

Aerodynamics & Flight Mechanics Research Group

Coleman Transformation for N Rotor Blades

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SCHOOL OF ENGINEERING SCIENCES

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Definition of the Transformation

The rotor consists of N blades, each of which is denoted by the suffix, k . Each individual blade then possesses a lag angle of ζ_k .

The Coleman transformation converts the N lag angles (ζ_k) to N variables, θ_j , via the following:

$$\theta_j = \sum_{k=0}^{N-1} \zeta_k e^{j \cdot \frac{2\pi i}{N} \cdot k} \quad (1)$$

For brevity, define:

$$\alpha = e^{\frac{2\pi i}{N}} \quad (2)$$

i.e. α is the primary N th root of unity.

Thus;

$$\theta_j = \sum_{k=0}^{N-1} \zeta_k \alpha^{j \cdot k} \quad (3)$$

for $j=1, 2, \dots, N-1$.



Therefore:

$$\begin{aligned}
 \theta_0 &= \zeta_0 + \zeta_1 + \zeta_2 + \dots + \zeta_{N-1} \\
 \theta_1 &= \zeta_0 + \zeta_1\alpha + \zeta_2\alpha^2 + \dots + \zeta_{N-1}\alpha^{N-1} \\
 \theta_2 &= \zeta_0 + \zeta_1\alpha^2 + \zeta_2\alpha^4 + \dots + \zeta_{N-1}\alpha^{2(N-1)} \\
 &\quad \Downarrow \\
 \theta_{N-1} &= \zeta_0 + \zeta_1\alpha^{N-1} + \zeta_2\alpha^{2(N-1)} + \dots + \zeta_{N-1}\alpha^{(N-1)(N-1)}
 \end{aligned} \tag{4}$$

The α terms represent the azimuthal spacing between the blades and are thus referenced to a rotating set of axes.

To cater for a non-rotating set of axes, the above equations are multiplied by:

$$e^{i\omega t} \tag{5}$$

also, denoting the blade azimuth via:

$$\psi_k = \omega t + \frac{2\pi}{N} \cdot k \tag{6}$$

from which:

$$e^{i\omega t} \cdot \alpha^k = e^{i\psi_k} \tag{7}$$



Since, α is an Nth root of unity we also have:

$$\alpha^{N-j} = \frac{\alpha^N}{\alpha^j} = \frac{1}{|\alpha|^{2j}} \cdot (\overline{\alpha})^j = \overline{\alpha^j} \quad (8)$$

Since:

$$|\alpha| = 1 \quad (9)$$

Rewriting the equations and premultiplying by powers of $e^{i\omega t}$ we find:

$$\begin{aligned} \theta_0 &= \zeta_0 + \zeta_1 + \zeta_2 + \dots + \zeta_{N-1} \\ e^{i\omega t} \cdot \theta_1 &= e^{i\omega t} \cdot \zeta_0 + \zeta_1 e^{i\omega t} \cdot \alpha + \zeta_2 e^{i\omega t} \cdot \alpha^2 + \dots + \zeta_{N-1} e^{i\omega t} \cdot \alpha^{N-1} \\ e^{2i\omega t} \cdot \theta_2 &= e^{2i\omega t} \cdot \zeta_0 + \zeta_1 e^{2i\omega t} \cdot \alpha^2 + \zeta_2 e^{2i\omega t} \cdot \alpha^4 + \dots + \zeta_{N-1} e^{2i\omega t} \cdot \alpha^{2(N-1)} \\ &\Downarrow \\ e^{-2i\omega t} \cdot \theta_{N-2} &= e^{-2i\omega t} \zeta_0 + \zeta_1 e^{-2i\omega t} \cdot \alpha^{N-2} + \zeta_2 e^{-2i\omega t} \cdot \alpha^{2(N-2)} + \dots + \zeta_{N-1} e^{-2i\omega t} \cdot \alpha^{(N-2)(N-1)} \\ e^{-i\omega t} \cdot \theta_{N-1} &= e^{-i\omega t} \cdot \zeta_0 + \zeta_1 e^{-i\omega t} \cdot \alpha^{N-1} + \zeta_2 e^{-i\omega t} \cdot \alpha^{2(N-1)} + \dots + \zeta_{N-1} e^{-i\omega t} \cdot \alpha^{(N-1)(N-1)} \end{aligned} \quad (10)$$

Since α is an Nth root of unity the last two equations can be rewritten:

$$\begin{aligned} e^{-2i\omega t} \cdot \theta_{N-2} &= e^{-2i\omega t} \zeta_0 + \zeta_1 e^{-2i\omega t} \cdot \alpha^{-2} + \zeta_2 e^{-2i\omega t} \cdot \alpha^{-4} + \dots + \zeta_{N-1} e^{-2i\omega t} \cdot \alpha^{-2(N-1)} \\ e^{-i\omega t} \cdot \theta_{N-1} &= e^{-i\omega t} \cdot \zeta_0 + \zeta_1 e^{-i\omega t} \cdot \alpha^{-1} + \zeta_2 e^{-i\omega t} \cdot \alpha^{-2} + \dots + \zeta_{N-1} e^{-i\omega t} \cdot \alpha^{-(N-1)} \end{aligned} \quad (11)$$



Construction of the Coleman Variables

The first equation

$$\theta_0 = \zeta_0 + \zeta_1 + \zeta_2 + \dots + \zeta_{N-1} \quad (12)$$

is left untouched.

The equations are then paired:

$$\begin{aligned} e^{i\omega t} \theta_1 &\Leftrightarrow e^{-i\omega t} \theta_{N-1} \\ e^{2i\omega t} \theta_2 &\Leftrightarrow e^{-2i\omega t} \theta_{N-2} \\ e^{3i\omega t} \theta_3 &\Leftrightarrow e^{-3i\omega t} \theta_{N-3} \end{aligned} \quad (13)$$

these are then added and subtracted to give real expressions in cosine and sine.

For instance, let us examine the first pairing:

$$\frac{(e^{i\omega t} \cdot \theta_1 + e^{-i\omega t} \cdot \theta_{N-1})}{2} = \frac{1}{2} \left(\begin{array}{l} e^{i\omega t} \cdot \zeta_0 + e^{-i\omega t} \cdot \zeta_0 \\ + \zeta_1 e^{i\omega t} \cdot \alpha + \zeta_1 e^{-i\omega t} \cdot \alpha^{N-1} \\ + \zeta_2 e^{i\omega t} \cdot \alpha^2 + \zeta_2 e^{-i\omega t} \cdot \alpha^{2(N-1)} \dots \\ + \zeta_{N-1} e^{i\omega t} \cdot \alpha^{N-1} + \zeta_{N-1} e^{-i\omega t} \cdot \alpha^{(N-1)(N-1)} \end{array} \right) \quad (14)$$

and since:

$$\begin{aligned} \alpha^N &= 1 \\ \alpha^{N-1} &= \alpha^{-1} \\ \alpha^{(N-1)(N-1)} &= \alpha^{-(N-1)} \end{aligned} \quad (15)$$



we find:

$$\frac{(e^{i\omega t} \cdot \theta_1 + e^{-i\omega t} \cdot \theta_{N-1})}{2} = \frac{1}{2} \begin{pmatrix} e^{i\omega t} \cdot \zeta_0 + e^{-i\omega t} \cdot \zeta_0 \\ + \zeta_1 e^{i\omega t} \cdot \alpha + \zeta_1 e^{-i\omega t} \cdot \alpha^{-1} \\ + \zeta_2 e^{i\omega t} \cdot \alpha^2 + \zeta_2 e^{-i\omega t} \cdot \alpha^{-2} \dots \\ + \zeta_{N-1} e^{i\omega t} \cdot \alpha^{N-1} + \zeta_{N-1} e^{-i\omega t} \cdot \alpha^{-(N-1)} \end{pmatrix} \quad (16)$$

