# The Host Galaxies of Rapidly Evolving Transients in the Dark Energy Survey

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#### ABSTRACT

Rapidly evolving transients (RETs), also termed fast blue optical transients (FBOTs), are a distinct class of astrophysical event. They are characterised by lightcurves that decline much faster than common classes supernovae (SNe), span vast ranges in peak luminosity and can be seen to redshifts greater than 1. Their evolution on fast timescales has hindered high quality follow-up observations, and thus their origin and explosion/emission mechanism remains unexplained. In this paper we define the largest sample of RETs to date, comprising 106 objects from the Dark Energy Survey, and perform the most comprehensive analysis of RET host galaxies. Using deep-stacked photometry and emission-lines from OzDES spectroscopy, we derive stellar masses and star-formation rates (SFRs) for 49 host galaxies, and metallicities for 37. We find that RETs explode exclusively in star-forming galaxies and are thus likely associated with massive stars. Comparing RET hosts to samples of host galaxies of other explosive transients as well as field galaxies, we find that RETs prefer galaxies with high specific SFRs, indicating a link to young stellar populations, similar to stripped-envelope SNe. RET hosts appear to show a lack of chemical enrichment, their metallicities akin to long duration gamma-ray bursts and superluminous SN host galaxies. There are no clear relationships between properties of the host galaxies and the peak magnitudes or decline rates of the transients themselves.

**Key words:** transients: supernovae – galaxies: star formation – galaxies: abundances – galaxies: photometry

#### 1 INTRODUCTION

In the standard paradigm of stellar evolution, stars with a zero-age main sequence (ZAMS) mass above  $8M_{\odot}$  are be-

lieved to explode as a result of a catastrophic collapse of their iron cores and are known as core-collapse supernovae (CC-SNe). CCSNe can be split into observationally-determined subclasses based on their lightcurve and spectral evolution:

SNe II display hydrogen features in their spectra (Minkowski 1941), and are thought to occur in stars that retain a large fraction of their hydrogen envelope. Conversely, SNe Ib and Ic do not show signatures of hydrogen (e.g. Filippenko 2002) and are thus referred to collectively as stripped-envelope SNe (SESNe), having undergone a partial removal of their outer atmospheres. The SN IIb subclass, which shows hydrogen only at early epochs that disappears after a few weeks (Filippenko 1988), is also commonly grouped along with SESNe. SNe IIn display much narrower hydrogen emission lines when compared to standard SNe II (Schlegel 1990). The narrow emission originates from the ejecta impacting on slow-moving circumstellar material (CSM). Since the turn of the century, observations of CCSNe, whose lightcurves are primarily powered by the radioactive decay of freshly synthesised Ni-56, have been supplemented by rarer, more exotic transient classes.

Long duration gamma-ray bursts (LGRBs), although first discovered in the 1960s (Klebesadel et al. 1973), were unequivocally linked to collapsing massive stars through their associations with broad-lined type Ic SNe (Galama et al. 1998; Hjorth et al. 2003). Thought to be caused by accretion onto a newly-formed black hole at the centre of a collapsing, rapidly-rotating massive star (e.g. Woosley 1993; Woosley & Bloom 2006; Woosley & Heger 2006), LGRBs comprise roughly 1% of all SNe Ic, which themselves make up only 15% of all CCSNe (Kelly & Kirshner 2012; Graham & Schady 2016). Another exotic class of SNe is the particularly bright superluminous supernovae (SLSNe; e.g. Quimby et al. 2011; Gal-Yam 2012). Originally grouped due to their slowly-evolving lightcurves and extreme luminosity (peaking at  $M_B < -21$  mag; 10-100 times brighter than typical CCSNe), recent observations have revealed a continuum of spectroscopically similar objects with peaks as faint as  $M_B \sim -19$  mag (De Cia et al. 2018; Lunnan et al. 2018; Angus et al. 2019), which overlaps with the bright end of the CCSN luminosity function (Li et al. 2011). The lightcurve evolution of SLSNe is not well described by models of Ni-56 decay, with the most popular alternative hypothesis being the magnetic coupling of the ejecta with the spin down of a newly formed, rapidly rotating magnetar.

Along with observations of the transients themselves, studies of host galaxies are frequently used to make strong inferences about the progenitor stars and explosion mechanisms. CCSNe are confined almost exclusively to galaxies hosting recent or ongoing star formation, due to their origin from massive stars. There are correlations between the expected progenitor mass of different sub-classes of CCSNe and host galaxy properties. On average, SESNe reside in galaxies with higher specific star-formation rates (sSFRs; James & Anderson 2006; Kelly et al. 2008), while studies of the local environments tend to show that SESNe explode closer to [HII] regions than SNe II, indicating that the progenitors are younger and more massive than the various subclasses of hydrogen-rich SNe II (e.g. Anderson et al. 2012; Galbany et al. 2018). More extreme events tend to occur in galaxies low in mass and high in sSFR, with both LGRBs (e.g. Fruchter et al. 2006; Le Floc'h et al. 2006; Levesque et al. 2010; Krühler et al. 2015; Vergani et al. 2015; Perley et al. 2016a; Palmerio et al. 2019; Taggart & Perley 2019) and to an even greater degree SLSNe (e.g. Neill et al. 2011; Lunnan et al. 2014; Leloudas et al. 2015; Angus et al. 2016;

Schulze et al. 2018; Taggart & Perley 2019) exhibiting this association.

The chemical composition of the interstellar medium (ISM) is an important consideration when comparing host galaxy properties. While it does not appear to play a significant role in the relative production of CCSNe (although there are some trends, with SESNe typically found in slightly less metal-rich environments than SNe II; Galbany et al. 2018), it appears to be vitally important in the production of LGRBs and SLSNe. Theory predicts that the production of a LGRB should only be possible in stars with a metallicity of  $Z/Z_{\odot} \leq 0.3$  (Woosley 1993) in order for the likely Wolf-Rayet or blue supergiant progenitors not to lose their outer atmospheres through metal-driven winds, thus conserving sufficient angular momentum to power the black-hole-driven jet or rapidly rotating magnetar. Many LGRB host galaxy studies have indeed revealed a metallicity threshold between 0.5 and 1 times the solar value (e.g. Stanek et al. 2006; Modjaz et al. 2008; Krühler et al. 2015; Perley et al. 2016a; Japelj et al. 2016; Vergani et al. 2017), although this is not a rigid threshold. SLSN host galaxies also appear to be lower in metallicity than would be expected for their stellar mass, with a suppression of SLSN production above a value of about half-solar (Lunnan et al. 2014; Chen et al. 2016; Perlev et al. 2016b). Like LGRBs, SLSNe also require a particularly high sSFR, suggesting that they are explosions of very young, rapidly rotating massive stars.

