Low Loss Kagome Hollow Core Fibers Operating from the Near- to the Mid-IR

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We report the fabrication and characterization of Kagome hollow core antiresonant fibers which combine low attenuation (as measured at ~30 cm bend diameter) with a wide operating bandwidth and high modal purity. Record low attenuation values are reported: 12.3 dB/km, 13.9 dB/km and 9.6 dB/km in three different fibers optimized for operation at 1 μm, 1.55 μm and 2.5 μm respectively. These fibers are excellent candidates for ultrahigh power delivery at key laser wavelengths including 1.064 μm and 2.94 μm, as well as for applications in gas-based sensing and non-linear optics. © 2017 Optical Society of America

Hollow core anti-resonant fibers (HC-ARFs) form a subset of the expanding family of hollow core-photonic crystal fibers (HC-PCFs). A HC-ARF confines light in an air core in specific wavelength regions through anti-resonance with the glass membranes which define the core boundary. The spectral regions where a HC-ARF can transmit light with low loss are therefore defined by the thickness of these membranes. As well as providing low loss and octave spanning spectral guidance in multiple spectral windows, HC-ARFs have been attracting increasing interest due to their large core diameters (~typically several tens of microns and up to ~120 μm), ultralow nonlinearity and low group velocity dispersion. Due to these properties, HC-ARFs are excellent candidates for high power beam delivery, in particular of ultra-short pulses, and gas-based non-linear optics applications [1-3].

Anti-resonant guidance opens up a high degree of flexibility in fiber design and several HC-ARF structures have been reported to date [4-6]. Kagome hollow core anti-resonant fibers (K-ARFs) are the most highly developed HC-ARF structure and were initially described in 2002 [7]; further improvements in fabrication and optical performance, as well as increased insight into this type of guidance mechanism, were provided in 2007 [8]. Until recently, K-ARFs have typically been associated with attenuation at the level of ~dB/m, however substantial efforts to reduce the overlap of the core guided modes with the core surround through engineering the shape of the core boundary have resulted in dramatic loss reduction [9]. In particular, focused efforts to achieve a core boundary shape with a highly negative curvature have led to remarkable low attenuation values of 17 dB/km [10] and very recently, 8.5 dB/km [11], both around the 1 μm wavelength region. As such, the attenuation of K-ARFs at wavelengths around 1 μm or below can now compete, or even be lower than, losses achieved in hollow core photonic bandgap fibers [12]. Yet, in these record low loss fibers it has so far proven extremely challenging to combine low loss with low bend sensitivity and wide operating bandwidth. For example, the 17 dB/km K-ARF [10] has a low loss bandwidth of only ~10 nm and therefore compromises on one of the key parameters that is typically exploited in a K-ARF and the 8.5 dB/km K-ARF in [11] has high bend loss, so that the low attenuation value is measured with a 70 cm bend diameter. As such, further progress is required to simultaneously achieve very low loss over a wide spectral range and with low bend sensitivity, while producing long, uniform lengths of fiber.

In this paper, we report the fabrication of three K-ARFs, designed for operation at 1 μm [13], 1.55 μm and 2.5 μm respectively, that all have record low attenuation (as measured at ~30 cm bend diameter) while maintaining broad operating bandwidths and therefore represent a new state of the art for K-ARFs. This is achieved, in part, by operating in the fundamental transmission window, which enables a wide, low loss operating bandwidth and has the added benefit of reducing the sensitivity of the fiber to bend loss; measurements on our 1.55 μm K-ARF show this fiber demonstrates record low bend loss for a K-ARF. We also report detailed measurements of the modal properties of our 1.55 μm K-ARF, which indicate that these new K-ARFs can have a high modal purity if appropriate launch conditions are used. The

Table I: Key structural and optical properties of fibers A, B and C.

	Core diameter (μm)	Core wall thickness (nm)	Minimum loss wavelength (nm)	Minimum loss (dB/km)	3 dB bandwidth (nm)	Curvature parameter 'b'
Fiber A	43	375	1010	12.3 ± 2.2	150	0.5
Fiber B	64	600	1566	13.9 ± 1.0	198	0.57
Fiber C	97	975	2462	9.6 ± 3.0	191	0.69

combination of ultra-low loss, low bend sensitivity, large mode field diameter and high modal purity render these fibers excellent for high power laser delivery at key wavelengths including 1.064 μm and 2.94 μm, where, to the best of our knowledge, our K-ARF demonstrates the lowest optical loss achieved in any hollow core fiber reported to date.

