



Observation of Suppression of Light Scattering Induced by Dipole-Dipole Interactions in a Cold-Atom Ensemble

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(Received 18 February 2014; published 26 September 2014)

We study the emergence of collective scattering in the presence of dipole-dipole interactions when we illuminate a cold cloud of rubidium atoms with a near-resonant and weak intensity laser. The size of the atomic sample is comparable to the wavelength of light. When we gradually increase the number of atoms from 1 to ~ 450 , we observe a broadening of the line, a small redshift and, consistently with these, a strong suppression of the scattered light with respect to the noninteracting atom case. We compare our data to numerical simulations of the optical response, which include the internal level structure of the atoms.

DOI: [10.1103/PhysRevLett.113.133602](https://doi.org/10.1103/PhysRevLett.113.133602)

PACS numbers: 42.50.Ct, 03.65.Nk, 32.80.Qk, 42.50.Nn

When resonant emitters, such as atoms, molecules, quantum dots, or metamaterial circuits, with a transition at a wavelength λ , are confined inside a volume smaller than λ^3 , they are coupled via strong dipole-dipole interactions. In this situation, the response of the ensemble to near-resonant light is collective and originates from the excitation of collective eigenstates of the system, such as super- and subradiant modes [1–3]. Dipole-dipole interactions affect the response of the system and the collective scattering of near-resonant light differs from the case of an assembly of noninteracting emitters [4]. It has even been predicted to be suppressed for a dense gas of cold two-level atoms [5].

Following the recent measurement of the collective Lamb shift [6] in a Fe layer [7], in a hot thermal vapor [8], and in arrays of trapped ions [9], it was pointed out [10] that the collective response of interacting emitters is different between ensembles exhibiting inhomogeneous broadening, such as solid state systems or thermal vapors, and those free of it, such as cold-atom clouds. In particular, inhomogeneous broadening suppresses the correlations induced by the interactions between dipoles, leading to the textbook theory of the optical response of continuous media [10,11]. In the absence of broadening, however, this theory fails and should be revisited to include the light-induced correlations [12–19]. Several recent experiments aiming at studying collective scattering with identical emitters used large and optically thick ensembles of cold atoms [20–23]. However, the case of a cold-atom ensemble with a size comparable to the optical wavelength has not been studied experimentally, nor has the transition between the well-understood case of scattering by an individual atom [24] to collective scattering. In particular, the suppression of light scattering when the number of atoms

increases in a regime of collective scattering has never been directly observed.

Here, we study—both experimentally and theoretically—the emergence of collective effects in the optical response of a cold-atom sample due to dipole-dipole interactions, as we gradually increase the number of atoms. To do so, we send low-intensity near-resonant laser light onto a cloud containing from 1 to ~ 450 cold ^{87}Rb atoms, with a size comparable to the wavelength of the optical transition at $\lambda = 780$ nm. Starting from one atom, we observe a broadening of the line as the number of atoms increases, as well as a small redshift and a strong suppression of the amount of scattered light with respect to the case of noninteracting atoms. We show that this suppression is consistent with the measured broadening and shift. We finally compare our measurements to a numerical simulation of the response of the system in the low excitation limit, including the internal level structure of the atoms.

The suppression of light scattering by resonant dipole-dipole interactions can be understood qualitatively as follows. Consider a laser radiation with frequency ω impinging on an ensemble of classical radiating dipoles with resonance frequency $\omega_0 = 2\pi c/\lambda$ (see Fig. 1). When the dipoles interact through the dipole-dipole potential

$$V_{ja}^{lb} = -V_{dd}[p_{\alpha\beta}(ikr - 1) + q_{\alpha\beta}(kr)^2]e^{ikr}, \quad (1)$$

the system features collective modes with various eigenfrequencies and decay rates. Here, j and l denote two dipoles separated by a distance r , $V_{dd} = 3\Gamma/4(kr)^3$, $k = 2\pi/\lambda$, Γ is the radiative decay rate in the absence of interactions, and the angular functions $p_{\alpha\beta}$ and $q_{\alpha\beta}$ depend on the polarizations α and β and the relative orientations of the dipoles j and l [25,26]. In our experiment, the geometry

