

Water Droplet Impact Energy Harvesting with P(VDF-TrFE) Piezoelectric Cantilevers on Stainless Steel Substrates

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Received xxxxxx

Accepted for publication xxxxxx

Published xxxxxx

Abstract

Harvesting energy from ambient environmental sources using piezoelectric transducers has seen a tremendous amount of interest from the scientific community in recent times. The practicality of energy scavenging technology looks set to see continued relevance, with decreasing power demands of electrical systems, such as Wireless Sensor Networks (WSN), allowing such technology to progressively act as an energy source to drive and sustain them independently. In light of this, there is a growing opportunity for piezoelectric materials to prolong, or even replace, battery powered sensor systems positioned in remote or difficult to reach areas. It has been demonstrated that falling water droplets of millimetric-scale diameter can impart forces of over a thousand times their resting weight upon surface impact. As such, the potential for utilising piezoelectric transducers to drive sensor systems, by converting the kinetic impact energy of falling water droplets into useful electrical energy, is investigated. The key parameters that affect the efficiency of energy transfer between incident water droplets and piezoelectric cantilever structures made of stainless steel foil coated with the lead-free piezoelectric material P(VDF-TrFE) are investigated. A peak power output of 28 nJ achieved from the impact of a 5.5 mm diameter droplet upon a single energy harvesting transducer illustrated that, for droplets of diameter 3.1 mm to 5.5 mm impacting from heights between 0.5 to 2.0 m, it is desirable to utilise piezoelectric transducer beams of bending stiffness in the range of 0.067 to 0.134 N/m in order to achieve good energy transfer efficiency. Although the active electrode area was constrained in order to maintain consistency between samples, reducing the peak energy output, the achieved results correspond to a 15.9 J/m³ energy density, representing the significant energy transfer efficiency achievable through appropriate transducer mechanical tailoring to the excitation source.

Keywords: Piezoelectric, Energy Harvesting, Droplet, Smart Materials, Microelectromechanical Systems, Rain Water

1. Introduction

Falling water droplets have been demonstrated to impart forces of over a thousand times their resting weight upon impact [1] Tropical countries regularly receive a deluge of large rainfall droplets falling upon natural and man-

made structures, representing an interesting opportunity for further investigation regarding kinetic to electrical energy harvesting. Such systems may have the potential to drive interesting applications such as smart city sensor technology.

Research into droplet impact energy harvesting has

49 seen limited progress to-date, with the majority of
50 investigations utilising commercially available sensors to
51 analyse droplet impact mechanics [2]–[4]. A non-
52 exhaustive review of investigations into water droplet
53 impact energy harvesting by piezoelectric transducers is
54 provided by Wong *et al* [5]. Whilst good insight has been
55 produced from these analyses, using commercial sensors
56 as the energy harvesting transducer limits the degree to
57 which the transducers can be modified to efficiently
58 harvest the droplet impact stimulus. Investigations into
59 water droplet energy harvesting appear to have been
60 initiated formally by Guigon *et al* in 2008, where both the
61 theory and experimental work contributing to efficient
62 mechanical energy harvesting were outlined in a two-part
63 study [6], [7]. The initial theoretical investigation
64 illustrated how the mechanical sizing of the structure
65 optimises the transfer of deformation energy from the
66 drop to the piezoelectric polymer, before investigating
67 the structure conversion efficiency through surface
68 contact electrode design. It was concluded that the
69 piezoelectric transducer material must be very thin (μm
70 scale), not be pre-stressed and be of a width slightly
71 smaller than the maximum diameter of the impacting
72 drop for efficient energy harvesting to take place.
73 Additionally, it was considered optimal for the
74 piezoelectric material to be entirely covered with
75 conducting electrodes. Simulations demonstrated a
76 theoretical energy output of $25\ \mu\text{J}$ and peak instantaneous
77 power output of $12\ \text{mW}$ from a “downpour” drop of $5\ \text{mm}$
78 diameter [6].

79 The corresponding experimental study demonstrated
80 how mono-stretched PVDF polymer bands of $10\ \text{cm}$
81 length, $3\ \text{mm}$ width and $25\ \mu\text{m}$ thickness demonstrated a
82 peak power output of approximately $1\ \text{nJ}$ electrical
83 energy and $1\ \mu\text{W}$ of instantaneous power from a single
84 droplet impact. It was observed that the recovery of
85 electrical energy was maximised when droplet impacts
86 were slightly off-centre from the beam material [7]. It is
87 noted that no hydrophobic encapsulation was
88 administered to the bands, therefore the droplet impacts
89 were treated as inelastic.

90 These findings were supported by research carried out
91 by Vasileiou *et al*, which explored how a reduction in
92 substrate area density and stiffness imparts immediate
93 acceleration and intrinsic responsiveness to impacting
94 droplets [8]. Although this research did not utilize
95 piezoelectric materials, and as such did not consider the
96 effect of the relevant parameters on power output, their
97 findings demonstrated efficient energy conversion from
98 droplet kinetic energy to substrate kinetic and strain
99 energy when the sample had low area density,
100 comparable stiffness to that of the water droplet’s surface
101 tension, and minimal damping.

102 The initial work of Guigon *et al* inspired a number of
103 follow-up investigations. Ilyas and Swingler analysed the

104 voltage output profiles produced from droplet impact
105 upon commercially available sensors in detail,
106 identifying two distinct phases in voltage and power
107 output; log growth at the initial impact, before
108 exponential decay occurring throughout the remainder of
109 the impact event. It was demonstrated that the log growth
110 stage significantly contributes to the overall power output
111 of the device [2], [9].

