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Lap time optimisation of a racing go-kart

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ABSTRACT

The minimum lap time optimal control problem has been solved for a go-kart model. The symbolic algebra software Maple has been used to derive equations of motion and an indirect method has been adopted to solve the optimal control problem. Simulation has been successfully performed on a full track lap with a multibody model endowed with seven degrees of freedom. Geometrical and mechanical characteristics of a real kart have been measured by a lab test to feed the mathematical model. Telemetry recorded in an entire lap by a professional driver has been compared to simulation results in order to validate the model. After the reliability of the optimal control model was proved, the simulation has been used to study the peculiar dynamics of go-karts and focus to tyre slippage dynamics, which is highly affected by the lack of differential.

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1. Introduction

Optimal control problems applied to racing vehicle are nowadays widely used to simulate and predict the performance on racing tracks. Examples can be found for rally cars, [1] Formula 1 cars [2] and motorcycles. [3,4] However, in the literature there is a lack of optimal control analysis regarding go-karts. Go-karts are four-wheeled vehicle which, unlike cars, do not have suspensions nor differential. Therefore, the dynamics of karts is significantly different from that of other four-wheeled vehicles. The absence of rear axle differential makes the rear wheels to have a strictly correlated longitudinal slip and this characteristic makes it more difficult for a kart to turn than cars. Indeed, when turning at a constant speed, the rear inner tyre has a positive slip, while the rear outer tyre has a negative one; thus, longitudinal forces that tend to oppose to cornering are generated. In order to reduce this effect, the inner rear wheel should be unloaded as much as possible: it is quite usual to see expert drivers lifting up such a wheel while cornering. The frame stiffness and steering geometry are the most important parameters affecting this phenomenon; moreover, the absence of suspensions makes the tyres and frame absorb all shocks induced by road unevenness. Thus, kart frame stiffness becomes a key feature that profoundly affects kart performance, as highlighted in [5] through a multibody simulated analysis and a comparison with experimental data; this study was focused on circular trajectories at a constant speed only.

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47 In this work, kart dynamics have been studied with the use of optimal control simula-
48 tions of an entire circuit lap. The mathematical model of the go-kart has been developed
49 with the use of a symbolic algebra software (MBSymba [6]), and then translated into
50 an optimal control problem with third-part libraries (XOptima and Mechatronix). Such
51 libraries use an indirect formulation to solve the optimal control problem, which has been
52 demonstrated to be effective for such purposes.[7,8]

53 Simulations have been carried out on the track 'Pista Azzurra', in Jesolo (near Venice),
54 and the results have been validated by comparison with experimental data acquired on the
55 same track. Once the reliability of the model was confirmed, we analysed the dynamics of
56 the kart when treading a turn with particular attention to tyre slippage and how the lack
57 of differential requires high lateral load transfer in order to minimise manoeuvre time.

58 In the following section, the mathematical model of the kart is described. Then, in
59 Section 3, the optimal control formulation is presented. In Section 4, one model valida-
60 tion is illustrated, including the kart analysis in the laboratory, the data acquiring and the
61 telemetry-simulation comparison. Finally, in Section 5 the analysis of kart dynamics on a
62 turn is presented.

64 **2. Mathematical model of the go-kart**

65 A go-kart is a four-wheeled, rear axle traction vehicle whose distinguishing features are the
66 absence of the suspensions system and differential. Since the vehicle has four contact points
67 with the road, it is hyperstatic. Therefore, the load distribution on each tyre depends on the
68 tyre radial compliance as well as the compliance of the chassis, which is quite deformable.
69 The absence of the differential on the traction axle makes any cornering manoeuvre prob-
70 lematic. Indeed, while cornering, the inner rear wheel has a smaller forward velocity, but
71 the same spin velocity than the outer one. Thus, the inner longitudinal slip is bigger than
72 the outer one, which leads to a yaw torque opposite to the yaw rate. Since the longitudinal
73 force is approximately proportional both to the longitudinal slip and to the tyre load, to
74 reduce such an undesired effect, the vertical load of the inner wheel should be as low as
75 possible and ideally null. This condition may be obtained by designing a proper combina-
76 tion of tyres and chassis stiffness, as well as steering system geometry. Steering geometry
77 indeed plays an important role in load transfers: due to the particular linkages between
78 front wheels and chassis, when steering the inner front tyre gets pulled down by a quantity
79 proportional to the steering angle, while the outer front tyre is lifted up by the same amount.
80 Thus, part of the vertical load is transferred to the front inner wheel and the rear inner one
81 gets more – if not completely – unloaded. Such essential features are fully embraced in the
82 mathematical model of the vehicle, as depicted in Figure 1. The inertia of the vehicle (chas-
83 sis, engine, tanks, steering, wheels, etc.) is modelled with a single rigid body. Indeed, we
84 are not interested in vehicle vibration nor in any high-frequency dynamics. The mass and
85 inertia of the rider are also incorporated in such a rigid body, since the seat is very stiff and
86 the rider body is constricted. Some racing riders are still able to partially move their torso
87 while cornering, however, this effect is not so relevant and it is completely neglected here.
88 To capture the actual distribution of vertical tyres load, the chassis compliance is modelled
89 by using four vertical translation springs that ideally connect any wheel centre to the chas-
90 sis, while other four vertical translation springs connect wheel centres to the ground and
91 represent tyres radial stiffness. Two additional springs connecting the right and the left
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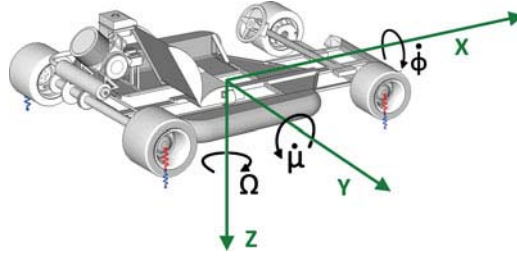


Figure 1. Go-kart model: the six degrees of freedom of the chassis are highlighted. The go-kart coordinate system is represented by the green axes. Black arrows and Greek letters indicate the rotational dof around the main axes. Red springs represent the reduced vertical stiffness of the chassis, while blue ones the tyre radial stiffness. Anti-roll bars (not represented) connect the left and right joint points between chassis and tyre springs.

wheel of each axle take into account for the anti-roll bars. The 125cc go-kart studied in this article was equipped with rear brake only, so the spin dynamics of the rear axle is included into the model, while the front axle is neglected. According to this description, the vehicle model has a total of seven degrees of freedom (dof), six of them associated with the vehicle gross motion and the latter associated with the rear axle spin rotation.

2.1. Newton's equations

The equations of motion have been derived by using MBSymba, which is an object-oriented language for the modelling of multibody systems and the automatic generation of the equations of motion.[6] The modelling procedure consists in the description of the multibody system by defining objects such as points, vectors, rigid bodies, forces and torques, and the relationships among them. In particular, it is possible to define which variables are small (infinitesimal) and which are not. Once the system has been modelled, the equations of motion are automatically derived, avoiding any human error. The Newton and Euler equations have been derived with respect to a coordinate system that is fixed to the go-kart body, except for having zero roll and pitch angles.

For a vehicle running on an horizontal, planar road surface, Newton's equations which describe the vehicle translation, respectively, along the longitudinal, lateral and vertical directions are

$$m(\dot{u} - \Omega v) = S_{rl} + S_{rr} - (F_{fl} + F_{fr}) \sin \delta - \frac{1}{2} \rho C_d A u^2, \quad (1a)$$

$$m(\dot{v} + \Omega u) = F_{rl} + F_{rr} + (F_{fl} + F_{fr}) \cos \delta, \quad (1b)$$

$$m\ddot{z} = mg - (N_{fl} + N_{fr} + N_{rl} + N_{rr}), \quad (1c)$$

where u and v are, respectively, the longitudinal and lateral speed of the vehicle centre of gravity (CoG), while z is the CoG vertical displacement with respect to the static equilibrium position, δ is the front wheel steering angle (assumed to be equal for both wheels), Ω is the yaw velocity, N_{as} , F_{as} , S_{as} are, respectively, the vertical, lateral and longitudinal force of each tyre, where $a \in \{r, f\}$ indicates either the rear axle, r , or the front one, f , and $s \in \{l, r\}$

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139 **Table 1.** Vehicle parameters.