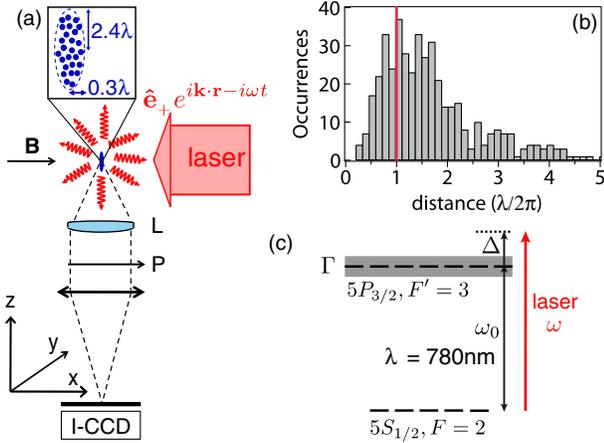


FIG. 1 (color online). (a) Experimental setup. The atoms are initially confined in a microscopic single-beam dipole trap (not shown) (wavelength 957 nm, depth 1 mK, and a waist $1.6 \mu\text{m}$, oscillation frequencies $\omega_x = \omega_y = 2\pi \times 62 \text{ kHz}$ and $\omega_z = 2\pi \times 8 \text{ kHz}$). The excitation laser propagates along the quantization axis x , set by a $B \sim 1 \text{ G}$ magnetic field. We collect the scattered light along z , after a polarizer P oriented at an angle of 55° with respect to x , using a lens L with a large numerical aperture ($\text{NA} = 0.5$) and an image intensifier followed by a CCD camera (I-CCD). (b) Simulation of the distribution of nearest neighbors for a single stochastic realization of a cloud of $N = 450$ atoms. (c) Structure of ^{87}Rb atoms relevant to this work. The excitation light at frequency ω is near resonant with the transition at $\lambda = 2\pi c/\omega_0 = 780 \text{ nm}$.

of the atomic system and its orientation with respect to the excitation are such that the incident laser best couples to only a few modes with decay rates Γ_c larger than Γ (superradiant modes), leading to a broader excitation spectrum. The excitation rate, and therefore the amount of scattered light, should thus be reduced by a factor $(\Gamma/\Gamma_c)^2$. The effect is stronger when the average distance between dipoles $\langle r \rangle$ is smaller than $\lambda/2\pi$.

To study the collective scattering by an ensemble of atoms coupled via resonant dipole-dipole interactions we use the setup depicted in Fig. 1(a). We prepare small clouds containing up to 450 atoms at a temperature $\sim 100 \mu\text{K}$, confined in a microscopic dipole trap [27], and illuminate them with laser light nearly resonant with the atomic transition at $\lambda = 780 \text{ nm}$. The Doppler width of the sample (150 kHz) is much smaller than the atomic linewidth $\Gamma/2\pi = 6 \text{ MHz}$, making inhomogeneous broadening negligible. The anisotropy of the trap results in an elongated cloud with calculated root-mean-square thermal sizes $\sigma_\rho = 0.3\lambda$ and $\sigma_z = 2.4\lambda$. The maximal density is $\rho = 2.5 \times 10^{14} \text{ at./cm}^3$ and the minimal average interatomic distance $\langle r \rangle = \rho^{-1/3} = 0.2\lambda$ [Fig. 1(b)]. In this regime, $k\langle r \rangle \sim 1$, leading to $V_{\text{dd}} \sim \Gamma$, and the resonant dipole-dipole interaction will therefore have an effect on the scattering.

