

# Numerical investigation of a pair of self-propelled AUVs operating in tandem

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## Abstract

This paper investigates the influence of the propeller race on upstream and downstream self-propelled AUVs. Initially simulations of a self-propelled hull are performed at the Reynolds Number  $3.2 \times 10^6$  with the commercial RANS code ANSYS CFX 12.1, utilising a body force model to replicate the impact of the propeller utilising momentum source terms. This is then extended to consider a fleet of two self-propelled vehicles operating at a range of longitudinal offset and transverse separations. The results highlight that operation in close proximity to another self-propelled vessel has a significant impact of both the flow around the hull and drag experienced by the vehicle. A propeller race deduction is proposed to account for the increase in vehicle drag due to the propulsors of other vehicles. The propeller race deduction is dependent upon both longitudinal and transverse separation. From a vehicle or mission design perspective, it is important to correctly understand the true propulsive energy budget of the vehicle and its impact on both range and endurance. This study highlights the importance of considering both thrust deduction and any propeller race deductions when calculating the propulsive power consumption of an individual or fleet of vehicles.

*Keywords:* co-operative AUVs, drag, RANS-SST, self-propelled model, body force model, propeller race deduction

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## 1. Introduction

The rapid growth in ocean operational services such as: offshore engineering, defence, shipping and marine ecosystem survey, challenge the performance and efficiency of manned and unmanned robotic systems (Woods, 1997). Since cold and deep ocean operational conditions are deemed too dangerous or impractical for humans (Caccia, 2006; Horgan and Toal, 2006; Smallwood and Whitcomb, 2004), unmanned submersibles are typically employed. Autonomous Underwater Vehicles (AUVs) are self-propelled robots which perform missions without requiring external powering or control. They are a cost-effective alternative compared to many current manned and tethered system, AUVs are becoming widely used to assist existing ocean science and engineering operations. Current AUVs are primarily designed to perform as a solo vehicle for specific missions, typically inspection or exploration tasks with an operational range up to one thousand kilometres and mission durations ranging from numbers of hours to several days.

The operational ability of an AUV depends on the required speed, range, depth, duration and manoeuvring performance of the vehicle. Since the range and endurance are limited by the cruise speed (Bingham et al., 2002; Furlong et al., 2007; McPhail, 2009; Stevenson et al., 2007; Tripp, 2006), the options to extend an endurance are: minimising drag (Huggins and Packwood, 1994; Parsons, 1972; Parsons et al., 1974), enhancing propulsive efficiency (Ageev, 2000, 1995; Stevenson et al., 2007), increasing specific energy of power sources (Ageev, 2000; Alers, 1981), and reducing hotel load (Furlong et al., 2007; Phillips et al., 2012).

Many simple observations suggest that the drag of an individual may be reduced in specific fleet configurations: examples such as, birds form an echelon formation for long distance migration (Andersson and Wallander, 2003), fish swim in shoals or schools (Alexander, 2004; Hanrahan and Juanes, 2001; Partridge et al., 1983). There are energy benefit to slipstreaming while cycling (Kyle, 1979), driving in convoy (Hucho and Ahmed, 1998) or swimming (Silva et al., 2008). There is a 60% energy saving for bottlenose dolphin calves when swimming in close proximity to their mothers (Weihs, 2004). Recent studies have also demonstrated energy savings for fleets of towed

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generic AUV hull forms at specific longitudinal and transverse offsets (Husaini et al., 2009; Rattanasiri et al., 2014). However, the propeller's influence has been excluded from these simulations. Axial and tangential accelerations, applied to fluid by the action of the propeller, lead to significant variations in the flow field around the stern of the vessel and downstream, see Figure 1.

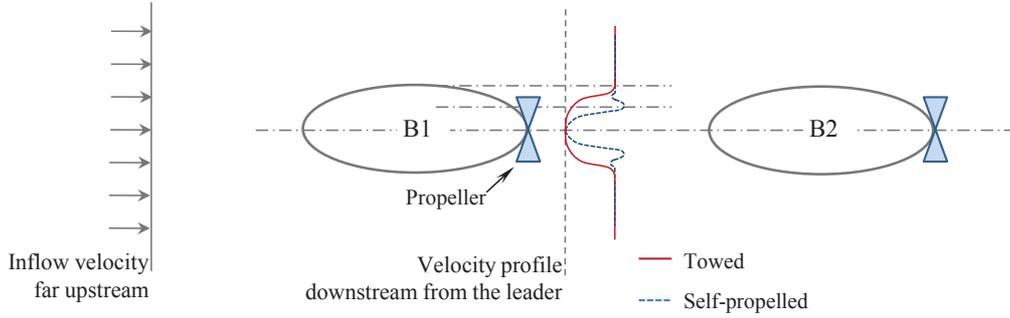


Figure 1: Velocity profile of fluid inflow past AUVs, B1 and B2 represents the leader and follower AUV, respectively.

Firstly, the accelerated flow due to the propeller results in increased skin friction resistance at the stern, as well as reduced local pressures which increase the form drag of the self-propelled body. For an AUV to move forward at the design speed, the required thrust,  $T_{SP}$ , must typically exceed its total towed drag,  $R$ , (Burcher and Rydill, 1994):-

$$T_{SP} = \frac{R}{1-t} \quad (1)$$

where  $R$  is the towed drag of the bare hull and propeller and  $t$  is defined as the thrust deduction which is a function of hull streamlining, propeller clearance and fullness (Burcher and Rydill, 1994).

Secondly, the propeller race from upstream or adjacent vessels will modify the flow field of nearby vehicles. For example a vehicle placed directly downstream of the propeller of an upstream vehicle will experience additional axial and tangential flow velocities in comparison to towed wake conditions, in this study, the increase in thrust required due to the propeller race from nearby vehicles is defined as a propeller race deduction,  $\gamma$ ,

$$T_{SP} = \frac{R}{1-t-\gamma}. \quad (2)$$

Investigation of self-propelled fleet configurations will provide more information on the true propulsive energy budget of a fleet of AUVs. The generated wake upstream/downstream could lead to a reducing/raising energy consumption, this leads to the underlying questions of:-

- does the impact of propeller provide the energy advantages/disadvantages to the fleet configuration?
- does a fleet configuration provide energy benefits for just an individual AUV or the whole fleet?
- what is the optimal configuration and optimal distances of the fleet?

The purpose of this paper is to provide guidance for operators on suitable spacing for multiple vehicles' missions. To achieve this aim, the two hydrodynamic processes of twin self-propelled AUVs: the body-to-body interference (or viscous interaction) and the increase of drag due to re-energised wake by a self-propelled vehicle must be numerically investigated.

## 2. Theoretical approach

Prediction of the hydrodynamic force acting on an AUV's hull can be modelled using a steady-state Reynolds Averaged Navier Stokes (RANS) simulation. This has proved to provide reasonably accurate results with modest computational cost (Jagadeesh et al., 2009; Karim et al., 2009; Phillips et al., 2008, 2007, 2010c; Sarkar et al., 1997). The commercial RANS solver ANSYS CFX 12.1 (ANSYS, 2010) is selected to perform in this study. By assuming the flow is incompressible, the continuity equation in tensor notation becomes:-

$$\frac{\partial U_i}{\partial x_i} = 0 \quad (3)$$

The momentum equation can be written as:-

$$\rho \left( \frac{\partial \overline{U}_i}{\partial t} + \frac{\partial \overline{U}_i \overline{U}_j}{\partial x_j} \right) = -\frac{\partial \overline{P}}{\partial x_i} + \frac{\partial}{\partial x_j} \left\{ \mu \left( \frac{\partial \overline{U}_i}{\partial x_j} + \frac{\partial \overline{U}_j}{\partial x_i} \right) \right\} - \rho \frac{\partial \overline{u'_i u'_j}}{\partial x_j} + \overline{F}_i \quad (4)$$

where the tensor  $x_i$  represents Cartesian co-ordinates ( $X, Y, Z$ ) and  $U_i$  are the Cartesian mean velocity components ( $\overline{U}_x, \overline{U}_y, \overline{U}_z$ ). The Reynolds stress tensor ( $\overline{\rho u'_i u'_j}$ ) is represented by the turbulence closure.  $\overline{F}_i$  represents additional momentum source terms ( $\overline{F}_X, \overline{F}_Y, \overline{F}_Z$ ).

The shear stress transport (SST) turbulence closure model (Menter, 1994) which blends  $k - \varepsilon$  and  $k - \omega$  was selected for this study. Previous investigations have shown that it is better able to replicate the flow around the ship and submarine hull forms than either  $k - \varepsilon$  or  $k - \omega$  model, notably with a moderate computational cost (Larsson and Baba, 1996; Phillips et al., 2010b).