Symbol	Value	Units	Description
h	0.250	m	^m Centre of gravity (CoG) height
a	0.645	m	^m Distance between the front axle and the vehicle CoG
b	0.400	m	^m Distance between the rear axle and the vehicle CoG
ρ	1.045		^m Wheelbase
$2t_f$	1.055	m	^m Front track
$2t_r$	1.200	m	^m Rear track
β	0.058	m/rad	^m Front wheel hub displacement to steering angle ratio
m	165	kg	^m Vehicle mass (rider included)
I_{xx}	20	kg m ²	^e Roll moment of inertia
I_{yy}	15	kg m ²	^e Pitch moment of inertia
I_{zz}	25	kg m ²	^e Yaw moment of inertia
I_{xz}	5	kg m ²	^e Mixed moment of inertia
I_r	0.2	kg m ²	^e Spin inertia of the rear axle
ρ	1.2	kg/m ³	Air density
C_dA	0.7	m ²	^e Drag surface coefficient
ks_r	60e3	N/m	^m Rear chassis stiffness (vertical displacements)
ks_f	17.7e3	N/m	^m Front chassis stiffness (vertical displacements)
kj_r	0	N/m	Rear anti-roll bar stiffness
kj_f	0	N/m	Front anti-roll stiffness
kt_r	61.3e3	N/m	^m Rear tyres radial stiffness
kt_f	64.5e3	N/m	^m Front tyres radial stiffness
ct_r	1.0e3	N s/m	^e Rear tyres radial damping
ct_f	1.0e3	N s/m	^e Front tyres radial damping
r_r	0.139	m	^m Rear tyres rolling radius
T_{\max}	17.6	N m	^m Maximum engine torque (@10,250 rpm)
P_{\max}	20.1	kW	^m Maximum engine power (@11,500 rpm)

160 Notes: ^m: Measured value, ^e: Estimated value, ^f: Fitted value.

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162 indicates either the left side, l , or the right one, r . Finally, model parameters are described
163 in Table 1.

164 Since tyres and chassis are stiff, it is reasonable to assume that the chassis pitch μ , roll
165 ϕ and vertical displacement z are infinitesimal, as well as their time derivatives. Therefore,
166 the equations of motion with respect to the roll, pitch and yaw rotation axis are
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$$168 \quad m[\dot{v}h + (h - z)\Omega u] + I_{xx}\ddot{\phi} - (I_{yy} - I_{zz} + I_{xx})\Omega\dot{\mu} - (I_{yy} - I_{zz})\Omega^2\phi \quad (2a)$$

$$169 \quad - (I_{xx} - I_{zz})\dot{\Omega}\mu - I_{xz}\dot{\Omega} = t_f(N_{fl} - N_{fr}) + t_r(N_{rl} - N_{rr}),$$

$$171 \quad m[(z - h)\dot{u} + h\Omega v] + I_{yy}\ddot{\mu} + (I_{yy} - I_{zz} + I_{xx})\Omega\dot{\phi} - (I_{xx} - I_{zz})\Omega^2\mu \quad (2b)$$

$$172 \quad + (I_{yy} - I_{zz})\dot{\Omega}\phi - I_{xz}\Omega^2 = a(N_{fr} + N_{fl}) - b(N_{rr} + N_{rl}) - \frac{1}{2}\rho C_d A u^2 (z - h),$$

$$175 \quad I_{zz}\dot{\Omega} - I_{xz}(\ddot{\phi} + 2\mu\dot{\Omega} + 2\dot{\mu}\Omega) = a(F_{fl} + F_{fr}) \cos \delta - b(F_{rr} + F_{rl}) \quad (2c)$$

$$176 \quad + t_f(F_{fr} - F_{fl}) \sin \delta + t_r(S_{rl} - S_{rr}).$$

177
178 It is worth pointing out that angular momenta have been evaluated with respect to the
179 vertical projection of the CoG onto the road plane. This choice simplifies the roll (2a) and
180 pitch (2b) equations by avoiding the presence of any term related to longitudinal or lateral
181 tyre forces.

182 Finally, the model is completed by the following rear axle spin equation of motion:

$$183 \quad I_r\dot{\omega}_r = T_r - S_{rr}(r_r - w_{rr}) - S_{rl}(r_r - w_{rl}), \quad (3)$$

184

185 where ω_r is the axle spin velocity, r_r is the rear tyre radius, w_{rr} and w_{rl} are rear right and
 186 rear left tyre deformations, respectively, and T_r is the axle torque due to either the engine
 187 ($T_r > 0$) or the brake ($T_r < 0$). The inertia I_r includes all parts connected to the shaft
 188 (wheels, disk brake, etc.) as well as the equivalent inertia of the engine, which is connected
 189 to the rear axle by a fixed ratio chain transmission.

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2.2. Tyres model

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Since the vehicle has four contact points with the road, it is not possible to determine the tyre vertical load unless the compliances of tyres and chassis are considered. In a kart, tyre stiffness is generally 3–5 times greater than chassis stiffness; however, in this model only their combined stiffness is considered. Indeed, free vibration frequencies of the chassis with attached wheels are far faster than the eigenfrequencies of vehicle gross motion, so chassis and wheels may be considered as a unique rigid body which is attached to the road by means of massless springs and dampers. According to these assumptions, for each wheel the load force due to tyre deformation has been made equal to the vertical force due to chassis deformation, as shown in the following formulas. Tyre deformations are given by

$$\begin{aligned}
 w_{fr} &= z + s_{fr} - a\mu + t_f\phi + \beta\delta, \\
 w_{fl} &= z + s_{fl} - a\mu - t_f\phi - \beta\delta, \\
 w_{rr} &= z + s_{rr} + b\mu + t_r\phi, \\
 w_{rl} &= z + s_{rl} + b\mu - t_r\phi,
 \end{aligned} \tag{4}$$

where s are frame vertical deformations and $\pm\beta\delta$ is the linear approximation of the kinematic relationship that links the vertical motion of the wheel centre to the steering angle (in other words β is the ratio between wheel vertical displacement and steering angle). Then, by making tyre loads equal to chassis deformation forces on each wheel:

$$\begin{aligned}
 s_{rr}k_{sr} + (s_{rr} - s_{rl})k_{jr} + w_{rr}k_{tr} &= 0, \\
 s_{rl}k_{sr} + (s_{rl} - s_{rr})k_{jr} + w_{rl}k_{tr} &= 0, \\
 s_{fr}k_{sf} + (s_{fr} - s_{fl})k_{jf} + w_{fr}k_{tf} &= 0, \\
 s_{fl}k_{sf} + (s_{fl} - s_{fr})k_{jf} + w_{fl}k_{tf} &= 0,
 \end{aligned} \tag{5}$$

where k_{sf} , k_{sr} are the chassis stiffness at the front and rear axles, and k_{jr} and k_{jf} are the rear and front anti-roll bar stiffness. In Equations (5), the first terms represent the force due to chassis deformations, the second terms are the force of the anti-roll bars and the last terms stand for tyres deformation loads. Once the expressions for s are obtained, they can be substituted into the tyre deformation expressions that become

$$\begin{aligned}
 w_{rr} &= \frac{k_{sr}(\mu b + \phi t_r + z)}{k_{sr} + k_{tr}} + \frac{2\phi k_{tr} t_r k_{jr}}{(k_{sr} + k_{tr})(2k_{jr} + k_{sr} + k_{tr})}, \\
 w_{rl} &= \frac{k_{sr}(\mu b - \phi t_r + z)}{k_{sr} + k_{tr}} - \frac{2\phi k_{tr} t_r k_{jr}}{(k_{sr} + k_{tr})(2k_{jr} + k_{sr} + k_{tr})},
 \end{aligned}$$

$$\begin{aligned}
231 \quad w_{fr} &= -\frac{k_{sf}(\mu a - \beta \delta - \phi t_f - z)}{k_{sf} + k_{tf}} + \frac{2k_{jf}k_{tf}(\beta \delta + \phi t_f)}{(k_{sf} + k_{tf})(2k_{jf} + k_{sf} + k_{tf})}, & (6) \\
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234 \quad w_{fl} &= -\frac{k_{sf}(\beta \delta + \mu a + \phi t_f - z)}{k_{sf} + k_{tf}} - \frac{2k_{jf}k_{tf}(\beta \delta + \phi t_f)}{(k_{sf} + k_{tf})(2k_{jf} + k_{sf} + k_{tf})}. \\
235
\end{aligned}$$

236 From these expressions, it can be noticed that tyre deformations are given by two terms:
 237 the first is that of a totally rigid chassis with tyre radial stiffness given by the series of chassis
 238 and tyre springs, and the second term is due to the anti-roll bar and proportional to the
 239 stiffness of the bars. Such expressions for tyre deformations have been substituted into tyre
 240 loads in order to algebraically remove the degrees of freedom of chassis deformations. As
 241 shown in Equation (7), a damping term proportional to the tyre deformation rate has been
 242 added in the calculus of tyre loads. Such a term, even if it is negligible for the dynamics of
 243 the vehicle, has been inserted to eliminate undesired oscillations in the simulation results.
 244 Finally, since while cornering the rear inner wheel may lift up and loose the contact with
 245 the ground, the rear tyre load expressions have been saturated accordingly. Thus:

$$\begin{aligned}
247 \quad N_{rr} &= \max(0, k_{tr}w_{rr} + c_{tr}\dot{w}_{rr}), \\
248 \quad N_{rl} &= \max(0, k_{tr}w_{rl} + c_{tr}\dot{w}_{rl}), & (7) \\
249 \\
250 \quad N_{fr} &= k_{tf}w_{fr} + c_{tf}\dot{w}_{fr}, \\
251 \quad N_{fl} &= k_{tf}w_{fl} + c_{tf}\dot{w}_{fl}, \\
252
\end{aligned}$$

