Near-Physical Point Lattice Calculation of Isospin-Breaking Corrections to $K_{\ell 2}/\pi_{\ell 2}$

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Andrew Zhen Ning Yong,[a,*] **Peter Boyle,**[a,b] **Matteo Di Carlo,**[a] **Felix Erben,**[a] **Vera Gülpers,**[a] **Maxwell T. Hansen,**[a] **Tim Harris,**[a] **Nils Hermansson-Truedsson,**[c] **Raoul Hodgson,**[a] **Andreas Jüttner,**[d,e] **Antonin Portelli**[a] **and James Richings**[a]

[a]*Department of Physics & Astronomy, University of Edinburgh,*
Peter Guthrie Tait Road, King's Buildings, Edinburgh, United Kingdom

[b]*Physics Department, Brookhaven National Laboratory, Upton, NY, USA*

[c]*Albert Einstein Center for Fundamental Physics, Institute for Theoretical Physics, Universität Bern,*
Sidlerstrasse 5, 3012 Bern, Switzerland

[d]*School of Physics and Astronomy, University of Southampton, University Road, Southampton SO17 1BJ,*
United Kingdom

[e]*CERN, Physics Department, 1211 Geneva 23, Switzerland*

E-mail: andrew.yong@ed.ac.uk

In recent years, lattice determinations of non-perturbative quantities such as f_K and f_π, which are relevant for V_{us} and V_{ud}, have reached an impressive precision of $O(1\%)$ or better. To make further progress, electromagnetic and strong isospin breaking effects must be included in lattice QCD simulations.

We present the status of the RBC/UKQCD lattice calculation of isospin-breaking corrections to light meson leptonic decays. This computation is performed in a (2+1)-flavor QCD simulation using Domain Wall Fermions with near-physical quark masses. The isospin-breaking effects are implemented via a perturbative expansion of the action in α and $(m_u - m_d)$. In this calculation, we work in the electro-quenched approximation and the photons are implemented in the Feynman gauge and QED$_L$ formulation.

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*Speaker

1. Introduction

One of the ongoing efforts to search for new physics beyond the Standard Model (SM) of particle physics is to test the unitarity of the Cabibbo-Kobayashi-Maskawa (CKM) matrix. Since all SM and potential beyond-the-SM particles contribute to hadronic decays via virtual corrections, any deviations from the theoretical expectation of unitarity may hint at inconsistencies between Nature and the SM in the flavor sector. In 2020, the PDG [1] reports the following tension with unity in the first row of the CKM matrix:

$$|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 0.9985 \pm 0.0005. \tag{1}$$

The determination of CKM matrix elements is an endeavour involving the joint effort of precise experimental measurements and predictions from theory. One of the experimental avenues of interest is the leptonic decay channels of light pseudoscalars, $P^{\pm} = K^{\pm}, \pi^{\pm}$. The tree-level expression is

$$\Gamma_P^{\text{tree}} = \frac{G_F^2}{8\pi} M_{P^+} m_{\ell^+}^2 \left(1 - \frac{m_{\ell^+}^2}{M_{P^+}^2}\right)^2 |f_P|^2 |V_{qq'}|^2, \tag{2}$$

where G_F is the Fermi constant and $V_{qq'}$ is the CKM matrix element between quark flavors q and q'. The pseudoscalar decay constant, f_P, is a quantity in the isosymmetric theory, where $\alpha = 0$ and $\delta m \equiv m_u - m_d = 0$. We note that since the decay constant encapsulates the non-perturbative effects of QCD, it is evaluated numerically using lattice QCD methods. Traditionally, these calculations were performed in the $\alpha = 0, \delta m = 0$ regime. However, since recent lattice determinations of f_K and f_π have attained percent-level precision [2], further progress will necessitate the inclusion of isospin-breaking (IB) effects since $\alpha \sim \frac{\delta m}{\Lambda_{\text{QCD}}} \sim 1\%$.

