

Nonlinear refractive index of ultrafast laser inscribed waveguides in gallium lanthanum sulphide

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We demonstrate ultrafast all-optical switching in femtosecond laser inscribed nonlinear directional couplers in gallium lanthanum sulphide operated at 1.55 μm . We report on the evaluation of the nonlinear refractive index of the waveguides forming the directional couplers by making use of the switching parameters. The nonlinear refractive index is reduced by the inscription process to about 4-5 times compared to bulk material. © 2015 Optical Society of America

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1. INTRODUCTION

Despite the continuing progress in electronics technology, there is a growing interest in the development of optical switching and processing devices due to their potential of operating at speeds and bandwidths much greater than those achievable with electronic devices. This rapid development of optical communications requires novel materials with large and ultrafast nonlinear optical responses in the femtosecond or picosecond domain for fabricating compact, low-threshold all-optical switching (AOS) and processing devices.

Chalcogenide glasses (ChGs) are a particularly appealing candidate for such devices as they present high third-order (Kerr) nonlinearities (up to a thousand times that of silica) [1], low multi-photon absorption [2] and high photosensitivity [1], properties which allow for switching devices with ultrafast response, shorter interaction lengths, lower switching threshold and higher figures of merit.

Of all the ChGs, gallium lanthanum sulphide (GLS) is a particularly attractive one and a much-awaited alternative to toxic arsenic-based glasses. It presents optical transparency from the visible wavelengths extending in the infrared up to 10 μm and thermal stability up to 550°C [3, 4]. In addition, GLS can be melted in a large scale without the requirement of a sealed ampoule environment making production and processing safer, easier and more economical [4].

In this work, we have utilized the ultrafast laser inscription (ULI) technique [5, 6] to fabricate optical directional couplers in polished GLS glass substrates. ULI is a novel fabrication technique which offers the promise of highly functional optical circuits in three-dimensional geometries [7, 8] that are not possible with commercial optical fiber or

planar lightwave circuit (PLC) fabrication methods. It relies on the nonlinear absorption of sub-bandgap photons to induce permanent structural changes to a material. These changes can manifest themselves in multiple ways, including a change in refractive index, which can be exploited to confine and guide light in ULI fabricated structures in bulk substrates [9, 10].

An evaluation of the linear refractive index profiles of ULI fabricated waveguides in GLS substrates was performed in a previous work [11] by our group. Here we expand our studies. We investigate the effect the inscription process has on the nonlinear refractive indices of the fs-laser inscribed directional couplers in GLS. We use a femtosecond OPA and the closed-aperture Z-scan technique [12] to measure the magnitude of the nonlinear refractive index coefficient, n_2 , of the bulk substrate. We then evaluate the nonlinear refractive index of the waveguides forming the directional couplers by making use of the switching parameters when the directional couplers are operated as all-optical switches and compare with bulk GLS.

All-optical switching devices which exploit a nonlinear refractive change have attracted much attention because of their ultrafast response. In this approach the irradiance-dependent refractive index, $n = n_0 + n_2I$, is used to demonstrate all-optical switching in the nonlinear medium. Here, n_0 is the linear component of the refractive index, n_2 is the nonlinear refractive index coefficient and I is the irradiance. A nonlinear directional coupler, intended to operate as an all-optical switch, is designed in such way that when a low irradiance input is applied, a maximum amount of cross coupling is observed. Then, as the input irradiance is increased, the refractive index of the

bar waveguide, following the above equation, increases, detuning the coupling coefficient between the two adjacent waveguides and causing most of the light to remain in the bar port in the output; therefore, achieving all-optical switching.

2. DEVICE FABRICATION

All structures reported in this letter were fabricated using a mode-locked Yb-doped fiber laser emitting 360 fs pulses at a central wavelength of 1045 nm and a pulse repetition rate of 500 kHz. The substrates were mounted on air bearing Aerotech stages and the pulses from the fabrication laser were focused inside the substrates to a depth of 240 μm from the top surface using a 0.4 NA aspheric lens. The pulse energy incident on the samples was chosen to be 58 nJ, whereas the translation speed was set to 12 mm/s. The substrate translation was perpendicular to the laser beam direction. The cross section of the inscribed waveguides forming the directional couplers was controlled using the well-known multiscan waveguide shaping technique [13, 14], yielding waveguides with rectangular cross-sections. For each of the couplers, 10 scans with a scan-to-scan separation of 0.36 μm were used. The above are the optimal ULI parameters for realising high quality single-mode waveguides at 1550 nm and were identified in a previous work by our group [11]. After fabrication, the input and output facets of the substrates were polished to optical quality.