The K-ARFs reported here all have a 7-cell core design and were fabricated using the conventional stack and draw technique. The capillary dimensions in the stack were designed so that in the final fiber, the target wavelength would be in the fundamental transmission window in order to maximize the low loss operating bandwidth and to reduce bend loss. Operation in this spectral window requires the silica membranes which surround the air core to be substantially thinner than if the fiber is designed for operation in a higher order spectral window, and this makes fabrication of a core membrane more challenging, particularly if targeting a high degree of negative curvature. Therefore, in the fiber draw process it is essential to carefully control the difference between the core and cladding pressures, as well as the draw tension, to obtain a core surround with a negative curvature shape.

Scanning electron microscope (SEM) images of the three fibers (referred to as A, B and C for low loss at 1 μm, 1.55 μm and 2.5 μm respectively) are shown in Fig. 1(a) – (c). A summary of the key properties, both structural and optical, is presented in Table I. These fibers were each drawn from separate canes (with very similar structures) and the core size scales with the wavelength of the fundamental guidance window. The attenuation spectra (calculated using cutback measurements) for the three fibers are shown in Fig. 2. For clarity, the attenuation of the fundamental
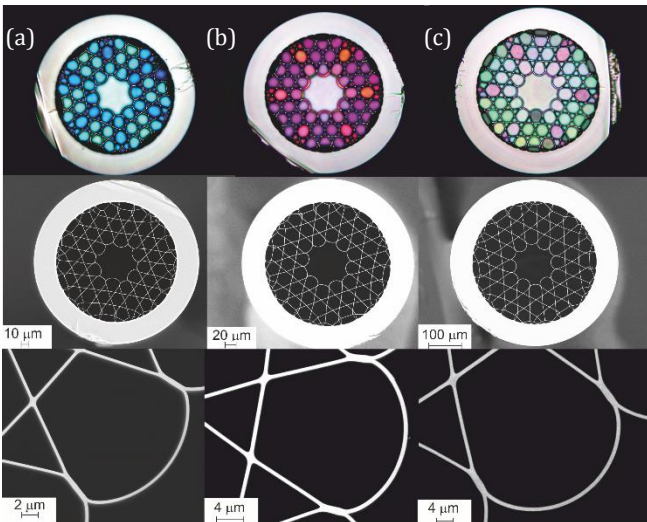

Figure 1: (a), (b) and (c) Optical microscope images (top row) and scanning electron microscope images of the full fiber (middle row) and the core surround (bottom row) of fibers A, B and C. respectively.

spectral window only is presented, but in each case the loss in the higher order windows was found to be higher than in the fundamental spectral window; for example, in fiber B, the minimum loss in the first higher order window is 62 dB/km at 878 nm. Fibers A and B were measured using a white light source (WLS) and optical spectrum analyzer (OSA). The WLS was coupled into these K-ARFs using a large mode area fiber with a similar mode field diameter to the fundamental mode of the K-ARF. The mode field diameters of our fibers were estimated to be ~73% of the core diameter [14]. Fiber C was measured using an in-house built mid-infrared supercontinuum source [15]. The fiber output of this source was butt-coupled directly to fiber C in order to carry out the loss measurements and due to a large mode field diameter mismatch (~6.5 μm for the supercontinuum output compared to ~75 μm for fiber C) this measurement does suffer from higher modal noise which is evident in the loss spectrum. A Yokogawa AQ6376 OSA was used for detection at mid-IR wavelengths. For all the loss measurements, the fibers were loosely coiled in ~30 cm diameter loops. From these measurements the minimum losses of fibers A, B and C were found to be 12.3 dB/km, 13.9 dB/km and 9.6 dB/km and each fiber has a low loss (3 dB) bandwidth of at least 150 nm. These fibers all represent record low loss values for K-ARFs (at 1 μm, considering the bend diameter in the measurement) in their respective spectral regions, and fibers B and C represent new record low loss values for any HC-ARF in their minimum loss windows. The measurement uncertainties in Table I were estimated by considering the variation in transmission spectra recorded using three different fiber cleaves performed at each length for each fiber and also the cutback fiber lengths used (63m to 5m, 148m to 15m and 100m to 10m for fibers A-C respectively). It is also noteworthy that the minimum loss of fiber C at 2.94 μm, which coincides with the Er:YAG high power laser source, is 18.6 dB/km, which we believe represents a new record low loss for any hollow core fiber at this wavelength. Furthermore, fiber C has an attenuation below 150 dB/km across a wide spectral region spanning from 2.25 to beyond 3.4 μm. The measurement of fiber C is truncated at the long wavelength edge due to the range limitation of the OSA used for detection, however measurements of shorter lengths using an FTIR show transmission up to 4.7 μm. The attenuation curve for each fiber is clearly structured as is common with K-ARFs as the resonances of the core guided modes with silica cladding nodes increase the fiber loss at specific wavelengths. The attenuation of fiber C also has loss peaks due to water vapor and carbon dioxide (~2.5–2.85 μm) and hydrogen chloride gas absorption (~3.3-3.4 μm); these species are known to be present in as-fabricated hollow core photonic crystal fibers and can be reduced or eliminated by purging the fiber with dry gas [16, 17].