Recently, inspection of high-cadence, large-area survey data sets have revealed more exotic transients that are difficult to explain with conventional models. Drout et al. (2014) presented a sample of rapidly evolving transients (RETs; also termed 'Fast Blue Optical Transients' - FBOTs or 'Fast Evolving Luminous Transients' - FELTs) in the Pan-STARRS survey (PS1). Arcavi et al. (2016) reported three rapidly rising, highly luminous objects in the Supernova Legacy Survey (SNLS) and one in the Palomar Transient Factory (PTF), of which one (SNLS04D4ec) declines rapidly like the PS1 sample. Pursiainen et al. (2018, hereafter P18) expanded the known number of RETs to beyond 80 with their sample from the Dark Energy Survey (DES), spanning a redshift range of  $\sim 0$  to > 1. A further sample of five objects has been discovered by the Hyper Suprime-Cam Subaru Strategic Program (SSP) Transient Survey (Tampo et al. 2020). RETs typically rise to peak brightness in less than 10 days, and decline to 10% of their peak brightness within 30 days, much faster than typical SNe. The photometric measurements of the PS1 and DES RETs seem to be well described by expanding blackbodies, although a handful show declining photospheric radii from the first detection. Due to the rapid nature of their lightcurves and location at high-redshift, spectral coverage is sparse and signal-to-noise ratio (SNR) is low, such that there has not yet been a conclusive detection of absorption or emission features from the transients and thus the physical mechanism responsible for their rapid evolution remains unexplained.

There are a limited number of events detected in the local Universe whose properties appear consistent with the RETs seen at cosmological distances in the PS1 and DES samples, the most widely studied of which is AT2018cow (e.g. Prentice et al. 2018; Margutti et al. 2019; Perley et al. 2019). The transient declined from its discovery, with constraints on the rise time of 1 day, and from X-rays through

to radio wavelengths did not resemble any known SN, GRB afterglow, or kilonova (KN; Ho et al. 2019). There are many diverse explanations for the power source of AT2018cow touted in the literature, including: magnetars (Mohan et al. 2020); electron capture collapse of merged white dwarfs (Lyutikov & Toonen 2019); a tidal disruption event (TDE) of a white dwarf (Kuin et al. 2019) or of a main sequence star by an intermediate mass black hole (Perley et al. 2019); common envelope jets supernova (CEJSN; Soker et al. 2019); or a wind-driven transient (Uno & Maeda 2020). Other nearby rapid transients include the local fast-declining SN-like transient SN2018kzr (McBrien et al. 2019) which is explained by the accretion-induced collapse of a white dwarf or a white dwarf-neutron star merger, and KSN-2015K (Rest et al. 2018) whose fast rise and decline is explained by the shock of an SN running into previously-expelled material. It is currently unclear whether these transients represent the local analogues of the DES and PS1 RETs.

In this paper, we present the first comprehensive study of the host galaxies of RETs. We make use of the final DES sample, which builds on P18 by adding the 5th and final season of DES-SN observations as well as more refined selection techniques. Using the deep DES photometry from Wiseman et al. (2020, hereafter W20) and spectra from OzDES (Lidman et al. 2020) we derive host galaxy properties and compare them to samples of host galaxies of CCSNe, LGRBs, and SLSNe, as well as the individual local rapid transients. For clarity, we will use the term RET to refer only to events in the high-redshift samples of DES and PS1.

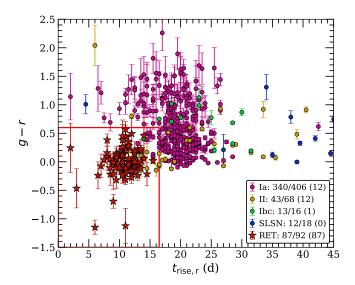
The order of the paper is as follows: in Section 2 we introduce the full DES RET sample and describe the host galaxy observations in Section 3. The analysis methods and results are described in Sections 4 and 5 respectively, before a discussion (Section 6.3) and conclusion (Section 7). Where applicable, we adopt a spatially flat  $\Lambda$ CDM cosmology with the parameters  $H_0 = 70~{\rm km~s^{-1}~Mpc^{-1}}$  and  $\Omega_{\rm M} = 0.3$ . Magnitudes are presented in the AB system (Oke & Gunn 1983), and values (uncertainties) are quoted at the  $50^{\rm th}$  ( $16^{\rm th}$  and  $84^{\rm th}$ , i.e.  $1\sigma$ ) percentiles of their probability density function.

### 2 SAMPLE SELECTION

We derive our sample from the 106 RETs discovered in the 5-year DES-SN transient survey. This number expands upon the 72 of P18. The first reason for the increased sample size is the use of the 5th year of DES-SN, as P18 were only able to make use of the first four years. By imposing the P18 selection criteria on season 5, the sample is increased to 92 objects. The second reason is an update to the sample selection technique, outlined in the following subsection, which adds a further 14 transients.

# 2.1 The DES supernova programme

The Dark Energy Survey (DES) makes use of the Dark Energy Camera (DECam; Flaugher et al. 2015) to survey 5000 deg<sup>2</sup> of the southern sky between 2013 and 2019. The supernova programme (DES-SN), designed primarily to detect type Ia SNe (SNe Ia) for cosmological measurements (Bernstein et al. 2012), consisted of five 6-month seasons of



**Figure 1.** Observer-frame g-r colour at maximum light and observer-frame r-band rise-time, derived from Gaussian Processed fits to DES-SN photometry for spectroscopically confirmed SNe and P18 RETs. The location of the red box is designed to maximise the completeness and purity of RETs. The fractions in the legend refer to the number of each class of transient passing the cuts of Section 2.2 compared to the total number of that class in DES-SN, while the number in parentheses refers to the number inside the red box defining RET parameter space.

approximately 7 day cadence in ten single pointings, known as the SN fields which cover a total area of  $\sim 27\,{\rm deg^2}.$  Of these fields, two were observed with longer exposure times and are referred to as deep fields with remaining eight known as shallow fields; the resulting approximate single-visit depths are 23.5 mag (shallow) and 24.5 mag (deep) in the DES r-band. Transient detection was performed by difference imaging using a custom pipeline (Kessler et al. 2015). Transient candidates were vetted via machine-learning techniques (Goldstein et al. 2015), leading to  $\sim 30,000$  viable supernova candidates over the full five seasons. DES-SN included an extensive spectroscopic follow-up program to identify transients and measure spectroscopic redshifts of host galaxies (D'Andrea et al. 2018). A full description of the search for rapid transients in DES can be found in P18.

#### 2.2 Improvements to the search method

The original method of finding RETs in the DES-SN data (and presented in P18) was designed to be simple and used light curve modelling with Gaussian and linear fits. The simplistic method made it possible to look for exotic transients without knowing their observed characteristics beforehand and resulted in a large sample of photometrically selected fast transients. However, as the search was simplistic and relied heavily on visual inspection of the available data (etc. images, light curves, host galaxy information), it is impossible to quantify the completeness of the sample. For instance, due to the large redshift range within the sample it is entirely possible that distant events with longer rise times could have been missed due to time dilation stretching their lightcurves, while faster events and low redshift may evolve on timescales

# 4 P. S. Wiseman

Table 1. Host galaxy information for the 106 RETs in the DES 5-year sample. A full-length, machine-readable version of this table is available in the online version of the manuscript.

