# High pulse energy, picosecond MgO:PPLN optical parametric oscillator using a single-mode fiber for signal feedback

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**Abstract:** We demonstrate a high-pulse-energy, synchronously-pumped (7.19MHz), 100ps, widely tunable MgO:PPLN OPO providing  $0.49\mu J$  pulses at  $1.5\mu m$  and  $0.19\mu J$  pulses at  $3.6\mu m$ . A single-mode fiber is employed in the OPO to keep the 42m-long cavity compact. **OCIS codes:** 190.0190, 190.4970, 140.7090, 060.2320, 140.2020

### 1. Introduction

Synchronously-pumped optical parametric oscillators (SPOPO) are useful sources for ultrashort, tunable optical pulses. Progress in high-power, compact, ultrafast pump sources means that they are becoming ever more attractive for a broad range of applications. Advances have included a high-average-power femtosecond OPO pumped by a thin-disk Yb:YAG laser [1], a high-pulse-energy picosecond OPO pumped by an Yb:fiber laser [2], and a controllable repetition rate picosecond OPO pumped by a fiber-amplified gain-switched laser diode [3].

In this contribution, we report a picosecond SPOPO outputting high-energy pulses by using a pump source consisting of a 100ps gain-switched laser diode operating at 7.19MHz and amplified in three Yb-doped fiber amplifiers, similar to that in [4], and by applying a single-mode fiber in the SPOPO cavity for signal feedback [1] to maintain a compact device.

# 2. Experimental setup

The 1.06 $\mu$ m pump source was similar in all respects to that described previously [4]. However, the chirped fiber Bragg grating, used to compensate for the predominantly linear chirp developed on the gain-switched seed pulse, was by-passed here leading to an increase of the pulse duration from 20ps to 100ps. The reduced peak power helped to avoid nonlinear effects in the fiber amplifiers enabling operation with pulse energies of up to 8.6 $\mu$ J, corresponding to an average power of 62W at the 7.19MHz repetition rate. The bandwidth of the output pulses was 0.25nm and was thus ~10 times the bandwidth limit. This highly efficient and compact pump source delivered near-diffraction-limited (M²  $\leq$  1.2 at any power level), linearly polarized output pulses with user-controllable repetition rate and a minimum of free-space optical components.

A schematic diagram of the SPOPO system is shown schematically in figure 1. The majority of the optical path length of the SPOPO ring cavity was provided by a ~27m piece of standard single-mode fiber (Corning SMF28). This allowed a small-footprint and compact overall resonator layout. A periodically-poled,  $40x10x1mm^3$  MgO:LiNbO<sub>3</sub> (MgO:PPLN) crystal was used as gain medium (Covesion Ltd.), where the 40mm length provided sufficient phase-matching bandwidth and relatively small temporal walk-off for the 100ps pump and signal pulses. The crystal had seven, 1mm by 1mm, periodically-poled gratings with periods ranging from 29.5 $\mu$ m up to 31.5 $\mu$ m, and was held in an oven at 150°C to eliminate any residual photorefractive effects. By using a lens with 250mm focal length to focus the pump beam into the crystal, a spot radius (1/e²-intensity) of 104 $\mu$ m was achieved giving a good match to the calculated 99 $\mu$ m signal spot radius obtained with 250mm radius of curvature mirrors CM1 and CM2.

A 90% transmission output coupler was used in order to achieve efficient output power extraction. The high output coupling also ensured that relatively low signal powers were returned through the optical fiber in order to avoid nonlinear distortions such as Raman scattering. Nevertheless, a fiber-coupling arrangement with a heat-sink was required to avoid gradual damage to the fiber end-facet. The MgO:PPLN crystal and the coupling lenses L1, L2 were broad-band anti-reflection coated at the signal wavelengths, and the MgO:PPLN coating was also low reflectivity for the pump and idler wavelengths. The cavity was singly-resonant for the signal with CM1 and CM2 being highly reflective for the signal wavelengths and all cavity mirrors being highly transmissive for the idler wavelengths. Furthermore, the fiber was not transmissive at mid-infrared wavelengths. The idler was extracted

through CM2, which had a transmission of 88% at the idler wavelengths. The transmission of CM1 for the pump wavelength of  $1.06\mu m$  was 92%.

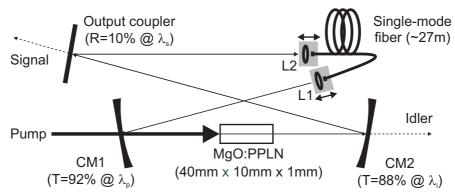


Fig. 1. Schematic diagram of the SPOPO fiber feedback cavity.