112 Research conducted by Viola *et al* has investigated
113 droplet energy output with commercially available
114 piezoelectric sensors, in addition to testing cantilever
115 beam, bridge and centralized floating diaphragm
116 harvester configurations [10]. It was found that the
117 cantilever beam harvester configuration achieved the best
118 response to impacting droplets, with a commercial
119 LDT1-028K MEAS piezoelectric sensor producing $17\ \text{V}$
120 output from droplets dispensed from a height of $2\ \text{m}$.

121 The performance of piezoelectric energy harvesters in
122 both simulated and actual rainfall was evaluated by
123 Wong *et al* [11], [12]. A spray-type rain simulator was
124 used to dispense a range of droplet diameters from a
125 height of $2.5\ \text{m}$ onto a Midé VultureTM commercial
126 piezoelectric sensor (model V25W), composed of two
127 PZT layers of $46\ \text{mm}$ length, $33\ \text{mm}$ width and $0.6\ \text{mm}$
128 overall thickness. With the PZT layers connected in
129 series across a $15\ \text{k}\Omega$ load resistor, simulated rainfall
130 rates of 33 , 40 , 62 and $99\ \text{mm/h}$ for a period of 300
131 seconds produced total power outputs of $0.074\ \mu\text{W}$,
132 $0.156\ \mu\text{W}$, $0.167\ \mu\text{W}$ and $0.207\ \mu\text{W}$ respectively.
133 Furthermore, the testing of the piezoelectric sensor in
134 three actual rainfall events of duration 250 , 204 and 301
135 minutes produced total harvested energies of $155.6\ \mu\text{J}$,
136 $438.9\ \mu\text{J}$ and $2076\ \mu\text{J}$ respectively. The significantly long
137 time scale required to capture these energies highlights
138 the inefficiency of attempting to harvest energy directly
139 from rainfall droplet impacts.

140 In this work, fundamental findings regarding droplet
141 impact energy harvesting are advanced by investigating
142 the key parameters dictating efficient energy transfer. In
143 order to encourage higher electrical power output, the
144 piezoelectric transducer’s bending stiffness k_b and
145 resonant frequency f_r are focused upon. COMSOL
146 Multiphysics simulations exploring simplified droplet
147 impact onto piezoelectric beams were carried out
148 initially. This was followed by an experimental study
149 utilising purpose-built samples consisting of
150 piezoelectric P(VDF-TrFE) polymer deposited onto
151 stainless steel foil substrate beams. Using these samples,
152 a variety of droplet diameters are dispensed onto the
153 transducer beam end from a range of heights in order to
154 explore the effect of beam stiffness on energy transfer
155 efficiency.

156 Finally, an investigation into the energy transfer
157 efficiency depending on droplet impact frequency is
158 carried out. Whilst it is rational to assume that driving an

159 energy harvester at its resonant frequency will result in
 160 maximised power output, it is proposed elsewhere that,
 161 depending on the relation between the beam resonant
 162 frequency and the natural vibration frequency of the
 163 impacting droplet (illustrated in Equation 1, where T is
 164 the natural oscillation period of a droplet [13]), the
 165 reactive transducer movement to impact can
 166 synergistically, passively or destructively contribute
 167 towards the droplet kinetic energy after recoil from the
 168 substrate [8]. As such, this research seeks to validate the
 169 effect of impact frequency upon the energy transfer
 170 efficiency, as this has yet to be explicitly demonstrated
 171 elsewhere.

$$173 \quad f_d = \frac{1}{T} = \frac{1}{\frac{\pi}{4} \sqrt{\frac{\rho D_0^3}{\sigma}}} \quad (1)$$

174
 175 It has been reliably demonstrated that droplet-surface
 176 interactions are not trivial, given the elastic nature of
 177 droplets upon impact. Substrate flexibility directly
 178 affects the droplet impact behaviour, with “softer”
 179 substrates suppressing droplet splashing [14]. It is
 180 possible to achieve a two-fold reduction in contact time
 181 during a droplet impact if the superhydrophobic substrate
 182 is tailored to an appropriate elasticity, directly
 183 influencing the transfer of energy from the droplet to the
 184 transducer [15]. Correspondingly, the wettability of the
 185 sample surface directly influences the energy output
 186 achieved. Gart *et al* studied the effect of different surface
 187 treatments on the generated torque produced by elastic
 188 cantilever beams from droplet impact [16]. It was found
 189 that hydrophilic surfaces, which encourage water droplet
 190 adhesion after impact, led to greater bending energy
 191 being produced in the beam than beams where a
 192 hydrophobic surface treatment had been applied. This
 193 was due to the additional mass of the water droplet
 194 sticking to the cantilever during the impact process,
 195 subsequently generating a higher torque over time. It is
 196 proposed that, for a limited number of impacts,
 197 hydrophilic surface treatments are desirable – however,
 198 such surface treatments encourage the accumulation of
 199 water upon the energy harvesting surface, leading to
 200 excessive damping and loss of electrical energy output
 201 [17].

202 In this study, a superhydrophobic surface treatment
 203 will be utilised to isolate the piezoelectric beam.
 204 Although it has been demonstrated that hydrophobic
 205 treatments reduce the bending energy experienced by
 206 piezoelectric transducers, as the droplets tend to roll-off
 207 the harvester surface during impact, such treatments
 208 reduce the accumulation of water and increase the energy
 209 transfer efficiency of the harvesters in the long term. For
 210 reference, the contact angle of droplets upon a P(VDF-
 211 TrFE) layer without surface treatment is approximately

212 75° [18], increasing to above 160° with a
 213 superhydrophobic treatment, representing a significant
 214 decrease in surface wettability [19].

