Optimizing Oxygen Reduction Catalyst Morphologies from First Principles

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Supporting Information

ABSTRACT: Catalytic activity of perovskites for oxygen reduction (ORR) was recently correlated with bulk d-electron occupancy of the transition metal. We expand on the resultant model, which successfully reproduces the high activity of LaMnO3 relative to other perovskites, by addressing catalyst surface morphology as an important aspect of the optimal ORR catalyst. The nature of reaction sites on low index surfaces of orthorhombic (Pnma) LaMnO3 is established from First Principles. The adsorption of O2 is markedly influenced by local geometry and strong electron correlation. Only one of the six reactions sites that result from experimentally confirmed symmetry-breaking Jahn–Teller distortions is found to bind O2 with an intermediate binding energy while facilitating the formation of superoxide, an important ORR intermediate in alkaline media. As demonstrated here for LaMnO3, rational design of the catalyst morphology to promote specific active sites is a highly effective optimization strategy for advanced functional ORR catalysts.

INTRODUCTION

Perovskites have been studied for their activity toward the oxygen reduction reaction (ORR) since the 1970s.1–3 Interest in alkaline fuel cells (AFCs) has been renewed in the past decade because of the development of suitable hydroxide-exchange membrane materials allowing the alkaline equivalent of the proton exchange membrane fuel cell (PEMFC).4 In AFC systems low-cost transition metals (TMs) both are stable and catalyze the ORR well, in contrast to the acidic PEMFC where only expensive noble metal catalysts show suitable activity and stability in spite of decades of research.5,6 LaMnO3 has been identified to be among the most catalytically active perovskites, and recent experiments on a range of perovskites suggest that the high performance is linked to the d4 character of the Mn3+ ion. This has been attributed to the single Mn eg electron allowing for an interaction with O2 molecules that is neither too strong nor too weak.7,8

The strength of the TM–O2 interaction has a significant effect on the activity of perovskite catalysts for the ORR. Previous analysis of the reaction mechanism has focused on the bulk electronic structure and simple models based on the covalent interaction of metals ions with adsorbed species. These models do not give full consideration to the important role of electron correlation in generating strong Jahn–Teller distortions of the MO6 octahedra in the active d4 (and d5) perovskites. This gives rise to a symmetry breaking orbital ordering of the bulk crystal and, as will be shown below, significant variations in surface reaction sites.

There have been previous studies of the adsorption of oxygen on the surface of LaMnO39–11 but the majority of studies has focused on solid oxide fuel cell (SOFC) applications and therefore consider only the highly symmetric, high temperature, cubic phase. For the lower operating temperature of AFCs, it is the orbitally ordered orthorhombic phase that is relevant.12–14 The schematic in Figure 1 shows that the Jahn–Teller distortion of the low-temperature phase leads to a set of crystallographically distinct surfaces, which are likely to exhibit different reactivities. The manufacturing procedure and resulting surface terminations have also shown to play a crucial role in determining catalytic activity.15 We use first-principles calculations, which reproduce the effects of strong electronic correlation, to demonstrate that local Jahn–Teller (J–T) distortions play an important role in determining the interaction of O2 molecules with LaMnO3 surfaces. The insight gained is of direct relevance to the design of new catalysts.

Previous studies of O2 adsorption on orthorhombic (Pnma) LaMnO3 have considered the (010) and (001) surfaces.10,16,17 The studies focus on a range of adsorption modes and binding sites, again with consideration to the SOFC reaction mechanism.10,16 Among the likely reaction sites in an AFC

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catalyst under reaction conditions as they provide examples of that prevail in the cubic high-temperature phase. As the O2 (equatorial 1.92/1.99 Å) relative to the undistorted octahedra contacts that are either elongated (apical 2.31 Å) or shortened orthorhombic LaMnO3 has yet to be fully explored, we provide listed in Table 1 and characterized by the surface on which they

Order of Increasing Surface Formation Energies (J)

Table 1. TM Adsorption Sites Available in Each Primitive Unit Cell of the Nonpolar, Stoichiometric, and Unreconstructed Low-Index Surfaces of LaMnO3 (Given in Order of Increasing Surface Formation Energies (E))^{16}