The flow around a rotating propeller is a complex transient flow, high mesh resolution is required around the blade in order to resolve the flow features. Small time steps are required to capture the transient flow behaviour. The problem considered in this paper does not concentrate on the propeller, but rather on a representative model of the velocity field downstream of the propeller. Therefore, this work utilises a body force approach proposed for self-propelled ship simulations. When using a body force model, the geometry of the propeller is not explicitly represented. The effect of the propeller on the flow is modelled for  $\overline{F}_i$  as distributed axial and tangential momentum source terms which induce axial and swirl accelerations in the fluid. The considered model prescribes a radial distribution of thrust and torque based on the Hough and Ordway (1965) circulation distribution, which has zero loading at the tip and root was shown to match Goldstein's optimum distribution (Goldstein, 1929). Coupling this distribution with a RANS simulation has been implemented by Paterson et al. (2003); Phillips et al. (2008, 2010a); Phillips (2009); Stern et al. (1988); Turnock et al. (2008). The non-dimensional thrust distribution ( $\overline{F}'_x$ ) and the non-dimensional torque distribution ( $\overline{F}'_\theta$ ) are given by:-

$$\overline{F}'_x = A_x r^* \sqrt{1 - r^{*2}} \quad (5)$$

$$\overline{F}'_\theta = A_\theta \frac{r^* \sqrt{1 - r^{*2}}}{(1 - Y_h) r^{*2} + Y_h} \quad (6)$$

where:

$$A_x = \frac{C_{th}}{\Delta x} \frac{105}{16(4 + 3Y_h)(1 - Y_h)} \quad (7)$$

$$C_{th} = \frac{8K_T}{\pi J^2} = \frac{2T}{\rho V_a^2 \pi R_p^2} \quad (8)$$

$$A_\theta = \frac{K_Q}{\Delta x J^2} \frac{105}{\pi(4 + 3Y_h)(1 - Y_h)} \quad (9)$$

$$J = \frac{V_a}{n(2R_p)} \quad (10)$$

$$V_a = \frac{(1 - w_t)V}{V} \quad (11)$$

where: both  $C_{th}$  and  $K_T$  are the thrust coefficient with different form as shown in Equation 8.  $K_Q$  is the torque coefficient.  $J$  is the advance coefficient,  $T$  is the thrust,  $w_t$  is the wake fraction,  $n$  is the revolutions per second,  $V$  is the vehicle speed and  $V_a$  is the advance speed. The momentum source terms are then applied to an annulus with finite thickness ( $\Delta x$ ) defined by:-

$$Y_h = \frac{R_h}{R_p}, \quad Y = \frac{r_p}{R_p}, \quad r^* = \frac{Y - Y_h}{1 - Y_h} \quad (12)$$

where  $R_p$  and  $R_h$  are the radius of propeller and hub, respectively.  $r_p$  is the local propeller radius. See Figure 2 for detail.

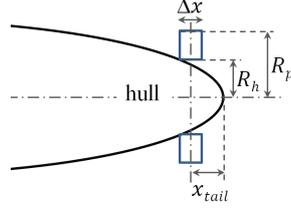


Figure 2: Body force propeller model

### 3. Case study

#### 3.1. Assumed hull shape

A shape of a 6:1 prolate spheroid with a length ( $L$ ) of 1200 mm and a diameter of 200 mm, is assumed to represent the AUV hull shape, giving a surface area ( $A_w$ ) of 0.601 m<sup>2</sup>. This investigation neglects the effect of appendages and other protrusions through the hull. Previously a series of wind tunnel experiments were performed on the flow past twin prolate spheroids of this shape to characterise the side-force and yawing moment interactions at different transverse spacings (Molland and Utama, 1997). Tests were carried out in the 7' × 5' (2.20 m × 1.57 m) low speed wind tunnel at the University of Southampton. The top spheroid (B1) was placed at the middle breadth and 1.07 m height from the floor. The lower spheroid (B2) was placed at the transverse separation ( $S/L$ ) of 0.27, 0.37 and 0.47 away from the centreline of B1. The noses of both spheroids are aligned with zero longitudinal offset ( $D/L = 0$ ) as shown in Figure 3.

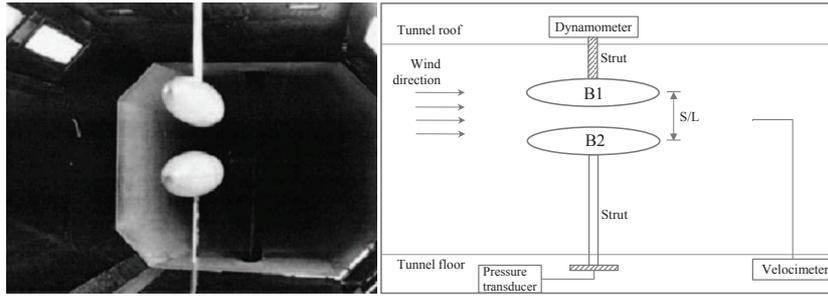


Figure 3: Twin prolate spheroids in the 14 m long, 7' × 5' wind tunnel, figure adapted from Molland and Utama (1997)

#### 3.2. Assumed propeller characteristics

The propeller performance is scaled from Autosub3 (Griffiths, 2003), the ratio of propeller diameter to maximum hull diameter is 0.8. The  $K_T$  and  $K_Q$  of Autosub3 related to  $J$  is predicted by the following equations (Phillips et al., 2009):-

$$K_T = 0.5106J^3 - 0.7700J^2 + 0.2017J + 0.0529 \quad (13)$$

$$10K_Q = 0.3920J^3 - 0.7499J^2 + 0.3218J + 0.0032 \quad (14)$$

For numerical stability, the ratio of finite length of the propeller disc and maximum diameter of the disc is suggested to be  $\geq 0.1$  (Phillips et al., 2010a). The other parameters are in the Table 1.

Table 1: Parameter of modelled propeller

Parameters	$R_p$ (m)	$R_h$ (m)	$\Delta x$ (m)	$x_{tail}$ (m)
Setting	0.080	0.0154	0.0154	0.007

In this study, the propeller rpm is iteratively varied to determine the appropriate advance coefficient (Equation 10), thrust coefficient (Equation 13) and torque coefficient (Equation 14) to achieve self-propulsion.

### 3.3. Present study

The configuration of a pair of AUVs is shown in Figure 4 where  $S/L$  is the transverse separation and  $D/L$  is the longitudinal offset, B1 and B2 represents the leader and follower, respectively. The action of the propeller is to re-energise the wake compared to a towed vessel. This propeller race will interact with any vehicles directly downstream impacting on their own drag. The propeller test case matrix in this study is shown in Table 2.

- Initial simulations are performed for a single hull. Case S1 is the benchmark towed single hull. Case S2 corresponds to a self-propelled vehicle where the resulting is modelled and torque is neglected. Case S3 corresponds to a self-propelled vehicle modelling thrust and torque. Comparing S2 and S3, the influence of the propeller torque and swirling flow may be examined.
- Subsequently, simulations are performed for a fleet of two vehicles. Cases F1 to F6 examine a fleet performance for various combinations of the towed and self-propelled vehicles at  $S/L = 0$  and  $D/L = 1.47$ .
- The impact of spacing on drag, thrust deduction and the propeller race deduction of a pair of self-propelled vehicles is then investigated by setting up Case F6 with varying  $1.17 \leq D/L \leq 1.77$  and  $0.0 \leq S/L \leq 0.37$ .
- Cases F1 to F3 are then extended to consider in detail; two self-propelled vehicles at a selection of longitudinal offsets for  $1.17 \leq D/L \leq 4.47$  and transverse separations for  $0.17 \leq S/L \leq 0.47$ .

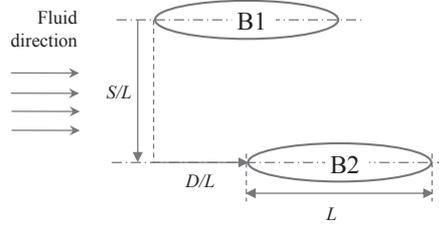


Figure 4: A pair of prolate spheroids in fleet configurations. Where  $S/L$  is the transverse separation and  $D/L$  is the longitudinal offset, B1 and B2 represents the leader and follower AUV, respectively.

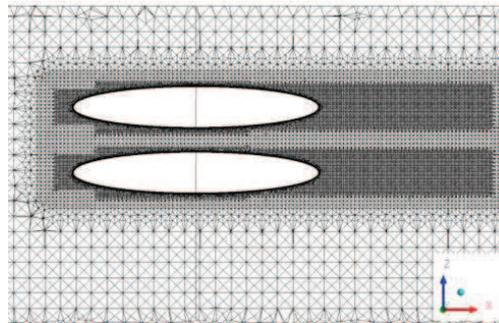
Table 2: Propeller test case matrix to simulate twin self-propelled vehicles in  $S/L = 0$  and  $D/L = 1.47$  configuration at various propulsive conditions.