Recalling (7), we have finally:

$$\frac{(e^{i\omega t} \cdot \theta_1 + e^{-i\omega t} \cdot \theta_{N-1})}{2} = \begin{pmatrix} \zeta_0 \cdot \cos \psi_0 \\ + \zeta_1 \cdot \cos \psi_1 \\ + \zeta_2 \cdot \cos \psi_2 \\ \dots \\ + \zeta_{N-1} \cdot \cos \psi_{N-1} \end{pmatrix} \quad (17)$$

and by a similar process:

$$\frac{(e^{i\omega t} \cdot \theta_1 - e^{-i\omega t} \cdot \theta_{N-1})}{2} = \begin{pmatrix} \zeta_0 \cdot \sin \psi_0 \\ + \zeta_1 \cdot \sin \psi_1 \\ + \zeta_2 \cdot \sin \psi_2 \\ \dots \\ + \zeta_{N-1} \cdot \sin \psi_{N-1} \end{pmatrix} \quad (18)$$

If the rotor has an odd number of blades (N) this process will complete the transformation. If the number of blades (N) is even then there will be one equation (θ_M) left. In this case the multiplication by $e^{Mi\omega t}$ does not take place and we have:

$$\theta_M = \zeta_0 + \zeta_1 \alpha^M + \zeta_2 \alpha^{2M} + \dots + \zeta_{N-1} \alpha^{M(N-1)} \quad (19)$$



here we have:

$$M = \frac{N}{2}$$

as can be seen the multiplying factors for the N lag angles, ζ_k , are powers of α^M .

Since:

$$\alpha = e^{\frac{2\pi i}{N}} \quad (20)$$

we find:

$$\begin{aligned} \alpha^M &= \left(e^{\frac{2\pi i}{N}} \right)^M \\ &= e^{\frac{2\pi Mi}{N}} \quad (21) \\ &= e^{\pi i} \\ &= -1 \end{aligned}$$

In other words the coefficients of the ζ_k terms are powers of -1, in other words of alternating sign.

Reference

R.P. Coleman, A.M. Feingold, Theory of self-excited mechanical oscillations of helicopter rotors with hinged blades, NACA Report TN 1351, 1958



