

Hydrogen desorption kinetics from H-Si (111) surfaces studied by optical sum frequency generation and second harmonic generation

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Q3 We have studied hydrogen desorption from a flat H-Si (111)1 × 1 surface at 711 K by observing sum frequency generation (SFG) and second harmonic generation (SHG) spectra. Flat H-Si (111) surfaces were prepared by dosing hydrogen molecules in a ultra-high vacuum chamber with a base pressure of 10⁻⁸ Pa. Combining the SFG and SHG methods, the desorption order has been clarified over the whole hydrogen coverage range from 1 monolayer (ML) to 0 ML. The hydrogen desorption was assigned as second order in the high coverage range of 1 ML–0.18 ML by using SFG spectroscopy and as first order in the coverage range of 0.18 ML–0.0 ML by using SHG spectroscopy. Copyright © 2016 John Wiley & Sons, Ltd.

Keywords: sum frequency generation (SFG); second harmonic generation (SHG); Si (111)1 × 1 surface; desorption order; hydrogen coverage

Introduction

Hydrogenated flat Si (111) surfaces are attractive substrates for silicon-based devices. The hydrogen desorption process from the Si surface plays an important role in the epitaxial growth of Si by chemical vapor deposition.^[1] For this reason, the hydrogen desorption process has been investigated by various methods such as laser-induced thermal desorption (LITD),^[1,2] second harmonic generation (SHG) spectroscopy,^[3] *in situ* surface infrared (IR) spectroscopy,^[4] scanning tunneling microscopy,^[5] and temperature-programmed desorption,^[6] and so on. The hydrogen desorption from a Si (111)7 × 7 surface was studied by using LITD by B.G. Koehler and his coworkers. They suggested that the desorption order of hydrogen molecules above 0.2 monolayer (ML) displayed second-order kinetics.^[1] Gupta *et al.* studied hydrogen desorption by temperature-programmed desorption, and they revealed that on the H-Si (111)1 × 1 surface hydrogen atoms started to desorb at 720 K and completely desorbed at 800 K.^[7] They also suggested that the desorption of hydrogen molecules for monohydride species follows second-order kinetics. These suggestions are consistent with other groups such as B.G. Koehler *et al.*^[1] Hien *et al.* studied isothermal hydrogen desorption from a H-Si (111)1 × 1 surface prepared by chemical etching by using sum frequency generation (SFG) spectroscopy at different temperatures. They suggested that the hydrogen desorbed homogeneously, but they could not decide the desorption order at hydrogen coverage below 0.2 ML,^[8] because at the coverage lower than 0.2 ML, the resonant SFG signal was indistinguishable from the non-resonance background. The surface coverage corresponding to ~1 ML is the maximum surface density of the hydrogen atoms on the Si (111) surface, namely, 1 ML = 7.8 × 10¹⁴ (atoms/cm²).^[9,10]

The aforementioned works show that hydrogen desorption from the single crystal Si (111) surface has been studied with interest for

a long time. However, some unanswered problems about the desorption order still remain, especially at low hydrogen coverage. The desorption order for coverage below 0.2 ML over a range of temperature of 680 K < T < 800 K was investigated using SHG by Reiders *et al.* only.^[3] They found that the result can be adequately characterized by an intermediate reaction order of 1.5 ± 0.2.

In our study, we investigated the hydrogen desorption mechanism from a flat H-Si (111)1 × 1 surface at 711 K by using both SFG and SHG spectroscopy. We have used SFG spectroscopy for investigating the hydrogen desorption at high hydrogen coverage, and SHG for low hydrogen coverage on the hydrogenated Si (111) surfaces. In our SHG experiment, the sample was heated for a much longer time than Reiders *et al.* in order that the fitting of the desorption curve should be much more exact.^[3]

Nonlinear optical spectroscopies such as SFG and SHG techniques are used as useful nondestructive and sensitive tools to study the properties of surfaces and interfaces. SFG is a second-order nonlinear optical process. In SFG, two photons with different frequencies of visible (ω_1) and IR (ω_2) interact with a nonlinear medium simultaneously to generate a photon with the sum of

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the two frequencies $\omega_1 + \omega_2$. SHG is a special case of SFG with $\omega_1 = \omega_2 = \omega$ and a photon with the frequency of 2ω is generated. In this study, SFG spectroscopy was used to observe the vibrational spectrum of an H-Si surface before and after heating at 711 K. A SFG signal could be obtained with hydrogen coverage from 1 ML to 0.2 ML. When hydrogen coverage became low (~ 0.2 ML), the SFG signal was close to the background, and the vibrational mode could not be seen. On the other hand, SHG was very sensitive to dangling bonds on the surface.^[3] Therefore, we applied SHG spectroscopy to measure the remaining hydrogen coverage when the coverage was low. Then, the desorption order could be clarified in the whole coverage range from 0 ML to 1 ML.