All the directional couplers used in this work had the same geometric layout (Fig. 1). They were formed from two waveguides, single-mode at 1550 nm. Each of the waveguides consisted of straight input and output sections linked to s-bends leading to a straight interaction region of length L . The input and output ports of the two waveguides forming each coupler were separated by more than 100 μm to avoid evanescent coupling in any other sections of the coupler than the interaction length. This was achieved by choosing the lateral displacement in the x axis (x -offset) to be 50 μm for each of the waveguides. The bend regions of the waveguides were inscribed by choosing a displacement of 5 mm along the propagation direction (y -offset) and a radius of curvature of 110 μm , resulting in waveguides which presented zero bend losses. The zero bend losses were verified by observing that waveguides of the above outlined geometry exhibited the same total throughput as a straight single-mode, at 1550 nm, waveguide of the same length inscribed with the same laser parameters. The propagation losses in straight ultrafast laser inscribed waveguides in GLS were previously found to be in the order of $\sim 0.8 \text{ dB} \cdot \text{cm}^{-1}$ [15].

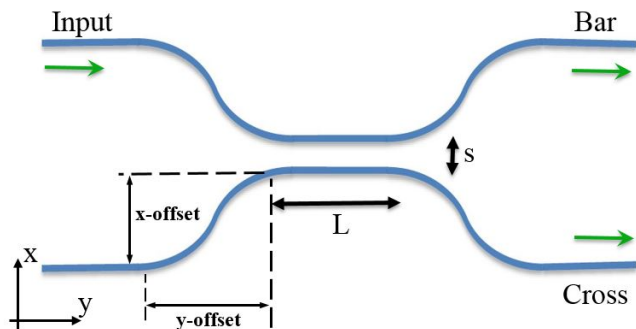


Fig. 1. Schematic of the directional coupler structure, where s is the core-to-core separation between the waveguides along the interaction length (L).

While keeping all other parameters constant, a series of directional couplers were inscribed using interaction lengths between 3.55 and

8.25 mm in steps of 50 μm . For all couplers, the core-to-core separation in the interaction region was set to 7.5 μm from the center of the cores.

3. ALL-OPTICAL SWITCHING

As discussed earlier, a directional coupler intended to be operated as an all-optical switch is designed such way that for a low irradiance input a maximum amount cross coupling in the output is observed. Therefore, prior to carrying out the switching experiment, the output relative transmission for the bar and cross ports for each of the couplers was measured by coupling light from a 1550 nm CW fiber laser to the input port. The use of the CW laser ensured that no nonlinearities were excited whilst propagating light in the directional couplers, verifying the output cross-to-bar ratio of the couplers at low irradiances before any switching occurs.

A single beam experiment was then performed to examine the operation of the couplers as all-optical switches. A regeneratively amplified Ti:Sapphire source delivering ~ 100 fs pulses, pumping an optical parametric amplifier (Newport Spectra-Physics OPA 800) at a repetition rate of 1 kHz was used as input. The output from the OPA was centered at 1550 nm and after being passed through a diamond pinhole held in vacuum, ensuring its Gaussian shape, was coupled into and collected from the directional couplers using two AR coated 16X (0.25NA) aspheric lenses. The input irradiance was measured before the sample using a pyro-electric power meter (Laser Probe Inc. RkP-575), whereas the output light from the couplers was collected using a large-area germanium photoreceiver (New Focus 2033). The output cross-to-bar ratio as a function of the input irradiance for two directional couplers inscribed with different interaction lengths while keeping all the other parameters constant is shown in figure 2.