Table I also shows the calculated core curvature parameter 'b' (as defined in [10]) for each fiber. This parameter increases as the core membrane thickness increases from fibers A to C, as the fundamental transmission window is shifted to longer
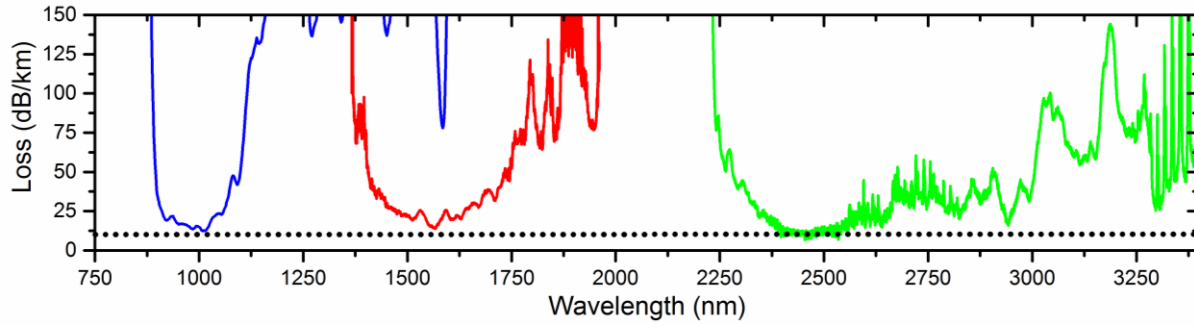
Figure 2: Attenuation spectra of fibers A, B and C (blue, red and green lines respectively) measured using the cutback technique. Dotted black line indicates loss equal to 10 dB/km.

wavelengths. In the previous state of the art K-ARFs, b is equal to 1 and 0.9 (in [10] and [11] respectively), and in these papers the authors argue that increasing b leads to reduced loss. For fibers A, B and C we can see that, although b is significantly lower than 1, the loss is still reduced. We believe that this is because the core curvature parameter cannot be considered alone when predicting the minimum loss structure; in fact the shape and size of the nodes in the core surround, in combination with the core curvature will define the minimum loss. Additionally, longitudinal uniformity along the fiber lengths is important to avoid defect-induced losses [18] and for our fibers, longitudinal uniformity was confirmed for fibers A and B using an OTDR.
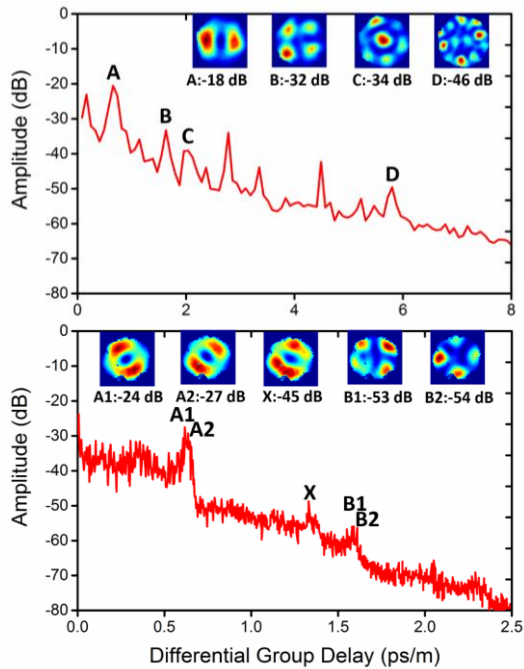