Case	B1	B2
S1	Towed Hull	-
S2	Thrust Only Propeller Model	-
S3	Thrust and Torque Propeller Model	-
F1	Towed Hull	Towed Hull
F2	Thrust Only Propeller Model	Towed Hull
F3	Thrust Only Propeller Model	Thrust Only Propeller Model
F4	Thrust and Torque Propeller Model	Towed Hull
F5	Thrust and Torque Propeller Model	Thrust Only Propeller Model
F6	Thrust and Torque Propeller Model	Thrust and Torque Propeller Model

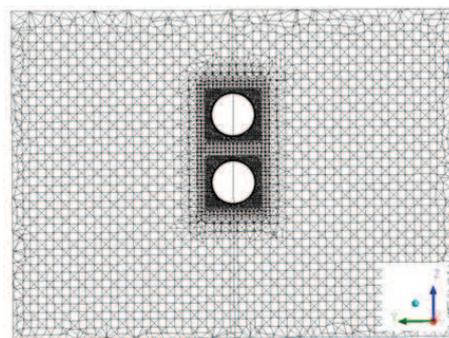
### 3.4. Numerical settings and mesh strategy

The fluid domain is modelled as  $1.4L \times 12L \times 1.8L$  to replicate the experimental domain of Molland and Utama (1997). Both hulls are modelled using a no slip wall condition. A free slip wall condition is used for the roof, floor and walls. The simulation is performed at a length Reynolds number ( $Re$ ) of  $3.2 \times 10^6$  (Typical AUV operation occurs at  $10^5 < Re < 10^7$ ) corresponding to an air inlet velocity ( $V$ ) of 40 m/s. The outlet is set at zero relative pressure. The air density ( $\rho_{air}$ ) and the air kinematic viscosity ( $\nu_{air}$ ) at room temperature are  $1.185 \text{ kg/m}^3$  and  $1.545 \times 10^{-5} \text{ m}^2/\text{s}$ , respectively. The computational parameters are provided in Table 3. Figure 5a and 5b show the sample of meshes cut in the  $ZX$  plane and the  $YZ$  plane, respectively. Figure 6 shows the influence on mesh size on the drag of a towed spheroid, eleven meshes where built ranging from 1.2 to 22.7 million elements. The

computational cost of using the medium mesh (**8.9 million elements**) is 10 times lower than that of the fine mesh (**22.7 million elements**) with a less than 2% variation in drag. Details of the fluid domain discretisation and mesh strategy's validation is completely presented in Rattanasiri et al. (2014). Results utilising these meshes exhibited good correlation with the pressure distribution, the side-force coefficient, form factor and the towed drag measured in: the wind tunnel experiments (Molland and Utama, 1997), previous numerical analysis (Molland and Utama, 2002) and empirical predictions (Hoerner, 1965). By implementing a propeller model, the investigation of the impact of propellers to a fleet of self-propelled twin hulls can be performed.



(a) ZX plane at the centre line at  $Y = 0$  with the fluid flow from left to right



(b) YZ plane at  $X = 0.6$  m from the noses

Figure 5: Mesh cut around a pair of spheroids for  $S/L = 0.27$  and  $D/L = 0$

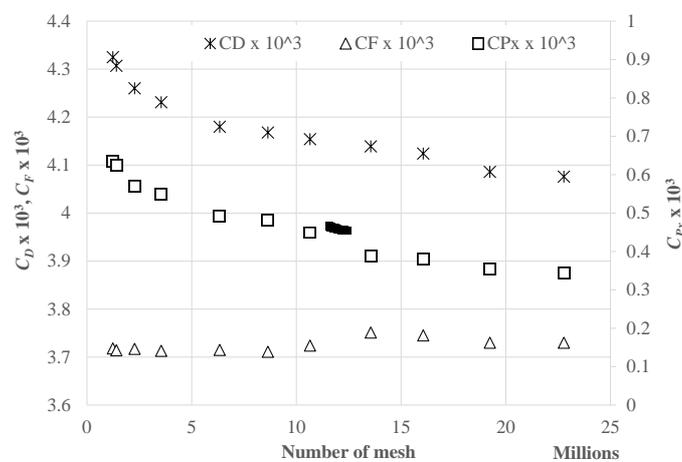


Figure 6: Mesh convergence of a pair of spheroids

Table 3: Computational parameters

Parameters	Setting
Global mesh size	0.20 m
Mesh type	Unstructured with local refinement around spheroids and in wake regions
$y^+$ average	30
No. of elements	8-15M with 15 prism layers in the boundary layer
Turbulence model	Shear Stress Transport
Inlet turbulent intensity	1%
Wall modelling	Automatic Wall Function (ANSYS, 2010)
Spatial discretisation	High Resolution (ANSYS, 2010)
Timescale control	Auto Timescale (ANSYS, 2010)
Convergence criteria	RMS residual $< 10^{-6}$
Computing	IRIDIS 3 Linux Cluster
Run type	Parallel (12 partitions run on 4×Dual core nodes, each with 2GB RAM)
Simulation Time	2.0-2.5 wall clock hours

## 4. Results

### 4.1. Single vehicle

Table 4 details the thrust deduction calculated for a single hull. The propeller revolution speed of each self-propelled AUV is iterated until the error between the self-propelled thrust ( $T_{SP}$ ) and the self-propelled drag ( $T_D$ ) are within 2%. Assuming no swirl is imparted to the flow ( $K_Q = 0$ ) results in a predicted thrust deduction of  $t = 0.13$ . Inclusion of a representative torque results in a prediction of  $t = 0.14$ . The numerical predictions of thrust deduction ( $t$ ) are within the ranges suggested in Burcher's experiment (Burcher and Rydill, 1994).

Table 5 shows a breakdown of the pressure and skin friction components of the drag for cases S1, S2 and S3. The thrust deduction is driven by a change to the pressure drag. Figure 7 shows an accelerated flow regime behind the propeller plane B1. These cases demonstrate the importance of considering the self-propelled drag when assessing total energy budget of an AUV.

Table 4: Thrust deduction and drag results of a single towed hull and a single self-propelled AUV with two different propulsive conditions

Body Force Model	Towed drag, $R$ (N)	Self-propelled thrust, $T_{SP}$ (N)	Self-propelled drag, $T_D$ (N)	Thrust deduction, $t$
S1: Towed hull (No propeller model)	2.3014	-	-	-
S2: Thrust Only Propeller	-	2.6393	2.6426	0.13
S3: Thrust and Torque Propeller	-	2.6830	2.6874	0.14

Table 5: Drag coefficient of a single towed hull and a single self-propelled AUV at two different propulsive conditions. Define the percent drag difference compare to case S1 as  $\% \Delta_{to S1} = \frac{drag_i - drag_{(S1)}}{drag_{(S1)}}$ , where  $i$  represents cases S2 and S3

Cases	Total drag coef. $C_D \times 10^3$	$\% \Delta_{(S1)}$	Skin friction drag coef. $C_F \times 10^3$	$\% \Delta_{(S1)}$	Pressure drag coef. $C_{Px} \times 10^3$	$\% \Delta_{(S1)}$
S1: Towed hull	4.053	-	3.643	-	0.347	-
S2: Thrust Only Propeller	4.632	14.7	3.647	0.1	1.012	191.6
S3: Thrust and Torque Propeller	4.709	16.6	3.648	0.1	1.089	213.8

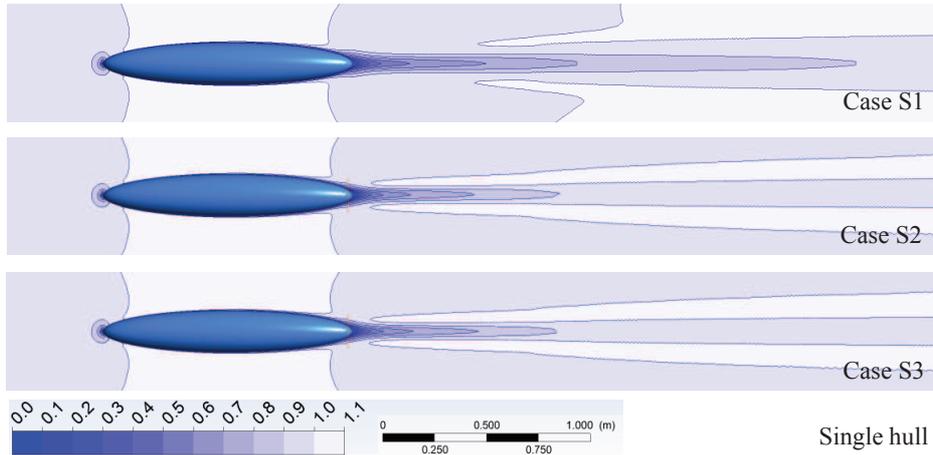


Figure 7: Comparison of velocity profile  $V/V_\infty$  of fluid inflow past a single AUV; (Top) Case S1: Towed hull, (Middle) Case S2: Self-propelled AUV by thrust only propeller and (bottom) Case S3: Self-propelled AUV by thrust and torque propeller.