_	RA	Dec	$m_r^{\rm a}$	z	Survey	Exposure time b
Transient name	(J2000	(J2000)	(AB)			(hours)
DES13C1acmt	54.32925	-26.83371	$23.18 \pm 0.04$	c	OzDES	-
DES13C1tgd	54.06436	-27.63867	$20.30 \pm 0.05$	0.19647	OzDES	10.17
DES13C3abtt	52.62108	-28.16151	$21.50 \pm 0.03$	d	-	-
DES13C3asvu	52.83670	-27.36071	$22.46 \pm 0.02$	d	-	-
DES13C3avkj	51.97076	-27.52792	$24.07 \pm 0.04$	c	OzDES	-
DES13C3bcok	53.02711	-28.62476	$18.61 \pm 0.02$	0.34577	LADUMA	
DES13C3nxi	51.96356	-28.35720	$24.99 \pm 0.05$	c	OzDES	-
DES13C3smn	51.97112	-28.08362	$25.30 \pm 0.06$	c	OzDES	-
DES13C3uig	52.94416	-27.58544	$22.17 \pm 0.03$	0.67346	ACES	
DES13E2lpk	10.09911	-43.53903	$20.49 \pm 0.03$	0.47541	OzDES	6.00
DES13S2wxf <sup>e</sup>	41.61268	-0.02634	$21.39 \pm 0.039$	0.56985	OzDES	5.42
DES13X1hav	35.03245	-5.11022	$23.63 \pm 0.07$	0.58236	OzDES	4.00
DES13X2oyb	35.31056	-5.67893	$22.91 \pm 0.06$	c	OzDES	-
DES13X2wvv <sup>e</sup>	34.86526	-6.71603	$21.89 \pm 0.02$	0.47503	OzDES	4.50
DES13X3aakf	35.71205	-4.69915	$25.47 \pm 0.09$	c	OzDES	_
DES13X3afjd	37.00386	-4.58049	$20.75 \pm 0.05$	d	-	-
DES13X3alnb	37.18496	-5.14232	$25.04 \pm 0.06$	c	OzDES	_
DES13X3gmd <sup>e</sup>	36.50409	-4.21949	$22.90 \pm 0.05$	0.78082	OzDES	13.92
DES13X3gms	35.80095	-4.49384	$23.02 \pm 0.06$	0.64792	OzDES	6.58
DES13X3kgm	36.50380	-4.86636	$26.46 \pm 0.18$	d	-	-
DES13X3npb	36.64218	-4.13381	$20.79 \pm 0.05$	0.49542	OzDES	5.17
DES13X3nyg	36.99228	-3.91327	$23.38 \pm 0.06$	0.71205	OzDES	45.50
DES13X3pby	36.33330	-5.31408	$23.82 \pm 0.04$	c	OzDES	-
DES14C1jnd	54.35153	-27.49300	$24.33 \pm 0.08$	c	OzDES	_
DES14C3gzj <sup>e</sup>	52.57168	-28.14019	$26.04 \pm 0.10$	d	-	_
DES14C3tnz <sup>e</sup>	52.86248	-28.51318	$22.26 \pm 0.03$	0.70452	OzDES	5.79
DES14C3tvw	53.32230	-27.90621	$21.31 \pm 0.02$	0.70390	ACES	0.10
DES14E1aqi <sup>e</sup>	8.73970	-43.40182	$25.61 \pm 0.02$	d	-	_
DES14E1aqr DES14E2xsm	9.67624	-43.58736	$23.39 \pm 0.04$	c	OzDES	
DES14S2ang	41.27776	-0.74529	$17.57 \pm 0.03$	0.05211	OzDES	0.75
DES14S2aliq DES14S2plb	41.85667	-1.61811	$17.37 \pm 0.03$ $18.38 \pm 0.03$	0.03211 $0.11531$	OzDES	2.00
DES14S2pli DES14S2pli	41.22837		$20.86 \pm 0.03$		OzDES	3.42
DES14S2pii DES14X1bnh		-1.09830 4.70254		0.35478	OzDES	3.50
DES14X10III DES14X3pkl	33.74902 37.21103	-4.79254 4.80741	$21.95 \pm 0.03$ $22.40 \pm 0.06$	0.82982 $0.29537$	OzDES	5.17
-		-4.80741		0.29551 d	OZDES	0.17
DES14X3pko	36.90824	-3.69642	$24.48 \pm 0.06$		- O-DEC	1.0.01
DES15C2eal	54.06180	-29.23000	$22.99 \pm 0.03$	0.22347 <sub>d</sub>	OzDES	16.91
DES15C3edw	52.55235	-27.71022	$22.69 \pm 0.03$		- O DEG	
DES15C3lpq	52.71204	-28.61319	$23.28 \pm 0.03$	0.61365	OzDES	6.17
DES15C3lzm	52.17460	-28.23186	$20.47 \pm 0.02$	0.32690	ATLAS	
DES15C3mem <sup>e</sup>	52.14216	-29.05851	$20.28 \pm 0.04$	0.61618	PRIMUS	
DES15C3mgq	52.76901	-28.20882	$22.97 \pm 0.03$	0.23031	OzDES	8.17
DES15C3nat	52.88528	-28.72370	$23.39 \pm 0.03$	0.83929	OzDES	18.62
DES15C3opk	51.66147	-28.34737	$23.05 \pm 0.03$	0.56984	OzDES	12.17
DES15C3opp	51.73962	-28.11496	$23.25 \pm 0.04$	0.44242	OzDES	16.75
DES15C3pbi	52.23620	-28.00223	$25.10 \pm 0.05$	d	-	-
DES15E2lmq	9.62005	-43.98734	$25.65 \pm 0.18$	d	-	-
DES15E2nqh	9.73174	-43.08707	$23.29 \pm 0.03$	0.51525	OzDES	6.83
DES15S1fli	43.18790	-0.88620	$20.98 \pm 0.02$	0.44739	OzDES	2.33
DES15S1fll	42.78846	-0.19747	$20.95 \pm 0.03$	0.22647	OzDES	0.67
DES15X2ead	36.48913	-6.45118	$20.08 \pm 0.03$	0.23175	OzDES	2.00
DES15X3atd	35.84011	-4.29142	$23.87 \pm 0.05$	c	OzDES	-
DES15X3kyt	36.27500	-5.41103	$24.74 \pm 0.07$	d	-	-

<sup>&</sup>lt;sup>a</sup> Apparent r-band Kron magnitude according to DES-SN deep coadds of W20, not corrected for Galactic foreground reddening.

<sup>&</sup>lt;sup>b</sup> Exposure time only given for spectra which we have used for line measurements rather than just redshift.

 $<sup>^{\</sup>rm c}$  Host targetted by OzDES but no redshift measurement possible.

<sup>&</sup>lt;sup>d</sup> Host not targetted by OzDES.

<sup>&</sup>lt;sup>e</sup> Found by updated search method (Section 2.2).

f Found in Season 5 using P18 method.

g Redshift updated since to P18.