253 where c_{tr} , c_{tf} are rear and front tyre radial damping and $\max(x, y)$ is a function that returns
 254 the maximum between x and y . For the sake of legibility, the tyre deformation expressions
 255 (Equation (6)) have not been substituted into previous equations. Lateral tyres forces and
 256 longitudinal tyres forces are calculated according to the Magic formula [9] as a function of
 257 tyre vertical load N , sideslip angle λ and longitudinal slip κ . In particular, sideslips may be
 258 calculated as follows:

$$\begin{aligned}
259 \quad \lambda_{rr} &= -\arctan\left(\frac{v - \Omega b}{u - \Omega t_r}\right), \\
260 \\
261 \quad \lambda_{rl} &= -\arctan\left(\frac{v - \Omega b}{u + \Omega t_r}\right), \\
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264 \quad \lambda_{fr} &= \delta - \arctan\left(\frac{v + \Omega a}{u - \Omega t_f}\right), & (8) \\
265 \\
266 \quad \lambda_{fl} &= \delta - \arctan\left(\frac{v + \Omega a}{u + \Omega t_f}\right). \\
267 \\
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\end{aligned}$$

269 Longitudinal slips are calculated for the rear axle only:

$$\begin{aligned}
270 \quad \kappa_{rr} &= \frac{\omega_r r_r}{u - \Omega t_r} - 1, \\
271 \\
272 \quad \kappa_{rl} &= \frac{\omega_r r_r}{u + \Omega t_r} - 1. & (9) \\
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\end{aligned}$$

275 Since the time lag which is present in the generation of lateral forces due to lateral compli-
 276 ance of tyre carcass is known to have a remarkable influence on the vehicle stability, such

277 a phenomenon is modelled as a first-order relaxation equation [9] for each tyre:

$$\begin{aligned}
 278 & \\
 279 & \frac{\sigma_r}{u - \Omega t_r} \dot{F}_{rr} + F_{rr} = F_{\text{magic},r}(N_{rr}, \lambda_{rr}, \kappa_{rr}), \\
 280 & \\
 281 & \frac{\sigma_r}{u + \Omega t_r} \dot{F}_{rl} + F_{rl} = F_{\text{magic},r}(N_{rl}, \lambda_{rl}, \kappa_{rl}), \\
 282 & \\
 283 & \frac{\sigma_f}{(u - \Omega t_f) \cos \delta + (v + \Omega a) \sin \delta} \dot{F}_{fr} + F_{fr} = F_{\text{magic},f}(N_{fr}, \lambda_{fr}, \kappa_{fr}), \\
 284 & \\
 285 & \frac{\sigma_f}{(u + \Omega t_f) \cos \delta + (v + \Omega a) \sin \delta} \dot{F}_{fl} + F_{fl} = F_{\text{magic},f}(N_{fl}, \lambda_{fl}, \kappa_{fl}). \\
 286 & \\
 287 &
 \end{aligned} \tag{10}$$

288 On the real axle, small spin inertia leads to fast dynamics and hence longitudinal forces are
 289 calculated without any time lag:

$$\begin{aligned}
 291 & S_{rr} = S_{\text{magic}}(N_{rr}, \lambda_{rr}, \kappa_{rr}), \\
 292 & S_{rl} = S_{\text{magic}}(N_{rl}, \lambda_{rl}, \kappa_{rl}), \\
 293 &
 \end{aligned} \tag{11}$$

294 where $S_{\text{magic},r}$, $F_{\text{magic},r}$ and $F_{\text{magic},f}$ are the functions returning, respectively, the rear lon-
 295 gitudinal, rear lateral and front lateral tyre forces (front and rear tyres have different
 296 parameters). The Pacejka Magic Formula Tyre Model ([9], chapter 4) has been used in
 297 $S_{\text{magic},r}$, $F_{\text{magic},r}$ and $F_{\text{magic},f}$, but the dependence on the tyre roll angle has been neglected.
 298 All the coefficients used in these formulas are listed in the appendix; the parameters that
 299 do not appear in the table in the appendix have been set to zero (or to one, in the case of
 300 scaling coefficients). This is because, as described in Section 4.1, only the tyres lateral stiff-
 301 ness have been measured in lab tests; thus, we used only the smallest number of parameters
 302 in Pacejka's magic formulas.

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2.3. Road and vehicle tracking

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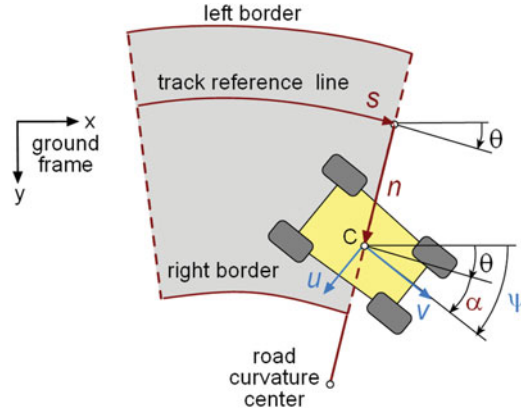
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Real roads are similar to strips: they are long and narrow. According to this idea, Figure 2
 illustrates a string-shaped road, which may be described by specifying the (x, y) coordi-
 nates of the middle line \mathcal{C} and the lateral width. As explained in [10], the most effective
 way to describe the road centre is the curvilinear abscissa approach, which makes it pos-
 sible to define \mathcal{C} by only giving the expression of the line curvature κ as a function of the
 line length s . Then, the road heading angle Θ as well as x, y coordinates may be calculated
 by integrating the curvature as follows:

$$\begin{aligned}
 \frac{d}{ds} \Theta &= \kappa(s), \\
 \frac{d}{ds} x &= \cos \Theta, \\
 \frac{d}{ds} y &= \sin \Theta.
 \end{aligned} \tag{12}$$

The advantage of this approach becomes evident while reminding that real roads and tracks
 are commonly designed in terms of straight segments (null curvature) which are connected

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335 **Figure 2.** Road tracking.
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to circular arcs (constant curvature) by means of clothoids (which curvature varies linearly). Indeed, while driving the steering angle is approximately proportional to the road and hence it is essential to avoid any discontinuity in the curvature.

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The advantage of the curvilinear coordinates approach becomes even more significant when it is applied to track the position and orientation of the vehicle as shown in Figure 2, where s and n are, respectively, the longitudinal and lateral position on the road strip, while α is the vehicle heading relative to the road (in other words, α is the angle between the road mid-line direction and the x -axis fixed on the vehicle). The calculus of the vehicle forward and lateral speed, as well as yaw velocity, yields

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2.4. First-order formulation

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Newton's equations (1) and (2) may be immediately reduced to the first order by introducing auxiliary variables for the relevant chassis speeds:

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$$\begin{aligned} \dot{z} &= z_{\text{dot}}, \\ \dot{\phi} &= \omega_{\phi}, \\ \dot{\mu} &= \omega_{\mu}. \end{aligned} \quad (14)$$

369 At this point, Equations (1)–(3), (10), (13) and (14) and their subordinate expressions
 370 completely describe the go-kart dynamics as a system of 17 first-order differential equa-
 371 tions with as many state variables and 2 inputs, respectively, the rear axle torque T_r , which
 372 (mainly) control the longitudinal dynamics, and the steering angle δ , which (mainly) con-
 373 trol the lateral dynamics. However, to avoid unrealistic jerky manoeuvres, it is necessary
 374 to consider that human drivers have a limited rate of change of control variables and
 375 experimental results show that humans optimise their driving actions minimising the lon-
 376 gitudinal and lateral jerks.[11] For this reason, it is assumed that the steering angle is not
 377 controlled directly, but via its time derivative, as follows:

$$\begin{aligned} 378 \quad \dot{\delta} &= \omega_\delta, \\ 379 \quad \dot{T}_r &= mr_r j_u, \end{aligned} \quad (15)$$

381 where ω_δ is the steering speed, which is approximately related to the lateral jerk and j_u
 382 (approximately) corresponds to the longitudinal jerk.