Experimental setup

Si (111) substrates ($25 \times 5 \times 0.1$ mm³ in dimension) were cut from N type flat Si wafers. The resistivity of the samples was $1 \sim 5$ Ω cm. Sample preparation and hydrogen termination were reported in detail by Hien *et al.*^[8] Here, we will describe in detail the SFG and SHG measurements. The optical system is shown in Fig. 1. For SFG measurement, we used doubled-frequency pulses at wavelength 532 nm with photon energy of ~ 2.33 eV generated by a mode-locked Nd³⁺:YAG laser operating at a repetition rate of 10 Hz and a pulse width of 30 ps as the incident visible light. Tunable infrared light pulses output from an optical parametric generator with an amplifier (optical parametric generator/optical parametric amplifier) system were used as the incident IR light. The incident visible light was passed through a Glan polarizer, a band pass filter, a lens with focal length $f = 300$ mm, and a CaF₂ window of the UHV chamber with a pulse energy of ~ 15 μ J/pulse. The IR light was focused by a CaF₂ lens with a focal length of $f = 250$ mm with the pulse energy of ~ 70 μ J/pulse at the sample. The angles of incidence of the visible and IR light beams were $\sim 45^\circ$ and 60° , respectively. A delay line was used to adjust the temporal overlap of the IR and the visible pulses at the sample. The SFG light generated from the sample in the reflective direction was passed through a glass window of the chamber, band pass filters (Asahi SV0490), a polarizer plate (Sigma Koki, SPF-30C-32), and finally the SFG signal was focused onto the monochromator entrance by a lens and was detected by a photomultiplier. The SFG signals were obtained

as a function of the IR light wavenumber. The SFG spectra were measured from 2060 to 2110 cm⁻¹ with a scanning step of 1 cm⁻¹. The acquisition time for one SFG data point was ~ 30 s and that for one SFG spectrum was ~ 25 min. The typical photon count rate at the SFG peak was ~ 10 photons per second, and the energy resolution for SFG spectra was ~ 2 cm⁻¹. Each measurement was conducted in the polarization combination of *ppp* (SFG in *p*-polarization, visible in *p*-polarization and IR light in *p*-polarization). After heating for 10 s at 711 K, the sample was cooled down to room temperature, and the SFG spectrum was taken. This procedure was repeated after total heating time of 20, 30, 40 s... and up to 230 s. The coverage of hydrogen was calculated from the resonant SFG signal as a function of the heating time.

In the SHG measurement, we used a fundamental wavelength of 1064 nm with photon energy 1.17 eV from the same Nd³⁺:YAG laser. The incident laser light pulse with an energy of ~ 380 μ J/pulse was passed through a half wave plate ($\lambda/2$), a Glan polarizer, a color filter, a lens with focal length $f = 250$ mm, and the CaF₂ window of the UHV chamber and finally reached the sample. A color glass 2ω cut filter was placed between the polarizer and the sample to block unwanted SHG background light from the optics generated prior to the interaction with the sample. The SHG light generated from the sample in the reflective direction was passed through a glass window of the chamber, a focusing lens with focal length $f = 300$ mm, and ω cut color glass filters to block the fundamental radiation light beams before coming into the monochromator. A polarizer plate was put before the double monochromator and a photomultiplier to select the polarization of the SHG signal. When the SFG signal became comparable with the background at the lower hydrogen coverage, we switched to SHG measurement and detected the Si dangling bonds. The sample was heated for 50 s at 711 K and then cooled down to room temperature, and the SHG signal was taken. This procedure was repeated after total time heating of 230, 280, 330, and 380 s...up to 3880 s. Then we heated the sample for different interval of times and measured SHG signal up to a total heating time of 18330 s.