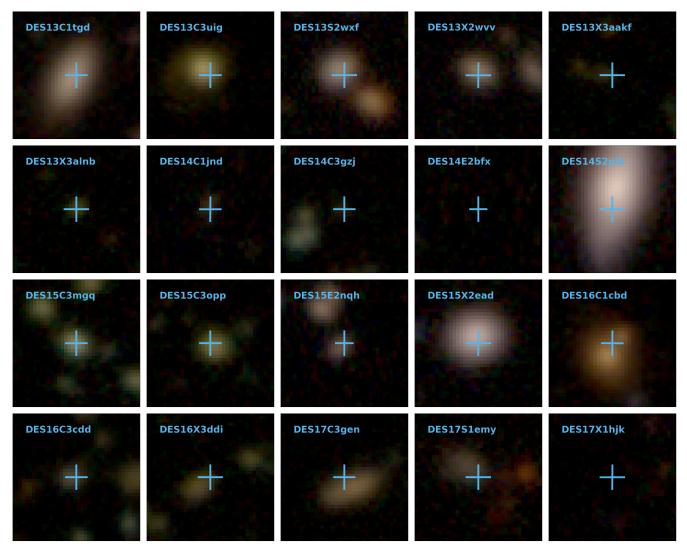


Figure 2. Selection of DES RET host galaxies in an RGB composite of the DES gri band deep coadds from W20. The locations of the transients are indicated with cyan crosses. The stamps have a size of 10" in each direction. DES14E2bfx and DES17X1hjk are considered hostless.

quicker than the survey cadence. Here, a more sophisticated search method is presented. As only a fraction of transients in DES-SN have redshift information from their host galaxies, we perform the search in the observer frame. The key features of RETs that separate them from most traditional SNe types are the fast light curve evolution (rise to peak in  $\lesssim$  15 days) and blue colour at peak  $(g-r\lesssim0.5)$ . Even though both of these quantities depend on the redshift of the transient, they still effectively distinguish the the fast events from traditional SN types. We thus attempt to select a sample based on observed rise times and colours at peak brightness.

#### 2.2.1 Gaussian processed lightcurves

Using observed photometric data points directly to infer rise time and colour has several problems that can be improved. For one, measuring peak colour is problematic: DES-SN did not always observe g- and r-bands on the same or even consecutive nights in the 'deep' fields, thus adding larger un-

certainty in the peak colour estimate. We do not have a light curve model for RETs, and therefore measuring a 10-15 day rise time is difficult with a one week cadence. Rather than fitting with a physically motivated model, we instead interpolate the lightcurves of all DES-SN candidates using Gaussian Processes (GP) as presented in Pursiainen et al. (2020). The interpolated lightcurves have a 0.5 day cadence and every epoch and band has a flux value and an associated uncertainty.

#### 2.2.2 Photometric definition of RETs

To make an improved selection of RETs we use a parameter space described by observed g-r colour and the time taken to rise from non-detection to peak r-band magnitude in the observer frame. Using rise times and colours from GP lightcurves, we populate this parameter space with the sample of 72 RETs from the P18 method, updated with 20 extra objects found using that method in the fifth season of DES. We add spectroscopically confirmed SNe of types Ia, Ibc, II,

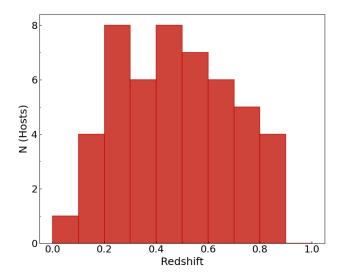


Figure 3. Redshift distribution for the host galaxies of RETs in DES for which a measurement was obtained.

and SLSNe observed by DES in order to verify that they are rejected by the search method. We keep objects passing the following selection requirements (cuts):

- (i) The transient was detected in only one DES-SN observing season.
- (ii) Maximum observed brightness in both g- and r-bands was brighter than 24 mag (in the eight 'shallow' DES-SN fields) or 25 mag (in the two 'deep' fields), as in P18.

SNe Ia and RETs populate two distinct regions of g-r vs.  $t_{\rm rise;r}$  parameter space (Fig. 1), where  $t_{\rm rise;r}$  is the time to rise from non-detection to peak r-band brightness, and the g-r colour is measured at peak brightness. RETs appear bluer and faster than the typical SNe. We define a region in this parameter space which minimises the contamination of non-RETs (purity) while maximising the total fraction of RETs (completeness). The resulting limits are -1.5 < g-r < 0.6 and  $t_{\rm rise;r} < 16.5$ , corresponding to the red box in Fig. 1.

#### 2.2.3 Removal of active galactic nuclei

In order to apply selection criterion i) from Section 2.2.2 it is necessary to distinguish between DES-SN candidates that are truly multi-season events (typically active galactic nuclei; AGN) and those that are single-season events with spurious detections in other seasons. To detect and filter out AGN we use a basic convolutional neural network (CNN) classifier. We train the CNN on spectroscopically confirmed SNe of all types as well as the 92 RETs identified using the P18 method (181 objects) and on spectroscopically typed AGN (182 objects), and use it to separate the sample into two photometric subtypes: AGN-like and SN-like. The classifier returns SNe-like objects with an accuracy of 0.992 on the test set (391 SNe and 79 AGN). No SNe were classified as AGN-like. The remaining AGN classified as SNe-like are removed by manual vetting later in the process. The CNN does not separate the SN-like objects into RETs and other SN subtypes: this is done in the subsequent processing step.

#### 2.2.4 Final DES RET sample

The GP lightcurves of all DES-SN candidates are classified as AGN-like or SNe-like by the CNN, and are then subject to the lightcurve quality cuts i) and ii) (Section 2.2.2), resulting in 2259 objects, of which 939 lie inside the colour and rise-time region which we defined for RETs. These objects are subject to a further set of cuts in order to remove remaining contaminants. We impose a cut based on a SN classifier (PSNID; Sako et al. 2008), that returns a normalised goodness of fit to different SN Ia and CCSN templates, along with a Bayesian probability of it being a SN Ia. To remove highly-probable SNe Ia, we use threshold probabilities of P(Ia) < 0.91 and P(Ia; Bayes) < 0.82 respectively to the above algorithms, which removes 46 objects from the RET parameter space. In order to further remove longer-lived SNe, the decline time to half of the peak brightness must be < 24 days. This removes 347 SNe, resulting in 546 objects remaining inside the parameter space. The final 546 transients have been visually inspected, with the majority rejected for clearly being spurious detections, obvious multi-season variability that was not picked up by the CNN, or showing a longer timescale decline.

Using the above method recovers 87 of the 92 RETs found using the P18 technique, and adds a further 14. The five were not recovered as their GP lightcurves were fainter than the limits given above in either g or r band. We refer to the resulting sample as DES RETs. Of the 106 objects in the sample, 96 have a host galaxy detected in deep host galaxy photometry of W20 when using the Directional Light Radius method (Sullivan et al. 2006) to associate hosts as per W20. Of these, 49 have a host galaxy spectroscopic redshift which we access through an internal release of the Oz<br/>DES Global Redshift Catalog (GRC; v.2020\_01\_04). The full OzDES redshift catalogue will be available alongside the public data release detailed in Lidman et al. (2020). A further three have redshifts obtained from narrow lines observed in spectra of the transients themselves. We do not consider these three objects for the analysis, since we are unable to separate transient and host contributions to the spectra. A selection of the host galaxies is shown in Fig. 2, centred on the location of the transient. The figure showcases the diversity of host galaxy morphologies and colours, while also displaying the limitations of ground-based observations of high-redshift, relatively small galaxies in terms of spatial resolution. Fig. 3 shows the distribution of redshifts amongst the 49 hosts for which such a measurement was possible. The effect that the redshift selection function has on the results is discussed in Section 6.1. The observational properties of the 96 detected hosts are displayed in Table 1. We highlight objects that were not presented in P18. We also highlight a small subset of objects for which the redshift has taken on a new value than that presented in P18 due to further OzDES observations leading to a more accurate determination.