Experimentally, we prepare the trapped atoms in the $F = 2$ hyperfine manifold with an efficiency better than 95%. We then release them in free space by switching off the

trapping light while exciting them with σ_+ polarized light at a frequency $\omega = \omega_0 + \Delta$ tuned near the ($5S_{1/2}, F = 2$) to ($5P_{3/2}, F' = 3$) transition (see Fig. 1). In this way we avoid extra light shifts induced by the trapping beam that would obscure the measurement of small collective shifts and broadening. Also, we choose the intensity saturation $I/I_{\text{sat}} = 0.1$ to be in the low excitation limit ($I_{\text{sat}} = 1.6 \text{ mW/cm}^2$). We interleave excitation pulses with duration 125 ns and recapture periods in the dipole trap with duration $1 \mu\text{s}$. This sequence is repeated 200 times using the same cloud of atoms, in order to improve the duty cycle of the experiment. Finally, we prepare a new atomic sample and repeat the set of excitation pulses a few hundred times. The scattered light that we collect in the z direction is therefore the result of an average over many spatial configurations of the atoms. The choice of the number of pulses (200) is a trade-off between getting a good signal-to-noise ratio and avoiding light-assisted losses [28] or heating of the cloud, both of which would lower the density. We checked that both effects do not exceed 5% over the entire set of pulses and that less than 5% of the atoms are depumped in the ($5S_{1/2}, F = 1$) hyperfine level during the excitation.

Figure 2 shows the number of photons $n_z(N, \Delta)$ detected by the I-CCD as a function of the detuning Δ of the excitation laser, for various numbers of atoms N . A Lorentzian fit agrees well with the data for the range of N explored here. As expected from the qualitative argument described above, we observe that the full width at half maximum (FWHM) increases with the number of atoms [see Fig. 3(a)], since the interatomic distance then decreases, leading to stronger dipole-dipole interactions. We also measure a small redshift $\delta\omega$ of the center frequency [Fig. 3(b)]. For $N = 1$, the FWHM is $1.35 \pm 0.15\Gamma$, in agreement with the short duration of the excitation pulses (125 ns), which broadens slightly the resonance. Figure 2

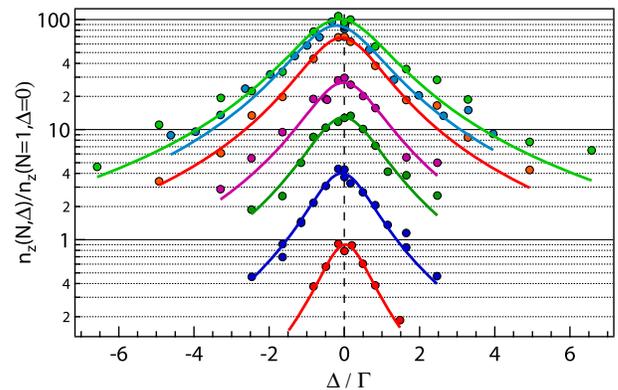


FIG. 2 (color online). Amount of scattered light detected $n_z(N, \Delta)$, versus the detuning Δ of the excitation light for numbers of atoms $N = 1, 5, 20, 50, 200, 325, 450$ (from bottom to top). The amplitudes of the curves are normalized to the amount of light detected at resonance for a single atom, $n_z(N = 1, \Delta = 0)$. Solid lines: Lorentzian fits to the data. Typical uncertainties: 10% (vertically) and 20% (horizontally).

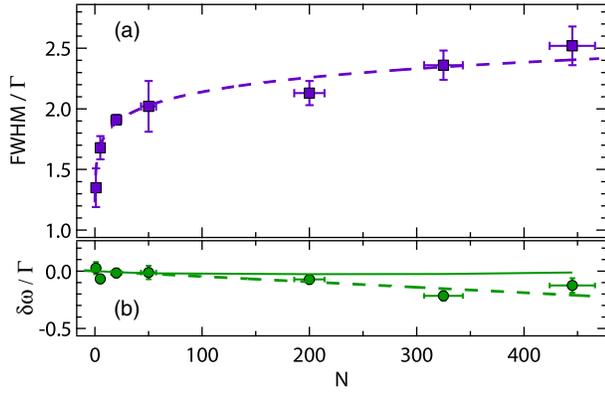


FIG. 3 (color online). (a) FWHM and (b) line shift $\delta\omega$ with respect to the atomic frequency for $N = 1$ atom, in units of Γ . Filled symbols: data extracted from the Lorentzian fits shown in Fig. 2, versus the number of atoms. Dashed lines: phenomenological fits of the FWHM and shift by, respectively, $\Gamma_c/\Gamma = 1.49(6) \times N^{0.08(1)}$, and $\delta\omega_c/\Gamma = 47(9) \times 10^{-5}N$. The error bars are from the fits of Fig. 2. Green solid line: results of the simulation (see text).