Results and discussions

Sum frequency generation response

In order to confirm the hydrogen desorption order of H-Si (111) 1×1 surfaces with high hydrogen coverage, we investigated the time dependence of isothermal desorption at temperature 711 K by taking the SFG spectra. Figure 2 shows a typical SFG spectrum of the H-Si (111) 1×1 surface at room temperature. The peak

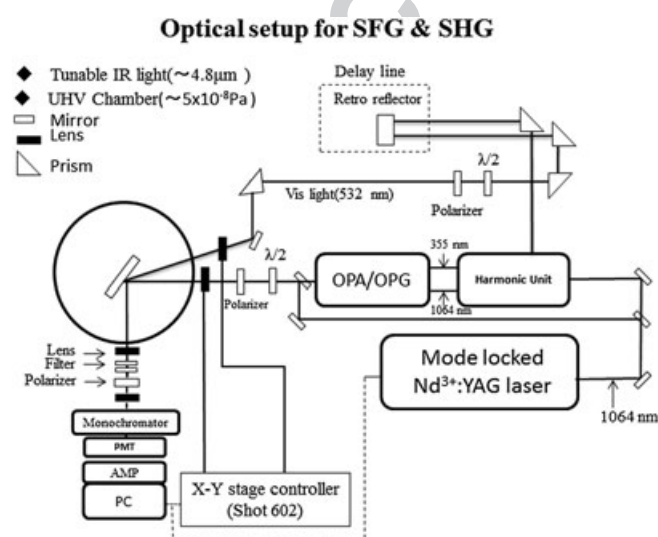


Figure 1. A schematic diagram of the sum frequency generation and second harmonic generation spectroscopic system.

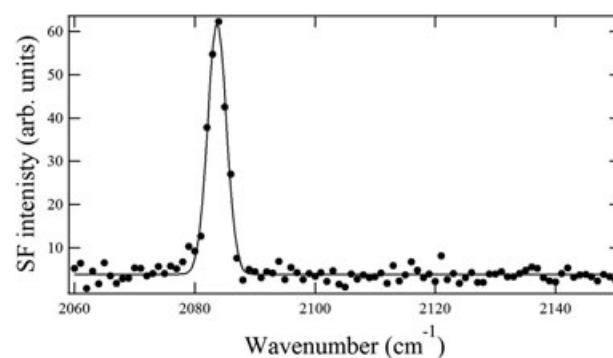


Figure 2. Sum frequency generation spectrum of the H-Si (111) 1×1 surface at room temperature. A sharp peak appears at 2083.7 cm⁻¹.

at 2083.7 cm^{-1} is assigned to the stretching vibration of monohydride on the Si surface. This result is consistent with the literature.^[8,11] In general, the hydrogen coverage θ can be estimated approximately from the peak height of SFG peaks by the following expression as

$$\theta \propto \sqrt{I_{\text{SFG}}} \propto \chi^{(2)} \quad (1)$$

Here, I_{SFG} is the peak height and $\chi^{(2)}$ is the nonlinear susceptibility. However, if there is interaction between Si-H oscillators, the coverage θ is not proportional to the value of $\chi^{(2)}$ anymore.^[12] Y. Miyauchi *et al.* proved that there was dipole coupling among Si-H oscillators on the flat Si (111) 1×1 and calculated the relation between θ and $\chi^{(2)}$ based on the coherent potential approximation method.^[13] In our paper, we will calculate the hydrogen coverage with respect to SFG signal ($\chi^{(2)}$) following Miyauchi's report.

The hydrogen coverage reduction on the Si surface in the desorption kinetics is given by using the Polanyi–Wigner desorption rate equation,^[1,8]

$$-\frac{d\theta}{dt} = \vartheta_d \theta^n e^{-E_d/RT} \quad (2)$$

Here, θ is the surface coverage, ϑ_d is the pre-exponential factor, E_d is the activation energy for desorption, R is the gas constant (8.31 J/mol K), and T_{surf} is the surface temperature. When $n = 1, 1.5$, and 2 , the solutions of desorption rate equation (1) becomes

$$\text{First-order desorption: } \theta_t = \theta_0 e^{-k_1 t} \quad (3)$$

$$\text{1.5th-order desorption: } \theta_t = \theta_0 \left(1 + \sqrt{\theta_0} k_{1.5} t\right)^{-2} \quad (4)$$

$$\text{Second-order desorption: } \theta_t = \theta_0 (1 + \theta_0 k_2 t)^{-1} \quad (5)$$