215 2. Sample Design Justification

216 Single-end clamped cantilever beams were selected as
 217 the test geometry, given the abundance of theoretical
 218 models which can be used to define and analyse such a
 219 system. The relationship expressing cantilever beam
 220 stiffness for a point load deflection applied at the free-
 221 end of the cantilever, perpendicular to the beam axis, is
 222 given in Equation 2. This relationship was used to
 223 describe the estimated stiffness of each test sample beam,
 224 where k_b is the beam bending stiffness, E is the beam
 225 elastic modulus, w is the beam width, t is the beam
 226 thickness and L is the beam length [20].

$$228 \quad k_b = \frac{Ewt^3}{4L^3} \quad (2)$$

229
 230 The relationship defining a cantilever beam’s first order
 231 resonant frequency is shown in Equation 3, where E , t , ρ
 232 and L relate to the cantilever’s Youngs modulus,
 233 thickness, density and length, respectively. The absence
 234 of a width term in this equation illustrates how the beam
 235 width can be varied in isolation to investigate the effect
 236 of beam stiffness variation on power output, without
 237 altering other mechanical parameters such as the beam’s
 238 resonant frequency.

$$239 \quad f_b = \left(1.875^2 \sqrt{\frac{\left(\frac{Et^2}{12}\right)}{\rho L^4}} \right) \cdot \frac{1}{2\pi} \quad (3)$$

240 Following the guidance in the literature outlining the
 241 need for thin, responsive substrates in order to achieve
 242 efficient droplet energy harvesting, stainless steel foil of
 243 25 μm thickness was used as a substrate, due to the
 244 material’s inherent robustness, corrosion resistance and
 245 flexibility at thin gauges. P(VDF-TrFE) was selected for
 246 usage as the piezoelectric polymer also due to its
 247 flexibility. Considering the thickness weighted material
 248 properties of both the substrate and the P(VDF-TrFE)
 249 layer, it was calculated that a P(VDF-TrFE) layer of
 250 approximately 7.2 μm in thickness would position the
 251 beam’s neutral axis to encourage an optimal piezoelectric
 252 voltage response, through facilitation of an advantageous
 253 bending regime [21]. In practice, it was found that such a
 254 thin layer of P(VDF-TrFE) on the selected stainless steel
 255 foil frequently resulted in breakdown during the poling
 256 process, so P(VDF-TrFE) layers of approximately 15 μm
 257 were fabricated.

258 3. Numerical Simulation of Stiffness Variation

259 Numerical simulation was carried out using COMSOL
260 Multiphysics® Modeling Software. The in-built library
261 materials “Steel AISI 4340” and “Polyvinylidene
262 fluoride (PVDF)” were used to simulate the foil substrate
263 and piezoelectric layer respectively. The properties of
264 these materials are tabulated in Table I.

265 Beams of length 43 mm were simulated, with one end
266 being held under fixed constraint. At the other end, a
267 pulsed load was applied using a circular work plane
268 parallel to the beam’s upper surface to represent the
269 impact of a droplet. The work plane diameter was
270 simulated at 3.1 mm, 4.4 mm and 5.5 mm dimensions to
271 represent the droplet diameter variations achievable from
272 the range of experimental syringe nozzle outlets
273 available.
274

TABLE I
SIMULATION MATERIAL PROPERTIES

Parameter	Value
Youngs Modulus Steel AISI 4340	205 GPa
Poisson’s ratio of Steel AISI 4340	0.28
Density of Steel AISI 4340	7850 kg/ m ³
Youngs Modulus Polyvinylidene fluoride (PVDF)	3 GPa
Poisson’s ratio of Polyvinylidene fluoride (PVDF)	0.18
Density of Polyvinylidene fluoride (PVDF)	1780 kg/ m ³
Polyvinylidene fluoride (PVDF) Piezoelectric Strain Constant d_{31} (shear mode direction 1)	13.6 pC/N
Polyvinylidene fluoride (PVDF) Piezoelectric Strain Constant d_{33} (thickness mode direction 3)	29.7 pC/N

275
276
277 The impact force for each diameter droplet was
278 calculated using Equation 4, which estimates the impact
279 force as a dynamic pressure $\rho_{water} v^2$ applied over a
280 surface area of πr^2 [1]. This force was then translated
281 into a force per unit area, and applied as a boundary load
282 to the work plane. The speed of the droplet upon impact
283 is assumed to be terminal velocity, and is calculated using
284 Equation 5, where ρ_{water} is the water density (1000
285 kg/m³), ρ_{air} is air density (1.225 kg/m³) and C_d is the
286 drag coefficient for a sphere (0.47). The duration of this
287 impact or “crash time” was calculated using Equation 6
288 [15], where τ is the crash time, D_o is the droplet diameter
289 before impact and v is the droplet’s impact speed,
290 assumed to be the droplet’s terminal velocity here.

291

292

293

$$F_0 = \pi \rho_{water} r^2 v^2 \quad (4)$$

294

295

$$v_t = \left(\frac{8 \rho_{water} r g}{3 \rho_{air} C_d} \right)^{1/2} \quad (5)$$

296

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$$\tau = \frac{D_o}{v} \quad (6)$$

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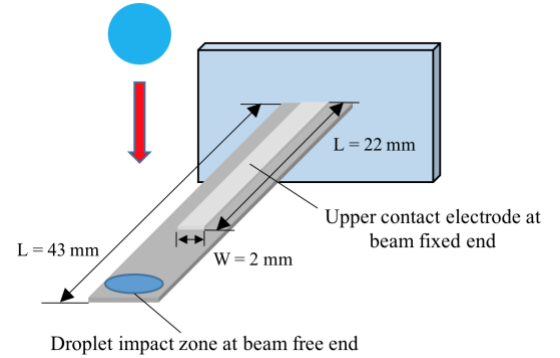
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An additional rectangular work plane of dimensions 22 mm length, 2 mm width was drawn on the upper surface of the beam to represent an upper contact electrode. An output load impedance sweep determined that the design was electrically impedance matched at 50 MΩ. The simulation time study period was 0.1 seconds, which ensured that the initial impact and subsequent initial beam displacement were captured. The simulation setup is illustrated in Fig. 1. The beam width was varied from 4 mm to 24 mm for each droplet diameter, with the power output results for each of the 3 droplet diameters tested illustrated in Fig. 2. The power output magnitude is over-estimated due to both the ideal behaviour of the simulated material, in addition to the sharp-impulse load duration used to represent non-trivial droplet impact dynamics.