# 2.3 Comparison samples

In order to compare the host galaxies of DES RETs to those discovered in other surveys as well as other types of explosive transient, we draw upon samples in the literature.

#### 2.3.1 RETs

Since the DES sample of RETs is by far the largest discovered to date, there is no other large sample of RETs with which to compare host galaxy properties. Drout et al. (2014) present host galaxies of 10 RETs discovered in the Pan-STARRS survey, with measurements of stellar masses, SFRs, and metallicities. We also compare with the host galaxy of SNLS04D4ec (Arcavi et al. 2016). To this we add the lowredshift transient AT2018cow (nicknamed "The Cow"). The host galaxy of AT2018cow has been studied with photometric measurements (Perley et al. 2019) as well as with an integral field spectrograph (Lyman et al. 2020), with consistent results. For our comparison, we use the galaxy-averaged stellar mass and SFR from Lyman et al. (2020) and the metallicity from the host nucleus as reported by both Morokuma-Matsui et al. (2019) and Lyman et al. (2020) as it best represents the method of obtaining spectra for our RET sample. We further compare to SN2018gep with data from Ho et al. (2019), and ZTF18abvkwla (nicknamed "The Koala") from Ho et al. (2020).

#### 2.3.2 SNe and GRBs

In compiling a set of comparison samples, we aim for the least biased selections possible. This requires surveys to be untargetted (they were not monitoring certain galaxies in order to search for SNe), ideally complete, and also covering a similar redshift range to the RETs. While in practice the second and third of these criteria are difficult to achieve, particularly with the fainter CCSNe, we are able to choose comparison samples from untargetted surveys to mitigate initial selection biases.

To compare with CCSNe, we draw on the untargetted sample of 47 SNe II from the Palomar Transient Factory (PTF; Stoll et al. 2013), which is likely complete in terms of hosts (all SNe have an associated host). While this sample lies at much lower redshift than the DES RETs (a maximum of 0.18 and mean of 0.05), redshift evolution is easier to account for in a less biased way than correcting for unknown incompleteness. We add to this the compilation of 56 untargeted SESNe from Sanders et al. (2012), with a maximum redshift of 0.26 and a mean of 0.05. Since Sanders et al. (2012) do not report host galaxy magnitudes, stellar masses or SFRs, we cross-match the SN positions with the the Sloan Digital Sky Survey (SDSS; York et al. 2000) Data Release 16 (DR16; Ahumada et al. 2019) and perform our own SED fit using the method outlined in Section 4.1. We are able to do this SED fit for 38 objects, with the others lying outside of the SDSS footprint.

We use the sample of GRB host galaxies of Krühler et al. (2015), using only galaxies with z < 1 in order to maintain completeness, resulting in a sample of 29 hosts with a mean redshift of 0.66. To investigate similarities with SLSNe, we use the host galaxy sample from PTF presented in Perley et al. (2016b) with a mean redshift of 0.24 and a maximum of 0.50.

The host galaxy properties of the above samples are not all derived using the same methods. In terms of SED fitting, the largest systematic offsets in derived properties are due to differences in the assumed initial mass function (IMF). Stoll et al. (2013) and Drout et al. (2014) assume a

Salpeter (1955) IMF whereas all other samples considered (including those calculated in this work in Section 4.1) are determined assuming a Chabrier (2003) IMF. Stellar masses and star-formation rates derived using a Salpeter IMF are roughly  $\sim 1.72$  times higher than those using a Chabrier IMF (Speagle et al. 2014), and we convert the Salpeter-derived values by this factor in order to compare them.

#### 2.4 Field Galaxies

To show how RETs compare to the galaxy population as a whole, we use a sample of  $\sim 800,000$  measurements from the MPA-JHU catalogues of stellar masses (based on the methods of Kauffmann et al. 2003; Salim et al. 2007), SFRs (based off Brinchmann et al. 2004), and metallicities (based off Tremonti et al. 2004) from the catalogues of SDSS Data Relase 7 (DR7; Abazajian et al. 2009). The mean redshift is 0.08 which is much lower than for the RETs, such that significant evolution in the galaxy population has happened between the majority of RET hosts and the SDSS sample. We do not correct for this, but take it into account when analysing our findings.

#### 3 HOST GALAXY OBSERVATIONS

#### 3.1 Photometry

The host galaxy photometry for the sample of RETs is taken from the catalogue of W20, which is based upon deep coadds reaching r-band limiting magnitudes of 26.5. The coadds were created using data from all five seasons of DES-SN, but by excluding one season at a time in order for that coadd not to include contamination from the transients in that season. For this sample, the limiting magnitude for obtaining a spectroscopic redshift (Section 3.2) is  $\sim 24.5$ , meaning that all hosts in the sample are detected with a high S/N.

# 3.2 Spectroscopy

Accurate redshifts for DES-SN were obtained by OzDES, a dedicated DES spectroscopic follow-up campaign based at the 3.9 m Anglo-Australian Telescope (AAT) using the AAOmega fibre-fed spectrograph and 2dF fibre positioner. The observation strategy of OzDES was to point at one of the ten DES-SN fields, and place fibres at the positions of transient hosts, continually coadding the spectra of a particular host until a redshift was obtained at which point the fibre could be allocated to a different transient. The spectra have a resolution of 1400-1700 and a wavelength range of 3700 – 8800 Å, and are reduced using the OzDES pipeline which makes use of a modified version of v6.46 of the 2dfdr (Croom et al. 2004) along with internal scripts. We use internal data release 7, a preliminary version of the public data release which is detailed in Lidman et al. (2020). Extensive description and discussion of OzDES can be found in Yuan et al. (2015); Childress et al. (2017); Lidman et al. (2020). We also obtained redshifts for some transient hosts serendipitously as part of the Looking at the Distant Universe with the MeerKAT Array (LADUMA) survey<sup>1</sup>, three of which

<sup>1</sup> http://www.laduma.uct.ac.za

are present in our sample. Objects for which the host already had a publicly available redshift were not observed with OzDES, but merged into the GRC nontheless. Surveys fulfilling this criteria include the Australia Telescope Large Area Survey (ATLAS; citealtMao2012, the Arizona CDFS Environment Survey (ACES; Cooper et al. 2012), and the PRIsm MUlti-object Survey (PRIMUS; Coil et al. 2011; Cool et al. 2013). Where the spectra from which those redshifts were derived are also public they are included in this analysis. These comprise the Galaxy and Mass Assembly survey (GAMA Driver et al. 2009; Baldry et al. 2018) and SDSS. In total we analyse 45 spectra, with a mean continuum signal-to-noise ratio (SNR) of 2.56/pixel. We stress that the emission lines are detected with a higher SNR than this.