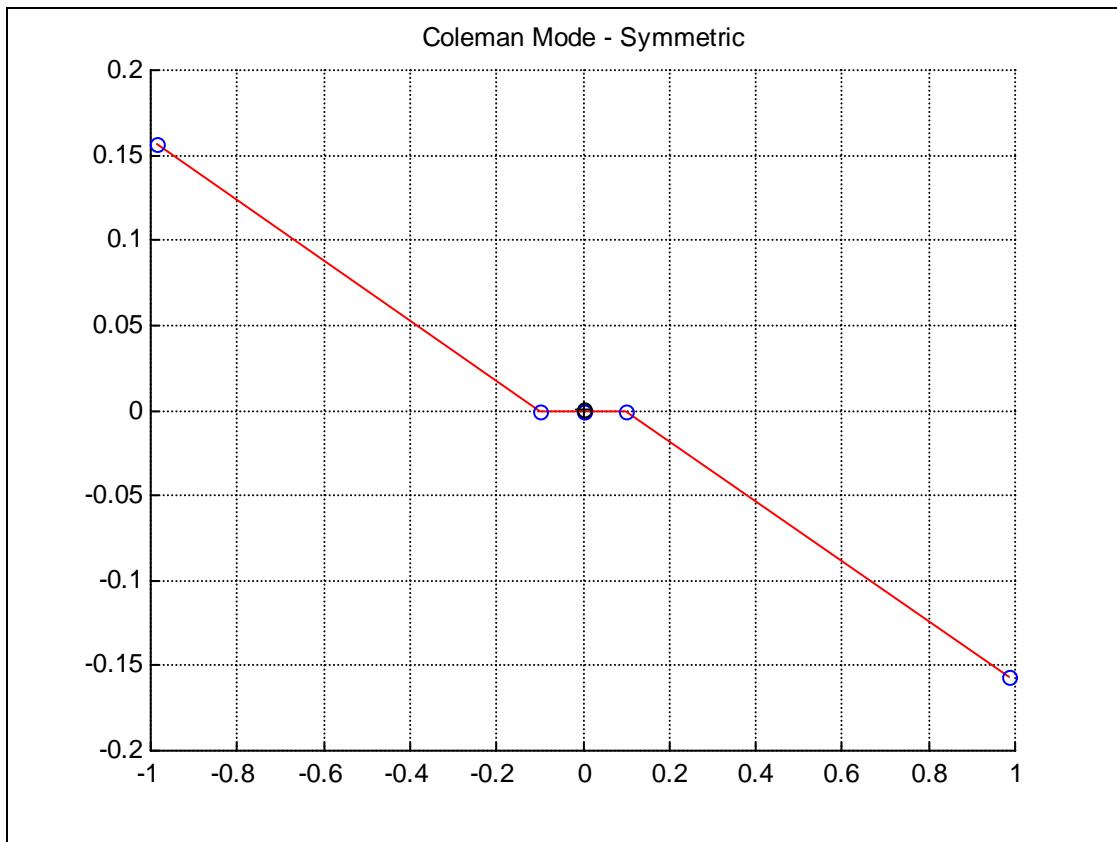
Examples

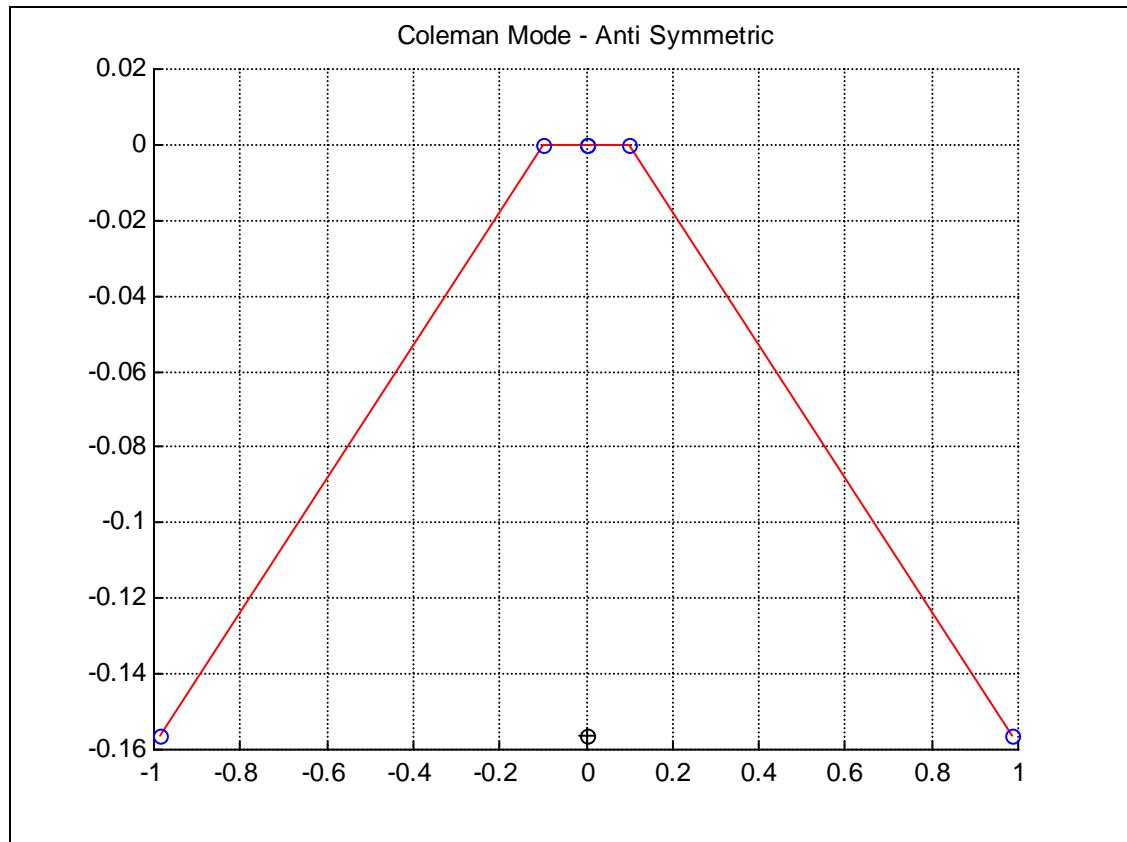
We therefore have the following examples of the transformation:

N=2

$$\Theta_0 = \zeta_0 + \zeta_1$$

$$\Theta_1 = \zeta_0 - \zeta_1$$





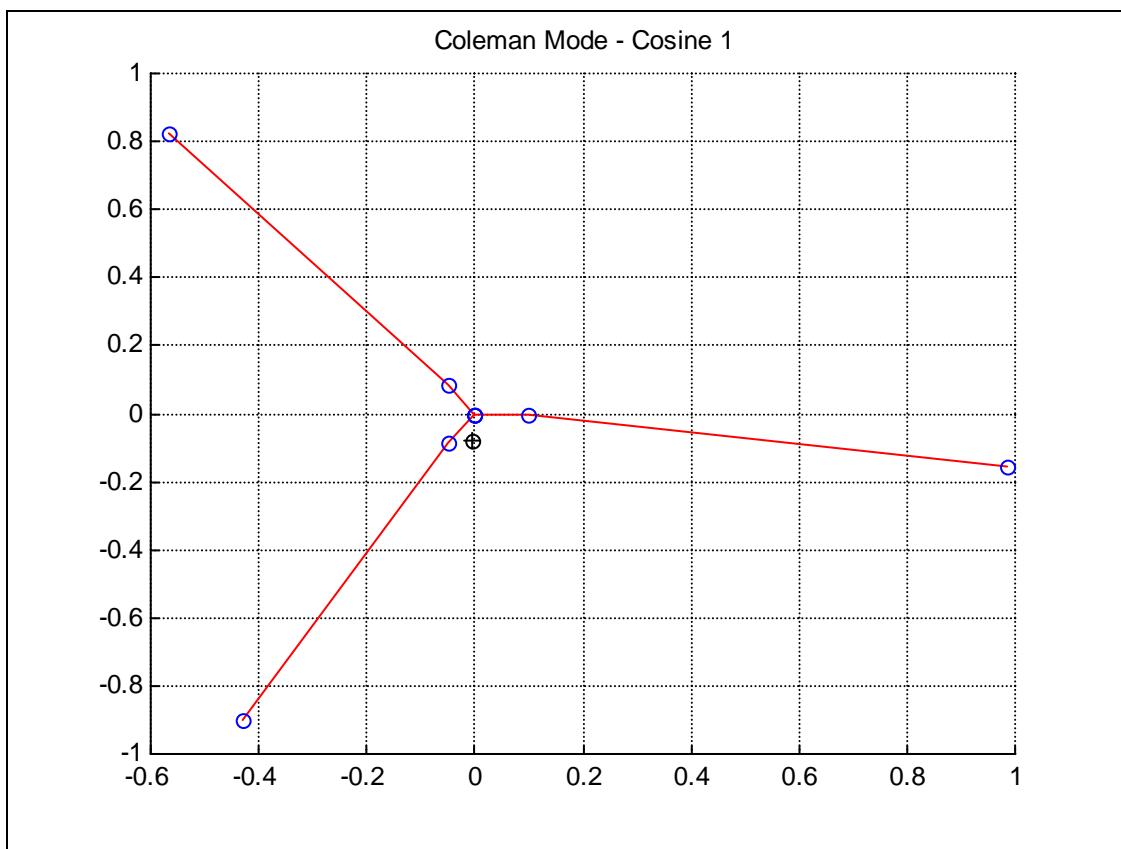
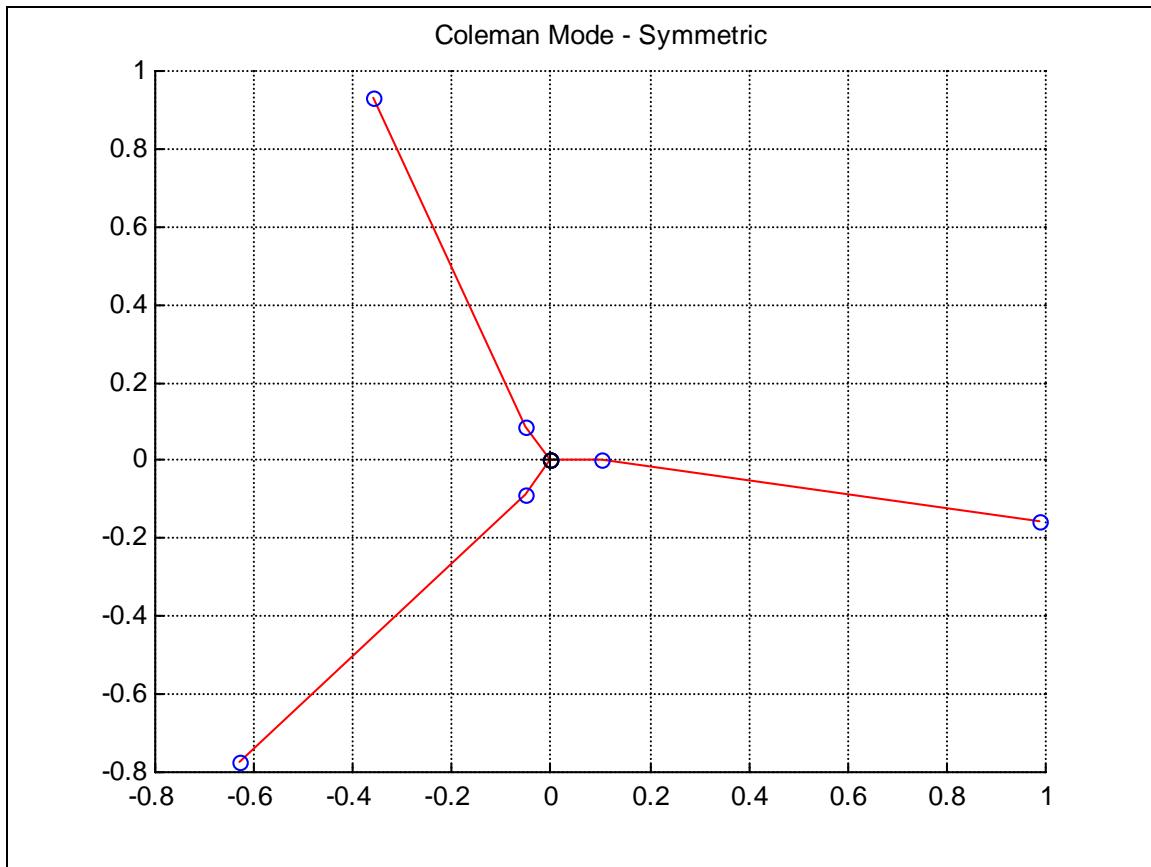
N=3