Due to the experimental challenge in distinguishing between final states with or without a soft photon, only the inclusive rates are measured in the case of pions and kaons. For the extraction of $V_{qq'}$ from experimental data, we combine the tree-level expression with its IB corrections with

$$\Gamma(P^+ \to \ell^+ \nu_\ell[\gamma]) = \Gamma_P^{\text{tree}}(1 + \delta R_P) + O(\alpha^2, \delta m^2, \alpha \delta m) \tag{3}$$

where δR_P contains the leading order IB contributions to the tree-level width.

In this proceeding, we report on our determination of $\frac{|V_{us}|}{|V_{ud}|}$, which can be extracted from a ratio of muonic inclusive rates via

$$\frac{|V_{us}|^2}{|V_{ud}|^2} = \frac{\Gamma(K^+ \to \mu^+ \nu_\mu[\gamma])}{\Gamma(\pi^+ \to \mu^+ \nu_\mu[\gamma])} \frac{M_{K^+}}{M_{\pi^+}} \frac{M_{\pi^+}^2 - m_{\mu^+}^2}{M_{K^+}^2 - m_{\mu^+}^2} \mathcal{F}^{-2}, \tag{4}$$

where

$$\mathcal{F} = \frac{f_K}{f_\pi} \sqrt{1 + \delta R_{K\pi}} \tag{5}$$

and $\delta R_{K\pi} \equiv \delta R_K - \delta R_\pi$. By using experimental measurements of the inclusive rate (*i.e.* branching ratios and mean lifetimes) and masses, the term $\delta R_{K\pi}$ encapsulates the ratio of IB correction to the hadronic decay amplitudes. The lattice determination of $\delta R_{K\pi}$ is thus the main focus of this calculation.

2. Calculation strategy

The methodology of separating the calculation of real and virtual IB corrections, along with the regulating and removal of finite volume effects (FVE) from the amplitudes computed on the lattice, was first proposed in [3]. Following this, a calculation of $\delta R_{K\pi}$ and $\frac{|V_{us}|}{|V_{ud}|}$ was accomplished in [4, 5]. Here, we adopt a similar strategy and write the inclusive rate as

$$\Gamma(P^+ \to \mu^+ \nu_\mu[\gamma]) = \lim_{L \to \infty}(\Gamma_0(L) - \Gamma_0^{(2)}(L)) + \lim_{\lambda \to 0}(\Gamma_0^{\mathrm{univ}}(\lambda) + \Gamma_1(\lambda, \Delta E_\gamma)), \qquad (6)$$

where the subscript denotes the number of real photons in the final state. The virtual corrections in Γ_0 are evaluated with numerical simulations since all momentum modes of the photon are involved in the interaction with the initial hadron. Thus, the lattice box size, L, is a natural choice of IR regulator. To remove the FVE of the lattice calculation, we introduced

$$\Gamma_0^{(n)}(L) = \Gamma_P^{\mathrm{tree}}\left(1 + 2\frac{\alpha}{4\pi}Y_P^{(n)}(L)\right) + O\left(\frac{1}{L^{n+1}}\right). \qquad (7)$$

The $O(L^{-1})$ FVE's were calculated in [6], and the $O(L^{-2})$ structure-dependent FVE's in [7]. Thus, the residual FVE coming from our lattice calculation now begins at $\Gamma_0(L) - \Gamma_0^{(2)}(L) \sim O\left(L^{-3}\right)$.

For the real photon contribution, we implement the analytic approach in [3]. In Equation (6),

$$\lim_{\lambda \to 0}(\Gamma_0^{\mathrm{univ}}(\lambda) + \Gamma_1(\lambda, \Delta E_\gamma)) = \Gamma_P^{\mathrm{tree}}\left(1 + \frac{\alpha}{4\pi}\delta\Gamma_{1,P}(\Delta E)\right), \qquad (8)$$

where ΔE is an energy threshold, below which the photon is sufficiently soft that it treats the initial hadron as a point-like particle. Here, for the IR regulator we use a fictitious photon mass, λ. Additionally, an intermediate 'universal' term, Γ_0^{univ}, is introduced to ensure the IR divergences in Equation (6) cancel numerically. In the following, we will discuss in detail the contributions going into $\Gamma_0(L)$.