#### 4 ESTIMATING HOST GALAXY PROPERTIES

#### 4.1 Photometric stellar parameters

To estimate the physical properties of the DES RET host galaxies, we generate synthetic photometry in the DES griz bands by combining the individual SEDs of simple stellar population models. We simulate a suite of synthetic galaxy star-formation histories from which we synthesise model SEDs using stellar population models from Bruzual & Charlot (2003) and a Chabrier (2003) initial mass function (IMF). The suite of models is drawn from the same distribution of parameters as used in Kauffmann et al. (2003) and similar papers (e.g. Gallazzi et al. 2005; Gallazzi & Bell 2009) and closely follows the method of Childress et al. (2013). From the synthetic SEDs we derive model magnitudes in the DES griz bands and compare them to the observed Kron magnitudes from the deep photometric catalogue of W20. For each set of model and observed magnitudes we calculate a  $\chi^2$ value, and from these estimate a probability density function (PDF) for key model parameters (mass-to-light ratio M/L, from which we derive  $M_*$ , and specific star-formation rate sSFR). To estimate uncertainties, we take the values at the 16th and 84th percentiles of the resulting PDF to be our  $1\sigma$ lower and upper bounds. The results are presented in Table 2. These parameters are estimated using global photometry of the entire galaxies, and thus represent the overall stellar population.

#### 4.2 Spectroscopic gas-phase parameters

To estimate parameters from the OzDES host galaxy spectra requires several processing steps. We first apply a flux calibration by 'mangling' the spectrum such that the integrated flux over the wavelength ranges of the DES photometric bands matches that measured in the photometry (further details on the mangling process are provided in Swann 2020). We use a circular aperture of diameter 2", matching the size of the spectrograph fibres. The resulting spectrum is a more accurate representation of the true spectrum at that point in the galaxy. This only holds, however, for the area covered by the fibre and we note that the resulting spectrum is not necessarily representative of the galaxy as a whole. There are several reasons this measurement may

differ from one made at the SN explosion site, such as metallicity and age gradients or structure such as bars and discs (see e.g. Iglesias-Páramo et al. 2013, 2016). We proceed with our analysis with this caveat acknowledged.

In order to subtract the stellar component of the host galaxy spectra, we use the Penalized PiXel-Fitting software (pPXF; Cappellari & Emsellem 2004; Cappellari & Michele 2012; Cappellari 2017), using the MILES library of single stellar populations (Vazdekis et al. 2010). By subtracting the best-fitting composite stellar spectrum from the pPXF fit, we are left with a 'gas' spectrum, comprising the emission lines. An example of this procedure is shown in Fig. 4. We fit the emission lines with Gaussian profiles. In order to estimate the uncertainty on the emission line fluxes, we fit  $10^4$  realisations of the line, each time adding perturbations to the line by drawing from a Gaussian distribution based on the variance spectrum. We take the mean and standard deviation of the resulting fits as our flux and its uncertainty, respectively. Line fluxes are presented in Table A1.

# 4.3 Estimating metallicities

The most common method used to estimate the metallicity of galaxies is to use emission line ratios that have been calibrated using theoretical or empirical models in order to approximate the gas-phase oxygen abundance in the interstellar medium. Emission lines originate from regions of ionised gas, but there are a number of possible causes of this ionisation. Using the Baldwin-Phillips-Terlevich diagram (Fig. 5; Baldwin et al. 1981), we demonstrate that the emission line ratios measured in RET hosts are consistent with ionisation caused by star-formation as opposed to AGN. Only 12 of the 45 RET host spectra have the necessary lines to plot an [NII]BPT diagram. This is also the case for the [SII] and [OI] versions of the diagram, and we find no evidence of AGN amongst the 12 hosts in those diagrams either.

Due to the low S/N of the spectra in this sample, we are constrained to a subset of metallicity diagnostics by the availability of only a handful of the strongest emission lines, namely H $\alpha$ , H $\beta$ , [OII]3727, [OIII]4959/5007, [NII]6548/6583, and [SII]6717/6731. Furthermore, for each host galaxy only a subset of these lines is detected - for example, H $\alpha$ , [NII] and [SII] are redshifted out of the spectral coverage at z>0.3, leaving only the oxygen and H $\beta$  lines available and thus the R23 diagnostic, which is based off the relative strengths of the [OII]and [OIII]lines. For hosts at z<0.3 we are able to use the [OIII]/[NII] (O3N2), [NII]/H $\alpha$  (N2), and [SII]/[NII] (S2N2) line ratios.

Due to the redshift range of our sample, and the limited wavelength coverage of the spectra (3000-8000 Å), we are unable to use a single line ratio to estimate the oxygen abundances. We thus determine a set of indicators for which to calculate abundances. For the O3N2 and N2 indicators we use the calibration of Pettini & Pagel (2004) (PP04), and if [SII] is detected we derive an abundance using the S2N2 diagnostic of Dopita et al. (2016) (D16). For the R23 indicator, we use the calibration of Kobulnicky & Kewley (2004) (KK04). At abundances around  $12 + \log (O/H) \sim 8.4$ , the R23 indicator becomes two-tailed, with a low and a high value of metallicity corresponding to a single R23 ratio. In cases where the lines are available, we break this degeneracy by cross-calibrating with the [NII]/[OII] ratio (Kewley

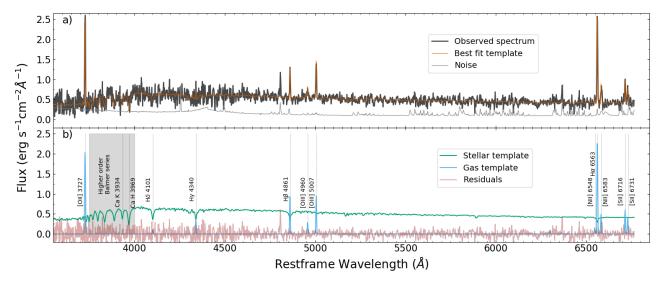
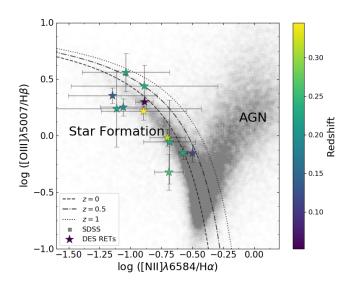


Figure 4. The spectrum of DES16C2ggt, decomposed into its constituent components according to the pPXF fit. a) the best fit superimposed on the observed spectrum. b) the constituent parts of the decomposition: the stellar template including absorption features, and the nebular gas emission. The grey area shows the location of higher-order ( $H\epsilon$ onwards) Balmer lines.



**Figure 5.** Baldwin-Phillips-Terlovich (BPT) diagram for RET hosts, showing that the emission lines are consistent with being generated by star formation rather than AGN activity. The three curves show the delimitation between star formation and AGN at z = 0, 0.5, 1 according to Kewley et al. (2013). The redshift of each RET is indicated by the colourbar on the right.