### 4.2. Velocity profile of drafting twin vehicles at $D/L = 1.47$

The flow velocity past two drafting hulls ( $S/L = 0$  and  $D/L = 1.47$ ) is demonstrated in Figure 8. Three cases are shown:

- case F1 where both hulls are towed (dotted line)

- case F2 where the leader B1 is the self-propelled vehicle and the follower B2 is towed (dashed line)
- case F3 where both vehicles are self-propelled (continuous line)

From a uniform inflow ( $D/L = -1.250$ ), the boundary layer develops over B1 ( $0 < D/L < 1.0$ ). At the propeller plane of B1 ( $D/L = 0.994$ ), a difference may be observed between cases F1 and F2&F3. For the towed case, F1, the advance velocity at the propeller plane is lower than the free-stream velocity due to the presence of the boundary layer the reduction in local velocity corresponds to a mean wake fraction ( $\bar{w}_i$ ) of 0.148 (calculated by Equation 11). Cases F2 and F3 experience an increase in the axial momentum on the propeller plane of B1, which is due to the action of the propeller. This axial increment increases with the slipstream contraction at  $D/L = 1.235$ .

The inflow to B2 ( $D/L = 1.470$ ) differs for the towed B1 (F1) where there is a wake deficit and the self-propelled B1 (F2&F3) where the propeller race has re-energised the wake, however, this is a non-uniform inflow. For all cases the boundary layer then develops around B2 at  $1.470 < D/L < 2.463$ . At the propeller plane of B2 ( $D/L = 2.463$ ), the mean wake fraction for B2 for case F2 is 0.143.

Downstream at  $D/L = 4.789$ , the three cases show different behaviour where F1 exhibits the largest wake deficit, F3 the smallest. Case F1 exhibits a wake deficit due to both vehicles' drag, F2 exhibits a wake deficit due to B2's drag and the re-energised wake of the propeller 1, and F3 exhibits a nearly completely uniform flow since both vehicles are self-propelled.

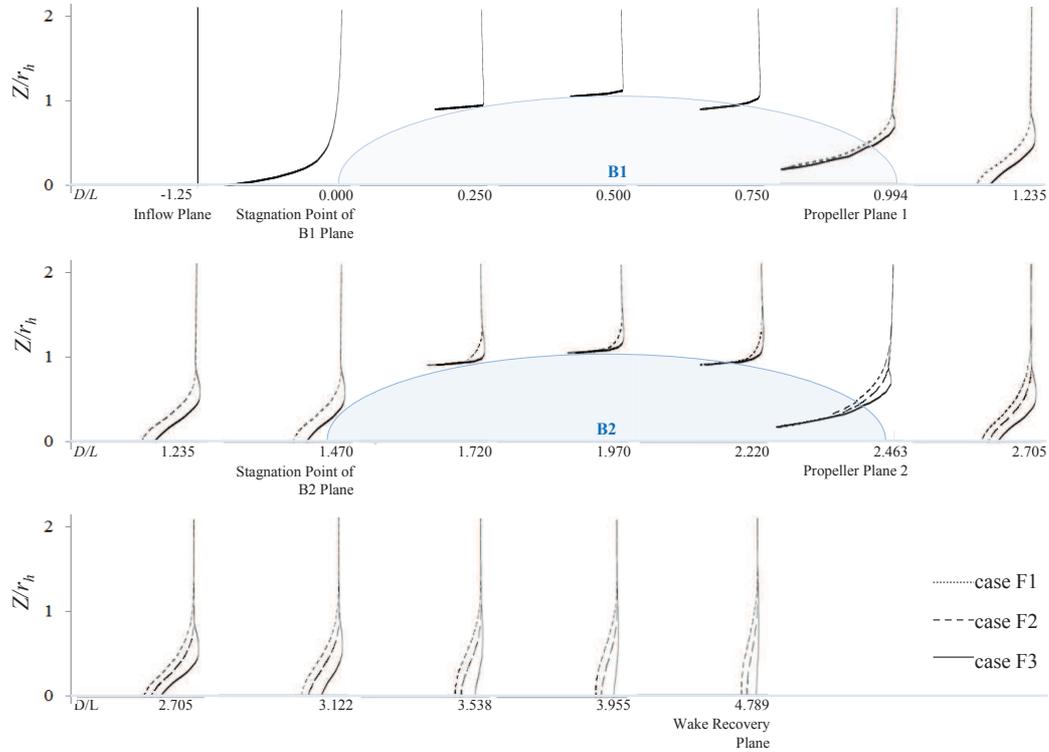


Figure 8: The velocity profile of the flow past two in-line hulls ( $S/L = 0, D/L = 1.47$ ). case F1 where both hulls are towed (dotted line), case F2 where the leader (B1) is the self-propelled vehicle and the follower (B2) is towed (dashed line), and a final case F3 where both vehicles are self-propelled (continuous line)

The influence of the propeller model on the self-propelled drag, the thrust deduction and the propeller race deduction are demonstrated in Figure 9 for a longitudinal separation of  $D/L = 1.47$ . The evaluated thrust deduction,  $t$ , of each case is within the empirical range of 0.13-0.14.

Starting with the twin towed case F1, the drag of B1 is reduced compared to a single towed vehicle. This is due to the bow stagnation of B2 aiding the pressure recovery at the stern of B1. For cases F2 to F6, B2 experiences higher drag than case F1 due to the propeller race at B1. This can be captured by using a propeller race deduction,  $\gamma$ . Thus the towed drag of the follower vehicle B2 (cases F2 and F4) may be calculated from:-

$$R_{B2}(\text{towed B1}) = (1 - \gamma) R_{B2}(\text{self-propelled B1}) \quad (15)$$

Hence, for a self-propelled follower B2 of cases F3, F5 and F6:-

$$R_{B2}(\text{towed B1}) = (1 - \gamma - t) T_{SP, B2}(\text{self-propelled B1}) \quad (16)$$

Comparing the results with and without representative torque, it is evident that the momentum source terms which represent propeller thrust are governing the predicted thrust deduction and the propeller race deduction.

Compare case  $i$  to case  $j$

$$\% \Delta = \frac{\text{drag}_i - \text{drag}_j}{\text{drag}_j} \times 100$$

Case	B1	B2	Comparison	$\% \Delta_{B1}$	$\% \Delta_{B2}$
S1	2.3014 / ---- / ----				
S2	2.6393 / 0.13 / ----	$K_T = 0.0201$	S2 to S1	14.7	
S3	2.6830 / 0.14 / ----	$10K_Q = 0.0215$ $K_T = 0.0203$	S3 to S1	16.6	
F1	2.2100 / ---- / ----	2.1305 / ---- / ----	F1 to S1	-4.0	-7.4
F2	2.5443 / 0.13 / ----	2.4583 / ---- / 0.13	F2 to F1	15.1	15.4
F3	2.5445 / 0.13 / ----	2.7927 / 0.11 / 0.13	F3 to F1	15.1	31.1
F4	2.5821 / 0.14 / ----	2.4609 / ---- / 0.13	F4 to S3	-3.8	6.9
F5	2.5824 / 0.14 / ----	2.7938 / 0.11 / 0.13	F5 to F4	0.0	13.5
F6	2.5824 / 0.14 / ----	2.8510 / 0.12 / 0.13	F6 to F4	0.0	15.9
			F6 to F3	1.5	2.1
			F6 to F1	16.9	33.8

Figure 9: Drag force results and the percentage drag differences in the drafting configuration at  $D/L = 1.47$ , where *drag* represent the towed drag ( $R$ ) for a towed case and the self-propelled thrust ( $T_{SP}$ ) for any propulsive conditions,  $t$  is the thrust deduction and  $\gamma$  is defined as the propeller race deduction.