also shows that the amount of light scattered in the z direction at resonance does not increase linearly with the number of atoms as one would expect for noninteracting atoms, but actually increases more slowly. Figure 4(a) indicates that this is also the case off resonance, where we plot $n_z(N, \Delta)/n_z(N = 1, \Delta)$ for different atom numbers and detunings. For noninteracting atoms this ratio is equal to the number of atoms N (and is thus independent of the detuning Δ), as we verified by collecting the scattered light after letting the atomic cloud expand in free space for a sufficiently long time [29]. By contrast, here we observe that the amount of scattered light is strongly suppressed on resonance as the number of atoms increases, and that we gradually recover the behavior of noninteracting atoms as we detune the laser away from resonance.

All the observations reported above can be reproduced by a single functional form:

$$n_z(N, \Delta) = C \frac{N}{\Gamma_c(N)^2 + 4[\Delta - \delta\omega_c(N)]^2}, \quad (2)$$

where C includes the detection efficiency of the imaging system. This is illustrated in Fig. 4(b): we find that the quantity $R(N, \Delta)/R(N = 1, \Delta)$, where $R(N, \Delta) = n_z(N, \Delta)[\Gamma_c^2 + 4(\Delta - \delta\omega_c)^2]$ and Γ_c and $\delta\omega_c$ are, respectively, the phenomenological fits of FWHM and the shift (see Fig. 3), collapses on a single curve whatever the detuning. For $N \lesssim 300$, this curve is linear with N with a slope of 1, in agreement with Eq. (2). It emphasizes that in this regime, the scattered intensity is suppressed by a factor $(\Gamma/\Gamma_c)^2$ at resonance, as expected from the qualitative discussion earlier. We note that this scaling cannot be explained by a model where the suppression would come from an incoherent superposition of the intensities scattered

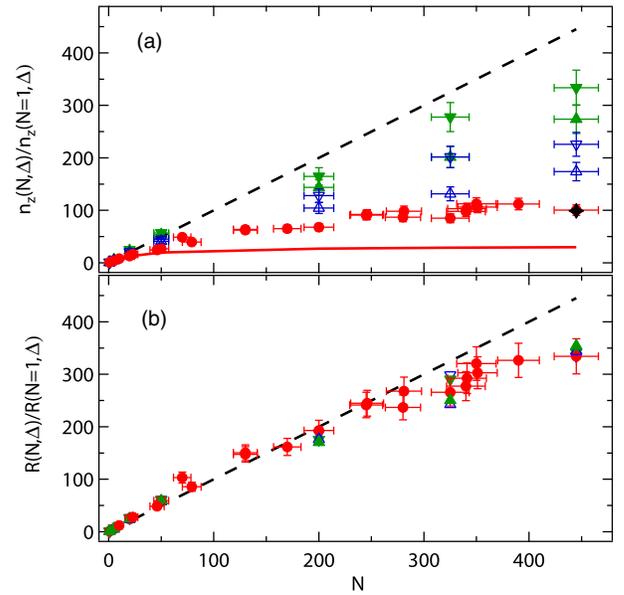


FIG. 4 (color online). (a) Scattered light detected in the z direction, versus the number of atoms, for different detunings of the laser: $\Delta = 0$ (red circles), $\pm\Gamma$ (up or down open triangles), and $\Delta = \pm 2.5\Gamma$ (up or down filled triangles). The intensity for each atom number is normalized to the single atom case at the same detuning. Red line: result of the simulation (see text) with widths of the cloud σ_ρ and σ_z . Black diamond: model with widths $2\sigma_\rho$ and $2\sigma_z$. (b) Ratio $R(N, \Delta)/R(N = 1, \Delta)$ versus the number of atoms (see text). Dashed line in (a) and (b): case of noninteracting atoms.