& Ellison 2008). In the cases where [NII] is not available, and there are no other diagnostics that can be used to inform the choice of branch, we use the host galaxy stellar mass to derive a crude metallicity estimate from the mass-metallicity relation (MZR) of Kewley & Ellison (2008) based upon the PP04 O3N2 diagnostic. For 12+log (O/H)<sub>MZR</sub> < 8.4 we chose the lower branch, while for higher MZR metallicities we choose the upper branch. We note that this is a rough estimation. If we leave the branch choice for those with no [NII] to be random, we find that the results are consistent to within uncertainties.

The samples to which we compare metallicities span different redshift ranges, were observed with different equipment, and in many cases were compiled before certain (particularly the D16) diagnostics were devised. Therefore, in order to compare oxygen abundances between different samples we transform all abundances onto the PP04 O3N2 scale using the conversion factors given in Kewley & Ellison (2008). This is not possible for the D16 diagnostic, so we discard it from the rest of our analysis, although for completeness we provide it for DES RET hosts where available. For samples that quoted multiple diagnostics, or for which sufficient line flux measurements were provided from which to calculate multiple diagnostics, we transform them all to the PP04 O3N2 scale. Following the prescription of Krühler et al. (2015) we simultaneously minimise the oxygen abundance against the PDFs of the various different diagnostics scaled to PP04 O3N2, resulting in a final 'best' PDF. We take  $1\sigma$  uncertainties from the  $16^{\rm th}$  and  $84^{\rm th}$  percentiles of these PDFs, and for DES RETs display the results in Table

#### 5 ANALYSIS

A summary of the masses, sSFRs and metallicities for the hosts of DES RETs and various comparison samples is presented in Table 3. In the following sections, we compare the host galaxy properties of RETs to each sample in detail.

#### 5.1 Star formation rate

Figs. 6 and 7, split in two for clarity, show the 'star formation main sequence' (SFMS) of RET host galaxies, as determined from photometric SED fitting along with that for the comparison samples and for the field galaxies of SDSS. RETs follow CCSNe, LGRBs and SLSNe in avoiding passive galaxies, evidence that RETs require the presence of star-formation and thus are linked to massive stars. One object

Table 2. Host galaxy properties for the 49 DES RET host galaxies with redshifts and host galaxy spectra. The table is available in the online version in a machine readable format.