Figure 3: S^2 measurements on 5 m (top) and 133 m (bottom) lengths of fiber B. Amplitude peaks corresponding to different modes propagating in the fiber are labeled; Mode groups A and B correspond to the LP_{11} and LP_{21} modes respectively. Unlabeled peaks correspond to double reflections in the measurement set-up [18]. Modal images for each fiber length are shown inset to each plot along with total power amplitude (MPI) in each mode compared to the fundamental mode power.

Fiber B was chosen for a more detailed study as its low loss operation at 1550 nm makes it amenable for measurement with our S^2 set-up [19], which enables analysis of the modal content of the fiber in near real time. The S^2 results for a 133 m length of fiber, cutback to 5 m, are shown in Fig. 3. These measurements were taken using a free space, lens optimized launch into the input of the fiber to ensure maximum coupling into the fundamental mode and as the measurements were taken as a cutback, the input coupling conditions for both measurements are the same. The S^2 measurement bandwidth (1555 – 1575 nm) was centered on the lowest transmission loss region for this fiber, and the laser source was swept in 1 pm steps at a rate of 500 pm/s. From Fig. 3 it is clear that, as expected [20], fiber B supports several modes. However, with appropriate coupling into the fiber, only a few higher order mode groups (HOMs) are present after a 5 m length and in total the HOM content is suppressed to approximately 18 dB below the fundamental mode (corresponding to ~1.6 % of the total power being transmitted in HOMs). Note that in both plots in Fig. 3, unlabeled peaks correspond to double reflections from different components in the measurement set-up, which are discussed in detail elsewhere [21]. Over a longer length of 133 m, the higher order mode extinction is ~24 dB; while this is an increase, this clearly shows that the differential loss between the LP_{11} mode group and the fundamental mode is not especially high; from this measurement the LP_{11} mode loss is estimated to be ~ 50 dB/km. This ratio of LP_{11} mode loss to fundamental mode loss (~ 3.6) is comparable with the results of our previous study of modal content in K-ARFs [20]. One feature to note is that over the 133m length, a mode (labeled 'X') is present with a differential group delay between that of the LP_{11} and LP_{21} mode groups; we believe this arises due to a discrete event where a very small amount of power is coupled between the LP_{11} and LP_{21} modes. As the core size of our fibers scales with wavelength, we believe that similar modal properties would be recorded for fibers A and C.

Figure 4 shows the bend loss measurements for fiber B at four different wavelengths (three in the fundamental transmission window and one (850 nm) in the first higher order transmission window of the fiber). These measurements were carried out using a white light source and by comparing the transmission of a very loosely coiled 5 m length of fiber B with the transmission as tighter bends were introduced along the fiber length, without changing the input or output coupling conditions. In the lowest loss spectral region around 1550 nm, the bend loss is very low for this fiber type, becoming significant compared with the measured cutback
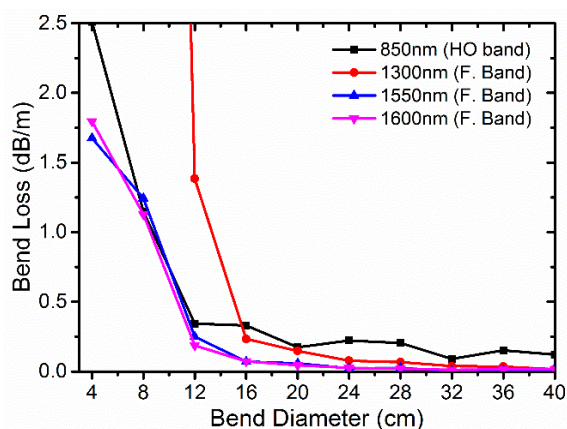
Figure 4: Bend loss measurements of fiber B at four different wavelengths indicating low bend sensitivity at 1550 nm and 1600 nm.