Despite the difference in power output between simulated and fabricated devices, the result trend indicating that beams with lower bending stiffness generate greater output power proves valid.

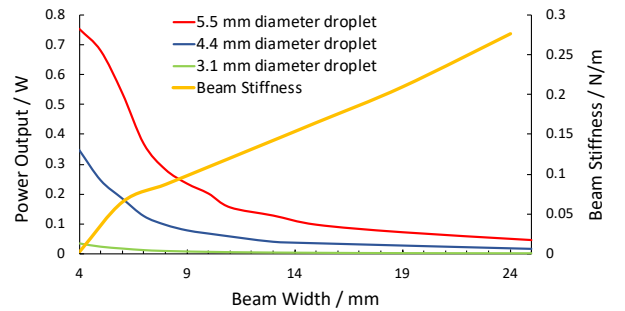


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Fig. 1. Illustration of numerical model geometry setup simulating water droplet impact upon cantilever beam end



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Fig. 2. Simulated effect of beam stiffness variation on peak power output from droplet impacts at terminal velocity.

323 **4. Fabrication and testing of P(VDF-TrFE) on**
 324 **stainless steel foil cantilever beams**

325 **4.1 Sample Fabrication**

326 A solution was prepared by dissolving 20% weight
 327 P(VDF-TrFE) co-polymer powder (70/30 mol ratio) in a
 328 solvent of dimethylformamide and acetone (volume ratio
 329 20/80). The solution was heated at 55°C in an oil bath
 330 and mechanically stirred for approximately 1 hour,
 331 before being degassed in an ultrasonic water bath for 1-2
 332 hours.

333 A micrometer adjustable applicator was used to
 334 uniformly spread the P(VDF-TrFE) solution on a
 335 stainless steel foil sheet. To accommodate for shrinkage
 336 during the annealing process, the applicator was set to
 337 produce 25 μm film thickness in order to achieve a final
 338 15 μm film thickness post heat treatment. After
 339 depositing, the sheet was placed in an oven at 100°C for
 340 5 minutes to allow the solution to dry. The sheet was then
 341 annealed for 2 hours at a temperature of 135°C in order
 342 to increase the piezoelectric material's crystallinity.

343 Following the heat treatment, the sheet was poled
 344 using a corona poling rig for 10 minutes at poling voltage
 345 of ~18 kV, before being cut into 6 samples of uniform
 346 length and thickness, but with varying widths of 6, 8, 10,
 347 24, 26 and 27 mm. In this instance, lengths of 63 mm
 348 were selected in order to achieve low sample resonant
 349 frequencies (~10 Hz), allowing for 20 mm of beam length
 350 for clamping and electrode connections. Relevant sample
 351 properties are given in Table II.

352 Silver electrodes of approximately 32 mm length, 2
 353 mm width, and 200 nm thickness were deposited through
 354 a shadow mask via e-beam evaporation. Wire
 355 connections were attached to the deposited electrode
 356 using silver electrode paste, before each sample was
 357 clamped using acrylic beams of 10 mm width, at a
 358 position of 10 mm away from the electrode connections
 359 area of the beams, as illustrated in Fig. 3. Finally, the
 360 samples were encapsulated using NeverWet@
 361 superhydrophobic surface treatment in order to isolate all
 362 electrical connections from water.

363 It was noted that samples suffered from slight initial
 364 deformation in some cases, due to both increasing mass
 365 and a residual stress gradient arising from the fabrication
 366 process. To negate any bias to testing results, two
 367 separate experiments were conducted – one which
 368 studied a series of samples, and a second which
 369 investigated a single sample in order to ensure
 370 consistency in experimental setup.

TABLE II
 LENGTH = 63 MM, THICKNESS = 40 μM SAMPLE EXPERIMENTAL VALUES

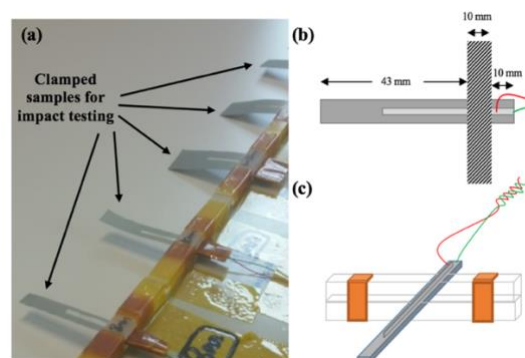
Sample Width / mm	Resonant Frequency / Hz	Electrically impedance matched load / MΩ	Measured beam stiffness / N/m
6	10	8	0.0943
8	10	8	0.1170
10	10	8	0.1290
24	11	5	0.3070
26	10	5	0.3067
27	10	8	0.3066

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372

TABLE III
 LENGTH = 63 MM, THICKNESS = 40 μM, PROGRESSIVE TRIM SAMPLE EXPERIMENTAL VALUES

Sample Width / mm	Resonant Frequency / Hz	Electrically impedance matched load / MΩ	Measured beam stiffness / N/m
4	10	8	0.067
7	10	8	0.094
11	10	8	0.134
15	10	8	0.188
19	10	8	0.188
25	10	8	0.235
27	10	8	0.235

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Fig. 3. Images to illustrate clamping set-up for 63 mm length samples. (a) depicts a photograph of samples used for testing. (b) illustrates sample dimensions - free beam length is 43 mm, with 10 mm being clamped securely with acrylic supports, and the last 10 mm of the beam reserved for electrode connections. (c) illustrates sample clamping set-up graphically.