#### 4.3. Impact of spacing on self-propelled vehicles

The results of the self-propelled thrust ( $T_{SP}$ ), thrust deduction ( $t$ ), propeller race deduction ( $\gamma$ ),  $K_T$  and  $10K_Q$  for a pair of self-propelled AUVs over a range of transverse and longitudinal spacings is shown in Figure 10. In a previous study of towed pairs of prolate spheroids, Rattanasiri et al. (2014) defined seven operating regions depending on the changes to individual drag and the total fleet, see Figure 11, where  $\%B1$  and  $\%B2$  is defined as the percentage difference of the individual drag of B1 and B2 referenced to a single towed hull drag, respectively.  $\%CB$  is defined as the combined drag which considered as a percentage difference of the fleet drag referenced to the sum of two single towed hull drags. Using a similar approach the results for the self-propelled case are presented in Figure 12. Where  $\%B1$  and  $\%B2$  is defined as the percentage difference of the individual self-propelled drag of B1 and B2 referenced to a single self-propelled resistance ( $T_D$ ), respectively.  $\%CB$  is defined as the combined self-propelled drag which considered as a percentage difference of the fleet drag referenced to the sum of two single self-propelled hull drags. Where  $T_{D(B1)}$  and  $T_{D(B2)}$  is the self-propelled drag of B1 and B2 in fleet, respectively. The sign + and - indicates increasing and decreasing of the towed and self-propelled drag, respectively.

$$\begin{aligned} \%B1 &= \frac{T_{D(B1)} - T_{D(s)}}{T_{D(s)}} \times 100, \\ \%B2 &= \frac{T_{D(B2)} - T_{D(s)}}{T_{D(s)}} \times 100, \\ \%CB &= \frac{T_{D(B1)} + T_{D(B2)} - 2T_{D(s)}}{2T_{D(s)}} \times 100 \end{aligned} \quad (17)$$

By varying the spacing of the pair, the results suggest some energy beneficial configurations exist for an individual but only very small gains are possible for the fleet. However, some configurations lead to a significant increase in the drag of both individual members and the fleet.

#### 4.3.1. Parallel region

The parallel region is where both hulls are parallel to the inflow, nose aligned with zero longitudinal offset ( $D/L = 0$ ), both hulls experience a propeller race deduction due to their close proximity; the propeller generates additional acceleration of the flow between the hulls. A comparison of velocity profile,  $V/V_\infty$ , for the flow past a pair of AUVs at  $S/L = 0.17$  and  $D/L = 0$  is shown in Figure 13. In an infinite domain, both hulls are expected to experience an equal propeller race deduction, however, due to the asymmetric finite numerical fluid domain, the tunnel wall results in slightly differing self-propelled thrust, thrust deduction and the propeller race deduction. As the transverse separation increases up to  $0.37L$ , the increase in accelerated flow due to the action of propeller reduces, resulting in no propeller race deduction. These results and the previous towed studies (Rattanasiri et al., 2014) suggested that the distance of  $0.5L$  is required as the minimum transverse separation to exclude body-to-body interaction and the propeller race deduction.

#### 4.3.2. Echelon region

With a transverse separation and a limited longitudinal offset within one body length, the configuration is similar to echelon or part of a vee formation. Figure 12 shows that while both hulls experience interaction, the follower B2 experiences a drag reduction while the leader B1 experiences a drag augment, due to the body-to-body interaction Rattanasiri et al. (2014). The comparison of results at  $S/L = 0.17$  and  $D/L = 0.57$  configuration as shown in Figure 14, implied that to maintain the speed of the follower in the fleet, the propeller has to operate to pull the AUV backward. As the transverse separation increases, the interactions between the hulls reduce. The information at  $S/L = 0.17$  and  $D/L = 0.27$  configuration is not available due to the high interaction between hulls which leads to the difficulty of obtaining the self-propelled follower.

#### 4.3.3. Push region

The push region is where the nose of follower is positioned close to the trailing edge of the leading AUV. The results show a significant decrease in the self-propelled drag of the leader, similar to the results for the towed case. A comparison of velocity profiles is shown in Figure 15. Vessels placed directly downstream of a propeller experience a significant propeller race deduction. A small transverse separation is sufficient to place the follower to one side of the propeller race where it experiences minimal propeller race deduction.

#### 4.3.4. Drafting region

By placing a towed B2 directly behind a towed B1 for  $D/L > 1.37$ , both individual AUVs and the fleet experience drag reduction (Rattanasiri et al., 2014), however this benefit has been removed by the action of the propeller for the self-propelled cases. Being placed in propeller race, the follower experiences a propeller race deduction of  $\gamma = 0.13$ . A comparison of velocity profiles is shown in Figure 16.

#### 4.3.5. No propeller race interaction zone

From Figure 10, there is minimal effect of propeller race deduction on the leader B1 or follower B2. A comparison of the velocity profile of a pair of AUVs at  $S/L = 0.17$  and  $D/L = 1.47$  configuration, Figure 17, shows a very small accelerated flow regime by propellers from leader B1 interact the flow regime around the follower B2. It could be concluded that the bigger the transverse separation and longitudinal offset, the smaller the accelerated flow regime between hulls.

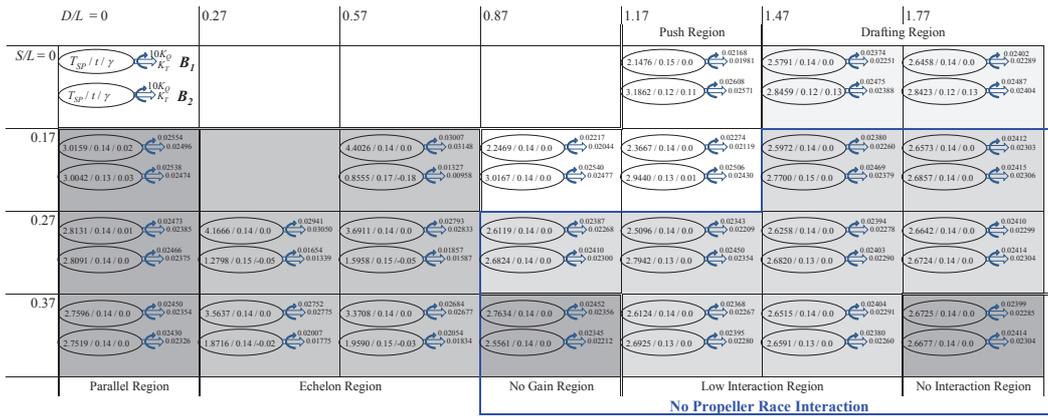


Figure 10: The results of individual self-propelled thrust ( $T_{SP}$ ), thrust deduction ( $t$ ), propeller race deduction ( $\gamma$ ),  $K_T$  and  $10K_Q$  at various spacings

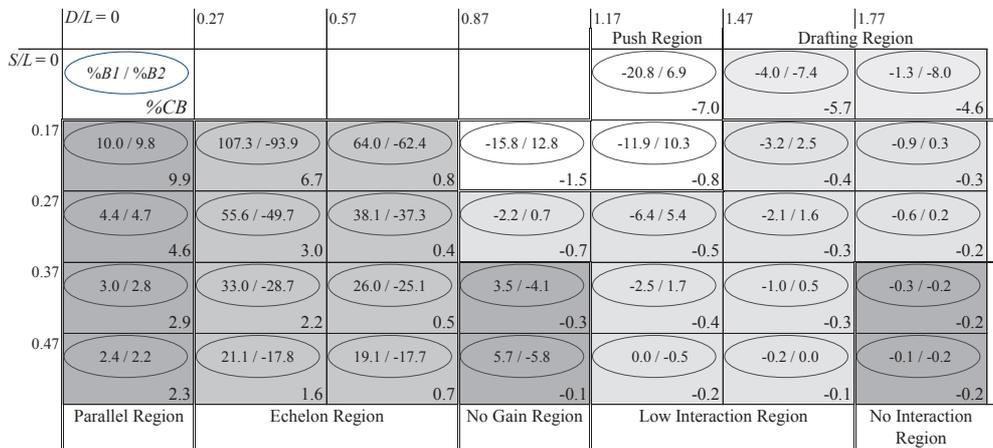


Figure 11: The effect of towed twin hulls' spacing on individual drags and combined drag of B1 and B2 at  $Re = 3.2 \times 10^6$  (Rattanasiri et al., 2014)

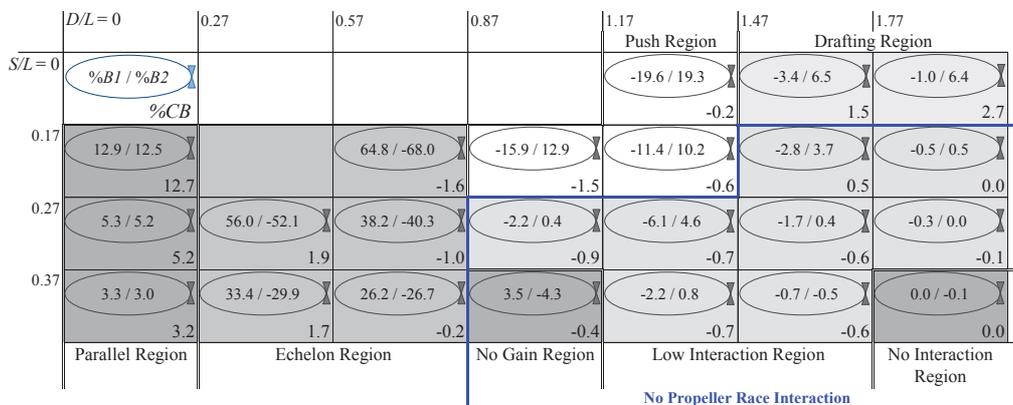


Figure 12: The effect of twin hulls' spacing on individual self-propelled thrust and combined drag of B1 and B2