F3 Figure 3 represents the time dependence of the hydrogen coverage at the heating temperature of 711 K . The hydrogen coverages during the isothermal desorption were fitted to the first-order and second-order theoretical curves. The solid dots are hydrogen coverage corresponding to the SFG intensities. The reduction of hydrogen coverage from 1 ML to 0.18 ML in Fig. 3 shows that the second order is the best fitted data with the coverage larger than 0.4 ML as in other reported results. This result is consistent with M. L. Wish *et al.*^[14] using LTD method with the surface temperature of 725 K . In that study, the hydrogen coverage reduced to 0.2 ML after 200 s . In the second-order process, one hydrogen atom leaves a Si atom and diffuses toward another Si-H site, and they combine

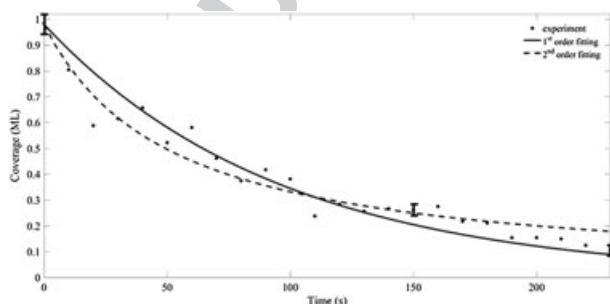


Figure 3. Isothermal hydrogen desorption from the H-Si (111) 1×1 surface at surface temperatures of 711 K heated for 10-s intervals. The solid dots are experimental results, and error bars represent the standard deviations. The solid and dashed lines correspond to the first-order and second-order desorption kinetics.

with each other to form a dihydride (Si-H_2). For a while, the dihydride state remains. Finally, the hydrogen atoms go beyond the highest potential barrier in the reaction coordinate, associate themselves with each other and desorb from the Si atom.^[15] In this way, the hydrogen re-combinative desorption occurs at a single Si atom. Our investigation confirmed that the hydrogen desorption can be assigned as second order as in the literature^[3] in the coverage range of 1.0 ML – 0.18 ML .

Second harmonic generation response

We have used SFG spectra to investigate hydrogen desorption at the high hydrogen coverage from 1 ML to 0.18 ML . Because SFG signal is unobservable at lower hydrogen coverage, the SHG spectroscopy is a powerful tool to detect hydrogen coverage below 0.18 ML . Reider *et al.* proved that SHG is sensitive to dangling bonds, especially when hydrogen coverage is lower than 0.3 ML .^[3]

After the SFG experiment, we continued the hydrogen desorption for the same sample as earlier and started the SHG measurement. Figure 4 shows the time dependence of SHG intensity of the H-Si (111) 1×1 surface when the surface was heated for 50-s intervals at 711 K . In this experiment, we used the polarization configuration $P_{\text{in}} P_{\text{out}}$. The SHG intensity initially increased rapidly as a function of heating time and then gradually saturated when the number of dangling bonds were saturated (Fig. 4).

At the lower coverage, the quenching of the surface susceptibility $\chi_{s(2)}^{(2)}(\theta)$ due to adsorbed hydrogen linearly depends on the coverage (θ) by the following equation.

$$\chi_{s(2)}^{(2)}(\theta) = \chi_{s(2)}^{(2)}(0) (1 - \alpha\theta), \text{ where } \theta \ll 1 \quad (6)$$

In the case of H-Si (111) 7×7 , the slope is $\alpha \approx 1.3$.^[16] Similarly, hydrogen adsorption on H-Si (100) 2×1 gives a proportionality constant of $\alpha \approx 3.1$.^[17] On the other hand, in our case of H-Si (111) 1×1 , the proportionality constant of $\alpha \approx 5.08$ was obtained by using Equation (6) and Fig. 4. Here, $\chi_{s(2)}^{(2)}(\theta)$ is proportional to the square root of SHG intensity $\sqrt{I_{\text{SHG}}}$, and $\chi_{s(2)}^{(2)}(0) \approx 16.47$ is the susceptibility at $\theta \sim 0\text{ ML}$ obtained from Fig. 4, and α was determined at $\theta = 0.18\text{ ML}$. Figure 5 shows the hydrogen coverage reduction with respect to the heating time from the H-Si (111) 1×1 surface, calculated via Equation (6). The initial coverage in the SHG measurement was 0.18 ML . The hydrogen coverage during the isothermal desorption was fitted with the first-order (1st), intermediate-order (1.5th), and second (2nd)-order theoretical