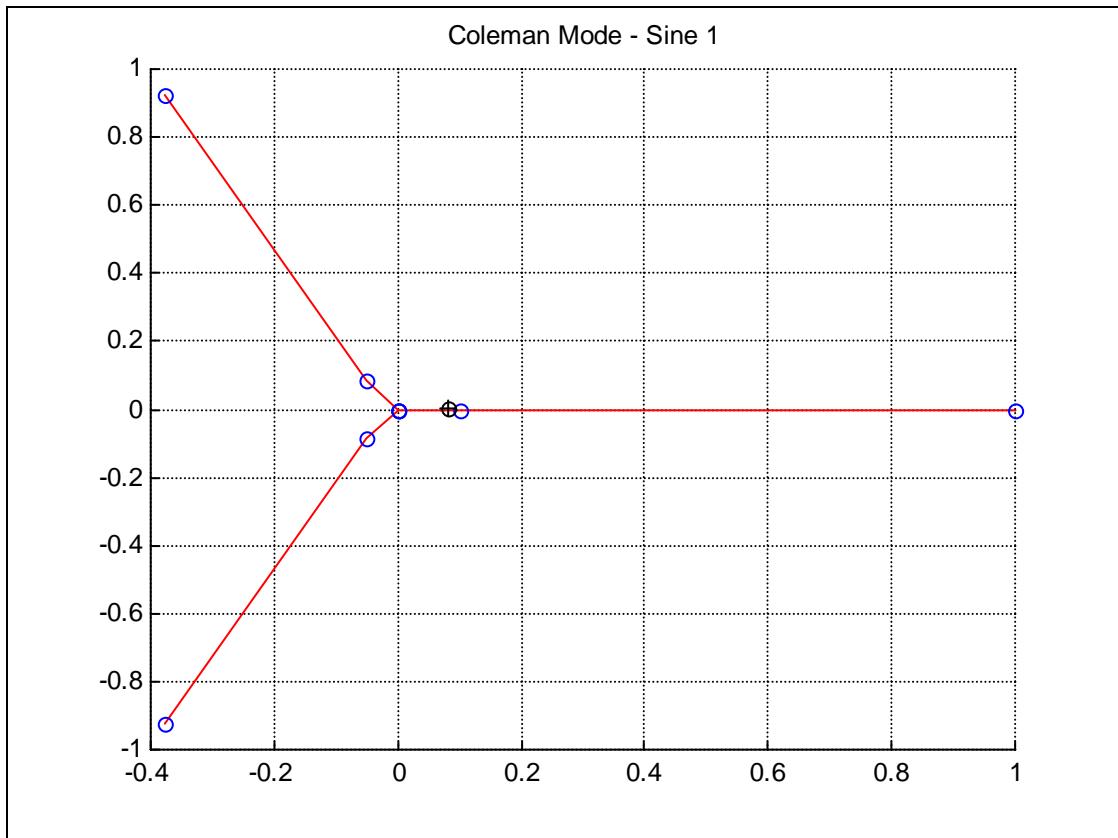
$$\Theta_0 = \zeta_0 + \zeta_1 + \zeta_2$$