3. Leptonic matrix elements from Euclidean correlation functions

The 4-fermion operator associated to this decay is

$$O_W = (\overline{\nu}\gamma_L^\tau \mu)(\overline{q}_1 \gamma_L^\tau q_2), \qquad (9)$$

where $\gamma_L^\tau = \gamma^\tau(\mathbb{1} - \gamma^5)$. For a two-body decay, the rate has a simple form:

$$\Gamma = \mathcal{K}\sum_{r,s}|\mathcal{M}_P^{r,s}|^2 \qquad (10)$$

where $\mathcal{K}$ contains the kinematic factors from the phase space integral and

$$\mathcal{M}_P^{r,s} \equiv \langle \mu^+, r; \nu_\mu, s|O_W|P^+\rangle, \qquad (11)$$

where r, s are the polarisations of the on-shell final state fermions. In the rest frame of the pseudoscalar, the amplitude in the isosymmetric theory of QCD is

$$\overline{\mathcal{M}}_P^{r,s} = \left(\overline{u}_\nu^r \Gamma_L^0 v_\mu^s\right) \overline{A}_P, \tag{12}$$

where the axial matrix element is

$$\overline{A}_P \equiv \langle 0|\overline{q}_2 \Gamma_L^0 q_1|\overline{P}\rangle = \overline{M}_P f_P, \tag{13}$$

with $\overline{M}_P$ the pseudoscalar mass in the $\alpha = 0$ isosymmetric theory[1]. At $O(\alpha, \delta m)$, all the possible virtual IB correction to the isosymmetric amplitude are presented in Figures 1 and 2, which we can write as

$$\frac{1}{2}\frac{\delta\Gamma_{0,P}}{\overline{\Gamma}_{0,P}} \equiv \frac{\sum_{r,s}\text{Re}\left[\overline{\mathcal{M}}_P^{r,s\,\dagger}\delta\mathcal{M}_P^{r,s}\right]}{\sum_{r,s}|\overline{\mathcal{M}}_P^{r,s}|^2} = \frac{\delta A_P}{\overline{A}_P} - \frac{\delta M_P}{\overline{M}_P} + \frac{\delta\mathcal{M}_{P\mu}}{\overline{\mathcal{M}}_P}, \tag{14}$$

where $\delta\mathcal{M}_{P\mu}$ is the amplitude correction corresponding to diagrams (e,f) in Figure 2. The contributions in Equation (14) are extracted from correlation functions generated in lattice simulations. Since $\alpha \sim \frac{\delta m}{\Lambda_{\text{QCD}}} \sim 1\%$ in the low energy regime, the QED and strong IB (SIB) corrections can be treated as a perturbation to our path integral expression [12]. The full QCD+QED expectation value for some observable O is

$$\langle O\rangle = \langle O\rangle_0 + \sum_q (m_q - \overline{m}_q)\frac{\partial}{\partial m_q}\langle O\rangle\bigg|_{m_q=\overline{m}_q} + \frac{1}{2!}e^2\frac{\partial^2}{\partial e^2}\langle O\rangle\bigg|_{e=0} + \dots, \tag{15}$$

where

$$\langle O\rangle = \frac{1}{\mathcal{Z}}\int \mathcal{D}[\psi]\mathcal{D}[\overline{\psi}]\mathcal{D}[U]\mathcal{D}[A]\, O[\psi,\overline{\psi},U,A]\, e^{-S_F[\psi,\overline{\psi},U,A]}e^{-S_G[U]}e^{-S_\gamma[A]} \tag{16}$$

is the path integral over the usual quark fields, $\psi, \overline{\psi}$; the $SU(3)$ gluonic fields, U; and the photon fields, A. Here, $\langle O\rangle_0$ is the QCD-only ($m_u = m_d$) expectation value. The SIB and QED corrections are then extracted from the slopes of the full correlation function, $\langle O\rangle$. Respectively, these derivatives generate correlation functions that contain scalar insertions (Figure 1) or two insertions of the electromagnetic (EM) current at $O(\alpha)$ (Figure 2). Together with the subtraction of FVE and the inclusion of the real photon contribution,

$$\delta R_{K\pi} = \left(\frac{\delta\Gamma_{0,K}(L)}{\overline{\Gamma}_{0,K}(L)} - \frac{\delta\Gamma_{0,\pi}(L)}{\overline{\Gamma}_{0,\pi}(L)}\right) - 2\frac{\alpha}{4\pi}\left(Y_K^{(2)}(L) - Y_\pi^{(2)}(L)\right) + \frac{\alpha}{4\pi}\left(\delta\Gamma_{1,K}(\Delta E) - \delta\Gamma_{1,\pi}(\Delta E)\right). \tag{17}$$