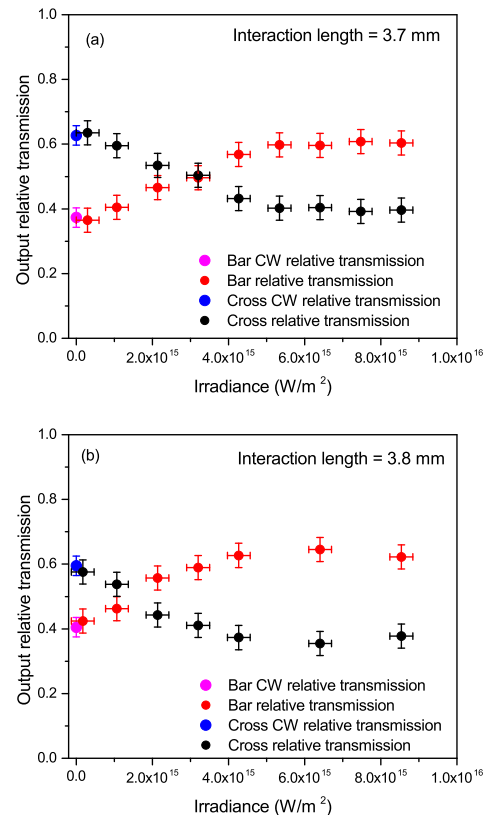


Fig. 2. Output relative transmission for the bar (red) and cross (black) ports as a function of the input irradiance for (a) a directional coupler with an interaction length of 3.7 mm and (b) a directional coupler with

an interaction length of 3.8 mm. The pink and blue circles in both graphs represent the bar and cross output ratio in the linear regime when using a 1550 nm CW fiber laser as input.

Coupler 1 was chosen to have an interaction length of 3.7 mm, whereas for coupler 2 the interaction length was set to 3.8 mm. The cross-to-bar output ratios, recorded with the 1550 nm CW laser as an input prior to the switching experiment was performed, are shown in the same axes system (at zero irradiance) with pink and blue circles.

As can be seen from the figure, for a low irradiance input about 60 % of the light is in the cross port at the output, whereas as the irradiance increases the bar-to-cross output ratio changes until it is completely reversed to about 60-40 % bar-to-cross at the highest irradiances used in this experiment. The incident irradiance required for coupler 1 (Fig. 2a) to completely switch from the cross to the bar state was about 5.34×10^{15} W/m², whereas the incident irradiance required for the second coupler (Fig. 2b) to completely switch from the cross to the bar state was about 4.27×10^{15} W/m². The lower irradiance required to switch coupler 2 is consistent with its longer interaction length outlining the common dilemma of the trade-off between power consumption and device dimensions when designing all-optical processing devices.

4. EVALUATION OF THE NONLINEAR REFRACTIVE INDEX FROM THE SWITCHING PARAMETERS

The phase change $\Delta\Phi$ due to nonlinear refraction is [16]:

$$\Delta\Phi = \frac{2\pi L n_2 I}{\lambda} \quad (1)$$

where L in this case is the interaction length of the directional couplers, λ is the operating wavelength and I the is the irradiance.

The phase change, $\Delta\Phi$, required to switch a nonlinear directional coupler is in the order of 4π [17]. Using this along with the corresponding interaction length and switching irradiance for each of the couplers of figure 2 in equation 1, and after accounting for the losses due to Fresnel reflection at the input facet, the n_2 of the modified region is calculated. The yielded n_2 for coupler 1 was 1.89×10^{-19} m²/W, whereas the yielded value for coupler 2 was 2.30×10^{-19} m²/W.

In a similar work demonstrating all-optical switching in an aluminum gallium arsenide (AlGaAs) directional coupler [18] the critical irradiance, I_c , was identified as the irradiance required to switch a half beat-length directional coupler to the condition where equal output powers emerge from the bar and cross waveguides. This critical irradiance was given as:

$$I_c = \frac{\lambda}{L n_2} \quad (2)$$

Adapting to our case, the critical irradiance for coupler 1 is 3.20×10^{15} W/m², whereas the critical irradiance for coupler 2 is 1.49×10^{15} W/m². Using these critical irradiances, again after accounting for the losses due to Fresnel reflection at the input, along with the interaction lengths for each of the couplers in the above equation the yielded n_2 values are 1.58×10^{-19} m²/W and 3.29×10^{-19} m²/W for couplers 1 and 2 respectively, in good agreement with the values calculated with equation 1.

The nonlinear refractive index coefficient of bulk GLS at 1550 nm was measured via the closed-aperture Z-scan technique. The output of the OPA outlined above was focused using a 200 mm focal length calcium fluoride lens after being passed through a diamond pinhole held in vacuum to ensure its Gaussian shape. A 1 mm thick GLS sample was then translated through the focus along the beam path whilst

recording the far field transmittance through an aperture using a large-area germanium photoreceiver (New Focus 2033). Figure 3 shows a closed-aperture Z-scan trace performed at a pulse energy of 100 nJ with an aperture transmission of 50%. The theoretical fit uses the simplified closed aperture fitting formula from Sheik-Bahae *et al.* [12]. The theoretical fit corresponds to a nonlinear refractive index coefficient of $(9.7 \pm 0.94) \times 10^{-19}$ m²/W.