loss reported above at bend diameters ~ 16 cm. From comparison with state of the art fibers in the literature [22] with a similar core size and which provide low loss at similar wavelengths, we believe this sets a record low bend sensitivity for a K-ARF. Note, we have also measured the bend sensitivity of fiber A (see [13]) and this fiber also shows reduced bend loss compared to the previous state-of-the-art K-ARF operating around 1 µm [11], however the bend loss becomes significant as compared to the cutback loss at a bend diameter ~ 24 cm. Considering that the structures of fibers A and B are very similar, the origin of this difference will be investigated further. From Fig. 4, it is clear that, as expected, the bend sensitivity increases at wavelengths closer to the edge of the transmission window, where the wavelength is closer to resonance, and also that the bend sensitivity increases in the higher order transmission window (as the wavelength is reduced relative to the core diameter). Even with this record low bend sensitivity, improvement in this optical property would be beneficial to many applications, such as the fabrication of compact fiber-based gas sensors.

The fibers presented here demonstrate low loss (measured at ~30 cm bend diameter), wide operating bandwidth and large mode field diameter. Furthermore detailed measurements of our fiber optimized for operation around 1.55 µm demonstrate record low bend loss for a K-ARF and high modal purity (with appropriate launch conditions). Therefore we believe these fibers represent a new state-of-the-art for K-ARFs. In the 1 µm spectral region our fiber presents comparable loss to state-of-the-art hollow core photonic bandgap fibers [12], and at 1.55 µm, our fiber has the lowest loss for any HC-ARF. At 2.94 µm our fiber represents the lowest loss hollow core fiber of any design. These fibers are highly suited for ultra-high power beam delivery from the near- to the mid-IR. Recently, several new low loss HC-ARFs have been reported [23 - 25], and in particular 'nodeless' designs hold high promise for further loss reduction as they have highly simplified structures and remove the attenuating resonances that are present in the cladding of a K-ARF. However, the optimal structure which will provide the best complete set of optical properties is still an open question and further work on both HC-ARF structures is ongoing.