$$\Theta_1 = \zeta_0 \cos \psi_0 + \zeta_1 \cos \psi_1 + \zeta_2 \cos \psi_2$$

$$\Theta_2 = \zeta_0 \sin \psi_0 + \zeta_1 \sin \psi_1 + \zeta_2 \sin \psi_2$$







N=4

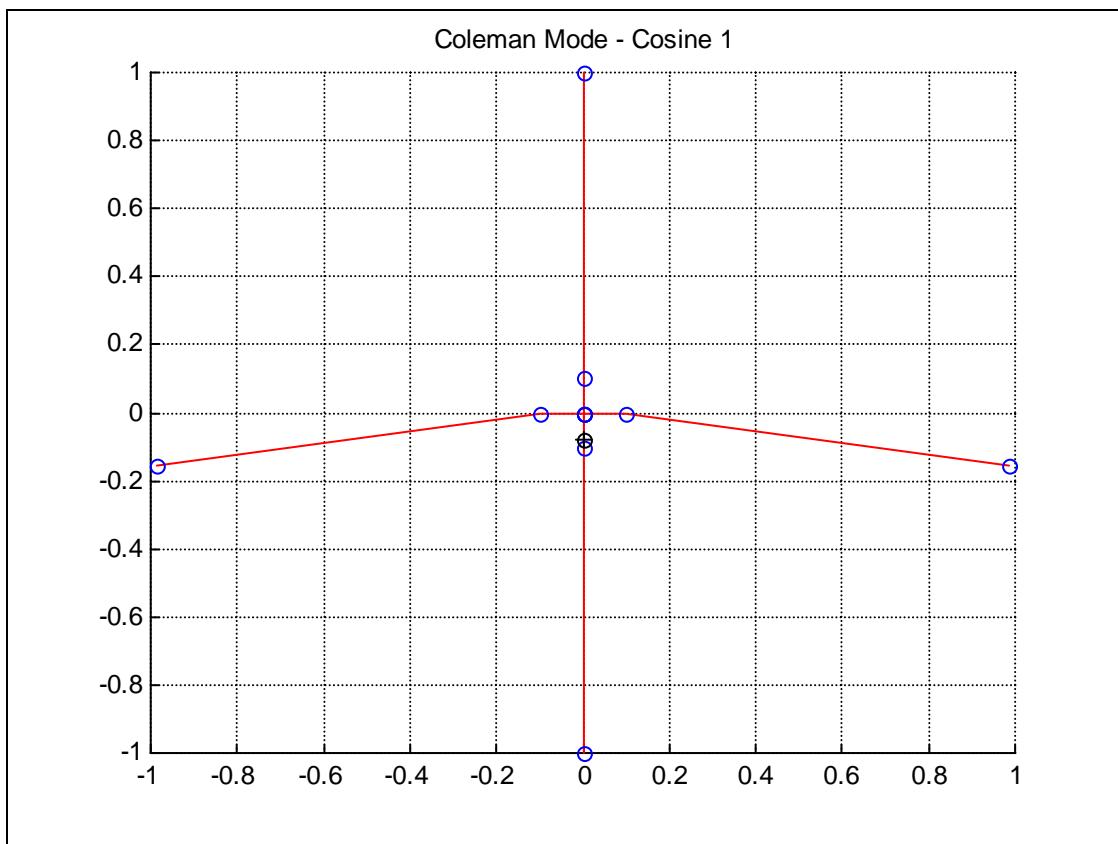
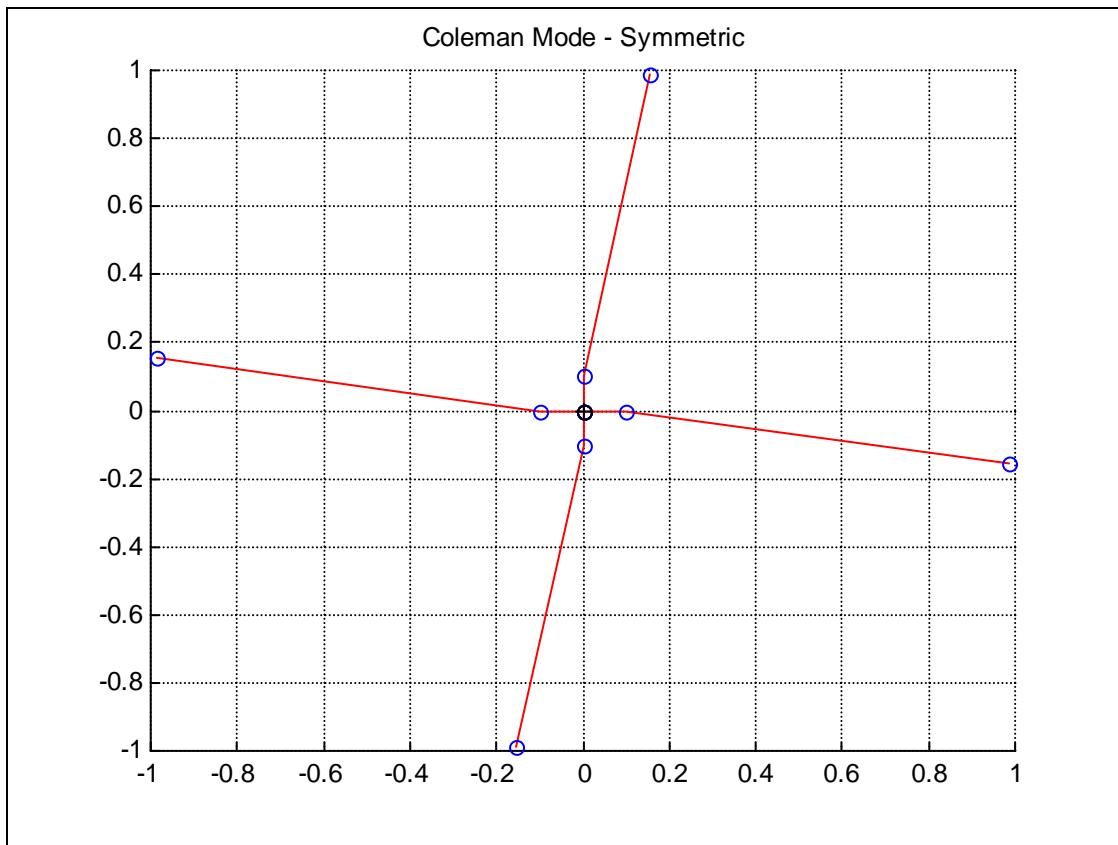
$$\Theta_0 = \zeta_0 + \zeta_1 + \zeta_2 + \zeta_3$$