### 3. Results

Figure 2 shows the signal and the idler output power of the SPOPO as a function of the input pump power by using the T=90% signal output coupler. All power values given here are measured external to the cavity. The oscillation threshold was achieved at an average pump power of 3.94W. Up to 3.51W of average signal power at 1.504 $\mu$ m and 1.37W of average idler power at 3.591 $\mu$ m was obtained using the 29.5 $\mu$ m-period grating and average pump powers up to 20.07W. The pump depletion was typically 35%. These values correspond to pulse energies of 0.49 $\mu$ J and 0.19 $\mu$ J for the signal and idler, respectively. The highest pulse energy of 0.65 $\mu$ J (signal only, idler not used) reported previously stems from a relay-imaging free-space OPO resonator by employing a cavity dumping technique [5].

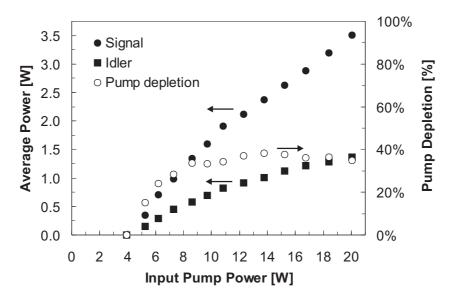


Fig. 2. Signal and idler output power as a function of input power. Pump depletion is shown on the right.

A manual fiber polarization controller was used in an attempt to control the polarization of the non-polarization-maintaining single-mode fiber and to deliver a linearly polarized beam with the correct orientation to the MgO:PPLN crystal. However, it was observed experimentally that this had no effect on the OPO output power. The parametric gain was in the high gain regime, where we estimated a required gain at the oscillation threshold of ~40 to overcome the losses due to the high output coupler transmission, the loss at fiber launching and the loss due to the incorrect polarization of the beam at the fiber output. In this regime the gain is given by 1/4-exp( $2\Gamma L$ ), where L is the crystal length and  $\Gamma$  is the gain coefficient, as defined in [6]. Hence, the threshold gain exponent  $2\Gamma L$  was ~5. At the maximum pump power, ~5 times above threshold, the exponent is  $\sqrt{5}$  times greater (because  $\Gamma^2$  is proportional to the pump power), i.e.  $2\Gamma L = 11.2$ , corresponding to a gain of ~43dB. An accurate *ab initio* calculation of the

effective gain exponent is complicated by the uncertainty over the effective pump intensity, since it has a Gaussian rather than a plane wave form.

We observed instability and a power roll-over of the output power when increasing the average pump power beyond 20W. The reason for the roll-over is not entirely clear but it appears to be not related to average power limitations for the MgO:PPLN crystal as a reduction of the pump average power and hence the thermal input to the crystal by placing a 50/50 duty cycle optical chopper in the pump beam did not improve the performance of the SPOPO. Physical damage, e.g. cracking or coating damage, of the crystal was not observed. There are indications that the limitation may lie in the pump source, especially in the optical isolator used between pump source and SPOPO. The crystal of the isolator may suffer from thermal effects, leading to a beam quality degradation and/or a spatial beam drift. Future work will investigate these limitations further.

The signal beam duration was measured with a 32GHz InGaAs detector and a 50GHz communication signal analyzer (experimentally determined minimum measurable pulse duration of ~30ps) to be 100ps and thus similar to that of the pump. Spectra were taken for both the signal (with an optical spectrum analyzer) and idler (with a monochromator) pulses giving FWHM values of 2.0nm and 12nm, respectively. These figures indicate strongly chirped pulses, which is again similar to the performance of the pump laser. The SPOPO signal and idler wavelength was tuned by accessing different poled gratings of the MgO:PPLN with the pump beam and a corresponding cavity length adjustment to account for the changing round-trip time of the signal pulses due to dispersion. A signal tuning range from  $1.5\mu m$  to  $1.7\mu m$  and an idler tuning range from  $2.9\mu m$  to  $3.6\mu m$  was achieved. The  $M^2$ -values were determined at the highest power to be 1.5 by 1.3 for the signal and 2.8 by 1.9 for the idler, where the higher value for both beams corresponds to the axis perpendicular to the nonlinear crystal plane and parallel to the ring cavity plane.

## 3. Summary

In summary, we have demonstrated a high pulse energy, synchronously-pumped, picosecond SPOPO operating at low MHz repetition rate. The use of a single-mode fiber to provide signal feedback is a key feature to maintain a compact resonator size despite its long cavity length. In combination with the efficient, compact, and power-scalable pump source consisting of a gain-switched laser diode and a chain of three Yb-doped fiber amplifiers, the overall system is attractive as a potential source for infrared materials processing [7,8]. Future work will concentrate on further power scaling into the  $\mu J$  pulse energy regime at pulse durations of a few picoseconds.

### 4. References

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