In the following, we classify Figure 1(a,b) and 2(a,b,c) collectively as **factorisable** diagrams and Figure 2(e,f) as **non-factorisable** diagrams. We do not consider the lepton self energy term (Figure 2(d)) as this is absorbed in the lepton renormalisation.

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[1]Such a theory is scheme-dependent and beyond the scope of the proceedings. Interested readers may refer to *e.g.* [8–11].
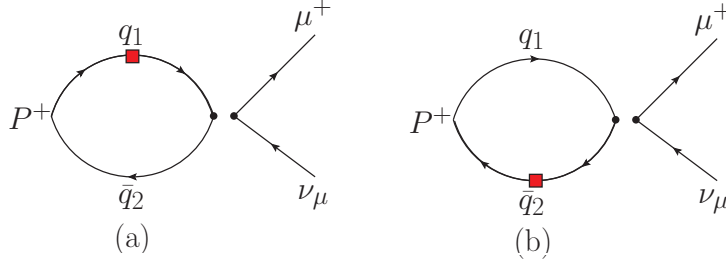
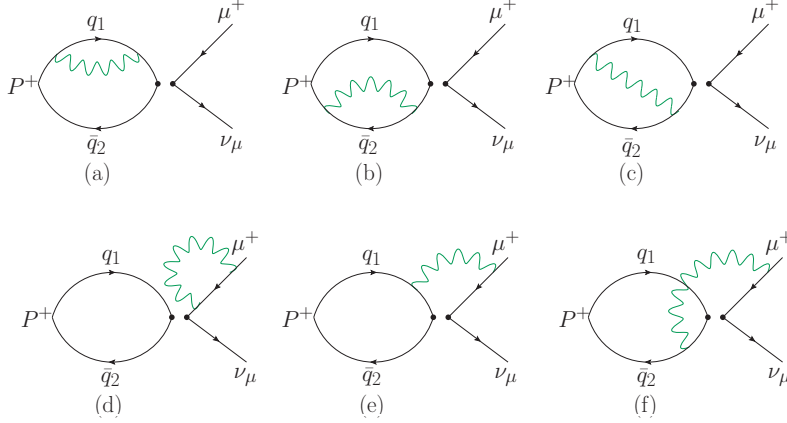**Figure 1:** Feynman diagrams of scalar insertions on quark legs (marked with red boxes).

Figure 2: Feynman diagrams of all possible insertions of the electromagnetic current (marked with green squiggle lines) at $O(\alpha)$. QED interactions with sea quarks are neglected.

4. Lattice methodology & implementation

For this calculation, we generate correlators in a $48^3 \times 96$ lattice using near-physical Möbius Domain Wall Fermions (DWF). The Domain Wall height and the length of the fifth dimension are $M_5 = 1.8$ and $L_s = 24$, respectively [13]. To reduce the computational cost of generating near-physical light quark propagators, we make use of ZMöbius fermions [14] together with the eigenvectors generated by the RBC/UKQCD collaboration for deflation. The 60 QCD gauge configurations used are also generated by the RBC/UKQCD collaboration using the Iwasaki gauge action [15]. The sea quark masses are $am_l^{sea} = 0.00078, am_s^{sea} = 0.0362$. We choose the valence up- and down-quark masses to have the same value as the sea, $am_u = am_d = am_l^{sea}$ and similarly for the valence strange quarks, $am_s = am_s^{sea}$. In this setup, the lattice spacing is $a^{-1} = 1.7295(38)$GeV and the ensemble pion mass is $M_\pi = 139.15(36)$ MeV.