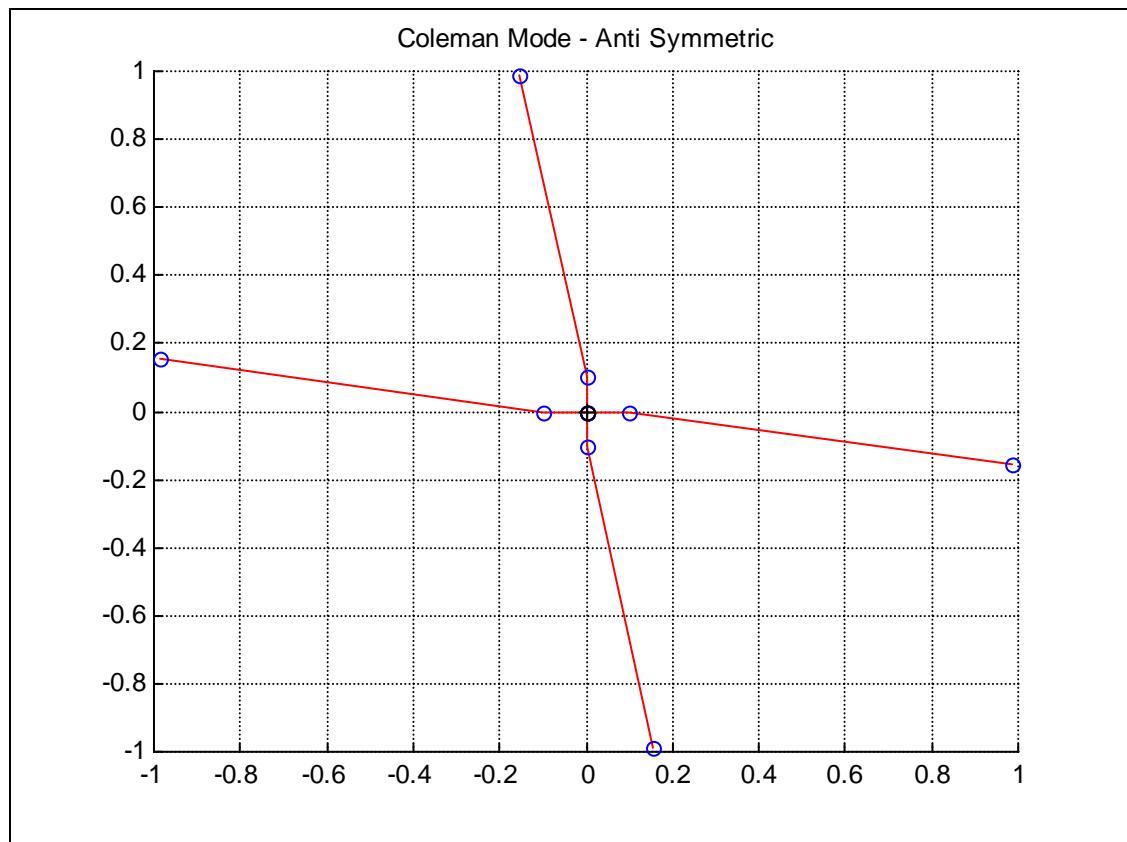
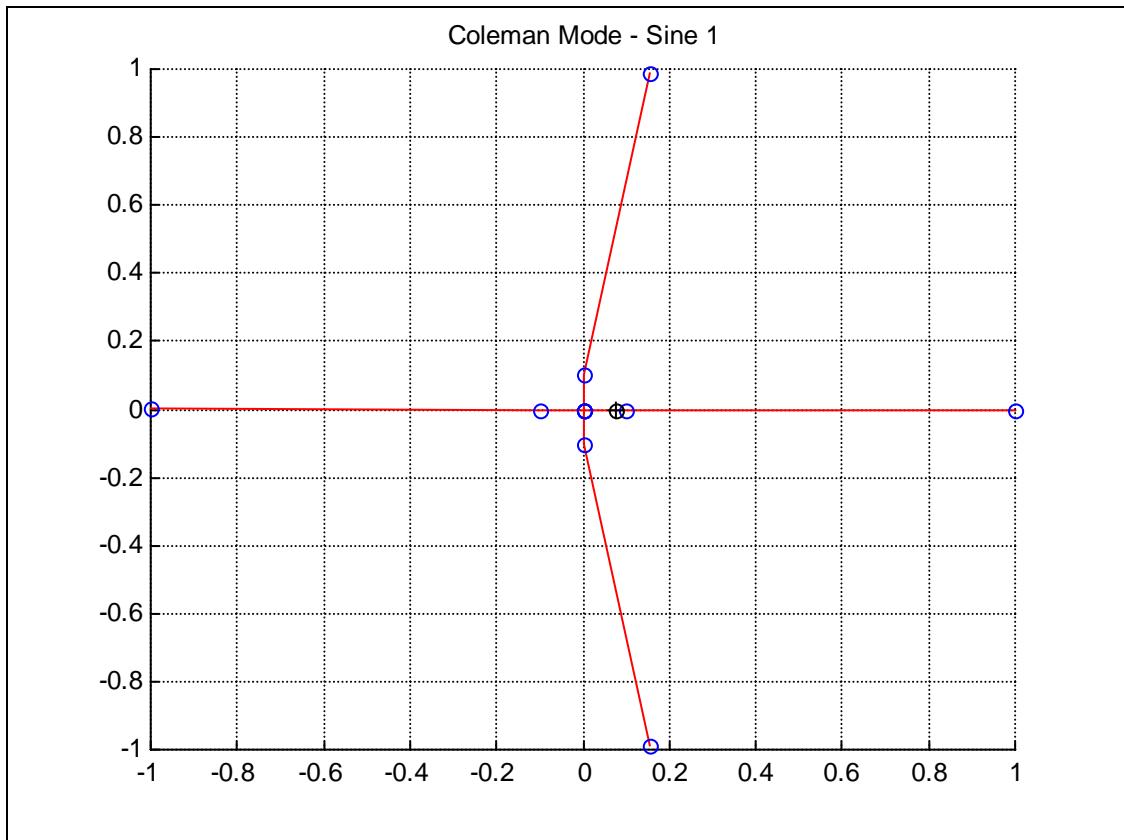
$$\Theta_1 = \sum_k \zeta_k \cos \psi_k$$

$$\Theta_2 = \sum_k \zeta_k \sin \psi_k$$

$$\Theta_3 = \zeta_0 - \zeta_1 + \zeta_2 - \zeta_3$$







N=5

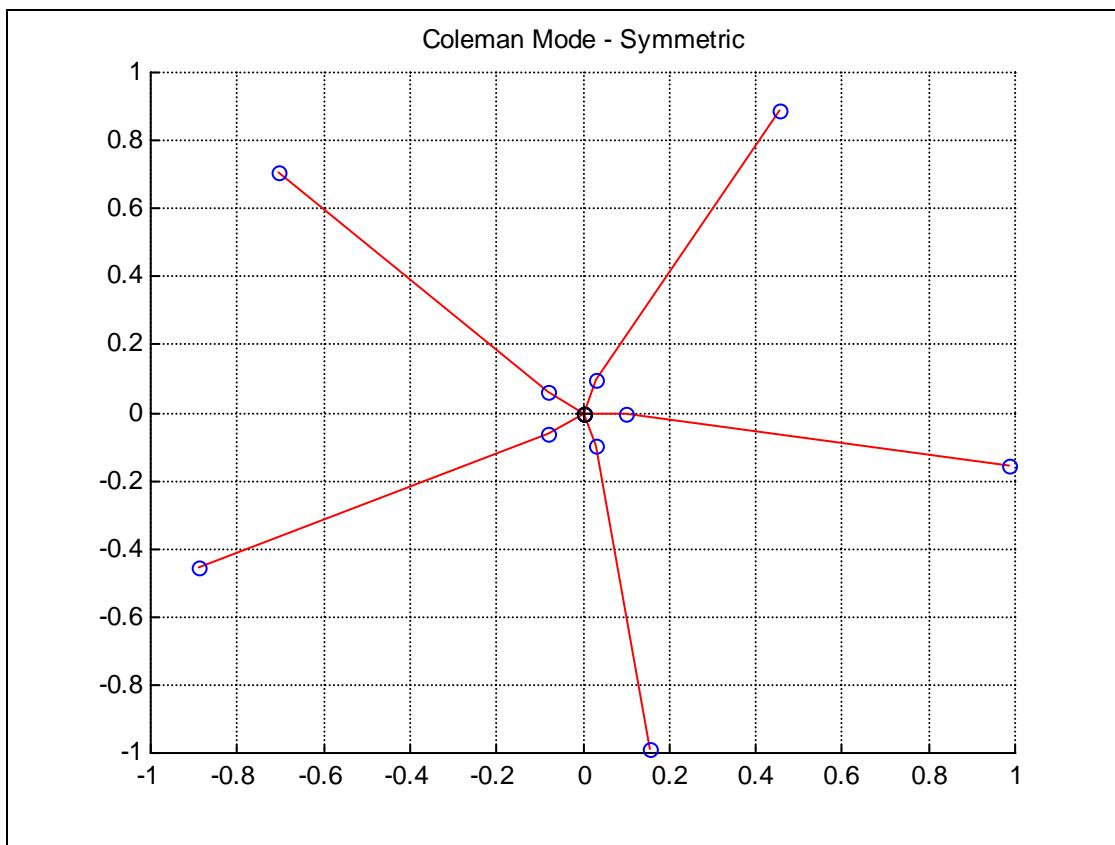
$$\Theta_0 = \zeta_0 + \zeta_1 + \zeta_2 + \zeta_3 + \zeta_4$$