by each atom with resonant frequencies inhomogeneously distributed over a distribution of width FWHM: that would lead to a suppression that would scale as Γ/Γ_c near resonance, instead of the $(\Gamma/\Gamma_c)^2$ scaling observed here. For $N > 300$, the departure from the linear law indicates that C depends on the number of atoms in this regime, and that the simple Lorentzian form (2) becomes inaccurate, as also found in the simulation (see below).

We have performed numerical simulations of the collective dynamic response of the atomic sample to near-resonant pulsed light in the low excitation limit. In this model, each atom, located at position \mathbf{r}_j ($j = 1, \dots, N$) and with dipole \mathbf{d}_j , is driven by the incident laser field and by the fields scattered by all the $N - 1$ other atoms, i.e., each dipole is coupled to the $N - 1$ other dipoles via the resonant interaction of Eq. (1). This classical electrodynamics simulation incorporates all the interactions between an ensemble of nonsaturated discrete dipoles. This approach has been used to study dielectric media comprising two-level or spatially averaged isotropic electric dipoles [10,14,16,18,30,31] as well as magnetodielectric circuit resonator systems [32]. Here, we also incorporate the Zeeman level structure of the atoms [13] and the shifts associated to the presence of the magnetic field. To calculate the dipoles \mathbf{d}_j in our experimental configuration,

we stochastically sample the positions of the atoms according to a 3-dimensional Gaussian density distribution with root-mean-square sizes given by the thermal sizes of the cloud along and perpendicular to the trap propagation axis; each atomic position is treated as an independent and identically distributed random variable. At each realization the N atoms are fixed at positions \mathbf{r}_j ($j = 1, \dots, N$) and we stochastically sample the magnetic quantum number of the Zeeman states m_j of each atom j . The probability of atom j being in state $|g, m\rangle$ ($m = \pm 2, \pm 1, 0$) is the initial population of that Zeeman state p_m ($0 < p_m < 1$; $\sum_m p_m = 1$). The optical pumping used in the preparation step before the excitation sequence, skews the initial populations; here we use the values $p_0 = p_1 = p_2 = 1/3$ and $p_{-1} = p_{-2} = 0$. We write the positive frequency component of the dipole produced by each atom j that oscillates at the laser frequency as $\mathbf{d}_j = \mathcal{D} \sum_{\sigma} \hat{e}_{\sigma} C_{m_j}^{(\sigma)} \mathcal{P}_{j\sigma}$, where the sum runs over the unit spherical polarization vectors $\sigma = \pm 1, 0$. The amplitude of the atomic dipole j associated to the optical transition $|g, m_j\rangle \rightarrow |e, m_j + \sigma\rangle$ is proportional to the reduced dipole matrix element \mathcal{D} , the atomic coherence $\mathcal{P}_{j\sigma}$, and the corresponding Clebsch-Gordan coefficient $C_{m_j}^{(\sigma)}$. The temporal evolution of the coherences is given by the set of coupled equations

$$\begin{aligned} \dot{\mathcal{P}}_{j\alpha} - i(\Delta_{j\alpha} + i\Gamma/2)\mathcal{P}_{j\alpha} \\ = -i\Omega_{j\alpha}(t) - i \sum_{l \neq j} \sum_{\beta} C_{m_l}^{(\beta)} C_{m_j}^{(\alpha)} V_{j\alpha}^{l\beta}(\mathbf{r}) \mathcal{P}_{l\beta}, \quad (3) \end{aligned}$$