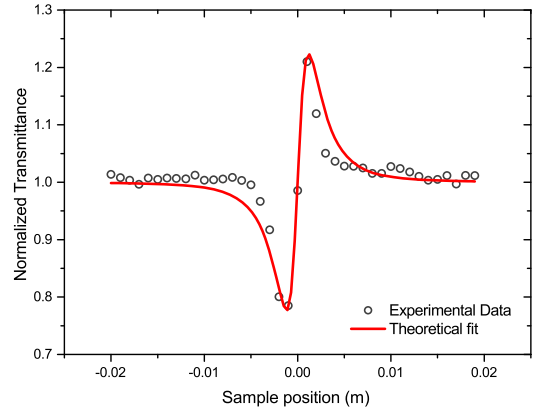


Fig. 3. Closed-aperture Z-scan trace of GLS at 1550 nm performed at a pulse energy of 100 nJ. The theoretical fit (solid red line) is calculated using an n_2 coefficient of 9.7×10^{-19} m²/W.

This value is about 4-5 times higher than the ones calculated for the modified area of the directional couplers by using the nonlinear phase change. By considering the values obtained using the critical irradiances, the n_2 values of the modified area are calculated to be 3-6 times lower than the one measured for the bulk. These findings indicate that the nonlinear refractive index of GLS is reduced by the inscription process.

5. DISCUSSION

This trend of decreased nonlinearity for the modified area was also observed in straight ULI GLS waveguides [10], in a work performed from previous members of our group. In that work the authors derived the nonlinear phase change from the self-phase modulation (SPM) peaks and calculated the n_2 coefficient of the modified area using equation 1. The n_2 coefficient they estimated for the modified area was approximate 5 times lower than the one they measured for the bulk, in very good agreement with the findings of this investigation. Moreover, it has been seen in studies on waveguides inscribed inside in fused silica substrates [19] that the ULI technique can cause changes to the nonlinear refractive index. It was shown that the nonlinear refractive indices of the waveguides were up to four times lower than that of bulk fused silica with their magnitude being very dependent on the inscription parameters.

Finally, it is worth mentioning that the CW cross-to-bar coupling ratio of about 60 % for the two couplers of Fig. 2 is the highest we could achieve. This coupling ratio would be unlikely to reach 100 % regardless of the sampling of the interaction length. This behavior is characteristic of differences in the propagation constants of the two waveguides forming the directional couplers [20], arising mainly from the fabrication procedure. In previous work by our group we showed that refractive index variations induced by ULI can extend beyond the geometrical structure of the waveguides up to a distance of ~ 16 μ m [11]. Therefore, in the case of the directional couplers used in this work, where the core-to-core separation was chosen to be 7.5 μ m, the optical properties (e.g. mode index) of the first waveguide for each coupler were directly affected by the inscription of the second

waveguide and vice versa leading to different propagation constants for the two waveguides. A possible solution to this problem would be to choose core-to-core separations bigger than $\sim 20 \mu\text{m}$. However, in this case the coupling coefficient would be greatly decreased and would therefore require very long interaction lengths for the maximum amount of power to transfer to the cross waveguide, which would cause other undesirable effects such as intermodal dispersion [21], completely jeopardizing the function of the directional couplers as optical switches. Other possible sources of the mismatch in the propagation constants of the two waveguides might be the drifts in laser power or temperature variations in the lab during the inscription process, where power fluctuations in the orders of only a few milliwatts or temperature variations as small as $\pm 1^\circ \text{C}$ might be sufficient to change the ULI parameters over the course of the experiment.

6. CONCLUSION

To summarize, we have investigated the nonlinear refraction properties of fs-laser inscribed structures. We have demonstrated ultrafast all-optical switching in nonlinear directional couplers in GLS at 1550 nm and by making use of the switching parameters we calculated the nonlinear refractive indices of the single-mode waveguides forming the directional couplers. It was found that the nonlinear refractive index of the modified, by the laser, area is reduced by the inscription process by about 4-5 times comparing to that of the bulk. We believe that our studies will contribute towards better understanding the nonlinear properties of ULI fabricated structures which will open the way for the optimum design of ultrafast laser inscribed integrated nonlinear optical devices. Finally, these studies are even more important in the case of chalcogenide glasses where their high Kerr nonlinearities make them very attractive candidates for such devices.

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