The correlators are built from quark propagators with Coulomb gauge-fixed wall sources. As such, we must generate correlators with both wall and point sinks in order to extract the axial matrix element. The implementation of QED on our lattice simulation is as follows: we remove the photon's spatial zero mode with the QED$_L$ formalism. As in the case of the scalar current, we sequentially insert a local EM current in Feynman gauge to obtain a sequential propagator with an $A\!\!\!/$ insertion. The correlators built from these propagators are then renormalised by appropriate factors of Z_V [16]. To build the hadron-leptonic correlators corresponding to diagrams (e,f) in Figure 2,

we include the charged lepton on the lattice. This is done by generating muon propagators with a free DWF action, using an input mass such that the pole mass of our DWF propagator matches with the experimentally measured value. The propagator is given twisted boundary conditions to conserve 4-momentum of this decay. The neutrino is a spectator fermion in this whole process, so we choose to omit it in the lattice simulation and include it in the analysis stage. We put the pseudoscalar interpolator at the origin and insert the 4-fermion operator on every timeslice, t_H, with the muon source-sink separation fixed at, $t_\mu - t_H = 12, 16, \ldots, 36, 40$.

In this calculation, we omit SIB contributions coming from sea quarks. We are also working in the electro-quenched approximation of QED - treating the sea quarks as electrically neutral.

5. Physical predictions from lattice calculations

Since the two sources of isospin-breaking are $O(1\%)$ effects, we can perform a linear expansion about the physical point and treat these effects as shifts from the isosymmetric point, where $\alpha = 0, \delta m = 0$. Let X be the observable of interest (*e.g.* hadronic mass). Then, at $O(\alpha, \delta m)$,

$$X = \overline{X} + \sum_q \Delta m_q \partial_{m_q} X|_{m_f = \overline{m}_f} + \frac{1}{2!} e^2 \partial_{e^2} X|_{e=0} + O(\alpha^2, (\Delta m_q)^2, \alpha \Delta m). \tag{18}$$

where $\partial_{y^n} X = \partial^n X / \partial y^n$. Working to this order, we can set the electric charge in terms of the fine structure constant in the Thomson limit, $e^2 = 4\pi \alpha_{\text{EM}}$[2]. The set of Δm_q are the isosymmetric-to-physical point bare quark mass shifts. For a theory of QCD with 3 flavors, we can solve for the three Δm_q's by imposing the following mass ratios:

$$\frac{(aM_P)^2}{(aM_{\Omega^-})^2} = \frac{\left(M_P^{exp}\right)^2}{\left(M_{\Omega^-}^{exp}\right)^2}, \tag{19}$$

where we choose $P = \pi^+, K^+, K^0$. Scale setting can also be done by considering the following dimensionful constraint:

$$a = \frac{aM_{\Omega^-}}{M_{\Omega^-}^{exp}}, \tag{20}$$

where we chose the omega baryon owing to the clean signal from lattice simulations.

An additional step is required of our setup. Namely, we are not simulating at but close to the desired isosymmetric point and thus we must correct for this mismatch before calculating the IB corrections. Here, we note that since there is no natural phenomena which interacts exclusively with the strong nuclear force, this unphysical definition of the isosymmetric point will depend on the separation scheme we prescribe to our calculation. To that end, consider the following mesonic quantities,

$$M_{ud}^2 = \frac{1}{2}\left(M_{\overline{u}u}^2 + M_{\overline{d}d}^2\right) \approx 2Bm_{ud} + \ldots$$
$$\Delta M^2 = M_{\overline{u}u}^2 - M_{\overline{d}d}^2 \approx 2B(m_u - m_d) + \ldots \tag{21}$$
$$M_{K\chi}^2 = \frac{1}{2}\left(M_{K^+}^2 + M_{K^0}^2 - M_{\pi^+}^2\right) \approx 2Bm_s + \ldots$$