where $\Omega_{j\alpha}(t)$ and $\Delta_{j\alpha} = \omega - \omega_{j\alpha}$ are, respectively, the time-dependent Rabi frequency and the detuning of the driving laser with respect to the Zeeman shifted transition of the α -polarized atom j with frequency $\omega_{j\alpha}$, and $\beta = \pm 1, 0$. Here, we deduce $\Omega_{j\alpha}(t)$ from the experimentally measured temporal profile of the excitation pulse. The last term in Eq. (3) couples the α -polarized dipole j to the β -polarized dipole l separated by $\mathbf{r} = \mathbf{r}_j - \mathbf{r}_l$ according to Eq. (1). We have solved Eqs. (3) numerically in the presence of a 1 G magnetic field to calculate the light field amplitude that is scattered into the solid angle encompassed by the aspherical lens in the far field. Finally, accounting for the polarization-sensitive detection scheme, we calculated the measured light intensity.

The simulation predicts that the spectra $n_z(N, \Delta)$ should present an increasing broadening and asymmetry (Fig. 5), a negligible shift [Fig. 3(b)], as well as a suppression of the scattered light [Fig. 4(a)] when the number of atoms increases. These features are in good agreement with our data for $N \lesssim 50$. In this range, the simulated spectra are well fitted by a Lorentzian for $N \lesssim 50$, thus justifying our fitting of the data by Eq. (2) and the collapse of the data shown in Fig. 4(b). For $N \gtrsim 50$, the agreement is only qualitative, as the effects are found to be less pronounced experimentally. We attribute these discrepancies to two possible reasons.

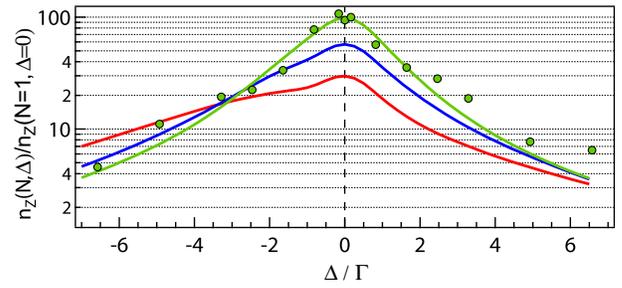


FIG. 5 (color online). Comparison between experiment and theory for the number of detected photons $n_z(N, \Delta)$ (normalized to the single atom case at resonance) for $N = 450$. The red, blue, and green lines correspond to samples with widths σ_{ρ} and σ_z , multiplied by a factor 1, 1.44, and 2, respectively.

First, forces induced by the dipole-dipole interactions may expel atom pairs with shortest interatomic distances, thus breaking down the assumption that atoms have frozen positions during the sequence of pulses excitations. This is all the more likely as the number of atoms is large, and could explain the evolution of the FWHM in Fig. 3(a). This effect is hard to check experimentally since the sample is smaller than the diffraction limit of our imaging system. We found numerically, however, that an increase by a factor 2 in the widths σ_{ρ} and σ_z already restores a nearly Lorentzian profile close to the measured spectra (see Fig. 5), and yields the observed suppression of light scattering [see Fig. 4(a)]. Second, the simulation predicts that the number of detected photons increases by a factor of 2 when the initial distribution of Zeeman state populations varies from $p_0 = p_1 = p_2 = 1/3$ to $p_2 = 1$. For large numbers of atoms, optical pumping during the set of excitation pulses may change the distribution of populations, an effect not accounted for in our model.

In conclusion, we have directly measured the suppression of light scattering induced by dipole-dipole interactions in an ensemble of cold atoms driven by a near-resonant weak laser field and compared it with a time-dependent model of coupled dipoles. The model reproduces the observed trends. In the future, we plan to investigate to what extent the observed collective scattering involves beyond-mean-field scattering processes, i.e., is cooperative in nature. Experimental investigations of the temporal response of the system, and comparisons to the case of a single atom [24], should also provide insight into the interplay between dipole-dipole interactions and collective scattering.

We acknowledge support from the E.U. through the ERC Starting Grant ARENA, from the Triangle de la Physique, EPSRC, and Leverhulme Trust. We thank P. Pillet, J.-J. Greffet, and J. Javanainen for fruitful discussions.

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