383 **4.2 Effect of beam stiffness variation on power**
 384 **output from droplet impact**
 385

386 The experimental test rig used to investigate beam
 387 stiffness variation is illustrated in Fig. 4. The rig
 388 consisted of an adjustable clamp, able to adjust to a height
 389 of approximately 2.4 m, a base stand and drip tray.
 390 Syringes placed in the clamp were used to dispense
 391 droplets of diameter 3.1 mm, 4.4 mm and 5.5 mm from
 392 different heights. Samples were fixed onto a single, solid
 393 acrylic board, which was positioned on top of the drip tray
 394 underneath the clamped syringe nozzle. The centroid
 395 location of droplet impact was no greater than 5 mm from
 396 the end of each beam's free end. The piezoelectric beam
 397 output was connected to an electrically impedance
 398 matched resistive load, with the generated voltage
 399 response across the resistor measured using an
 400 MDO3000 series Tektronix® Oscilloscope.

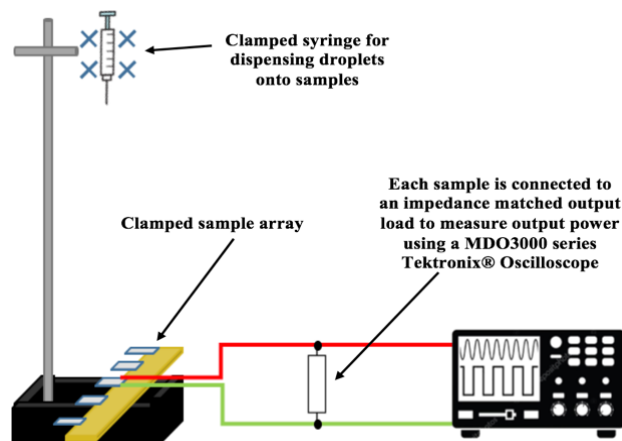
401 Two rounds of testing took place – the first measured
 402 the power output generated from the impact of the 3.1
 403 mm, 4.4 mm and 5.5 mm diameter droplets dispensed
 404 from heights of 0.5 m, 1 m, 1.5 m and 2 m for each of
 405 the 6 samples. It was observed that larger samples tended
 406 to suffer initial displacement due to their increased mass.
 407 In order to validate that this did not detrimentally effect
 408 any result trend achieved, a second test round was
 409 carried out.

410 This test round focused on progressively trimming the
 411 width of the 27 mm wide sample in isolation, illustrated
 412 in Fig. 5, whilst impact testing with the same three
 413 different droplet diameters from heights of 0.5 m, 1 m
 414 and 1.5 m. Testing was carried out in this manner to
 415 ensure that any physical differences between samples
 416 were negated. Furthermore, sample characteristics, such
 417 as initial displacement, were kept constant between width
 418 alterations adding further control to the experiment.
 419 Testing multiple samples, however, demonstrated that
 420 the stiffness values achieved were not unique to the
 421 individual sample used in the progressive trim testing.
 422 The properties of the sample used at each interval during
 423 the progressive trim testing round are displayed in Table
 424 III.

425 4.3 Droplet impact frequency effect on power 426 output

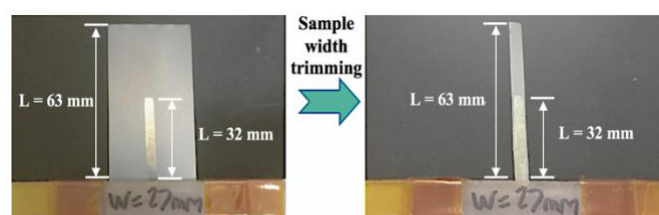
427 In order to test the effect of droplet impact frequency
 428 on power output, an Alaris IVAC P7000 syringe driver
 429 was used to dispense 3 mm-diameter droplets at a range
 430 of drip frequencies onto each sample beam end. This test
 431 round utilized the samples used in previous tests (widths
 432 of 4, 6, 8, 10, 24 and 26 mm). The properties of these
 433 samples are shown in Tables II and III. Fig. 5 presents the
 434 photos to illustrate progressive trim test sample
 435 modification. An IR detector was used to measure the
 436 drip frequency rate of the droplets, outputting a negative
 437 spike each time a droplet broke the IR beam, as shown in
 438 Fig. 6. This signal was measured using an oscilloscope
 439 and used to calculate drip frequency. Measured drip
 440 frequencies ranged from 1 Hz to at least 30 Hz for all
 441 samples.

442



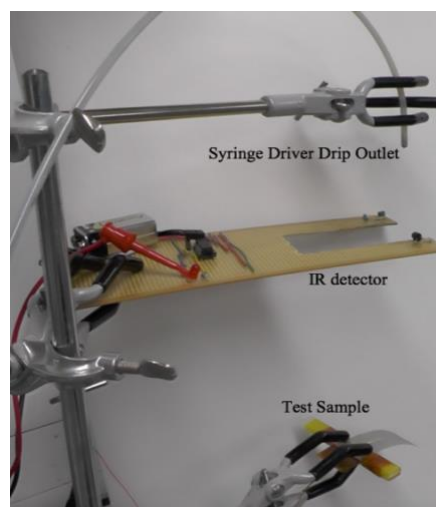
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Fig. 4. Diagram to illustrate droplet impact test setup



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Fig. 5. Photo to illustrate progressive trim test sample modification



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Fig. 6. Photograph to illustrate droplet impact frequency effect on power output test setup

451 5. Results and Discussion

452 5.1 Effect of beam stiffness variation on power 453 output – multiple samples results

454 The results in Fig. 7 show that beams with bending
 455 stiffness in the region of 0.1 N/m are better for droplet
 456 energy harvesting. The output energy per impact
 457 produced by the beam with 8 mm width (bending
 458 stiffness of 0.1170 N/m) produced the highest output
 459 energy levels throughout a range of different diameter

460 droplet impacts, peaking at an output of 28 nJ for a single
461 5.5 mm diameter droplet impact from a release height of
462 1.5 m. Fig. 8 displays the averaged energy outputs of all
463 release heights for each droplet diameter tested, in order
464 to highlight the peak energy output achieved by the 8 mm
465 width beam for each droplet diameter.