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[2]$\alpha_{\text{EM}} = 7.2973525693(11) \times 10^{-3}$[17]

where $\overline{q}q$ are neutral pseudoscalars made from connected-only propagators. Here, we have used the fact that, at leading order partially-quenched χPT [18, 19], these squared masses are proportional to the quark masses we are interested in, with B the chiral condensate. Thus, we can tune our setup to the isosymmetric point by setting $\alpha = 0$ in Equation (18) and determine another set of bare quark mass shifts, $\{\Delta m'_q\}$, with the following constraints:

$$M_{ud}^2 = \left(M_{\pi^0}^{exp}\right)^2, \quad \Delta M^2 = 0, \quad M_{K\chi}^2 = \frac{1}{2}\left(\left(M_{K^+}^{exp}\right)^2 + \left(M_{K^0}^{exp}\right)^2 - \left(M_{\pi^+}^{exp}\right)^2\right). \tag{22}$$

Most notably, by fixing $\Delta M^2 = 0$, we emulate the constraint in which $\delta m = 0$.

6. Constructing the amplitude correction

6.1 Factorisable IB correction

The extraction of IB correction to the amplitude from correlators corresponding to Figure 1(a,b) and Figure 2(a,b,c) proceeds in a manner analogous to the standard two-pt analysis, *i.e.* we need the pseudoscalar (p) and axial-pseudoscalar (a) interpolators:

$$\phi_p = \overline{q}_2\gamma^5 q_1 \quad \text{and} \quad \phi_a = \overline{q}_2\gamma^0\gamma^5 q_1. \tag{23}$$

Excluding backward propagating effects, the QCD-only pp/pa−correlator is

$$\overline{C}_P^{pj}(t) \equiv \int d^3x \, \langle 0|T\left\{\phi_j(\vec{x},t)\phi_p^\dagger(0)\right\}|0\rangle \, e^{-i\vec{k}\cdot\vec{x}} = c_P^{pj} e^{-\overline{M}_P t}, \tag{24}$$

with $c_P^{pj} = \langle 0|\phi_j|P\rangle \langle P|\phi_p^\dagger|0\rangle / 2\overline{M}_P$. For pp/pa−correlators containing either electromagnetic or scalar currents, we have

$$\partial_{g_k} C_P^{pj}(t) = \left(\partial_{g_k} c_P^{pj} + c_P^{pj} \partial_{g_k}\overline{M}_P t\right) e^{-(\overline{M}_P + \Delta g_k \partial_{g_k}\overline{M}_P)t}, \tag{25}$$

where $g_k = \{e^2, m_l, m_s\}$ and $\Delta g_k = g_k - \overline{g}_k$. Taking the ratio of Equation (25) and (24), we have

$$R_{P,g_k}^{pj}(t) = \frac{\partial_{g_k} c_P^{pj}}{\overline{c}_P^{pj}} - \partial_{g_k}\overline{M}_P t. \tag{26}$$

Since the QED and SIB correlator ratios share a common mass parameter, $\overline{M}_P$, we include R_{P,g_k}^{pj} for all j and g_k in a combined fit. In Figure 3, we present a preliminary result of this combined fit.

6.2 Non-factorisable IB correction

The hadron-leptonic correlators on the lattice are generated with a modified 4-fermion operator:

$$\widetilde{O}_{W,\beta} = (\gamma_L^\tau \mu)_\beta (\overline{q}_1\gamma_L^\tau q_2), \tag{27}$$
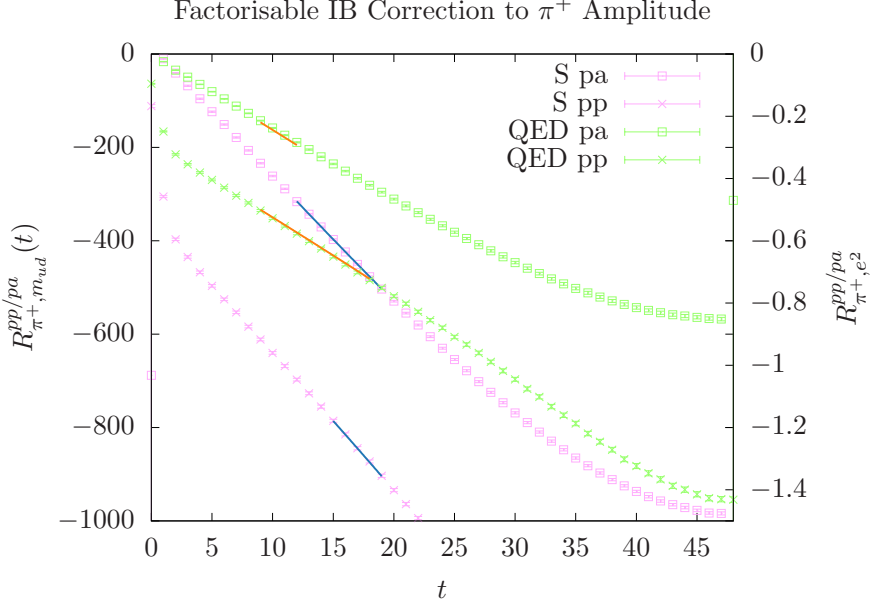
Figure 3: A combined fit of all correlator ratios defined in Equation (25). The pink points are the m_{ud} scalar insertion data and should be read with the left y-axis, while the green points are the QED data and should be read with the right y-axis. The blue and orange lines are the fits to the scalar insertion and QED data, respectively. The error bands are not visible. For this combined fit, $\chi^2/dof = 70/52 = 1.35$, p−value= 0.099.

where β is an open spinor index. The spectral representation of these novel correlators is

$$
\begin{aligned}
\overline{C}_P(t_H, t_\mu - t_H)_{\beta_1 \beta_2} &\equiv \int \prod_{j=\mu,H} d^3 x_j \, e^{-i \vec{p}_j \cdot \vec{x}_j} \langle 0| \left(\overline{\mu}(\vec{x}_\mu, t_\mu) \Gamma_L^0 \right) \widetilde{O}_W(\vec{x}_H, t_H) \phi^\dagger(0) |0\rangle_{\beta_1 \beta_2} \\
&= -\frac{\langle P|\phi_P^\dagger|0\rangle}{4 E_\mu \overline{M}_P} \left[\overline{\mathcal{M}}_P \cdot (\not{p}_\mu - i m_\mu) \right]_{\beta_1 \beta_2} e^{-\overline{M}_P t_H} e^{-E_\mu(t_\mu - t_H)}
\end{aligned}
\tag{28}
$$

where $\beta_{1,2}$ are spinor indices. The correlator containing electromagnetic currents has an analogous functional form. In order to extract the non-factorisable amplitude correction, we saturate the spinor indices by including the missing neutrino leg, giving us a trace over the correlator. Then, taking the ratio of the QED and QCD-only hadron-leptonic correlator, we find that at large time separations,

$$
\begin{aligned}
R_{P\mu}(t_H, t_\mu - t_H) &= \frac{\mathrm{Tr}\left[\not{p}_\nu \partial_{e_q} \partial_{e_\mu} C_P(t_H, t_\mu - t_H) \gamma_L^0 \right]}{\mathrm{Tr}\left[\not{p}_\nu \overline{C}_P(t_H, t_\mu - t_H) \gamma_L^0 \right]} \\
&\xrightarrow{t_\mu \gg t_H \gg 0} \frac{\mathrm{Tr}\left[\not{p}_\nu \partial_{e_q} \partial_{e_\mu} \mathcal{M}_P(\not{p}_\mu - i m_\mu) \gamma_L^0 \right]}{\mathrm{Tr}\left[\not{p}_\nu \overline{\mathcal{M}}_P(\not{p}_\mu - i m_\mu) \gamma_L^0 \right]}.
\end{aligned}
\tag{29}
$$