383 In conclusion, the vehicle dynamics is described by means of a set of two inputs:

$$384 \quad \mathbf{u} = \{\omega_\delta, j_u\}^T, \quad (16)$$

385 plus 19 state variables:

$$386 \quad \mathbf{x} = \{s, n, \alpha, z, \phi, \mu, u, v, z_{\text{dot}}, \Omega, \omega_\phi, \omega_\mu, \omega_r, F_{rr}, F_{rl}, F_{fr}, F_{fl}, \delta, T_r\}^T \quad (17)$$

387 and as many implicit first-order differential equations, which may be abbreviated to:

$$388 \quad \hat{\mathbf{A}}(\mathbf{x})\dot{\mathbf{x}} = \hat{\mathbf{f}}(\mathbf{x}, \mathbf{u}), \quad (18)$$

389 where matrix $\tilde{\mathbf{A}}$ is invertible provided that $u > 0$, i.e. the vehicle never stops, and
 390 $n < 1/\kappa(s)$, i.e. the vehicle never passes over the local curvature centre of the road.

391

392 **3. Optimal control problem**

393 **3.1. Vehicle dynamics in curvilinear abscissa domain**

394 The minimum lap time problem consists, as previously said, in finding the vehicle control
 395 inputs that minimise the time T necessary to move the vehicle along the track from the
 396 starting line to the finish one. Therefore, the curvilinear abscissa s varies between fixed
 397 initial point $s = 0$ and end point $s = L$, while the final value T of the independent variable
 398 t is unknown. For this reason, it is convenient to change the independent variable from
 399 t to s in the equations of motion (18). Such a variable change is based on the following
 400 derivation rule:

$$401 \quad \dot{\mathbf{x}} = \frac{d\mathbf{x}}{dt} = \frac{d\mathbf{x}}{ds} \frac{ds}{dt} = \mathbf{x}'\dot{s} = \mathbf{x}'v. \quad (19)$$

402 The time domain equations (18) are then transformed into the space domain as follows:

$$403 \quad v\tilde{\mathbf{A}}\mathbf{x}' = \tilde{\mathbf{f}}(\mathbf{x}, \mathbf{u}). \quad (20)$$

404 However, the variable change (19) transforms the differential equation (13a) into an alge-
 405 braic one, that should be eliminated from state equations (20) together with the variable s

415 which can no more belongs to the state vector \mathbf{x} . At this point, the variable t is no longer
 416 present in the model; however, it can be easily re-introduced by writing Equation (13a) as
 417 follows:

$$418 \quad \frac{dt}{ds} = t' = \frac{1}{v} = \frac{1 - n\kappa(s)}{u \cos \alpha - v \sin \alpha}. \quad (21)$$

420 Summarising, the s -domain state space model has $n = 19$ state variables:

$$421 \quad \mathbf{y} = \{t, n, \alpha, z, \phi, \mu, u, v, w, \Omega, \omega_\phi, \omega_\mu, \omega_r, F_{rr}, F_{rl}, F_{fr}, F_{fl}, \delta, T_r\}^T \quad (22)$$

422 and $p = 2$ inputs already defined in Equation (16) but now depend on s , while model
 423 equations may be summarised as a set of implicit differential equations:

$$424 \quad \nu A \mathbf{y}' = \mathbf{f}(\mathbf{y}, \mathbf{u}, s). \quad (23)$$

425 Equations (23) are not singular only and if only $\nu > 0$, i.e. the s -domain formulation
 426 requires that the vehicle never stops (a condition already required in the t -domain
 427 formulation) nor can it revert its direction of travel on the track.

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429 **3.2. The minimum lap time problem**

430 The Optimal control problem (OCP) consists in finding the vehicle control inputs that
 431 move the vehicle from the starting line $s = 0$ and $t(0) = 0$ to the finish one $s = L$ in the
 432 minimum time $T = t(L)$, while satisfying the mechanical equations of motion as well as
 433 other inequality constraints (tyres adherence, max power, track width, etc.) and may be
 434 formulated as follows:

$$435 \quad \text{find : } \min_{\mathbf{u} \in U} t(L) \quad (24a)$$

$$436 \quad \text{subject to : } \nu A \mathbf{y}' = \mathbf{f}(\mathbf{y}, \mathbf{u}, s), \quad \mathbf{y} \in \mathbb{R}^n, \quad \mathbf{u} \in \mathbb{R}^p \quad (24b)$$

$$437 \quad \boldsymbol{\psi}(\mathbf{y}, \mathbf{u}, s) \leq \mathbf{0}, \quad \boldsymbol{\psi} \in \mathbb{R}^m \quad (24c)$$

$$438 \quad \mathbf{b}(\mathbf{y}(0), \mathbf{y}(L)) = \mathbf{0}, \quad \mathbf{b} \in \mathbb{R}^{r \leq 2n}, \quad (24d)$$

439 where Equation (24b) is the state space model already discussed, Equation (24c) are
 440 algebraic inequalities that may bound both the state variables and control inputs and
 441 Equation (24d) is the set of boundary conditions used to (partially) specify the vehicle
 442 state at the beginning and at the end of the manoeuvre.

443 Inequalities (24c) are used to keep the vehicle inside the admissible range of operat-
 444 ing conditions. First of all, all wheels must remain inside the track (see Figure 2). Such a
 445 condition for the front right, front left, rear right and rear left wheel can be expressed as

$$446 \quad \begin{aligned} 447 \quad n + a \sin(\alpha) + t_f \cos(\alpha) &\leq w_r, \\ 448 \quad n + a \sin(\alpha) - t_f \cos(\alpha) &\geq -w_l, \\ 449 \quad n - b \sin(\alpha) + t_r \cos(\alpha) &\leq w_r, \\ 450 \quad n - b \sin(\alpha) - t_r \cos(\alpha) &\geq -w_l, \end{aligned} \quad (25)$$

451 where w_l and w_r are the distance of the left and right border from the track reference line
 452 (that might vary along the track), respectively. The propulsive torque on the rear axle T_r is

461 limited by the maximum engine torque, which is a function of the engine speed $\omega_e = \tau_e \omega_r$,
 462 as follows:

$$463 \quad T_r \leq \tau_e T_e(\tau_e \omega_r). \quad (26)$$

464
 465 It is not necessary to explicitly restrict tyre longitudinal and lateral forces within their
 466 adherence limits; indeed, such constraints are implicitly contained in the saturated Magic
 467 formula used in Equations (11) and (10). On the contrary, it is necessary to explicitly
 468 prevent vehicle roll-over by imposing the following constraints on the front vertical loads:

$$469 \quad N_{fr} \geq 0, \quad (27)$$

$$470 \quad N_{fl} \geq 0.$$

471
 472 Finally, the control inputs are limited in magnitude as follows:

$$473 \quad -\omega_L \leq \omega_\delta \leq \omega_L, \quad (28a)$$

$$474 \quad -j_L \leq j_u \leq j_L. \quad (28b)$$

475
 476 Equations (25)–(28) form a set of $m = 11$ unilateral constraints of type (24c). The pres-
 477 ence of such inequality constraints makes the OCP problem (24) particularly hard to solve.
 478 However, it is possible to convert such inequality constraints into penalty terms [7] to be
 479 included in the optimality criterion (24a). Each penalty term should be very small (ide-
 480 ally null) when a constraint is satisfied and suddenly should become large as the constraint
 481 limit is approached and possibly reached. In conclusion, the OCP problem (24) is replaced
 482 with the (almost) equivalent one which does not contain inequality constraints:

$$483 \quad \text{find : } \min_{u \in U} t(L) + \sum_{j=1}^m \int_0^L W_j(\psi_j(\mathbf{y}, \mathbf{u}, s)) ds \quad (29a)$$

$$484 \quad \text{subject to : } \nu A \mathbf{y}' = \mathbf{f}(\mathbf{y}, \mathbf{u}, s) \quad (29b)$$

$$485 \quad \mathbf{b}(\mathbf{y}(0), \mathbf{y}(L)) = \mathbf{0}. \quad (29c)$$

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 487
 488 To complete the problem definition, it is necessary to specify boundary conditions (24d).
 489 As the optimisation is made on a closed loop track, it is natural to impose cyclic boundary
 490 conditions for all state variables $\mathbf{y}(s)$, except for $t(L)$ which is of course free and under opti-
 491 misation. The OCP formulation (29) is general and the problem may be solved by using
 492 different approaches [12] such as non-linear programming, dynamic programming, and
 493 Pontryagin's indirect method, which is the one that has been used in the present paper.
 494 Such a method,[13] using Lagrange's multipliers technique, eliminates differential con-
 495 straints (29b) and transforms the OCP formulation (29) into a functional unconstrained
 496 minimisation problem. Then, according to the variational first principle, the stationarity
 497 condition of the functional to be minimised leads to a two-point boundary value prob-
 498 lem (BVP). Despite the elegance of this well-known technique, it is not straightforward to
 499 apply this to complex models as the one presented here. However, the utilisation of com-
 500 puter algebra tools (we used Maple) makes it possible to symbolically derive BVP equations
 501 and automatically convert them into a C++ code ready to be compiled, avoiding human
 502 errors. Finally, the BVP problem is discretised and numerically integrated using `XOPTima`,
 503 a specialised solver.[8]

507 4. Experimental test and model validation

508 In order to provide a validation of the mathematical model by comparison between
509 simulation results with experimental telemetry, a precise measurement of the go-kart char-
510 aracteristic is first necessary. Indeed, the reliability of the simulation is highly affected by the
511 agreement of the data used with real kart properties.