466 *5.2 Effect of beam stiffness variation on power* 467 *output – progressive trim testing of single 27 mm* 468 *width sample results*

469 When a single sample had its stiffness modified in the
470 case of the progressive trim tests, the results illustrated in
471 Fig. 9 again demonstrate that beams with bending
472 stiffness of approximately 0.1 N/m are better for droplet
473 harvesting, with the beam of width 7 mm, bending
474 stiffness 0.094 N/m, producing the highest energy output
475 levels per impact for the range of droplet diameters and
476 release heights tested. The peak output power for the 7
477 mm beam undergoing single droplet impact was 14 nJ,
478 generated by a 5.5 mm diameter droplet falling from a
479 release height of 1.5 m. Fig. 10 displays the averaged
480 energy outputs of all release heights for each droplet
481 diameter tested, in order to highlight the peak energy
482 output achieved by the 7 mm width beam for each droplet
483 diameter.

484 *5.3 Droplet impact frequency effect on power* 485 *output results*

486 The impact frequency testing results in Fig. 11 indicate
487 that, despite the non-trivial behaviour of droplet-surface
488 impact interactions, sample output power was better
489 when the frequency of incident impacting droplets onto
490 the sample beam ends' was close to/at the samples
491 resonant frequency. It was noted that the simulated power
492 output levels were a couple of magnitudes greater than
493 the practical results. It is proposed that this is due to the
494 simplified loading condition used in the simulation
495 environment, which neglects the complex fluid-structure
496 interaction of droplet impact in reality. Nonetheless,
497 these results support the conclusion that driving the
498 samples at their resonant frequency achieves optimal
499 power output. The slight difference in power output
500 peaks between the simulated and practical results are
501 likely due to mechanical imperfections in the
502 experimental sample set.

503 This is considered to be due to the disparities between
504 the simulation and experimental materials, which will
505 likely differ due to real-world processing requirements.
506 In each case, the power output was maximised when the
507 simulated/practical sample was driven by impacts
508 matching its resonant frequency.

509 *5.4 Result Analysis & Discussion*

510 Previous research carried out in the area of droplet
511 energy harvesting using piezoelectric materials has found
512 that the energy transfer efficiency of such systems is
513 typically very low, in the order of approximately 0.12%
514 [2]. This figure was achieved using a commercially
515 available piezoelectric sensor, the Pro-Wave (FS-2513P).
516 Although the theoretical energy transfer efficiency of the
517 results reported here was approximately 0.0013%,
518 representing the ideal energy available from a 5.5 mm
519 diameter droplet falling from 1.5 m as 1.28 mJ, it is
520 important to consider that the electrode area was
521 constrained in order to retain experimental consistency as
522 beam width variation was carried out.

523 Furthermore, a comparison of transducer energy
524 densities highlights the benefits of mechanical tailoring
525 to the excitation source. The results from the cited
526 literature utilising the commercial Pro-Wave (FS-2513P)
527 sensor [2] indicate an energy output of no more than 90
528 nJ from droplet impact. With an active volume of
529 approximately $0.975 \mu\text{m}^3$, the energy density is
530 calculated as 0.092 J/m^3 . In comparison, the peak energy
531 output of 28 nJ achieved here, across an active sample
532 volume of 1.76 nm^3 , represents an energy density of 15.9
533 J/m^3 .

534 As such, whilst the peak energy output achieved in this
535 research is not the highest possible, it demonstrates
536 significant energy transfer efficiency for the active
537 electrode areas used. It is proposed that the energy
538 density of the commercial sensors typically utilised in
539 other studies is notably lower due to the relatively stiff
540 Mylar coating used to encapsulate the sample.
541 Furthermore, with the sample clamped in a cantilever
542 beam orientation, the impact zone of the droplet upon the
543 transducer was targeted at the beam center, instead of the
544 beam end. This would have increased the transducer's
545 apparent stiffness to the droplet impact, decreasing the
546 responsiveness of the beam and resulting in further
547 degradation of the energy transfer efficiency. The
548 simulation and experimental results of this study indicate
549 that both low beam stiffness in the region of 0.1 N/m, in
550 addition to transducer resonant frequency being close
551 to/at the driving frequency of the energy source, are
552 required for efficient energy transfer to occur between the
553 impacting droplet and the piezoelectric transducer.

554 Furthermore, the results illustrate that a relatively
555 small variation in beam width, relating to a consequent
556 change in bending stiffness, can significantly affect the
557 energy transfer efficiency. For example, when impacting
558 a 3.1 mm diameter droplet from 0.5 m onto the
559 progressive trim test sample, an energy output of 5.1 nJ
560 was achieved with a 7 mm width beam, compared to a 1
561 nJ output obtained from the beam at 27 mm width. This
562 represents an approximate 5-times increase in energy
563 transfer efficiency, in relation to a variation of 20 mm in
564 sample width, underlining the importance of transducer