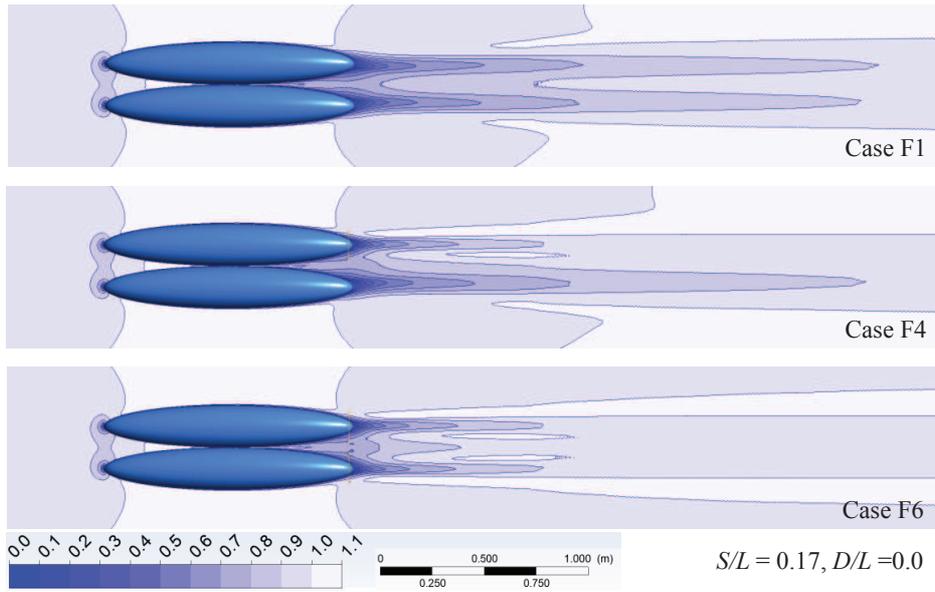


Figure 13: Comparison of velocity profile  $V/V_\infty$  of fluid inflow past a pair of AUVs at  $S/L = 0.17$  and  $D/L = 0.0$  configuration; (Top) Case F1: towed B1 and towed B2, (Middle) Case F4: self-propelled B1 and towed B2 (bottom) Case F6: self-propelled B1 and self-propelled B2.

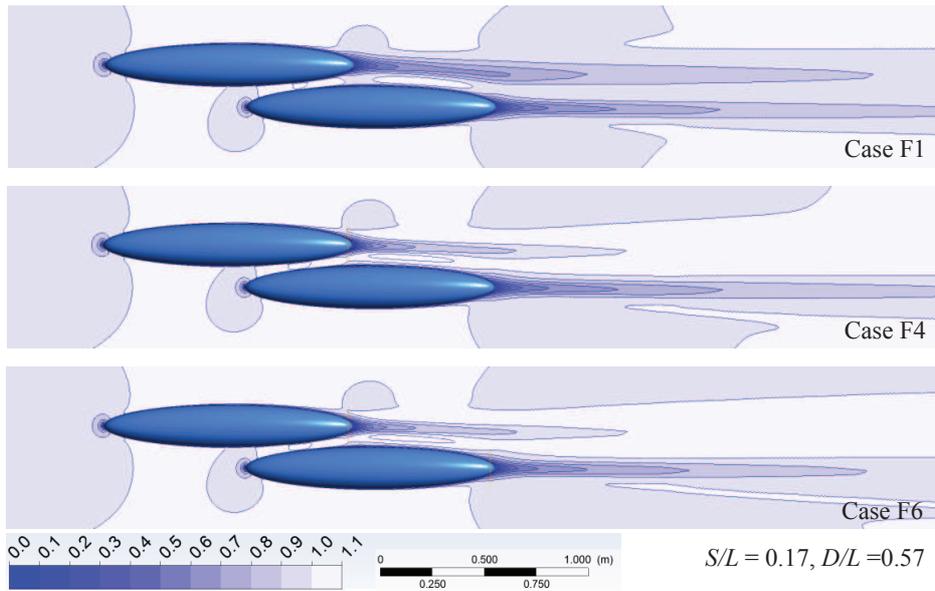


Figure 14: Comparison of velocity profile  $V/V_\infty$  of fluid inflow past a pair of AUVs at  $S/L = 0.17$  and  $D/L = 0.57$  configuration; (Top) Case F1: towed B1 and towed B2, (Middle) Case F4: self-propelled B1 and towed B2 (bottom) Case F6: self-propelled B1 and self-propelled B2.

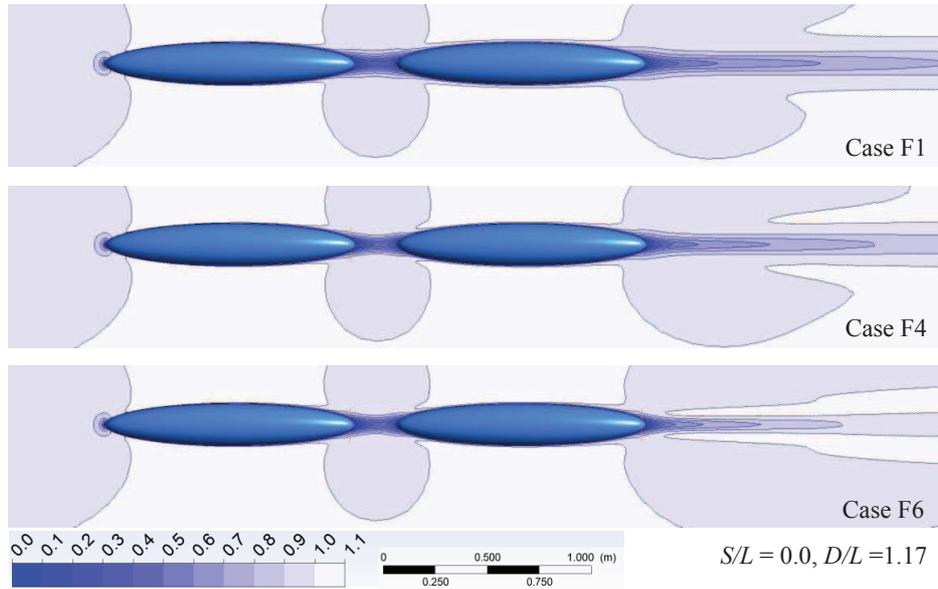


Figure 15: Comparison of velocity profile  $V/V_{\infty}$  of fluid inflow past a pair of AUVs at  $S/L = 0.0$  and  $D/L = 1.17$  configuration; (Top) Case F1: towed B1 and towed B2, (Middle) Case F4: self-propelled B1 and towed B2 (bottom) Case F6: self-propelled B1 and self-propelled B2.

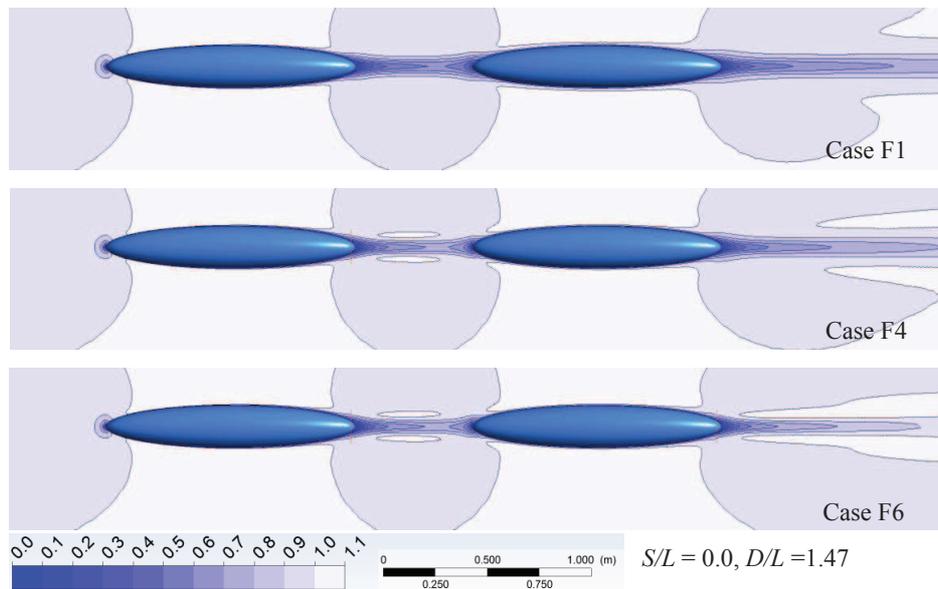


Figure 16: Comparison of velocity profile  $V/V_{\infty}$  of fluid inflow past a pair of AUVs at  $S/L = 0.0$  and  $D/L = 1.47$  configuration; (Top) Case F1: towed B1 and towed B2, (Middle) Case F4: self-propelled B1 and towed B2 (bottom) Case F6: self-propelled B1 and self-propelled B2.

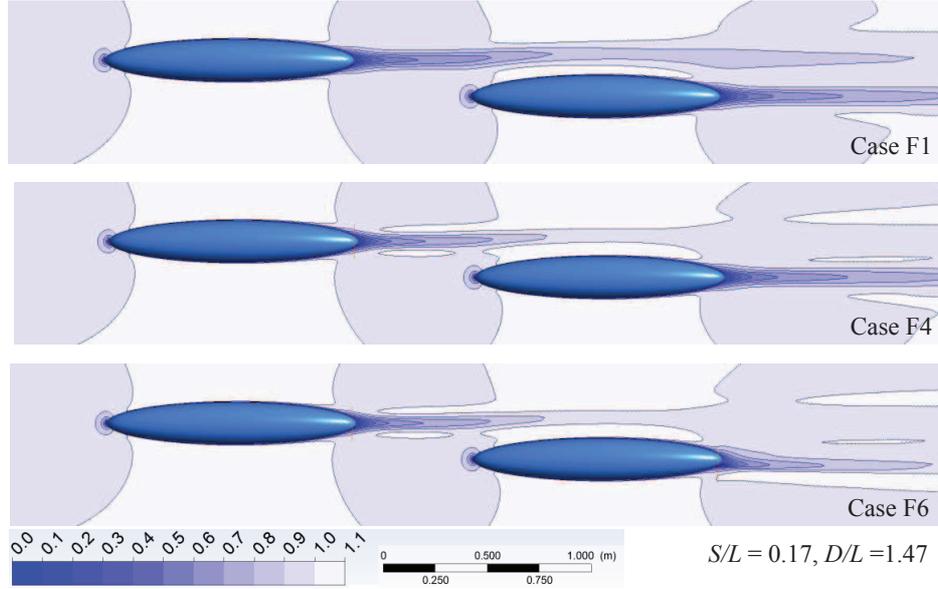


Figure 17: Comparison of velocity profile  $V/V_\infty$  of fluid inflow past a pair of AUVs at  $S/L = 0.17$  and  $D/L = 1.47$  configuration; (Top) Case F1: towed B1 and towed B2, (Middle) Case F4: self-propelled B1 and towed B2 (bottom) Case F6: self-propelled B1 and self-propelled B2.

#### 4.4. Twin vehicles at various drafting and parallel configurations

Thus we can conclude that the parallel, push and the drafting regions experience the greatest interactions, these will be examined in more detail in this section.

##### 4.4.1. Parallel configurations

Figure 18 illustrates the influence of the transverse separation in the parallel configuration on the components of drag. Implementing the propeller on both hulls increases the drag on both hulls. For case F1, at  $S/L = 0.17$ , the total drag of both towed vehicles is increased by 10%, it is due to a change in the pressure drag rather than a skin friction drag. As the separation increases, the drag of both hulls tends towards the drag of a single free stream towed resistance.

For cases F2, the total drag of self-propelled B1 is increased by 15% compared to the towed B1 (F1) and then tends toward the single self-propelled thrust as the separation increases. From Figure 18b and 18c, at the separation within  $0.27L$ , the propeller race deduction results from the self-propelled B1 is driven the pressure drag of towed B2 to be increased by 15% when compared to the towed B2 (F1), consequently increase a 4% of its total drag.

For case F3, at  $S/L = 0.17$ , the self-propelled B1 (F3) experiences a 4% drag augment compared to that of F2. This suggests that the propeller race deduction from both self-propelled AUVs interacts with each other, subsequently disturbing the pressure distribution around the hulls. For  $S/L > 0.30$ , the drag of both self-propelled vehicles tends towards the free stream thrust requirements.

##### 4.4.2. Drafting configurations

Figure 19 illustrates the influence of the longitudinal offset in the drafting configuration on the components of drag. Starting with case F1, at  $D/L = 1.17$ , the total drag of B1 is increased by 8% whilst the total drag of B2 is reduced by 21% (Figure 19a). As the longitudinal offset increases, the drag of both towed hulls tends towards the free stream drag as wake recovery occurs. This hydrodynamic behaviour is similar to cars in a convoy with a slipstream at  $S/L > 1.5$  (Hucho and Ahmed, 1998). This goes some-way to explain the energy benefit of cycling and swimming in the slipstream (Kyle, 1979; Silva et al., 2008). According to Figure 19b and 19c, the discrepancy of the total drag of both hulls is driven by a change to the pressure drag. Significantly, a change to the skin friction drag of B2 occurs when placed directly behind towed B1, this dominates the B2's drag reduction. Then at  $D/L = 4.47$ , the total drag of both hulls tends towards the single free stream towed resistance.

For case F2, at  $1.17 \leq D/L \leq 4.47$ , the total drag of self-propelled B1 is increased by 12% compared to that of towed B1 (F1). The drag of self-propelled B1 tends towards the single vehicle self-propelled thrust as the offset

increases. Considering Figure 19b and 19c, at  $D/L = 1.17$ , the drag of B2 is increased by 12% compared to that of towed B2 (F1), this is due to an addition to the skin friction drag results from a propeller race deduction. As the offset increases, the extra skin friction drag is reduced which results in a decrease of the total drag augment towards the single towed resistance at  $D/L > 5.0$ .

For case F3, the drag of self-propelled B1 (F3) shows the same results as that of case F2. Apparently, placing either a towed B2 (F2) or self-propelled B2 (F3) directly behind the self-propelled B1 has no impact to the B1's drag. At  $D/L = 1.17$ , the drag of the self-propelled B2 (F3) is increased by 18% compared to the towed B2 (F2). The drag of self-propelled B2 (F3) is then reduced towards the free stream thrust when the offset increases. Figure 19b and 19c show that a change of the total drag augment is influenced by the pressure drag rather than the skin friction drag.

For these drafting cases, in general, for vehicles following a self-propelled leader within  $5.0L$ , the skin friction drag of the follower is increased by the propeller race deduction, consequently, a rise in the total drag. For longitudinal offsets  $D/L > 5.0$ , the drag of both B1 and B2 tend towards the single free stream value, which are at a towed resistance and a self-propelled thrust for the towed vehicle and self-propelled vehicle, respectively.

For all drafting and parallel cases, with both vehicles in close proximity, significant change in the pressure drag dominates an increment in the total drag. The propeller race deduction influences a significant change in the skin friction drag of any vehicle placed directly behind a self-propelled vehicle. On the other hand, in very close proximity in the parallel configuration, the propeller race drives a 25% increase in the pressure drag for both vehicles; it is evident that a change of the skin friction drag has a minimal change of less than 1%.

## 5. Conclusion

This paper investigates the influence of the propeller race on upstream and downstream self-propelled AUVs. Initially simulations of a self-propelled hull are performed, utilising body force model to replicate the impact of the propeller by modelling the momentum source terms. This is then extended to consider a fleet of two self-propelled vehicles operating at a range of longitudinal offset and transverse separation.

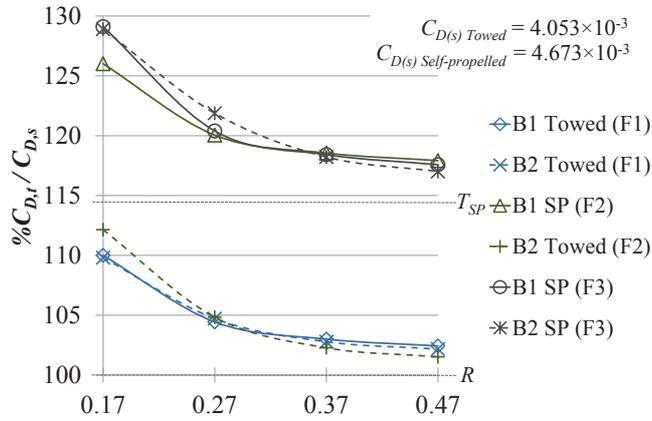
The results highlight that operation in close proximity to another self-propelled vessel has a significant impact of both the flow around the hull and drag experienced by the vehicle. A propeller race deduction is proposed to account for the increase in vehicle drag due to the propulsors of other vehicles. The propeller race deduction is dependent on both longitudinal and transverse separation.

The self-propelled results using a virtual propeller show no significant propulsive energy benefit to the total fleet energy budget. In-line operation which has previously been shown to be beneficial for towed vehicles is shown to give no or minimal benefit to self-propelled vehicles due to the upstream propeller re-energising the wake. An echelon configuration does allow a transfer of drag from the follower to the leader, but would require additional course keeping due to the hydrodynamic moment induced on the vehicle. A parallel configuration results in an increased propulsion cost for both vehicles. However, based on this numerical information, operators can determine the optimal configurations in transverse separation and longitudinal offset based on energy considerations. However, this information is based on a homogeneous fleet of torpedo style vessels, by adopting a heterogeneous fleet or a non-conventional body shape it may be possible to reduce the resistance of following vessels if they are designed such that the follower(s) operate in the wake of a leading vehicle but outside of its propeller race.

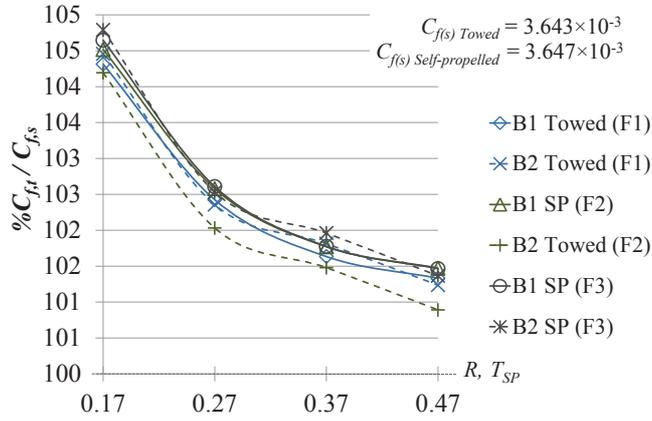
From a vehicle or mission design perspective, it is important to correctly understand the true propulsive energy budget of the vehicle and its impact on both range and endurance. This study highlights the importance of considering both thrust deduction and any propeller race deductions when calculating the propulsive power consumption of an individual or fleet of vehicles.

## Acknowledgment

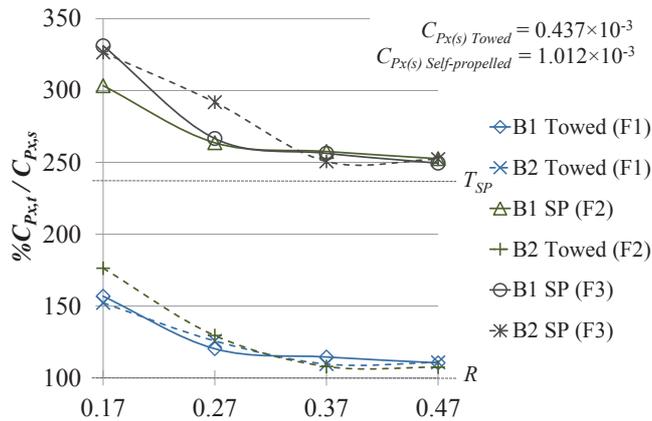
The authors acknowledge the use of the IRIDIS High Performance Computing Facility, and associated support services at the University of Southampton, in the completion of this work. The PhD studentship of Rattanasiri was financed by the Royal Thai Government.



(a) Total drag coefficient

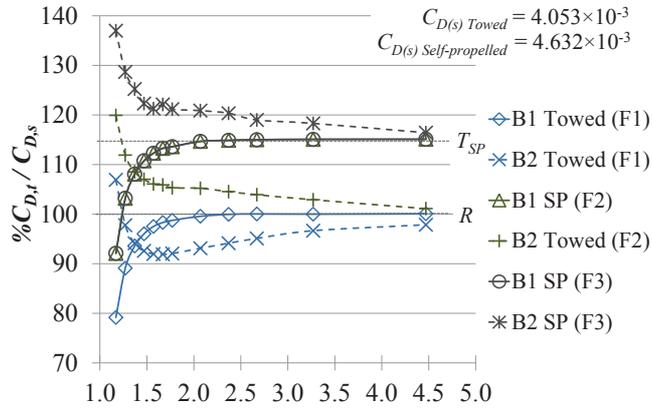


(b) Skin friction drag coefficient

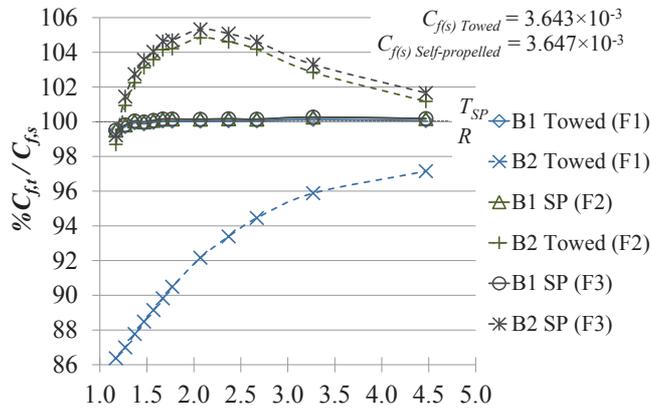


(c) Pressure drag coefficient

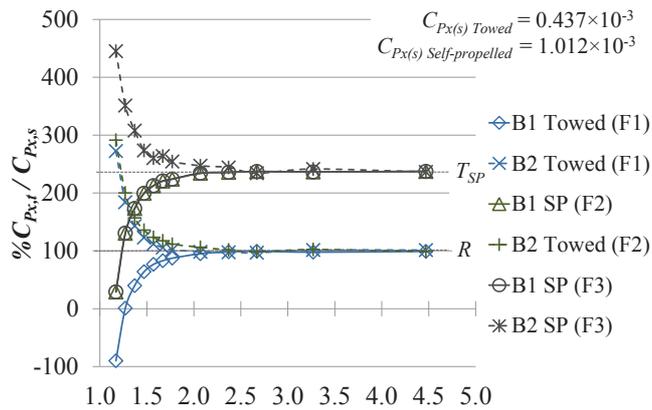
Figure 18: Effect of propeller on forces acting on the parallel configuration B1 and B2 at various transverse separations



(a) Total drag coefficient



(b) Skin friction drag coefficient



(c) Pressure drag coefficient

Figure 19: Effect of propeller on forces acting on drafting configuration B1 and B2 at various longitudinal offsets

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## Nomenclature

$1 + k$	The form factor of a single spheroid
$1 + \beta k$	The form factor of twin spheroids
$A_w$	Wetted surface area (m <sup>2</sup> )
B1 and B2	Leader spheroid and follower spheroid
%B1 and %B2	Individual drag of B1 and B2 referred to a single hull drag
%CB	Combined drag refer to sum of two single hull drags
$C_F$	Skin friction drag coefficient, $\frac{\text{Skin friction drag}}{0.5 \rho V^2}$
$C_{P_x}$	Pressure drag coefficient, $\frac{\text{Pressure drag}}{0.5 \rho V^2}$
$C_{th}$ and $K_T$	Thrust coefficient
$d_m$	Maximum diameter of the body of revolution (m)
$D$	Longitudinal offset (m)
$D/L$	Non-dimensional longitudinal offset
$d_h$	Maximum hull diameter (m)
$F_i$	The momentum source terms ( $\bar{F}_x, \bar{F}_y, \bar{F}_z$ )
$\bar{F}_x$	Non-dimensional axial momentum source term
$\bar{F}_\theta$	Non-dimensional tangential momentum source term
$\bar{F}_r$	Non-dimensional radial momentum source term
$J$	Advance coefficient
$K_Q$	Torque coefficient
$n$	Revolutions per second of propeller (s <sup>-1</sup> )
$S/L$	Transverse separation
$t$	Thrust deduction
$T_D$	The self-propelled drag (N)
$T_{D(B1)}$ and $T_{D(B2)}$	Self-propelled drag of B1 and B2 in fleet
$T_{SP}$	The require thrust (N)
$L$	Length of the body from nose to tail (m)
$R$	Towed drag (N)
$R_{(s)}$	towed drag of a single hull
$Re$	Length Reynolds number, $\frac{VL}{\nu}$
$R_p$ and $R_h$	Radius of propeller and hub (m)
$r_p$	Local propeller radius (m)
$V$	Vehicle speed (m/s)
$V_a$	Advance speed (m/s)
$U_i$	Cartesian mean velocity components ( $\overline{U_x}, \overline{U_y}, \overline{U_z}$ )
$w_t$	Wake fraction
$\bar{w}_t$	Mean wake fraction over propeller disc
$x_i$	Represents Cartesian co-ordinates ( $X, Y, Z$ ) (m)
$\nu$	Fluid kinematic viscosity, $\mu/\rho$ (m <sup>2</sup> /s)
$\rho$	Fluid density (kg/m <sup>3</sup> )
$\mu$	Fluid dynamic viscosity (kg/m.s)
$k$	Fluid turbulent kinetic energy (m <sup>2</sup> /s <sup>2</sup> )
$\epsilon$	Rate of dissipation of turbulent energy (m <sup>2</sup> s <sup>-3</sup> )
$\bar{\nu}$	Viscosity like variable
$\gamma$	The propeller race deduction
$\Delta$	The difference between drags