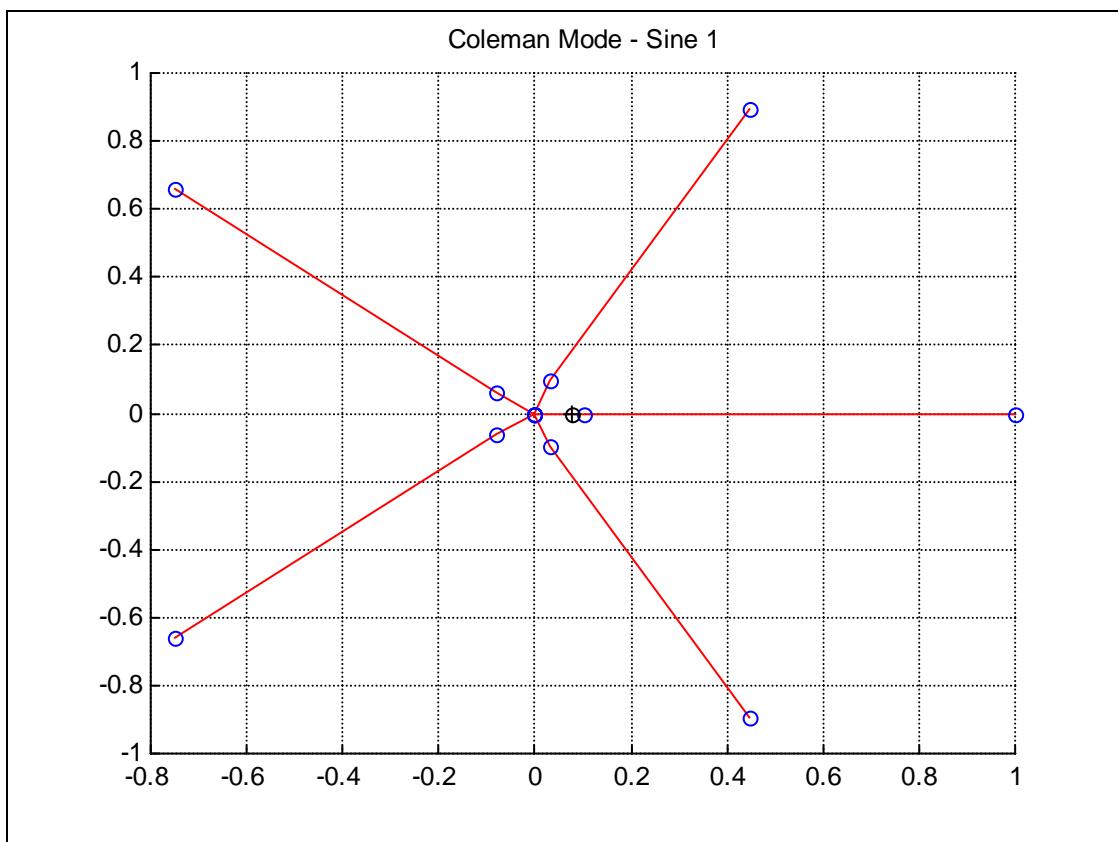
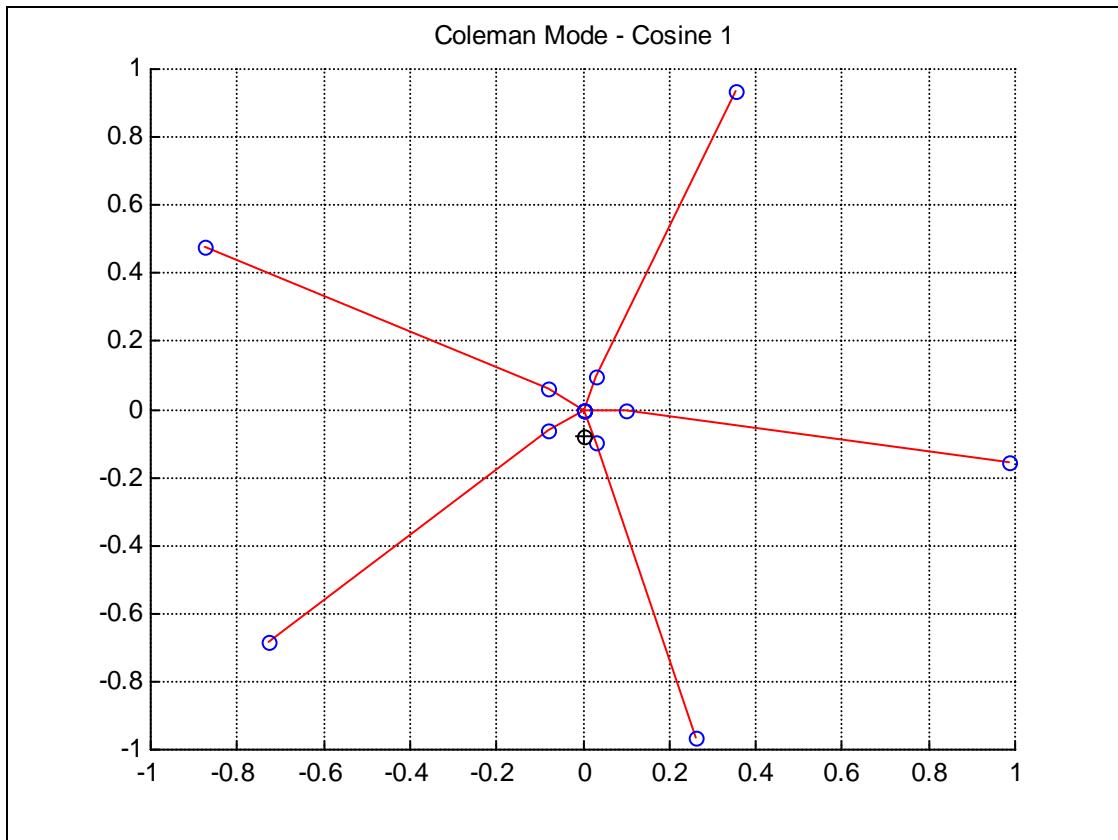
$$\Theta_1 = \sum_k \zeta_k \cos \psi_k$$