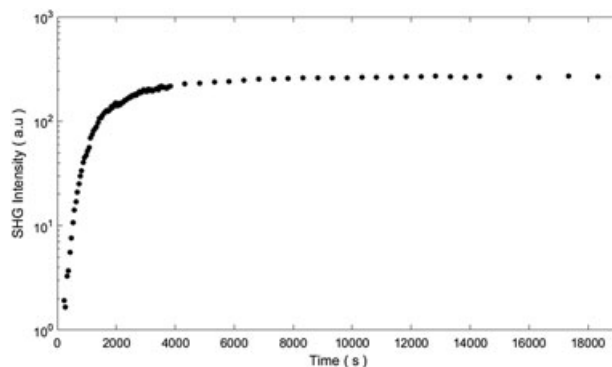


Figure 4. The second harmonic generation (SHG) intensity of the H/Si (111) 1×1 surface with respect to heating times. Excitation light wavelength is 1064 nm , the polarization of incident light and SHG light was $P_{\text{in}}/P_{\text{out}}$. The solid dots are experimental SHG intensities.

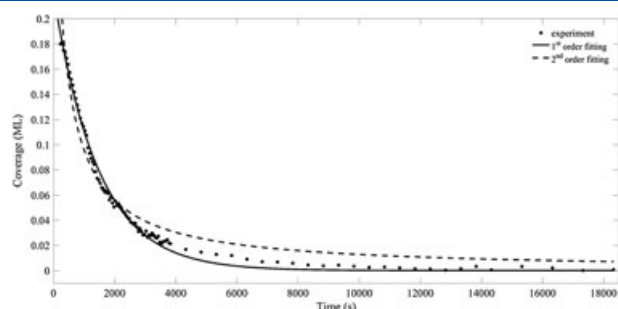


Figure 5. Isothermal hydrogen desorption from the H-Si (111)1 × 1 surface at surface temperatures of 711 K by using second harmonic generation spectroscopy. The solid dots are experimental hydrogen coverage and solid and dashed lines correspond to the first-order and second-order desorption kinetics.

curves using Equations (3), (4), and (5), respectively. The reduction of hydrogen coverage in Fig. 5 fits the first-order curve best. In the next section, we will try to explain the mechanism of the first-order desorption.

Hydrogen desorption model

Isothermal measurements for H desorption from Si (111) surfaces have been studied by several authors using different methods.^[1,7] They proposed a model that H desorption from Si (111) was re-combinative second-order desorption for monohydride Si (111) surfaces. Indeed, our results confirmed that the hydrogen desorbed in second-order kinetics at the coverage range from 1 ML to 0.18 ML. We agree that at the higher coverage of hydrogen than 0.18 ML, the hydrogen desorption from the order phase of 1 × 1: H obeys the second-order desorption. Namely, one hydrogen atom leaves a Si atom and diffuses toward another Si-H site and then they combine to form a dihydride (Si-H₂). Finally the molecular hydrogen desorbs from the Si atom.

For the first-order desorption, S. Ciraci *et al.* proposed that the hydrogen desorption occurs from dihydride species on the Si (100) surface. Two hydrogen atoms re-combined from the two adjacent silicon dihydride species.^[18] M. L. Wish *et al.* proposed another model for the first-order desorption from Si (100)2 × 1 surfaces. In this case, hydrogen desorption occurs from two hydrogen atoms paired on the same single dimer of Si-Si.^[14]

As seen in Fig. 5, the hydrogen desorbed in the first-order kinetics at the lower coverage from 0.18 ML–0.0 ML as a function of the heating time. In order to explain the mechanism of the hydrogen desorption, we suggest two candidate models:

Similarly to Y. Morita *et al.*, let us assume that there exist small islands of Si atoms and hydrogen atoms on the surface. After heating for several hundreds of seconds, these islands become crystallized into two-dimensional ((2D) islands containing one double layer of Si (111)1 × 1 atoms terminated by monohydride.^[5] During the hydrogen desorption from the surface, each Si atom from the 2D island carrying one hydrogen broke its three back bonds (called ≡Si-H species) and immigrated out of the islands and combined with three Si atoms on the surface in order to reduce the number of dangling bonds. Let us imagine now that there were two types of Si-H bonds on the surface. One is remaining monohydride Si-H from 1 × 1 phase created by hydrogen dosing. The other is ≡Si-H species coming from 2D islands. When the number of monohydrides Si-H on the 1 × 1 phase is large (1 ML), the hydrogen desorption on the 1 × 1 phase is dominant because

the distance between H-H atoms on the 1 × 1 phase is shorter 3.8 Å. Therefore, the second-order desorption is reasonable in spite of the co-existence of small amount of ≡Si-H species. After most of the monohydride Si-H desorbed, the coverage becomes lower and the surface contains H atoms on the 1 × 1: H phase and from the ≡Si-H species.

As in the paper of Y. Morita *et al.*,^[5] we define the coverage of ≡Si-H species as d_{adatom} (ML) and the coverage of remaining monohydride Si-H as $d_{\text{SiH}} = \theta$ (ML). The total remaining hydrogen coverage is $d_{\text{adatom}} + d_{\text{SiH}} = \varnothing$ (ML). Then, the desorption order of hydrogen at low coverage becomes smaller than second order following the equation

$$\frac{d\theta}{dt} = -\frac{k_d \theta^2}{(1 - 3\varnothing) + 3\theta} \quad (7)$$

From this equation, if $\varnothing = 1/3$, Equation (7) becomes Equation (3), first-order desorption. This suggestion of Morita *et al.* is consistent with our experimental results. In our case, when the total coverage of monohydride and ≡Si-H species is 0.18 ML, the experimental data is best fitted with the first-order desorption. Therefore, when the surface was heated at 711 K, very small amount of 1 × 1: H atoms existed within the disordered phase, the number of ≡Si-H species became large, and we expect the high possibility of the H-H recombination there.

We suggest another model for the first-order desorption. After hydrogen adsorption, the surface might have the remaining local areas of Si (111)7 × 7 or there might be small surface defects, and thus Si-H₂ and Si-H₃ might exist. G. Schulze *et al.* suggested that the dihydride and trihydride should desorb from 573 to 770 K while monohydride should desorb from 680 to 900 K.^[19] In our case, the surface was heated at 711 K. In that case, at very low hydrogen coverage, two hydrogen atoms of a Si-H₂ can combine with each other and desorb in the first-order kinetics. Previously, we proved in the SFG spectrum that there was no dihydride and trihydride on the surface.^[8] But in that case, the amount of these SiH₂ and SiH₃ was too small, and SFG could not have detected their vibration. This is just like our result in Fig. 3 that, when the number of monohydride becomes too small due to desorption, SFG could not have detected its vibration. Nevertheless, this suggestion is consistent with another report of A. Ichimiya *et al.*,^[20] studied by reflection high energy positron diffraction.

Conclusions

In this study, we have investigated the hydrogen desorption mechanism from the H-Si (111)1 × 1 by using SFG spectroscopy and SHG spectroscopy. Isothermal desorption was observed at a temperature of 711 K. We suggested that the hydrogen desorption was assigned as second order in the coverage range 1.0 ML–0.18 ML by SFG, and the hydrogen desorption was assigned as first order in the coverage range 0.18 ML–0.0 ML by SHG. We have not yet calculated the hydrogen desorption activation energy for the first-order desorption from 0.18 ML to 0.0 ML in this experiment. So further experiments at different temperatures are required. From that result, the mechanism of the first-order desorption will be determined.