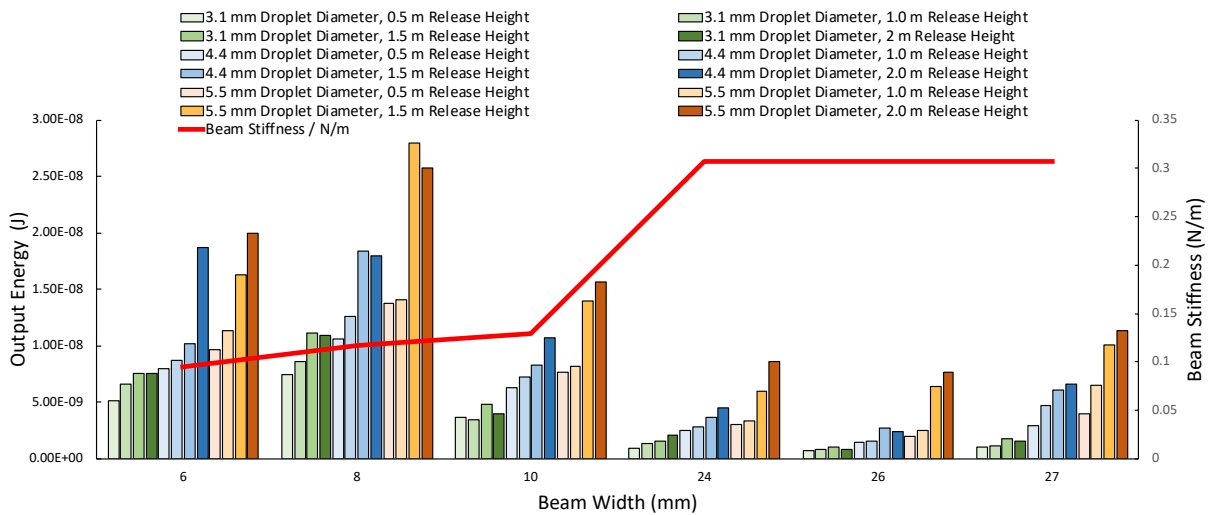
565 parameter tailoring depending on the expected droplet
 566 size and impact speed.

567 Similarly, in situations where continuous droplet
 568 impact is expected, the results achieved here underline
 569 the corresponding importance of transducer resonant
 570 frequency matching. For example, the sample of width 8
 571 mm produced approximately 3.6 nW peak output power
 572 when driven with droplets at a frequency of 1 Hz.
 573 However, when driven with droplets at its resonant
 574 frequency of 10 Hz, a peak output power of 12 nW was
 575 achieved, representing a 4-time increase in output power.

576 Given the significant effect of such parameter tailoring
 577 on energy transfer efficiency, should energy harvesting
 578 from rainfall be a target application, the results suggest
 579 that greater energy efficiency can be achieved through
 580 accumulating the rainfall in a storage tank initially,
 581 before dispensing droplets using a tiered outlet system.
 582 This would serve to control the droplet size and impact
 583 frequency, in order to better match the key transducer
 584 parameters outlined in this report, resulting in much

585 greater power output than perhaps otherwise achievable
 586 through attempting to harvest energy from direct rainfall
 587 impact. Should a sufficient amount of water be
 588 accumulated, an array of energy harvesting devices
 589 fabricated using the design guidelines outlined in this
 590 research could drive low power systems by storing and
 591 supplying the accumulated energy via an efficient energy
 592 management system, such as an ultra-sharp transistor
 593 based switching circuit [22]. For example, such energy
 594 could be used to drive medium to short range wireless
 595 sensor applications, utilising components such as
 596 Microsemi's ZL70550 RF transceiver [23] supported by
 597 an ultra-low power management IC such as the
 598 S6AE10xA energy harvesting PMIC from Cypress
 599 Semiconductor [24].

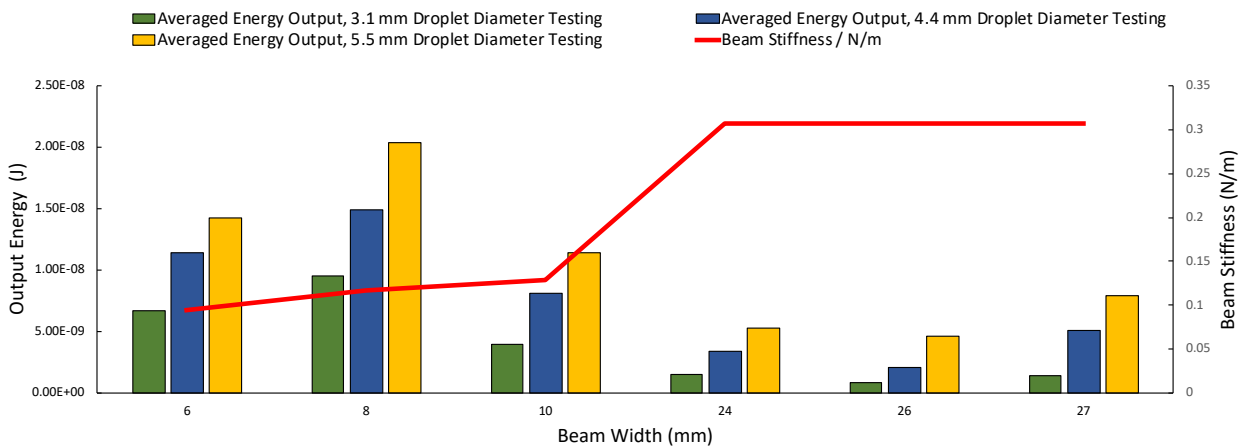
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Fig. 7. Average output energy per beam width as a function of impacting droplet diameter and release height, multiple test samples