512 In the following section, the description of the procedure used to measure go-kart
513 characteristics is reported, while in Section 4.2 simulation results are compared to track
514 data.

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4.1. Lab tests

519 The geometry, inertia and compliance of the go-kart, as well as tyre properties, have been
520 measured to feed the simulation model. The geometry of the chassis has been carefully
521 measured by using a 3D coordinate-measuring machine. Weighing balances have been
522 used to measure the go-kart mass and its longitudinal position, the CoG vertical position
523 was estimated by tilting the vehicle until it reached its (unstable) equilibrium position on
524 the two lateral wheels. Such geometric and inertial data was used to build a virtual pro-
525 totype of the go-kart, which has been used to estimate the moment of inertia as well as
526 to estimate the frame stiffness by means of a FEM analysis. The frame is complex and
527 its representation in terms of few lumped springs used in the vehicle dynamic illustrated
528 Section 2 disregards many degrees of freedom and the correlation of such lumped stiffness
529 to FEM analysis is not trivial. In this case, the focus is on the correct estimation of tyres
530 vertical loads, so the hub points stiffness have been estimated by simultaneously locking
531 three wheels hub and by applying a vertical force on the last one. Vehicle parameters are
532 collected in Table 1. Tyres properties have been measured on the rotating tyre test rig which
533 is available at the department.[14] The identification of the vertical structural stiffness k_{tr} ,
534 k_{tf} has been carried out by applying different vertical loads in the range of 0–1500 N both
535 front and rear tyres showed a good linear behaviour. Tyre adherence was measured with a
536 vertical load of 560 N, which is approximately the load of a rear tyre in static conditions.
537 Figure 3 shows the ratio between the lateral force and the vertical load as a function of the
538 sideslip angle, for both the front and the rear tyre. The tyre behaviour in the range of 0–6°
539 of sideslip has been accurately identified, but unfortunately the adherence peak was not
540 reached during measurements; indeed, the friction between the tyre and the test rig were
541 exaggeratedly high (as the tyre test rig was not specifically designed to measure kart tyres).
542 It is worth pointing out that in actual driving conditions the adherence limit remarkably
543 depends on many parameters and in particular on the asphalt characteristics. Hence, the
544 correlation with lab measurement would not be very easy in any case. On the contrary,
545 the cornering stiffness, i.e. the tyre behaviour at low sideslip angles, mainly depends on
546 the tyre carcass characteristics and it is not so much influenced by the tyre–road friction
547 properties.[9,15] For this reason, tyres adherence peak have been adjusted to fit the exper-
548 imental maximum accelerations (that can be read in Figure 6). Tyre forces that were used
549 for the simulation are reported in Figure 4. Rear tyre normalised lateral force is shown for
550 different values of longitudinal slip to highlight the coupling between lateral and longitu-
551 dinal forces. For the front tyre, the force is shown only for $\kappa = 0$ because the front tyre's
552 longitudinal slip is not modelled. Tyres adherence values, together with all coefficients of

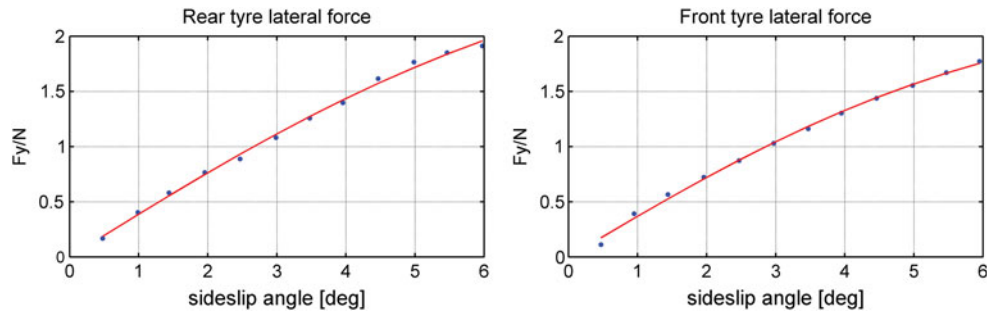


Figure 3. Lateral force of front and rear tyres measured on the test rig. Continuous red lines represent data fit.

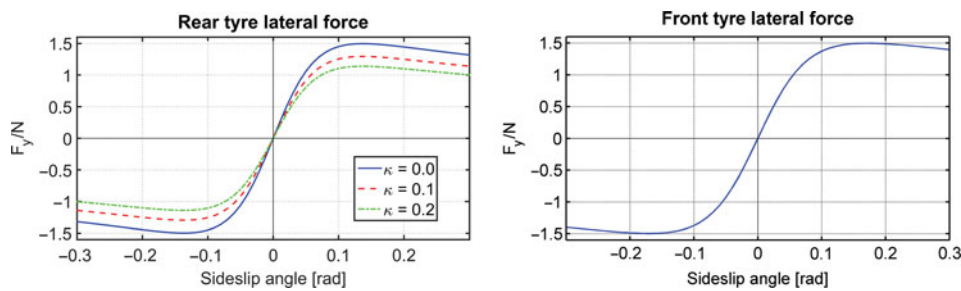


Figure 4. Tyre forces: normalised lateral forces are shown for rear and front tyres at a vertical load of 560 N. For rear tyre the force is shown at different values of longitudinal slip.

Pacejka's magic formula, are summarised in Table A1. The lateral friction peak coefficient (1.50) is significantly higher than longitudinal one (0.90) because in kart tracks thin layers of rubber get often deposited on the asphalt in curves due to high tyre slippage. Thus, turns are characterised by a higher friction coefficient than straights, and this effect can be modelled by considering tyres with greater lateral adherence. Such adherence conditions are confirmed in Section 4.3 by the validation of the model with experimental data. The engine torque curve used in this simulation is the one provide by Iame (www.iame.it).

4.2. Track tests

The go-kart has been equipped with an inertial measurement platform (IMU) composed of three accelerometers and three gyrometers. Moreover, a Hall sensor has been installed on the rear axle to measure the spin velocity and a rotational potentiometer has been installed on the steering system to measure the steering angle. All sensors have been connected to the data logger by means of a dedicated CAN bus. Tests have been carried out at Pista Azzurra (Jesolo), which is a national category race-track 1051 m long with an average width of 8 m. Tests were carried out on a cold day in November, with adherence conditions far from being optimal, by an expert driver who participates in go-kart races at the Italian level. Figures 6(b) and 6(c) highlighted the bounds of both the lateral accelerations $|a_y| < 15 \text{ m/s}^2$ and longitudinal one $-7 \text{ m/s}^2 < a_x < 5 \text{ m/s}^2$. While the traction acceleration is limited by the engine power, lateral and braking accelerations are limited by tyre adherence,

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Figure 5. Satellite view of Pista Azzurra with simulation trajectory overlaying (yellow line). Numbers inside boxes show the value of the curvilinear abscissa s along the track every 100 m. The turn inside the black rectangle is the one analysed in the next chapter.

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where lateral adherence is much bigger than the longitudinal one, thanks to the rubber deposits that tyres leave in curve. It is worth pointing out that, most likely, the driver was not able to drive very close to the tyre adherence limit and hence the acceleration bounds are a measure of driver's capability rather than tyre adherence. However, this is not a big issue for the validation of the simulation software, as the focus will be posed on the comparison between speed profile of real and virtual driver, constrained within the same acceleration bounds.

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4.3. Validation

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The simulation took approximately 230 s on a common laptop equipped with an Intel Core-i7 640M processor. Figure 5 shows a satellite view of the track with the overlying trajectory resulting from the optimal manoeuvring method. The trajectory confirms that the kart remains always within borders of the real track, suggesting that the road has been accurately reproduced. Simulated and experimental time lap are, respectively, 53.575 and 53.580 s. However, this is not an indication of the quality of the simulation since the information on the time lap has been used to fine tune tyre adherence, which has been estimated from measured lateral and longitudinal accelerations. The accuracy of the simulation should be instead evaluated by comparing measured data with simulated ones as they both vary along the track, as reported in Figure 6: dash-dot red lines refer to telemetry data, continuous blue ones to simulation. The simulated speed profile reproduces faithfully the experimental one; there are some slightly discrepancies in the speed at the middle of some turns, where the simulation is sometime faster, other slower. In addition to the simulation errors that are certainly present because of model approximation and parameter inaccuracy, it should be kept into consideration that asphalt was not perfectly dry so some parts of the circuit had a different adherence compared to others. Moreover, the driver

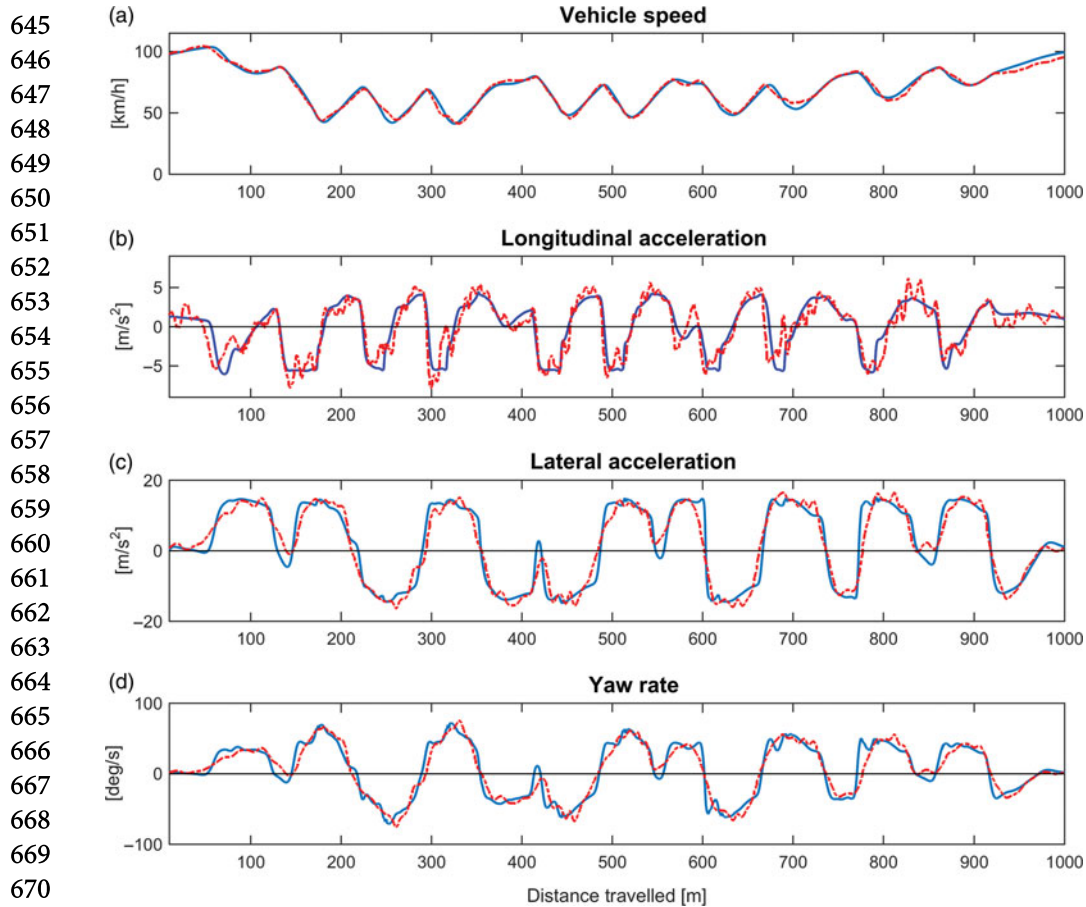
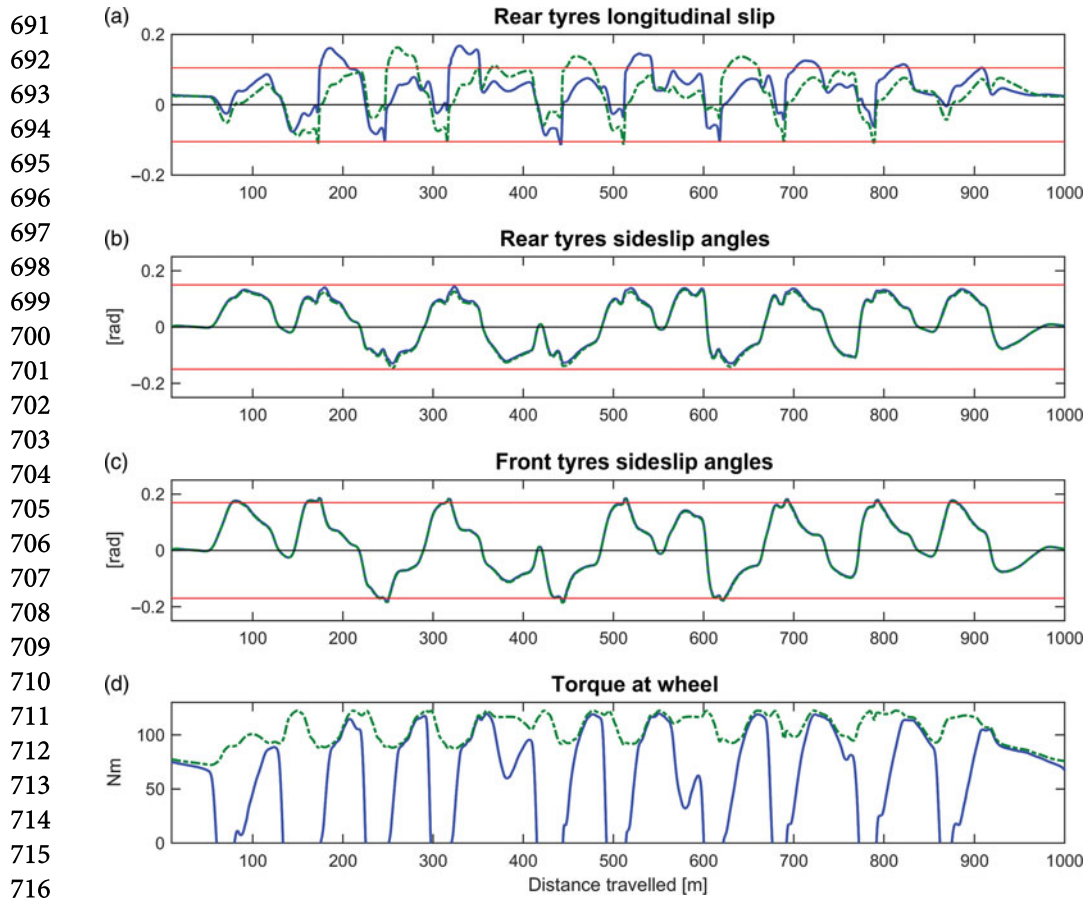


Figure 6. Vehicle gross motion for the full lap: dash-dot red lines refer to telemetry data, continuous blue ones to simulation results.

was very skilled but not professional. So, we think that the accordance between simulation and experiments is excellent. Longitudinal and lateral accelerations also reproduce faithfully the measured trend. Simulated signals are obviously smoother than the experimental ones, due to the absence of every source of unwanted vibrations. For this reason, the peak value of experimental accelerations are higher than the simulated one. The experimental lateral acceleration and yaw rate signal are less affected by external noise with respect to longitudinal acceleration and they have a better correspondence with the corresponding simulated quantities.

Figure 7 shows additional variables that are available for the simulation only and which are quite useful to analyse the go-kart behaviour. Rear tyre longitudinal and lateral slips are reported, respectively, on the first and the second line. As shown in the figure, sideslip angles of rear tyres are quite identical for left and right sides, while for the longitudinal slips there are significant differences. Due to the absence of a differential, such slippage differences are consequences of the yaw motion only. Thus, at a given time only one tyre can have the combination of κ and λ that produces the maximum force: this means that it is



717 **Figure 7.** Simulation details: in tyre slip charts, continuous blue lines refer to right tyre, dash-dot green
 718 lines to left one. Horizontal red lines represent the value at which the maximum of tyre adherence is
 719 reached for a non-combined tyre force. In the chart *d*, the continuous blue line represents the used
 720 torque, the dash-dot green one the maximum available torque.