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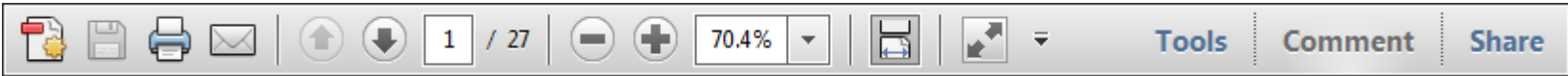
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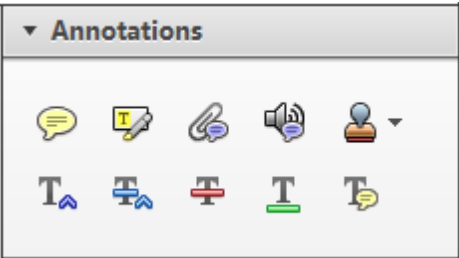
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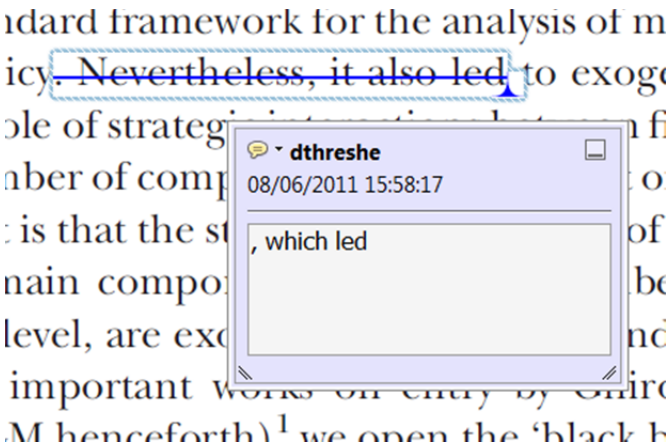
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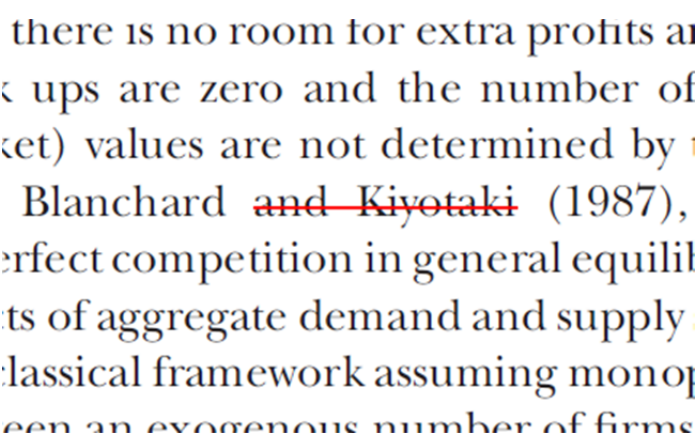
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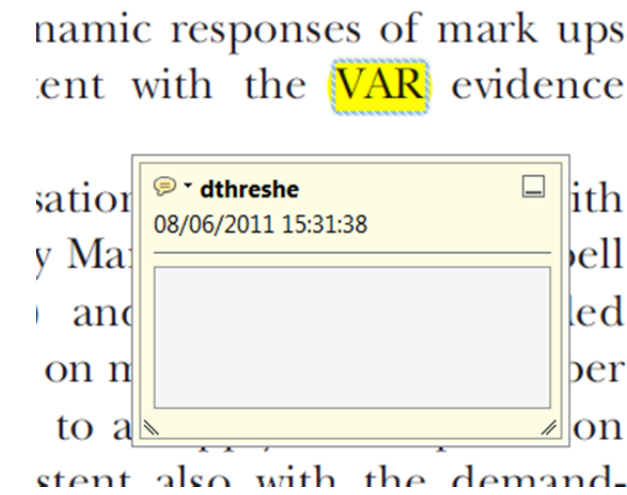
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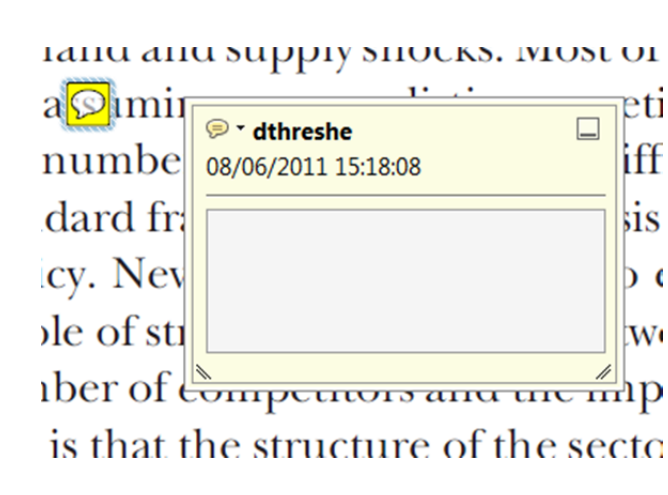
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
Marks a point in the proof where a comment needs to be highlighted.

How to use it

- Click on the [Add sticky note](#) icon in the Annotations section.
- Click at the point in the proof where the comment should be inserted.
- Type the comment into the yellow box that appears.

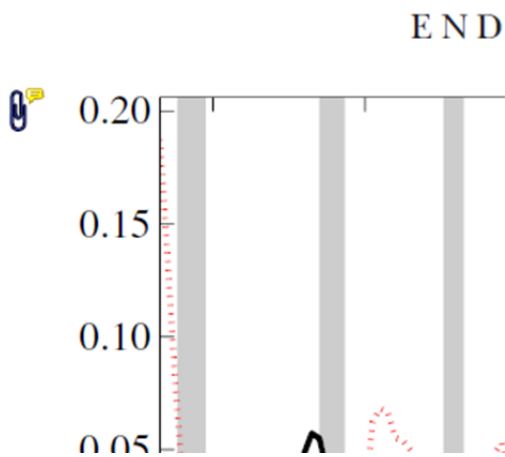


5. **Attach File** Tool – for inserting large amounts of text or replacement figures.


 Inserts an icon linking to the attached file in the appropriate place in the text.

How to use it

- Click on the **Attach File** icon in the Annotations section.
- Click on the proof to where you'd like the attached file to be linked.
- Select the file to be attached from your computer or network.
- Select the colour and type of icon that will appear in the proof. Click OK.



6. **Add stamp** Tool – for approving a proof if no corrections are required.

 Inserts a selected stamp onto an appropriate place in the proof.

How to use it

- Click on the **Add stamp** icon in the Annotations section.
- Select the stamp you want to use. (The **Approved** stamp is usually available directly in the menu that appears).
- Click on the proof where you'd like the stamp to appear. (Where a proof is to be approved as it is, this would normally be on the first page).

of the business cycle, starting with the
on perfect competition, constant returns
production. In this environment goods
extra profits and the structure of market
he number of firms in the individual firm
etermined by the model. The New-Key
otaki (1987), has introduced product
general equilibrium models with nominal
ed and supply shocks. Most of this literat

APPROVED

▼ Drawing Markups

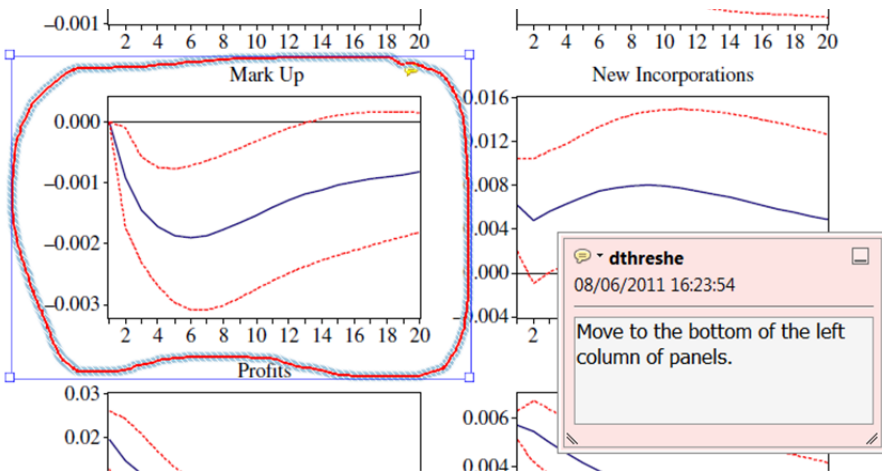


How to use it

- Click on one of the shapes in the **Drawing Markups** section.
- Click on the proof at the relevant point and draw the selected shape with the cursor.
- To add a comment to the drawn shape, move the cursor over the shape until an arrowhead appears.
- Double click on the shape and type any text in the red box that appears.

7. **Drawing Markups** Tools – for drawing shapes, lines and freeform annotations on proofs and commenting on these marks.

Allows shapes, lines and freeform annotations to be drawn on proofs and for comment to be made on these marks..



For further information on how to annotate proofs, click on the **Help** menu to reveal a list of further options:

