) \cdot e^{\frac{\lambda U \tau}{b}} d \tau \tag{12.}
\end{align*}
$$

If the forcing function, $f$, satisfies:

$$
\begin{equation*}
f(0)=0 \tag{13.}
\end{equation*}
$$

Equation (12) simplifies to:

$$
\begin{equation*}
\frac{C_{L}}{C_{L Q S}}=f(t)-A e^{-\frac{\lambda U t}{b}} \int_{0}^{t} f^{\prime}(\tau) \cdot e^{\frac{\lambda U \tau}{b}} d \tau \tag{14.}
\end{equation*}
$$

## Particular Forcing Function Sinusoidal Ramp

The function for this is piecewise and defined thus:

$$
f(t)=\left\{\begin{array}{c}
\theta \cdot \sin ^{2}\left(\frac{\pi t}{2 T}\right) ; 0 \leq t \leq T  \tag{15.}\\
\theta ; t>T
\end{array}\right.
$$

This satisfies the condition:

$$
\begin{equation*}
f(0)=0 \tag{16.}
\end{equation*}
$$

The profile of this function is shown in Figure 2:


Figure 2
The derivative of this function with respect to $t$ is given by:

$$
f^{\prime}(t)=\left\{\begin{array}{c}
\frac{\pi \theta}{2 T} \cdot \sin \left(\frac{\pi t}{T}\right) ; 0 \leq t \leq T  \tag{17.}\\
0 ; t>T
\end{array}\right.
$$

If we define two parameters thus:

$$
\begin{align*}
\alpha & =\frac{U \lambda}{b}  \tag{18.}\\
\beta & =\frac{\pi}{T}
\end{align*}
$$

Then a typical integral is:

$$
\begin{align*}
& \int_{0}^{t} e^{\alpha \tau} \cdot \sin \beta \tau \cdot d \tau  \tag{19.}\\
&=\frac{\beta+e^{\alpha t}(\alpha \sin \beta t-\beta \cos \beta t)}{\alpha^{2}+\beta^{2}}
\end{align*}
$$

In the case of the response after the forcing function has achieved its final value, the integral of (19) is the special case of:

$$
\begin{aligned}
& \int_{0}^{T} e^{\alpha \tau} \cdot \sin \beta \tau \cdot d \tau \\
& \quad=\frac{\beta+e^{\alpha T}(\alpha \sin \beta T-\beta \cos \beta T)}{\alpha^{2}+\beta^{2}}
\end{aligned}
$$

Noting that:

$$
\begin{gather*}
\beta T=\pi \\
\sin \beta T=0  \tag{21.}\\
\cos \beta T=-1
\end{gather*}
$$

Equation (20) simplifies to:

$$
\begin{equation*}
\int_{0}^{T} e^{\alpha \tau} \cdot \sin \beta \tau \cdot d \tau=\frac{\beta\left(1+e^{\alpha T}\right)}{\alpha^{2}+\beta^{2}} \tag{22.}
\end{equation*}
$$

From the above results, the final equation for the lift coefficient ratio is given by:

$$
\frac{C_{L}}{C_{L Q S}}=\left\{\begin{array}{c}
\frac{\pi \theta}{2 T} \cdot \sin \left(\frac{\pi t}{T}\right)-\frac{A \beta \theta}{2}\left[\frac{\beta e^{-\alpha t}+\alpha \sin \beta t-\beta \cos \beta t}{\alpha^{2}+\beta^{2}}\right] ; 0 \leq t \leq T  \tag{23.}\\
\frac{\pi \theta}{2 T} \cdot \sin \left(\frac{\pi t}{T}\right)-\frac{A \beta^{2} \theta}{2}\left[\frac{e^{-\alpha t}\left(1+e^{\alpha T}\right)}{\alpha^{2}+\beta^{2}}\right] ; t>T
\end{array}\right.
$$

## Results

The analysis was used to evaluate the lift variation for a sinusoidal-squared input.

The incidence variation is shown in Figure 3:


Figure 3 - Pitch Angle Variation
The Kussner function is shown in Figure 4:

## Southâmpoon



Figure 4 - Kussner Function

The resulting lift variation ratio is shown in Figure 5:


Figure 5 - Unsteady and Steady Lift Variation

## Matlab File

```
%
% Unsteady Aerodynamics -Exponential Lift Decrement
%
% SJN 9/3/08
%
clear
colordef white
lambda=.1;
A=1;
U=100;
C=1;
dclda=5.8;
thetadeg=10;
T=.5;
tmax=1;
nt=101;
b=c/2;
thetamax=thetadeg*pi/180;
alf=U* lambda/b;
bet=pi/T;
den=alf^2+bet^2;
t=linspace(0,tmax,nt);
clqs=zeros(1,nt);
clus=clqs;
thetdeg=clqs;
for i=1:nt
    if t(i) < T
        f=thetamax*sin(bet*t(i)/2)^2;
        num=(bet*thetamax/2)*(alf*sin(bet*t(i)) -bet*cos(bet*t(i)) +bet*exp(-
alf*t(i)));
    else
        f=thetamax;
        num=(bet^2*thetamax/2)*(exp(-alf*(t(i)-T))+exp(-alf*t(i)));
    end
    clqs (i)=dclda*f;
    clus(i)=dclda*(f-A*num/den);
    thetdeg(i)=(180/pi)*f;
end
clf
plot(t,thetdeg,'b','LineWidth', 2);
grid on
title(['Pitch Angle Variation T = ',num2str(T),' 0_M_A_X =
',num2str(thetadeg),'\circ']);
xlabel('Time');
ylabel('0');
figure
clf
plot(t,clqs,'b','LineWidth', 2);
hold on
```

```
plot(t,clus,'r','LineWidth',2);
grid on
title(['Lift Coefficient Ratio A = ',num2str(A),' \lambda =
',num2str(lambda),' T = ',num2str(T),' 0_M_A_X =
',num2str(thetadeg),'\circ']);
legend('Quasi Steady','Unsteady');
xlabel('Time');
ylabel('\DeltaC_L_U_S / \DeltaC_L_Q_S');
figure
clf
s=linspace (0,50,201);
phi=1-A* exp(-lambda*s);
plot(s,phi,'b','LineWidth',2);
grid on
title(['Indicial Response A = ',num2str(A),' \lambda =
',num2str(lambda)]);
xlabel('s');
ylabel('\phi');
```