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724 not possible to bring both tyres at adherence limits in the same manoeuvre, and the optimal
 725 control obviously tends to engage to the limit (at a longitudinal slip κ of ≈ 0.105) the
 726 tyre that has more load, i.e. the outer one. When it reaches such conditions, the inner one
 727 consequently is forced to have a greater slippage and reaches the high slip values (≈ 0.17)
 728 that are beyond the maximum of Pacejka's magic formula. Regarding sideslip angles, tyres
 729 are never pushed beyond the maximum of the magic formula thanks to the optimality condi-
 730 tion. Moreover, only the front tyres go slightly beyond the maximum of lateral adherence
 731 (at $\lambda \approx 0.17$ rad). This is not in contradiction with the optimality condition because lat-
 732 eral force does not decrease suddenly after the peak (see Figure 4); moreover, keeping the
 733 front tyres at a higher steering angle helps the vehicle to brake. The rear tyres instead are
 734 limited to lower sideslip angles because their lateral slippage is due only to kart drift angle.
 735 However, thanks to the higher sideslip stiffness, the rear tyres get extremely close to their
 736 maximum lateral adherence (at $\lambda \approx 0.15$ rad). Finally, the last graph reported in Figure 7

737 shows that engine torque represents the limiting factor when traction is required and never
 738 exceeds the maximum available power.

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741 5. Turn analysis

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In order to better understand kart dynamics, this section focuses on in-depth analysis of the vehicle behaviour while running along a particular curve of the entire track, the clockwise one comprised between 120 and 215 m of the curvilinear abscissa s . This is the turn within black box in Figure 5; the details of the track and the optimal trajectory are shown in Figure 8: the kart reaches the external track borders at the entrance of the curve, then it passes at the apex in the middle of the turn, finally, after exiting the curve, it remains in a central position before approaching the next, counterclockwise turn.

The speed profile in Figure 9 shows that the kart starts from a speed of ≈ 85 km/h and brakes to ≈ 45 km/h with a maximum longitudinal deceleration $a_x \approx -5.5$ m/s²: even if the estimated longitudinal tyre friction is equal to 0.91, the maximum deceleration is significantly lower than 0.91 g because the kart has brakes on the rear axle only. Braking deceleration remains nearly constant to its maximum until $s \approx 165$ m, then for $165 < s < 185$, a_x grows up to its maximum value of 4 m/s². Such acceleration is kept nearly constant while exiting the turn. Lateral acceleration also is kept almost constant at ≈ 13 m/s² while turning ($150 < s < 200$). Steering angle, drift angle and yaw rate trends highlight that a small *pendulum* manoeuvre is exploited before the beginning of the turn. Summarising, the manoeuvre can be dividend into three phases that can be

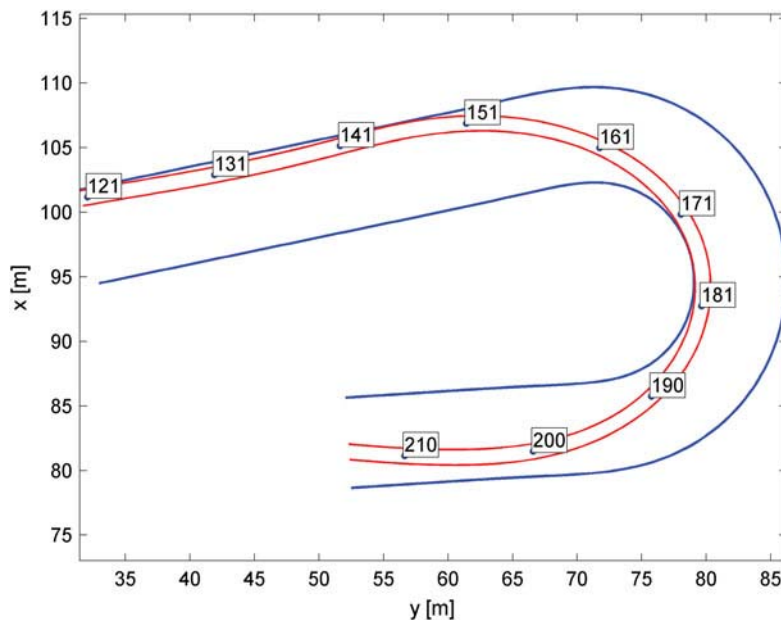
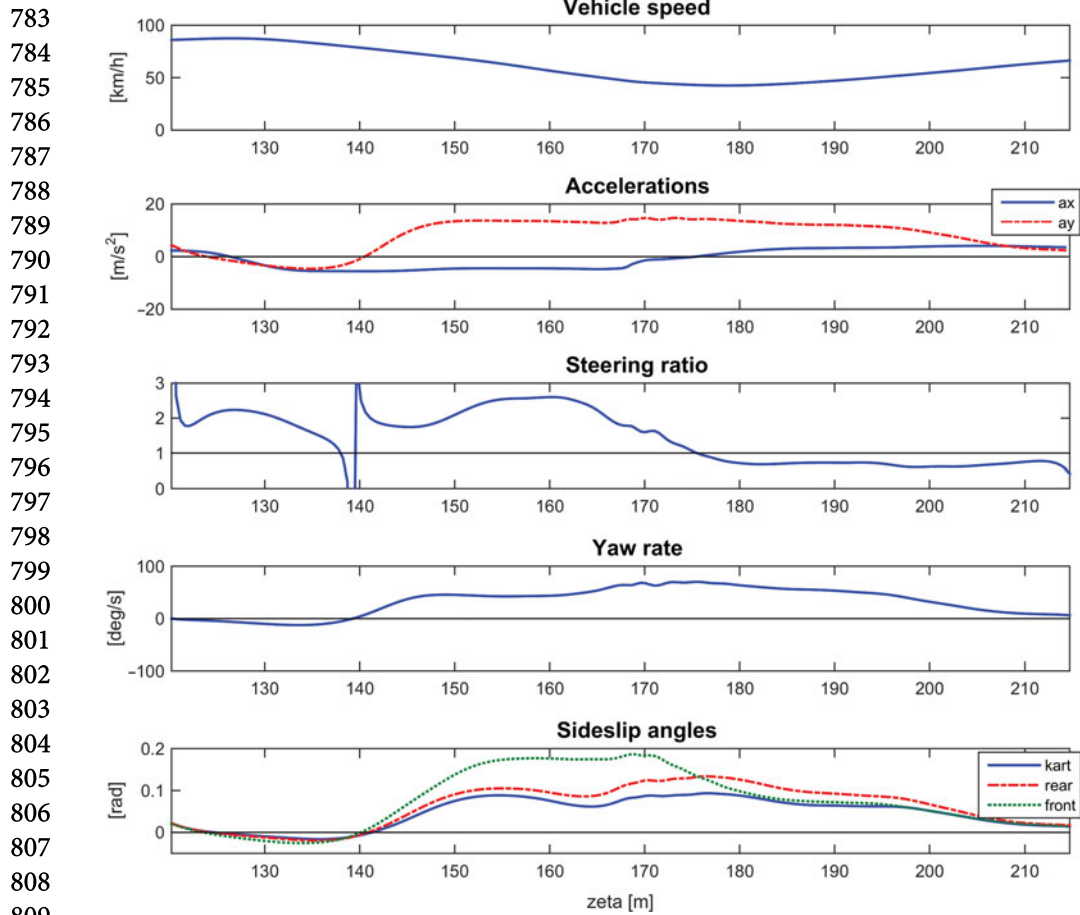


Figure 8. Turn analysis: optimal trajectory. Red lines represent kart edge trajectories, numbers in rectangular box show the value of the curvilinear abscissa s . The location of the turn in the circuit is shown in Figure 5.



810 **Figure 9.** Turn analysis: gross-motion. Steering ratio shown in the chart *d* is the ratio between the real
811 curvature radius and the kinematic one (that is the radius of curvature at zero slip). Values higher than 1
812 indicate that the vehicle is under-steering, lower than 1 over-steering. Rear (front) sideslip angle in figure
813 (e) is the average of the rear (front) left and right tyres.

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816 roughly associated with braking ($125 < s < 165$), cornering ($165 < s < 185$) and accel-
817 erating ($185 < s < 210$). The ratio between the real curvature radius and the kinematic
818 one (steering ratio) is lower than one after the second half of the manoeuvre, when the
819 vehicle is accelerating, and it is greater than one during all the rest of the time except for a
820 short transient at turn entrance in correspondence with the *pendulum* manoeuvre. When
821 the yaw rate passes through zero the steering ratio diverges to infinity since the kinematic
822 curvature radius is zero. The under-steering behaviour of the kart at the entrance of the
823 turn is confirmed also by Figure 9(e), which shows that the front tyres sideslip angle is
824 almost the double of that of rear ones. The tyre slip and forces that are generated in this
825 manoeuvre are shown in Figure 10. In the braking phase, longitudinal slips are obviously
826 negative for both tyres and they are near the value at which the maximum of the (uncoupled)
827 longitudinal force is exerted. The reason they do not reach such a value is because
828 the longitudinal force is coupled with the lateral one that is also exerted at the same time.