For each pseudoscalar, we have eight of these correlator ratios, corresponding to the different muon source-sink separations, $t_\mu - t_H$. We perform a 2D combined fit in the t_H and $(t_\mu - t_H)$ direction and Figure 4 shows the result of this fit analysis.
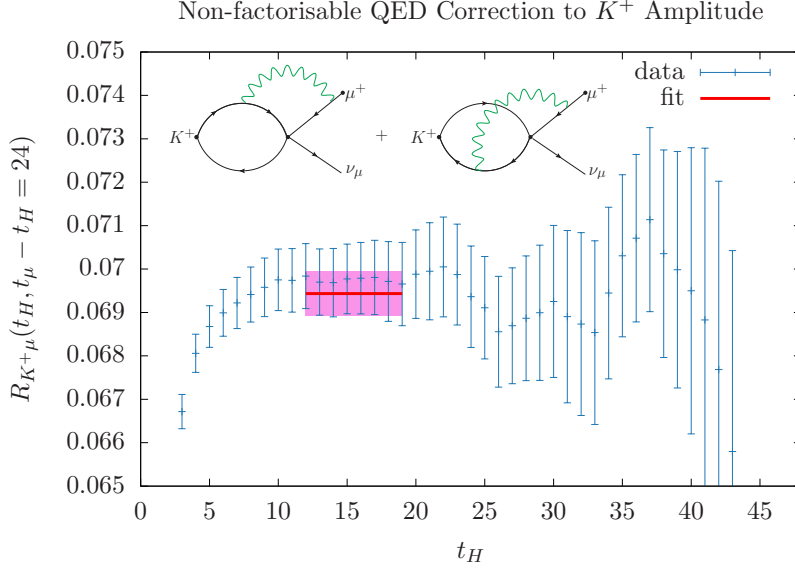
Figure 4: The non-factorisable correlator ratio defined in Equation (29). The red line and pink error band is the result of a combined fit of the hadron-leptonic correlators, in which only $t_\mu - t_H = 24$ is shown here. For this combined fit, $\chi^2/dof = 29/31 = 0.94$, $p-$value$= 0.812$.

6.3 Combining lattice and analytic results

We are now in the position to construct $\delta R_{K\pi}$ in Equation (17), which we recast here for convenience

$$\delta R_{K\pi} = \left(\frac{\delta \Gamma_{0,K}(L)}{\overline{\Gamma}_{0,K}(L)} - \frac{\delta \Gamma_{0,\pi}(L)}{\overline{\Gamma}_{0,\pi}(L)} \right) - 2\frac{\alpha}{4\pi} \left(Y_K^{(2)}(L) - Y_\pi^{(2)}(L) \right) + \frac{\alpha}{4\pi} \left(\delta \Gamma_{1,K}(\Delta E) - \delta \Gamma_{1,\pi}(\Delta E) \right). \tag{30}$$

The output of the combined fits discussed in the previous subsections, properly tuned to the isosymmetric point using Equation (18) and the appropriate bare quark mass shifts, $\{\Delta m_q\}$, will give the contribution of the first parentheses on the RHS of the above equation. Combining with the analytic results from [7] and [3], corresponding to the second and third parentheses, respectively, we obtain $\delta R_{K\pi}$.

7. Conclusion & Outlook

A high precision test of the unitarity of the CKM matrix is made possible, in part, with recent improvements in lattice simulations. Through the inclusion of IB effects, we are now in the position to predict light CKM matrix elements at percent-level precision or better. In the RBC&UKQCD collaboration, we have a lattice setup that allows us to extract the amplitude correction, $\delta R_{K\pi}$, from a near-physical point simulation. At the time of writing, we are estimating the systematics on the prediction of $\frac{|V_{us}|}{|V_{ud}|}$. Indeed, we expect to provide an update on the ratio of CKM matrix elements shortly after the publication of this proceeding.

Our immediate plan for the leptonic decay sector is as follows: progress is under way to renormalise the weak operator, O_W. This will enable us to obtain $|V_{ud}|$ and $|V_{us}|$ and, in turn, provide an update to the unitarity tension seen in Equation (1). Further improvements are possible with an unquenched calculation, *i.e.* including disconnected contributions and QED interactions with sea quarks. Additionally, we envision in the near future a departure from the point-like approximation to a first principle lattice calculation of the real photon contribution. In the long term, it is our objective to further constrain CKM matrix elements by studying semi-leptonic decays, *e.g.* $K^{\pm} \to \pi^0 \ell^{\pm} \nu_\ell$.

Acknowledgements

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