<table>
<thead>
<tr>
<th>plane</th>
<th>E\textsubscript{s} relaxed (J/m\textsuperscript{2})</th>
<th>Mn sites</th>
</tr>
</thead>
<tbody>
<tr>
<td>(100)</td>
<td>0.98</td>
<td>2 MnO\textsubscript{6}</td>
</tr>
<tr>
<td>(001)</td>
<td>1.39</td>
<td>2 MnO\textsubscript{5}</td>
</tr>
<tr>
<td>(101)</td>
<td>1.36</td>
<td>2 MnO\textsubscript{5}</td>
</tr>
<tr>
<td>(110)</td>
<td>1.98</td>
<td>1 MnO\textsubscript{5} 1 MnO\textsubscript{5}</td>
</tr>
</tbody>
</table>

The sites are J−T distorted octahedra (MnO\textsubscript{6}) from which O ions have been cleaved (MnO\textsubscript{5} or MnO\textsubscript{4}) in the formation of the surface. The superscripts Ap (apical) and Eq (equatorial) indicate the type of Mn−O bond cleaved from the octahedron to form the site.\textsuperscript{18}

occur, the number of O atoms cleaved to form the site, and the J−T distortion of the Mn−O bond cleaved when forming the site. These sites are indicative of those present in an active catalyst under reaction conditions as they provide examples of apically and equatorially distorted 4- and 5-coordinate Mn sites. The strength of O2 binding is relatively easily computed and is likely to be an excellent proxy for ORR activity as the formation of reduced O2 at the surface is known to be an important intermediate in the high pH environment of the AFC.\textsuperscript{19−21} The use of binding energies to determine reactivity has been exploited in previous studies where a middle ground is sought between overly strong or weak binding of adsorbates for higher catalytic activity.\textsuperscript{7,8} We also use the charge transfer to the adsorbed molecule as an indication of the likelihood of reaction steps that depend on the oxidation of the metal site and therefore assert that the combination of binding energy and charge transfer provides a strong indication of the catalytic activity of the sites considered.

RESULT

To verify the existence of J−T distortion at the orthorhombic LaMnO3 surface, TEM was performed on crystalline powder. A typical image is shown in Figure 2 in which the predicted J−T distortion of Mn−O octahedra in the bulk crystal is clearly visible and is perpetuated from the bulk to the surface. Heuristically, one would be inclined to propose a model for adsorption of O2 on these distorted surfaces where binding is stronger at lower coordinated Mn sites and at the shorter Mn−O bond provided by equatorial sites relative to the apical sites. While the former is confirmed by our detailed study, the complex nature of the interaction between O2 and the active sites leads to a nonintuitive dependence of adsorption strength with bond length.

The computed electronic and structural data for the adsorption of O2 at each site are presented together with binding energies in Table 2. Note that in the DFT calculations of the bulk-like regions of the LaMnO3 lattice the charge of an O ion estimated by Mulliken population analysis (MPA) is −1.3 lel. Since O\textsuperscript{2−} is the formal oxidation state in bulk LaMnO3, a value of −0.65 lel for the adsorbed molecule (O2 \textsubscript{q} in Table 2) indicates a formal charge of −1 lel. The computed occupations of the surface Mn 3d orbitals are documented in Table 3. These reflect changes in the oxidation state of the surface sites.

It is evident that O2 adsorption to MnO\textsubscript{3} sites results in strong change transfer to the molecule and oxidation of Mn
from $d^3$ to $d^1$. This is accompanied by a large surface relaxation that may be interpreted as an anti-J-T distortion as the $e_g$ orbital is no longer occupied. This is illustrated very clearly in the spin density at the surface sites before and after $O_2$ adsorption (Figure 3). The final geometry of each MnO$_5$ site is therefore rather similar, and the critical role that the apical and equatorial bond lengths of the initial geometry play in the adsorption energetics is not reflected in the final Mn−O$_2$ distances. This results in a rather unexpected trend in adsorption energy with Mn−O$_2$ distance, as typically one would expect stronger binding to be correlated with a shorter interatomic distance.