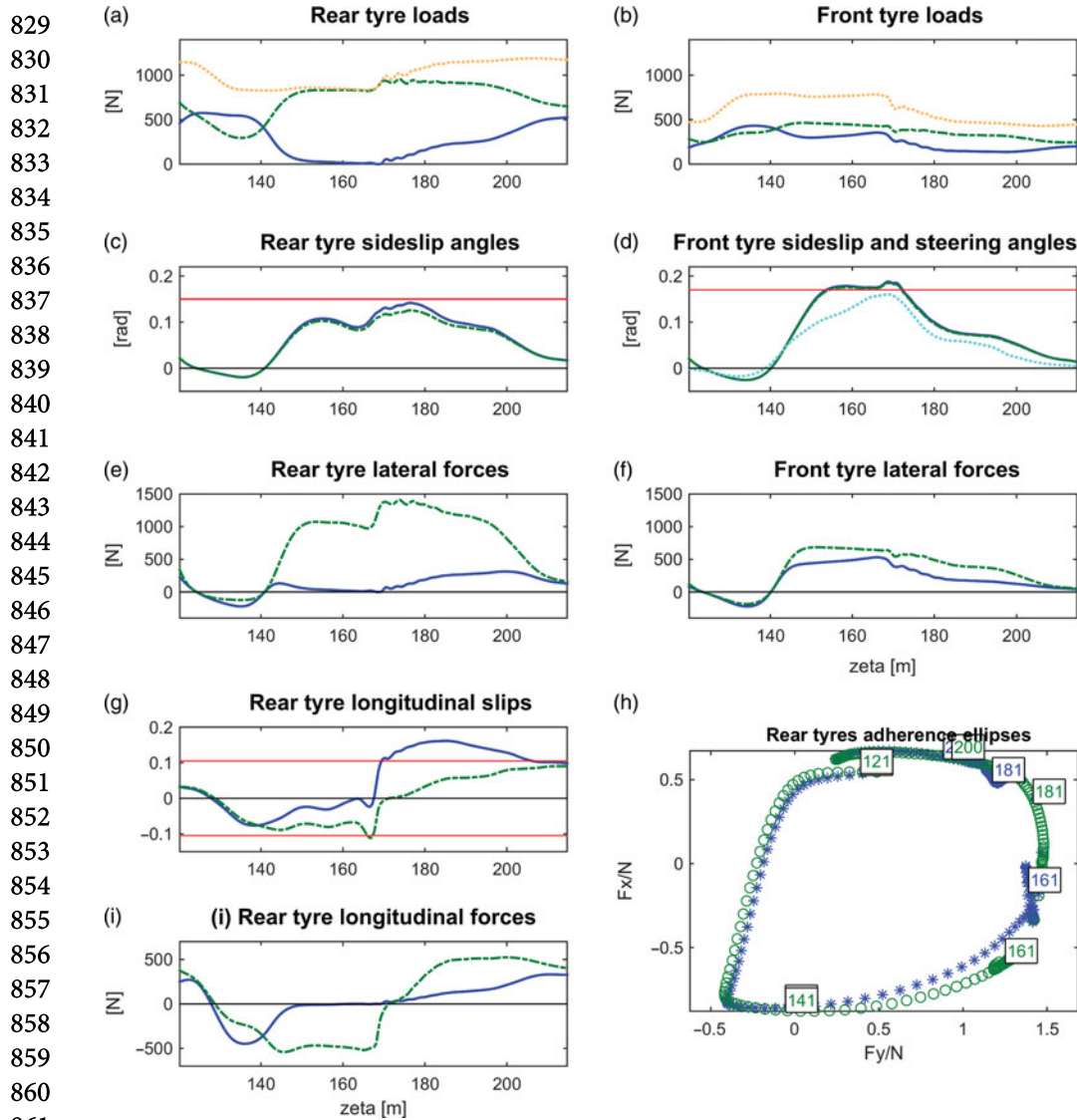


Figure 10. Turn analysis: tyres kinematic. Continuous blue lines refer to right tyres, dash-dot green ones to left tyres. Dotted orange lines, when present, show the sum of left and right tyres. Horizontal red lines in figure (c), (d) and (g) represent the value at which the maximum of tyre adherence is reached for a non-combined tyre force at nominal load. In plot (d), the dotted cyan line represents the steering angle. In figure (h), asterisks refer to right tyre, circles to left one; numbers within a rectangular box show the value of the curvilinear abscissa s of the point. Turn manoeuvre begins from the top of the picture and then continues counterclockwise. Tyre vertical loads in (a) and (b) clearly highlight how the right rear tyre (inner tyre) gets completely unload during the manoeuvre.

Then, while turning, the longitudinal slip of the two tyres shows significant differences due to the yaw motion (the absence of differential forces the wheels to have the same spin) and, since the yaw is positive, the right tyre has always a higher slip. When the kart has exited from the turn and the yaw rate is nearly zero, the longitudinal slip of the two tyres becomes equal again. Regarding the lateral slippage, the sideslip angles of rear tyres start increasing

875 from zero when the kart enters the curve, then they reach the maximum at the middle
876 of the turn and then decrease and become zero again when the kart has exited the turn.
877 The rear tyres never reach the maximum of pure lateral adherence. Concerning front tyres,
878 it can be noticed that the difference between right and left load is significant but as huge
879 as for rear tyres. This is a clear evidence that the steering effect transfers a considerable
880 load to the front inner wheel. The sideslip angles of front tyres have a trend that is almost
881 proportional to that of the steering angle (Figure 10(d), cyan line).

882 Adherence ellipses shown in the bottom-right corner of Figure 10 confirm what has
883 been said in the previous chapter: in every cornering only the outer tyre has the combina-
884 tion of longitudinal slip and sideslip angle engaging it to the limits of adherence ellipses.
885 Indeed, red points are arranged over a bigger ellipsoidal shape.

886 Finally, from the comparison of rear tyres longitudinal forces in Figure 10 with the steer-
887 ing ratio shown in Figure 9, it can be noticed that, at least as rough approximation, the
888 vehicle is under-steering when braking, and over-steering when accelerating. Keeping in
889 mind the longitudinal tyre force trend, this behaviour is reasonable: since most of the lon-
890 gitudinal force is exerted by the outer tyres (inner is less loaded), when braking such force
891 generates an aligning yaw torque, and a turning torque when accelerating.

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894 **6. Conclusions**

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895 In this work, we demonstrated the development of a multibody model for go-karts. First,
896 we presented the mathematical formulation of the multibody system together with optimal
897 control formulation. An indirect method was used to solve the optimal control problem
898 which allowed us to perform simulations. The model, fed with kart characteristics mea-
899 sured by lab tests, has been demonstrated to reproduce enough accurately the telemetry
900 recorded during the experimental track test with an expert driver. Speed and accelerations
901 measured resulted in very good agreement with simulated signals. Simulation results have
902 been exploited to study the dynamics of karts with particular attention to tyres dynamics.
903 We have been able to reproduce racing manoeuvres that bring to the lifting of the rear inner
904 wheel when turning. The importance of such manoeuvres due to the lack of differential has
905 been highlighted too.

906 With this work we developed a powerful tool that can be used to optimise kart charac-
907 teristic (such as frame and anti-roll bar stiffness) when designing such vehicles ; it allows
908 us to investigate how they affect vehicle manoeuvrability before to execute any track test.
909 Moreover, this model can be modified and used as a starting point for other four-wheeled
910 vehicles for optimal control problem. Even if in the literature other authors have presented
911 OCP simulations for cars [1] and Formula1 cars,[2] this work is among the first to calculate
912 the tyre loads from the vehicle dynamics and tyre deformations.

913 Further perspective can be the parametric analysis and optimisation of structural
914 characteristic.

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917

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the experimental activities.

921 **Disclosure statement**

922 No potential conflict of interest was reported by the authors.

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Appendix. Tyre parameters**Table A1.** Pacejka's tyre magic formula coefficients.

Coefficient	Rear tyre	Front tyre
p_{Cx1}	2.3	
p_{Dx1}	0.9	
p_{Ex1}	0.95	
p_{Kx1}	20	
p_{Kx2}	1	
p_{Kx3}	−0.5	
p_{Cy1}	2.3	2.13
p_{Dy1}	1.5	1.5
p_{Ey1}	0.9	0.8
p_{Ky1}	−37.6	−34.1
p_{Ky2}	1.6	1.6
p_{rBx1}	14	
p_{rBy1}	12	
p_{rCy1}	0.6	
λ_{Fz0}	1.6	1.6

All Pacejka's coefficients ([9], chapter 4) that are not listed here have been set to their neutral value (either 0 or 1).

Q2

Q3