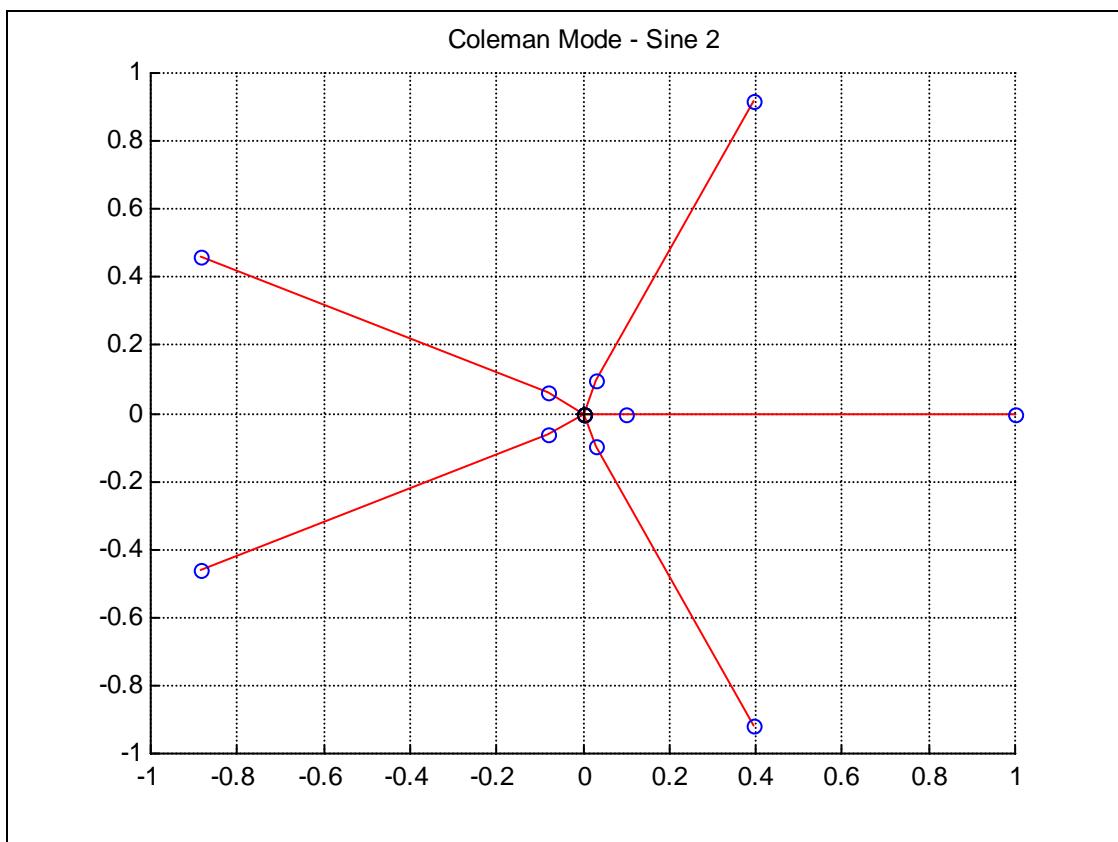
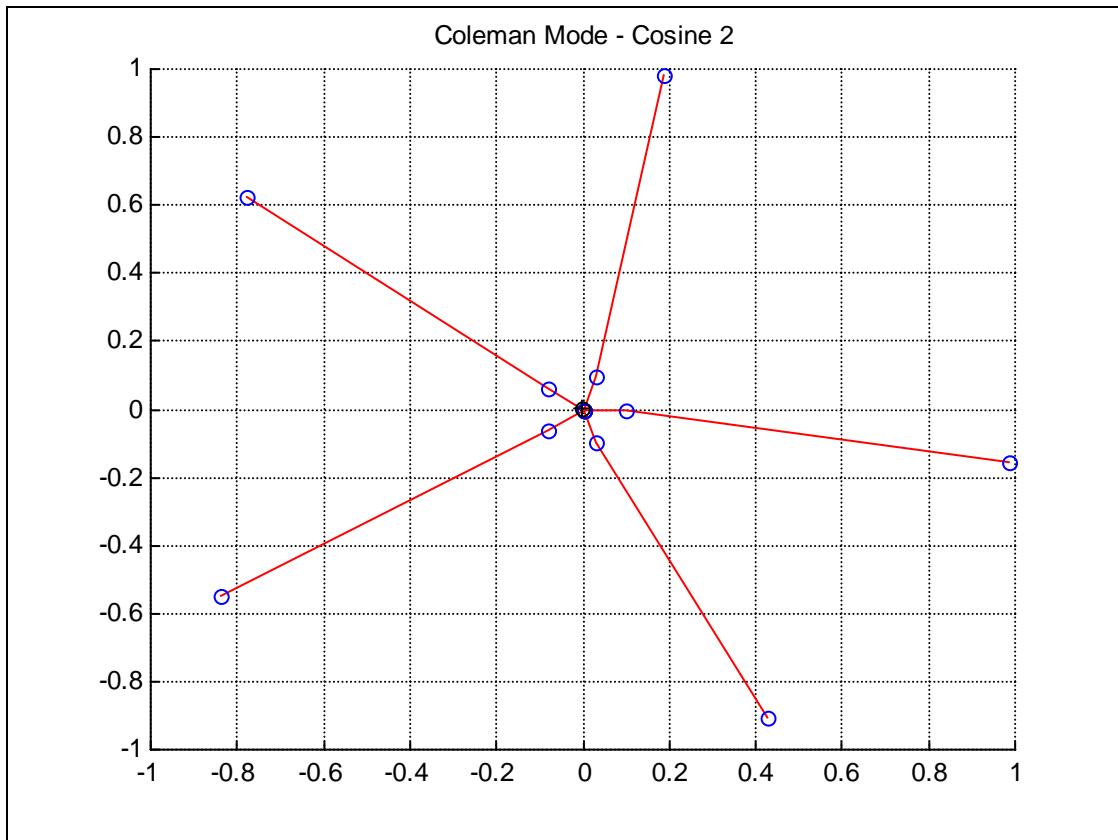
$$\Theta_2 = \sum_k \zeta_k \sin \psi_k$$

$$\Theta_3 = \sum_k \zeta_k \cos 2\psi_k$$

$$\Theta_4 = \sum_k \zeta_k \sin 2\psi_k$$







Or in tabular form:

N	$\Sigma+$	$\Sigma+-$	$\text{Cos}\psi$	$\text{Cos } 2\psi$	$\text{Cos } 3\psi$	$\text{Cos } 4\psi$	$\text{Sin}\psi$	$\text{Sin } 2\psi$	$\text{Sin } 3\psi$	$\text{Sin } 4\psi$
1	✓									
2	✓	✓								
3	✓		✓				✓			
4	✓	✓	✓				✓			
5	✓		✓	✓			✓	✓		
6	✓	✓	✓	✓			✓	✓		
7	✓		✓	✓	✓		✓	✓	✓	
8	✓	✓	✓	✓	✓		✓	✓	✓	
9	✓		✓	✓	✓	✓	✓	✓	✓	✓
10	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓