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References

1. C. Saraceno, F. Emaury, A. Diebold, I. Graumann, M. Golling, and U. Keller, in *SPIE Defense+ Security. International Society for Optics and Photonics 2016*, 9835, p. 98350X (2016).
2. B. Debord, M. Alharbi, L. Vincetti, A. Husakou, C. Fourcade-Dutin, C. Hoenninger, E. Mottay, F. Gerome and F. Benabid, Opt. Express, **22** 10735 (2014).
3. C. Wang, N. V. Wheeler, C. Fourcade-Dutin, M, Grogan, T. D. Bradley, B. R. Washburn, F. Benabid and K. L. Corwin, Appl. Opt., **52**, 5430, (2013).
4. F. Couny, F. Benabid, and P. S. Light, Opt. Lett., **31,** 3574 (2006).
5. F. Yu, W. J. Wadsworth, and J. C. Knight, Opt. Express, **20,** 11153 (2012).
6. A. N. Kolyadin, A. F. Kosolapov, A. D. Pryamikov, A. S. Biriukov, V. G. Plotnichenko, and E. M. Dianov, Opt. Express, **21**, 9514 (2013).
7. F. Benabid, J. C. Knight, G. Antonopoulos, and P. S. J. Russell, Science, **298**, 399 (2002).
8. F. Couny, F. Benabid, P. J. Roberts, P. S. Light, and M. G. Raymer, Science, **318**, 1118 (2007).
9. Y. Y. Wang, N. V. Wheeler, F. Couny, P. J. Roberts, and F. Benabid, Opt. Lett., **36**, 669 (2011).
10. B. Debord, M. Alharbi, T. Bradley, C. Fourcade-Dutin, Y. Y. Wang, L. Vincetti, F. Gerome and F. Benabid, Opt. Express, **21**, 28597 (2013).
11. B. Debord, M. Maurel, A. Amsanpally, M. Adnan, B. Beaudou, J. M. Blondy, L. Vincetti, F. Gerome and F. Benabid, in *Proc. of SPIE*, 10094, p.100941M (2017).
12. Y. Chen, H.C.H. Mulvad, S. R. Sandoghchi, E. Numkam, T.D. Bradley, J. R. Hayes, N.V. Wheeler, G.T. Jasion, S. U. Alam, F. Poletti, M.N. Petrovich and D. J. Richardson, *Conference on Lasers Electro-Optics 2016*, (Optical Society of America, 2016), STu4P.1.
13. N. V. Wheeler, T. D. Bradley, J. R Hayes, M. A. Gouveia, Y. Chen, S. R. Sandohgchi, F. Poletti, M. N. Petrovich and D. J. Richardson, in *Advanced Photonics 2016* (Optical Society of America, 2016) SoM3F.2.
14. E. A. J. Marcatili and R. A. Schmeltzer, Bell Syst. Tech. J. **43**, 1783 (1964).
15. A. M. Heidt, J. H. V. Price, C. Baskiotis, J S. Feehan, Z. Li, S. U. Alam and D. J. Richardson, Opt. Express, **21**, 24281 (2013).
16. N. V. Wheeler, A. M. Heidt, N. K. Baddela, E. N. Fokoua, J. R. Hayes, S. R. Sandoghchi, F. Poletti, M. N. Petrovich and D. J. Richardson, Opt. Lett., **39**, 295 (2014).
17. N. V. Wheeler, M. N. Petrovich, N. K. Baddela, J. R. Hayes, E. N. Fokoua, F. Poletti and D. J. Richardson, in *Conference on Lasers Electro-Optics 2012*, (Optical Society of America, 2012), CM3N.5.
18. S. R. Sandoghchi, M. N. Petrovich, D. R. Gray, Y. Chen, N. V. Wheeler, T. D. Bradley, N. H. L. Wong, G. Jasion, J. R. Hayes, E. Numkam Fokoua, M. Botelho Alonso Gouveia, S. M. Abokhamis Mousavi, D. J. Richardson and F. Poletti, Opt. Express, **23,** 27960 (2015).
19. D. R. Gray, M. N. Petrovich, S. R. Sandoghchi, N. V. Wheeler, N. K. Baddela, G. T. Jasion, T. Bradley, D. J. Richardson and F. Poletti, IEEE Photonics Technol. Lett., **28**, 1034 (2016).
20. T. D. Bradley, N. V. Wheeler, G. T. Jasion, D. R. Gray, J. R. Hayes, M. Botelho Alonso, S. M. Sandoghchi, Y. Chen, F. Poletti, D. J. Richardson, M. N. Petrovich, Opt. Express, **24**, 15798 (2016).
21. D. R. Gray, S.R.Sandoghchi, N.V.Wheeler, N.K.Baddela, G.T.Jasion, M.N.Petrovich, F.Poletti and D.J.Richardson, in *OFC 2015*, (Optical Society of America, 2015), W4I.6.
22. M. Alharbi, T. Bradley, B. Debord, C. Fourcade-Dutin, D. Ghosh, L. Vincetti, F. Gerome and F. Benabid, Opt. Express, **21**, 28609 (2013).
23. J. R. Hayes, F. Poletti, S. Abokhamis Mousavi, N. V. Wheeler, N. K. Baddela and D. J. Richardson, J. Light. Technol., **35**, 437 (2017).
24. B. Debord, A. Amsanpally, M. Chafer, A. Bas, M. Maurel, J-M. Blondy , E. Huggonot, L. Vincetti, , F. Gerome and F. Benabid, Optica, **4**, 209 (2017).
25. M. Michieletto, J. K. Lyngso, C. Jakobsen, J. Laesgaard, O. Bang and T. T. Alkeskjold, Opt. Express, **24**, 7103 (2016).

Full References

[1] C. Saraceno, F. Emaury, A. Diebold, I. Graumann, M. Golling, and U. Keller, "Trends in high-power ultrafast lasers," in *SPIE Defense+ Security. International Society for Optics and Photonics*, 2016, vol. 9835, pp. 98350X–1–98350X–7.

[2] B. Debord, M. Alharbi, L. Vincetti, A. Husakou, C. Fourcade-Dutin, C. Hoenninger, E. Mottay, F. Gerome and F. Benabid, "Multi-meter fiber-delivery and pulse self-compression of milli-Joule femtosecond laser and fiber-aided laser-micromachining.," *Opt. Express*, vol. 22, no. 9, pp. 10735–46, 2014.

[3] C. Wang, N. V. Wheeler, C. Fourcade-Dutin, M, Grogan, T. D. Bradley, B. R. Washburn, F. Benabid and K. L. Corwin, "Acetylene frequency references in gas-filled hollow optical fiber and photonic microcells," *Appl. Opt.*, vol. 52, no. 13, 2013.

[4] F. Couny, F. Benabid, and P. S. Light, "Large-pitch kagome-structured hollow-core photonic crystal fiber," *Opt. Lett.*, vol. 31, no. 24, pp. 3574–3576, 2006.

[5] F. Yu, W. J. Wadsworth, and J. C. Knight, "Low loss silica hollow core fibers for 3-4 μm spectral region.," *Opt. Express*, vol. 20, no. 10, pp. 11153–8, 2012.

[6] A. N. Kolyadin, A. F. Kosolapov, A. D. Pryamikov, A. S. Biriukov, V. G. Plotnichenko, and E. M. Dianov, "Light transmission in negative curvature hollow core fiber in extremely high material loss region," *Opt. Exp.*, vol. 21, no. 8, pp. 9514–9519, 2013.

[7] F. Benabid, J. C. Knight, G. Antonopoulos, and P. S. J. Russell, "Stimulated Raman Scattering in Hydrogen-Filled Hollow-Core Photonic Crystal Fiber," *Science*, vol. 298, no. 5592, pp. 399–402, 2002.

[8] F. Couny, F. Benabid, P. J. Roberts, P. S. Light, and M. G. Raymer, "Generation and photonic guidance of multi-octave optical-frequency combs," *Science (80-.).*, vol. 318, no. 5853, pp. 1118–1121, 2007.

[9] Y. Y. Wang, N. V. Wheeler, F. Couny, P. J. Roberts, and F. Benabid, "Low loss broadband transmission in hypocycloid-core Kagome hollow-core photonic crystal fiber," *Opt. Lett.*, vol. 36, no. 5, pp. 669–671, 2011.

[10] B. Debord, M. Alharbi, T. Bradley, C. Fourcade-Dutin, Y. Y. Wang, L. Vincetti, F. Gerome and F. Benabid, "Hypocycloid-shaped hollow-core photonic crystal fiber Part I: Arc curvature effect on confinement loss," *Opt. Express*, vol. 21, no. 23, pp. 28597– 608, 2013.

[11] B. Debord, M. Maurel, A. Amsanpally, M. Adnan, B. Beaudou, J. M. Blondy, L. Vincetti, F. Gerome and F. Benabid, " Ultra-low loss (8.5 dB/km @ Yb-laser wavelength range) inhibited-coupling Kagome HC-PCF for laser beam delivery applications," Proc. SPIE 10094, Frontiers in Ultrafast Optics: Biomedical, Scientific, and Industrial Applications XVII, 100941M (2017); doi:10.1117/12.2251889

[12] Y. Chen, H.C.H. Mulvad, S. R. Sandoghchi, E. Numkam, T.D. Bradley, J. R. Hayes, N.V. Wheeler, G.T. Jasion, S. U. Alam, F. Poletti, M.N. Petrovich and D. J. Richardson, "First Demonstration of Low Loss, Bend Insensitive 37-Cell Hollow-Core Photonic Bandgap Fiber at 1μm for High Power Delivery Applications," *Conference on Lasers Electro-Optics (CLEO) 2016,* STu4P.1, 2016.

[13] N. V. Wheeler, T. D. Bradley, J. R Hayes, M. A. Gouveia, Y. Chen, S. R. Sandohgchi, F. Poletti, M. N. Petrovich and D. J. Richardson, "Low Loss Kagome Fiber in the 1 μm Wavelength Region," in Advanced Photonics 2016 (IPR, NOMA, Sensors, Networks, SPPCom, SOF), OSA Technical Digest (online) (Optical Society of America, 2016), SoM3F.2.

[14] E. A. J. Marcatili and R. A. Schmeltzer, "Hollow Metallic and Dielectric Waveguides for Long Distance Optical Transmission and Lasers," Bell Syst. Tech. J. 43(4), 1783–1809 (1964).

[15] A. M. Heidt, J. H. V. Price, C. Baskiotis, J S. Feehan, Z. Li, S. U. Alam and D. J. Richardson, "Mid-infrared ZBLAN fiber supercontinuum source using picosecond diode-pumping at 2 μm," *Opt. Express*, vol. 21, no. 20, p. 24281, 2013.

[16] N. V. Wheeler, A. M. Heidt, N. K. Baddela, E. N. Fokoua, J. R. Hayes, S. R. Sandoghchi, F. Poletti, M. N. Petrovich and D. J. Richardson, "Low-loss and low-bend-sensitivity mid-infrared guidance in a hollow-core-photonic-bandgap fiber.," *Opt. Lett.*, vol. 39, no. 2, pp. 295–8, 2014.

[17] N. V. Wheeler, M. N. Petrovich, N. K. Baddela, J. R. Hayes, E. N. Fokoua, F. Poletti and D. J. Richardson, "Gas Absorption between 1.8 and 2.1 μm in Low Loss (5.2 dB/km) HC-PBGF," *Conference on Lasers Electro-Optics (CLEO), 2012* CM3N.5, 2012.

[18] S. R. Sandoghchi, M.N.Petrovich, D.R.Gray, Yong Chen, N.V.Wheeler, T.D.Bradley, N.H.L.Wong, G.Jasion, J.R.Hayes, E.Numkam Fokoua, M.Botelho Alonso Gouveia, S.M.Abokhamis Mousavi, D.J.Richardson and F.Poletti, "Optical side scattering radiometry for high resolution, wide dynamic range longitudinal assessment of optical fibers," *Opt. Express*, vol. 23, no. 21, pp. 27960– 27974, 2015.

[19] D. R. Gray, M. N. Petrovich, S. R. Sandoghchi, N. V. Wheeler, N. K. Baddela, G. T. Jasion, T. Bradley, D. J. Richardson and F. Poletti, "Real-Time Modal Analysis via Wavelength- Swept Spatial and Spectral (S2) Imaging," *IEEE Photonics Technol. Lett.*, vol. 28, no. 9, pp. 1034– 1037, 2016.

[20]   T. D. Bradley, N. V. Wheeler, G. T. Jasion, D. R. Gray, J. R. Hayes, M. Botelho Alonso, S. M. Sandoghchi, Y. Chen, F. Poletti, D. J. Richardson, M. N. Petrovich, "Modal content in hypocycloid Kagome hollow core photonic crystal fibers," *Opt. Express*, vol. 24, no. 14, pp. 15798–15812, 2016.

[21]   D. R. Gray, S.R.Sandoghchi, N.V.Wheeler, N.K.Baddela, G.T.Jasion, M.N.Petrovich, F.Poletti and D. J. Richardson, "Mitigating Spectral Leakage and Sampling Errors in Spatial and Spectral (S 2) Imaging," in *Optical Fiber Communications (OFC) Conference 2015*, W4I.6, 2015.

[22]   M. Alharbi, T. Bradley, B. Debord, C. Fourcade-Dutin, D. Ghosh, L. Vincetti, F. Gerome and F. Benabid, "Hypocycloid-shaped hollow-core photonic crystal fiber Part II: cladding effect on confinement and bend loss.," *Opt. Express*, vol. 21, no. 23, pp. 28609–16, 2013.

[23]   J. R. Hayes, F. Poletti, S. Abokhamis Mousavi, N. V. Wheeler, N. K. Baddela and D. J. Richardson, "Antiresonant Hollow Core Fiber with Octave Spanning Bandwidth for Short Haul Data Communications," *J. Light. Technol.*, 35 (3), pp. 437, 2017

[24]   B. Debord, A. Amsanpally, M. Chafer, A. Bas, M. Maurel, J-M. Blondy , E. Huggonot, L. Vincetti, , F. Gerome and F. Benabid, Optica, 4(2), pp. 209, 2017.

[25]   M. Michieletto, J. K. Lyngso, C. Jakobsen, J. Laesgaard, O. Bang and T. T. Alkeskjold, "Hollow core fibers for high power pulse delivery," *Opt. Express,* vol. 24, no. 7, pp. 7103- 7119, 2016.