The differences in the initial coordination of the reaction sites, on the other hand, are reflected in the final Mn−O$_2$ distances. $O_2$ adsorption on the MnO$_4$ sites results in significantly lower $d$(Mn−O$_2$), as listed in Table 2. The end-on adsorption mode on the MnO$_5$ site proves to be an exception to this, as in this mode the adsorption results in formation of a superoxide (O$_2^-$) rather than the peroxide (O$_2^2-$). In the following we analyze the nature of the adsorption at each of the six surface reaction sites (Figures 4–7) in detail.

(100) MnO$_5^{6p}$. The lowest binding energy and charge transfer (−0.15 eV and 0.48 lel) are observed for $O_2$ adsorption on this site. Adsorption takes place at an elongated five-coordinate apical site, and the $d$(O−O) of 1.31 Å is the least perturbed from that computed for the free $O_2$ molecule. This is consistent with a general trend that $d$(O−O) increases in proportion to the binding energy.$^9,11,22$

(001) MnO$_5^{6p}$. This site is similar to that on the (100) surface but is formed from the cleavage of an equatorial bond. Adsorption is therefore far stronger with a binding energy of −0.61 eV, with significantly higher charge transfer and, as expected, a somewhat longer $d$(O−O) of 1.34 Å.

(101) MnO$_5^{6p}$ > La Bridge. The geometry of the (101) surface results in a unique adsorption mode for the $O_2$ molecule where it is not bound to the MnO$_5^{6p}$ site. Instead, as shown in Figure 6, the minimum energy geometry is in a bridging position between the surface La ions. The binding energy at this site is −1.17 eV compared to the −0.15 we observe for the MnO$_5^{6p}$ site. The higher binding energy and

Table 2. Computed Binding Energies (BEs) and Bond Distances ($d$, in Å) of $O_2$ Adsorption on the LaMnO$_3$ Surfaces, Indicating the Initial > Final Adsorption Site Where Different $q(O_2)$ is the Charge Transferred in lel to $O_2$ and Resulting Species (Where 0.65 lel by MPA is Equivalent to One Electron Based on the Bulk Charges)$^a$

<table>
<thead>
<tr>
<th>adsorption mode</th>
<th>d(Mn−O$_2$)</th>
<th>d(O−O)</th>
<th>BE (eV)</th>
<th>$q$(O$_2$)</th>
</tr>
</thead>
<tbody>
<tr>
<td>free $O_2$</td>
<td>-</td>
<td>1.23</td>
<td>-</td>
<td>0</td>
</tr>
<tr>
<td>(100) MnO$_5^{6p}$</td>
<td>2.00</td>
<td>1.31</td>
<td>-0.15</td>
<td>0.48 (O−O$^−$)</td>
</tr>
<tr>
<td>(001) MnO$_5^{6p}$</td>
<td>2.07</td>
<td>1.34</td>
<td>-0.61</td>
<td>0.61 (O−O$^−$)</td>
</tr>
<tr>
<td>(101) MnO$_5^{6p}$ &gt; La bridge</td>
<td>-</td>
<td>1.36</td>
<td>-1.17</td>
<td>0.74 (O−O$^−$)</td>
</tr>
<tr>
<td>(110) MnO$_4^{6p}$ &gt; MnO$_4^{6p}$</td>
<td>1.93</td>
<td>1.48</td>
<td>-2.04</td>
<td>1.24 (O−O$^−$)</td>
</tr>
<tr>
<td>(110) MnO$_4^{6p}$ &gt; MnO$_4^{6p}$</td>
<td>2.03</td>
<td>1.34</td>
<td>-2.32</td>
<td>0.70 (O−O$^−$)</td>
</tr>
<tr>
<td>(110) MnO$_4^{6p}$ &gt; MnO$_4^{6p}$</td>
<td>1.94</td>
<td>1.46</td>
<td>-2.67</td>
<td>1.04 (O−O$^−$)</td>
</tr>
</tbody>
</table>

$a_{eg}$ refers to the formal $e_g$ electrons at the active site before $O_2$ adsorption.

Table 3. Occupations of the Mn 3d Orbitals of the Adsorption Sites Estimated by MPA before and after $O_2$ Adsorption

<table>
<thead>
<tr>
<th>adsorption mode</th>
<th>initial 3d (lel)</th>
<th>final 3d (lel)</th>
</tr>
</thead>
<tbody>
<tr>
<td>(100) MnO$_5^{6p}$</td>
<td>3.75</td>
<td>3.06</td>
</tr>
<tr>
<td>(001) MnO$_5^{6p}$</td>
<td>3.75</td>
<td>2.91</td>
</tr>
<tr>
<td>(101) MnO$_5^{6p}$ &gt; La bridge</td>
<td>-</td>
<td></td>
</tr>
<tr>
<td>(110) MnO$_4^{6p}$ &gt; MnO$_4^{6p}$</td>
<td>3.69, 4.52</td>
<td>2.70, 3.77</td>
</tr>
<tr>
<td>(110) MnO$_4^{6p}$ &gt; MnO$_4^{6p}$</td>
<td>3.69, 4.52</td>
<td>3.21, 3.70</td>
</tr>
<tr>
<td>(110) MnO$_4^{6p}$ &gt; MnO$_4^{6p}$</td>
<td>3.69, 4.52</td>
<td>3.17, 3.99</td>
</tr>
</tbody>
</table>

Figure 3. Calculated spin density of (a) the unoccupied Mn$^{6p}$ site on the (100) surface and (b) after $O_2$ adsorption. (c) Model of the idealized $d_{xy}$, $d_{xz}$, $d_{yz}$, and $d_z^2$ suborbitals for reference.
charge transfer (−0.74 $\text{e}^-$) are reflected in the longer $d(\text{O}−\text{O})$ of 1.36 Å. In this case the charge is transferred from fully coordinate MnO$_6$ octahedra at the surface rather than the surface MnO$_5$ sites (Figure 6, charge transfer quantified in Supporting Information Table S2). The adsorption geometry and charge transfer here suggest an alternative pathway for the ORR, via electron transfer from the surface rather than from a specific site, which has not been discussed in previous literature. Studies of O$_2$ adsorption on the La site have only been carried out for the high-temperature cubic LaMnO$_3$ surfaces, where it is reported that binding is generally weaker than at Mn sites. The (101) surface results in this adsorption mode as access to the Mn site in the second surface layer is sterically inhibited, and the surface layer contains two La ions in close proximity (Figure 6).

(110). This surface presents three distinct types of adsorption sites: MnO$_4^{\text{Eq}}$, MnO$_4^{\text{Ap}}$, and MnO$_5^{\text{Eq}}$. We investigated these sites independently by adsorbing an O$_2$ molecule at an initial geometry above the sites. However, as observed in Figure 7, three metastable adsorption modes are obtained that center on the MnO$_4^{\text{Eq}}$ site: MnO$_4^{\text{Eq-Lateral}}$, MnO$_4^{\text{Eq-End-on}}$, and MnO$_4^{\text{Eq-Bidentate}}$. The corresponding initial adsorption modes are provided in Table 2.
MnO$_{4}^{\text{Eq-Bidentate}}$, resulting in binding energies of −2.04, −2.32, and −2.67 eV, respectively. The $d(O-O)$, which has been linear with binding energy for the other surfaces, appears to be insensitive to the binding energies on the MnO$_{4}^{\text{Eq}}$ site. The trend in $d(O-O)$ can instead be rationalized more accurately by the variations in charge transfer since increased filling of the antibonding $O_2\pi^*$ orbital destabilizes the $O-O$ bond. The O$_2$ molecules in the MnO$_{4}^{\text{Eq-Lateral}}$ and MnO$_{4}^{\text{Eq-Bidentate}}$ modes accept −1.24 $e$ and −1.04 $e$, while in the MnO$_{4}^{\text{Eq-End-on}}$ mode it accepts −0.70 $e$, resulting in $d(O-O)$ of 1.48, 1.46, and 1.34 Å, respectively. The differences in binding energy between the adsorption modes on the MnO$_{4}^{\text{Eq}}$ site occur in part due to variations in structural relaxation stemming from the orientation of the O$_2$ molecule. End-on adsorption is known to be more energetically favorable than lateral adsorption, perhaps due to lower interference with lattice O ions.\(^{10}\) For the MnO$_{4}^{\text{Eq-Bidentate}}$ mode, the high binding energy can mostly be attributed to the stabilization of both adsorption sites of the MnO$_4$ octahedron simultaneously, which effectively cancels out its undercoordination and massively stabilizes the surface.

Binding on the MnO$_{4}^{\text{Eq}}$ site of the (110) surface results in shorter $d$(Mn–O$_2$) and significantly larger binding energies than observed on the MnO$_{5}^{\text{Eq}}$ of the (001) surface. This can be attributed to undercoordination. The MnO$_{4}^{\text{Eq}}$ site goes from 4 to 5 coordination upon O$_2$ adsorption but still does not achieve the 6 coordination of the bulk, thus retaining a stronger attraction between the Mn center and the O$_2$ molecule than observed on a MnO$_5$ site after O$_2$ adsorption.

**DISCUSSION**

Our current understanding of optimal ORR activity in d$^4$ and d$^\pi$ perovskites is based on the experimental study of a wide range of perovskites.\(^7\) The data suggest that the requirement for the ORR mechanism to proceed optimally is the existence of a single $e_\pi$ electron in the TM ion that is able to form a covalent interaction with the adsorbate. As a consequence, variations in the ORR kinetics between manganites (d$^4$) and nickelates (d$^\pi$), for instance, are seen to be a result of the changes in TM–O hybridization that modify the covalency of the TM–O$_2$ (site–adsorbate) bond. However, this simplified concept of a single $e_\pi$ electron site and a covalent interaction is insufficient to describe the data presented herein. For example, the Mn$_X^{\text{pp}}$ site that shows single $e_\pi$ occupancy produces weak charge transfer and binding when compared to the doubly $e_\pi$ occupied MnO$_{4}^{\text{Eq}}$ site, which generates strong charge transfer and binding.

After consideration of the computed binding energies and charge transfer at the various sites, we propose that the ORR mainly proceeds via the disproportionation reaction of adsorbed superoxide O$_2^{\text{-}}$ with water to form the readily desorbed hydroperoxide HOO$^-$ and hydroxide OH$^-$.\(^{6,20,21}\) This pathway only proceeds via the mediation of catalyst sites that lead to the formation of O$_2^{\text{-}}$. The sites which result in the formation of strongly bound peroxide O$_2^{2-}$ and superoxide species, but due to the strong distortions of the truncated surface octahedra and the variations in site coordination upon cleavage, only one of these sites is found to be particularly suitable for the ORR: the MnO$_X^{\text{pp}}$ site on the (001) surface. At all other sites the binding energy is either too weak (MnO$_X^{\text{pp}}$ sites) or too strong (MnO$_X^{\text{pp}/3\text{-}}$ sites).

The predicted relative inactivity of many sites on the catalyst may seem to be a strong contraindication for its use in alkaline fuel cells as only 20% of the catalyst surface is predicted to be the (001) surface.\(^{18}\) This result does, however, suggest a route to optimization of the catalyst through control of the crystallite morphology to maximize the area of the (001) facets. As the data presented make it clear that the binding energy of O$_2$ is strongly affected by the distortion of the surface MnO$_5$ sites, we further propose optimization routes that vary the Mn–O distances of the surface sites. This includes doping in nanocrystalline powders or strain engineering in epitaxial thin films. These strategies are highly likely to be successful in catalyst optimization, and their implementation has been shown to be achievable in recent experiments on LaMnO$_3$ thin films.\(^{23}\)

**CONCLUSIONS**

This work details the structure and energetics of the adsorption sites on the (100), (001), (101), and (110) surfaces of orthorhombic (Pnma) LaMnO$_3$ by adsorption of molecular oxygen (O$_2$) in hybrid-exchange (B3LYP) density functional theory calculations. The binding energies of O$_2$ indicate that MnO$_{5}^{\text{Eq}}$ adsorption sites formed through cleavage of J−T distorted apical bonds have a weak interaction with O$_2$, while equatorially cleaved sites bind more strongly. In both instances, O$_2$ is reduced to superoxide O$_2^{\text{-}}$. Lower coordination sites (MnO$_{5}^{\text{pp}}$) have overly strong binding and allow for the formation of peroxide O$_2^{2-}$, which is not likely to be effective for the ORR in alkaline fuel cells (AFCs).

Only the MnO$_{4}^{\text{Eq}}$ site of the (001) surface satisfies the conditions of binding O$_2$ with an energy neither too weak nor too strong and results in the formation of a superoxide species necessary to proceed with the oxygen reduction reaction (ORR) in AFCs. This is shown to be due to the local distortion of the site which is caused by the J−T distortion of bulk MnO$_6$ octahedra. This model suggests that manipulation of the site activity could be achieved by controlling bulk Mn–O distances through doping and lattice strain. In addition, the total activity can be increased by maximization of the active surface area (surfaces that exhibit MnO$_X^{\text{Eq}}$ sites) through control of the crystallite morphology. This work also establishes the need to consider the effects of strong on-site electron correlation in determining the geometry and chemical activity of reaction sites in transition metal oxide catalysts. It is therefore of direct relevance to the discovery and optimization of functional transition metal oxides in many applications.

**METHODOLOGY**

Since previous studies have indicated that the ORR is catalyzed via adsorption at a surface Mn site,\(^{7,25}\) we adopt an initial geometry for O$_2$ adsorption with the molecule in an “end-on” orientation at the Mn site at a position reflecting the type of Mn–O bond cleaved when forming the surface (Figure 8). Symmetric adsorption on both sides of a periodic slab is found to converge the adsorption energy rapidly as a function of slab thickness as it prevents the generation of an unphysical electrostatic dipole perpendicular to the surface. The predicted adsorption geometry is then generated by an unconstrained
optimization of all atomic coordinates. Details of the surface construction are provided in the Supporting Information.

**Binding Energies.** In general computed surface binding energies are corrected for basis set superposition error (BSSE) using a counterpoise correction (CPC). For systems studied in this work the CPC method is found to be inapplicable as strong charge rearrangements occur at the adsorption sites during relaxation. The BSSE is therefore estimated at the unrelaxed geometry. The details of this approach and the full list of binding energies obtained using the different methods of estimating BSSE are provided in the Supporting Information.

**TEM.** The LaMnO₃ single-crystal particles were synthesized using a coprecipitation method as described elsewhere. High-resolution transmission electron microscopy (TEM) was performed on a spherical aberration (Cs) corrector at the image plane FEI Titan 80-300 TEM operated at 300 kV. For enhancement of the contrast in the micrographs, the corrector was set at negative spherical aberration imaging conditions, and it was tuned to −13 μm. Simulations of the atomic positions to assist the interpretation of the high-resolution micrographs were done using the multislice method in JEMS software.

**Computational Details.** We adopt the hybrid-exchange B3LYP functional for our DFT calculations, as it provides a description of strong on-site correlations and thus an accurate description of the energetics and electronic structure of systems in which strong electron–electron interactions result in electron localization and orbital ordering (particularly for transition metal oxides). It has also recently been shown to provide quantitative agreement between calculated and experimental formation energies for a wide range of Mn-oxides including LaMnO₃.

All calculations have been performed using the CRYSTAL09 code, based on the expansion of the crystalline orbitals as a linear combination of a local basis set (BS) consisting of atom centered Gaussian orbitals. The Mn and O ions are described by a triple valence all-electron BS: an 86-411d(411) contraction (one s, four sp, and two d shells) and an 8-411d(1) contraction (one s, three sp, and one d shell), respectively; the most diffuse sp(d) exponents are αₘₙ = 0.4986(0.249) and αₒ = 0.1843(0.6) Bohr⁻². The La basis set includes a pseudopotential to describe the core electrons, while the valence part consists of a 411p(411)d(311) contraction scheme (with three s, three p, and three d shells); the most diffuse exponent is α₂⁺ = 0.15 Bohr⁻² for each s, p, and d shell.

The internal coordinates of the slab have been determined by minimization of the total energy within an iterative procedure based on the total energy gradient calculated with respect to the nuclear coordinates. Convergence was determined from the root-mean-square (rms) and the absolute value of the largest component of the forces. The thresholds for the maximum and the rms forces (the maximum and the rms atomic displacements) have been set to 0.00045 and 0.00030 (0.00180 and 0.0012) in atomic units. Geometry optimization was halted when all four conditions were satisfied simultaneously.

**REFERENCES**


**ASSOCIATED CONTENT**

*Supporting Information*

**Procedure for calculations and the full Mn 3d orbital occupations of surface Mn ions. The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acs.jpcc.5b05460.**

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**Notes**

The authors declare no competing financial interest.

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