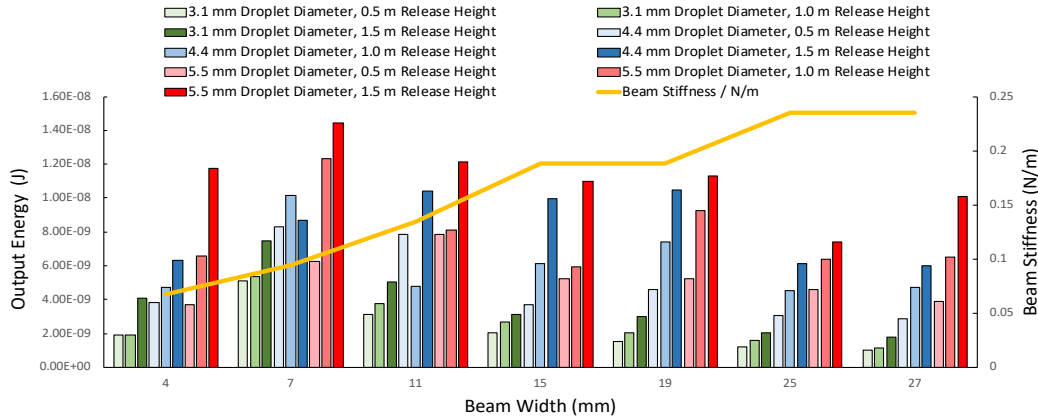


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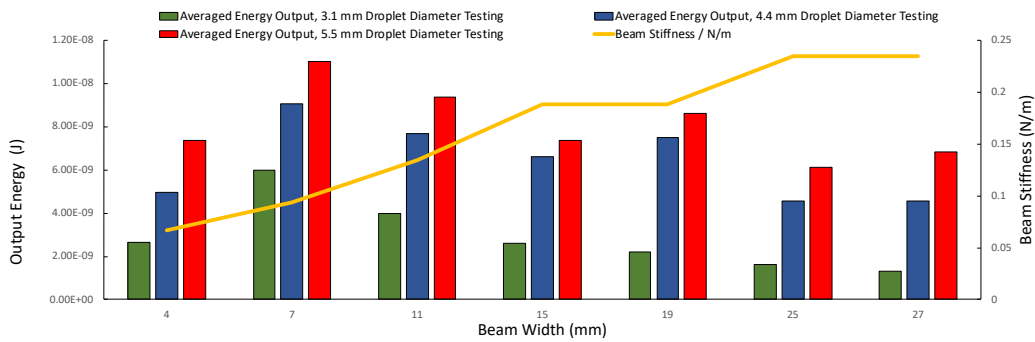
Fig. 8. Graph displaying all droplet release height results from Fig. 7. (multiple test samples) averaged to give a single output energy, per beam width, for each droplet diameter tested



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Fig. 9. Average output energy per beam width as a function of impacting droplet diameter and release height, progressive trim testing

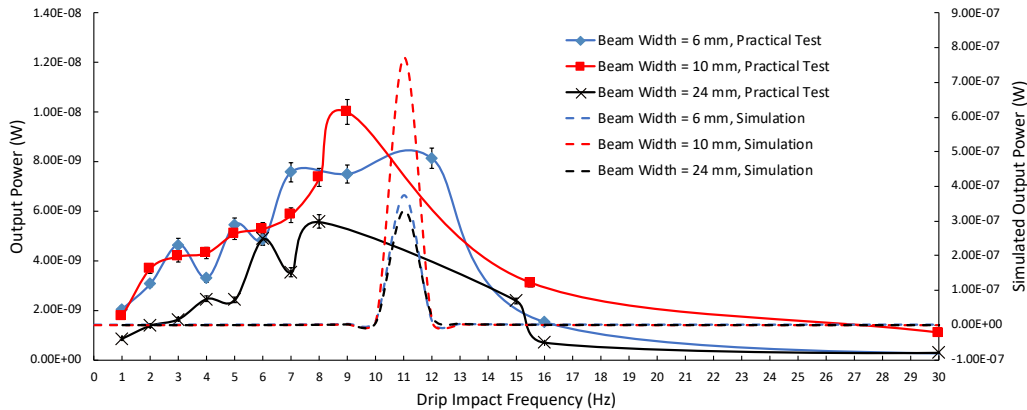


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Graph displaying all droplet release height results from Fig. 9 (progressive trim testing) averaged to give a single output energy, per beam width, for each droplet diameter tested



612

613

Sample output power as a function of droplet impact frequency

614 **6. Conclusion**

615 In this paper, we investigated the key parameters that affect
 616 the efficiency of energy transfer between incident water
 617 droplets and piezoelectric cantilever structures made of
 618 stainless steel foil coated with piezoelectric P(VDF-TrFE).
 619 The experimental results and analyses achieved underline the
 620 importance of appropriate transducer stiffness and resonance
 621 frequency matching, depending on the expected impact
 622 behaviour of incident droplets. The resulting enhanced
 623 power output was demonstrated to be of magnitude scale with

624 optimal piezoelectric transducer stiffness and resonance
 625 frequency matching.

626 From the results of both impact tests, it was demonstrated
 627 that for droplets of diameter 3.1 mm to 5.5 mm, impacting
 628 from heights between 0.5 to 2.0 m, it is desirable to utilise
 629 piezoelectric transducer beams of bending stiffness in the
 630 range of 0.067 to 0.134 N/m in order to achieve good energy
 631 transfer efficiency, resulting in enhanced electrical power
 632 output. Additionally, it was found that the power output was
 633 further increased when the impact frequency was close to/at

634 the transducers resonant frequency. Our investigation res 688
635 provide valuable guidelines for piezoelectric transdu 689
636 design for effective mechanical energy harvesting from 690
637 droplets. 691

638 **Acknowledgements** 692

639 S. C. J. Jellard would like to sincerely thank the IMRE lab 693
640 staff for their support and guidance throughout this research 694
641 in particular Dr. Weng Heng Liew, Dr. Meysam Sharifzadeh 695
642 Mirshekarloo and Dr. Lei Zhang. The authors at IMRE 696
643 acknowledge the partial support by Singapore Maritime 697
644 Institute under the Maritime Sustainability R&D 698
645 Programme, Project ID: SMI-2015-MA-07 (IMRE/16 699
646 7P1125). S. C. J. Jellard would also like to thank the 700
647 University of Southampton's Faculty of Engineering and 701
648 Physical Sciences for the Tizard studentship and A*STAR 702
649 for the ARAP Scholarship. 703

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