Transient Name	$\log{(M_*)}$	log (SFR)	log (sSFR)	$12 + \log{(O/H)}$						
	$(M_{\odot})$	$\left( M_{\odot} \mathrm{yr}^{-1} \right)$	$\left( yr^{-1}\right)$	$\mathrm{Best^a}$	D16	$PP04~\mathrm{N2}$	PP04 O3N2	KK04 R23	Average O3N2 <sup>b</sup>	
DES13C1tgd	10.240.08	$0.10^{0.38}_{0.47}$	$-10.14_{0.39}^{0.30}$	8.670.18	$8.62_{0.20}^{0.18}$	$8.72^{0.16}_{0.16}$	-	-	8.80_0.18	
DES13C3bcok	$11.57_{0.11}^{0.11}$	$1.20_{0.50}^{0.49}$	$-10.37_{0.38}^{0.38}$	$8.67_{0.19}^{0.18} \\ 8.99_{0.29}^{0.08}$	- 0.20	- 0.10	_	$8.99_{0.29}^{0.08}$	$8.66_{0.31}^{0.10}$	
DES13C3uig	$11.57_{0.11}^{0.11} \\ 10.28_{0.10}^{0.23}$	$1.20_{0.50}^{0.49} \\ 0.59_{0.20}^{0.40}$	$-9.68_{0.10}^{0.39}$	- 0.29	_	-	_	- 0.29	- 0.31	
DES13E2lpk	$10.66_{0.09}^{0.10}$	$0.59_{0.29}^{0.39}$ $0.46_{0.36}^{0.45}$	$\begin{array}{c} -10.37_{0.38}^{0.38} \\ -9.68_{0.19}^{0.16} \\ -10.20_{0.27}^{0.35} \\ -9.57_{0.04}^{0.04} \end{array}$	$8.89_{0.14}^{0.09}$	-	-	_	$8.89_{0.14}^{0.09}$	$8.54_{0.15}^{0.11}$	
DES13S2wxf	$9.86_{0.03}^{0.06}$		$-9.57_{0.04}^{0.04}$	$7.62_{0.70}^{0.62}$	-	-	_	$7.62_{0.70}^{0.62}$ $8.62_{0.46}^{0.24}$	- 0.13	
DES13X1hav	0 160.38	$0.29_{0.07}^{0.10}$ $-0.44_{0.37}^{0.93}$	$-9.60^{0.55}_{0.21}$	$8.62_{0.46}^{0.70}$	-	-	-	$8.62_{0.46}^{0.74}$	$8.28_{0.24}^{0.24}$	
DES13X2wvv	$9.79_{0.14}^{0.16}$	$0.19_{0.43}^{0.37}$	$-9.60_{0.28}^{0.21}$ $-9.59_{0.25}^{0.42}$	$8.62_{0.46}^{0.24} \\ 8.68_{0.27}^{0.15}$	-	-	_	$8.68_{0.27}^{0.46}$	$8.33_{0.17}^{0.16}$	
DES13X3gmd	10 330.44	0 = 10.36	$-9.59_{0.25}^{0.26}$	-	-	-	-	- 0.27	- 0.17	
DES13X3gms	$9.41_{0.16}^{0.27}$	$0.74_{0.75}^{0.80}$ $-0.06_{0.51}^{0.78}$	$-9.39_{0.25}^{0.12}$ $-9.47_{0.35}^{0.52}$	$8.51_{0.96}^{0.50}$	-	-	_	$8.51_{0.96}^{0.50}$	$8.21_{0.48}^{0.48}$	
DES13X3npb	11 020.18	1 1/11.33		$8.72_{0.37}^{0.23}$	_	-	_	$8.72_{0.37}^{0.23} \\ 7.69_{0.72}^{0.61} \\ 8.38_{0.85}^{0.42}$	$8.37_{0.22}^{0.46}$	
DES13X3nyg	0.210.42	$0.05_{0.53}^{0.73}$	~ ~ -0.33	$7.69_{0.72}^{0.37}$	_	-	_	$7.69_{0.72}^{0.31}$	- 0.22	
DES14C3tnz	$9.31_{0.26}^{0.42} \\ 10.21_{0.22}^{0.41}$	$0.78_{0.60}^{0.33}$ $1.02_{0.46}^{0.39}$	0.440.47	$7.69_{0.72}^{0.61} \\ 8.38_{0.85}^{0.42}$	_	_	_	8.380.42	$8.15_{0.29}^{0.29}$	
DES14C3tvw	11 000.09	$1.02_{0.39}^{0.39}$		-	_	_	_	- 0.83	- 0.29	
DES14S2ang		0 240.54		$8.34_{0.15}^{0.02} \\ 8.65_{0.05}^{0.29} \\ 8.95_{0.08}^{0.06}$	$8.16^{0.09}_{0.11}$	$8.34^{0.04}$	$8.35^{0.02}_{0.02}$	$8.36^{0.01}$	$8.29^{0.02}$	
DES14S2plb	$9.41_{0.10}^{0.23}$ $10.36_{0.23}^{0.25}$	$-0.34_{0.43}^{0.33}$ $0.89_{0.64}^{0.77}$	$-9.75_{0.33}^{0.32}$ $-9.47_{0.40}^{0.52}$	$8.65_{0.29}^{0.15}$	$8.65_{0.07}^{0.11}$	$8.34_{0.04}^{0.04} \\ 8.63_{0.05}^{0.05}$	$8.35_{0.02}^{0.02} \\ 8.62_{0.03}^{0.03}$	$8.36_{0.01}^{0.01}$ $8.95_{0.05}^{0.04}$	$8.29_{0.02}^{0.02}$ $8.64_{0.05}^{0.04}$	
DES14S2pli	$10.27_{0.11}^{0.23}$	$0.50_{0.38}^{0.64}$	$-9.47_{0.40}^{0.32}$ $-9.77_{0.27}^{0.27}$	$8.95^{0.05}_{0.06}$	- 0.07	- 0.05	0.03	$8.95_{0.08}^{0.05}$	$8.64_{0.05}^{0.05}$ $8.61_{0.09}^{0.07}$	
DES14X1bnh		1.530.52		- 0.08	_	_	_	- 0.08	- 0.09	
DES14X3pkl	$9.60_{0.34}^{0.25}$	$1.53_{\substack{0.52\\0.95\\0.26_{\substack{0.85\\0.85}}}}^{0.38}$	$-10.18_{0.25}^{0.25}$ $-9.34_{0.51}^{0.46}$	8.380.20	8.220.23	$8.48_{0.16}^{0.15}$	_	_	$8.52_{0.18}^{0.18}$	
DES15C2eal		$-1.29_{0.30}^{0.85}$		$8.38_{0.29}^{0.20}$ $8.30_{0.31}^{0.25}$	$8.22_{0.30}^{0.23} \\ 8.04_{0.51}^{0.32}$	$8.32_{0.19}^{0.16}$	$8.34_{0.20}^{0.15}$	8 390.50	$8.28_{0.16}^{0.18}$	
DES15C3lpq	$8.42_{0.09}^{0.18}$ $9.32_{0.20}^{0.39}$	$0.86_{0.73}^{1.08}$	- 1-0.21	$8.30_{0.58}^{0.31}$	-0.51	0.320.19	0.540.20	$8.39_{0.34}^{0.50} \\ 8.30_{0.58}^{0.33}$	$8.15_{0.14}^{0.16}$	
DES15C3lzm	$10.06_{0.20}^{0.28}$	$1.08_{0.68}^{0.75}$	$-8.45^{0.68}_{0.53}$ $-8.98^{0.47}_{0.48}$ $0.720.27$	0.500.58	_	_	_	0.50	0.150.14	
DES15C3mem	10 760.37	4 0 40 64	$-9.72_{0.31}^{0.48}$	_	_	_	_	_		
DES15C3mgq	8.41 <sup>0.29</sup>	$1.04_{0.37}^{0.04}$ $-1.12_{0.37}^{0.55}$	$-9.72_{0.31}^{0.31}$ $-9.53_{0.25}^{0.26}$	$8.37_{0.27}^{0.29}$	$8.45_{0.50}^{0.48}$	$8.43_{0.23}^{0.24}$	$8.34_{0.17}^{0.11}$	$8.34_{0.34}^{0.39}$	$8.32_{0.14}^{0.20}$	
DES15C3nat	$10.49_{0.17}^{0.12}$	$1.14_{0.60}^{0.37}$	$-9.34_{0.41}^{0.25}$	0.370.27	0.450.50	0.43 <sub>0.23</sub>	0.340.17	0.34	0.520.14	
DES15C3nat DES15C3opk	0.47 <sub>0.17</sub> 0.80 <sup>0.22</sup>	$0.89_{1.20}^{0.67}$	0.010.45	8 570.33		_	_	8 570.33	$8.24_{0.31}^{0.31}$	
DES15C3opp	$9.89_{0.44}^{0.22}$ $9.15_{0.21}^{0.32}$	$0.13_{0.78}^{0.84}$	$-9.01_{0.75}^{0.15}$ $-9.02_{0.57}^{0.52}$	$8.57_{0.80}^{0.33}$ $8.45_{0.59}^{0.32}$		_	_	$8.57_{0.80}^{0.33}$ $8.45_{0.59}^{0.32}$ $8.83_{0.52}^{0.15}$	$8.17_{0.24}^{0.24}$	
DES15E2nqh	0.200.28	0 110 94		8.83 <sup>0.15</sup>	-	-	-	9.43 <sub>0.59</sub> 9.930.15	8.48 <sup>0.17</sup>	
DES15E2liqli DES15S1fli	$10.24_{0.27}^{0.26}$	1.540.67	0.43	9 960.06	-	-	-	$8.86_{0.08}^{0.52}$		
DES15S1fil		$0.11_{0.72}^{0.72}$ $1.54_{1.01}^{0.67}$ $-0.14_{0.64}^{0.54}$	$-8.70_{0.75}^{0.41}$ $-9.34_{0.43}^{0.22}$	$8.86_{0.08}^{0.06}$ $8.26_{0.20}^{0.23}$ $8.45_{0.23}^{0.23}$	o 270.36	0 200.13	$8.22_{0.14}^{0.10}$	0.00034	$8.52_{0.08}^{0.07}$ $8.22_{0.12}^{0.12}$	
DES15X2ead	$9.20_{0.20}^{0.20}$ $9.92_{0.08}^{0.20}$	$-0.14_{0.64}^{0.34}$ $0.22_{0.41}^{0.47}$	0.000.27	0.20 <sub>0.20</sub> 0.45 <sup>0.23</sup>	$8.37_{0.45}^{0.36} \\ 8.22_{0.34}^{0.25}$	$8.28_{0.16}^{0.13} \\ 8.47_{0.16}^{0.15}$	$8.22_{0.14}^{0.14}$ $8.52_{0.15}^{0.14}$	$8.26_{0.22}^{0.32}$ $8.61_{0.77}^{0.38}$	$8.43_{0.09}^{0.24}$	
DES15X2ead DES15X3mxf	0.020.32	$0.22_{0.41} \\ 0.78_{1.18}^{0.90}$	$-9.69_{0.32}^{0.32}$ $-9.15_{0.88}^{0.58}$	0.43 <sub>0.33</sub> 0.730.15	6.22 <sub>0.34</sub>	0.470.16	6.32 <sub>0.15</sub>	0.01 <sub>0.77</sub> 9.720.15	0.43 <sub>0.09</sub> 9.260.16	
	$9.93_{0.30}^{0.32}$ $10.77_{0.15}^{0.22}$		-9.13 <sub>0.88</sub>	$8.72_{0.25}^{0.15}$ $8.79_{0.37}^{0.20}$	-	-	-	$8.72_{0.25}^{0.15}$ $8.79_{0.37}^{0.20}$	$8.36_{0.18}^{0.16}$ $8.44_{0.22}^{0.22}$	
DES16C1cbd	$9.86_{0.30}^{0.15}$	$-1.43_{1.94}^{1.48}$ $0.13_{0.46}^{0.94}$	$-12.20_{1.79}^{1.26}$ $-9.73_{0.18}^{0.59}$	$8.79_{0.37}^{0.37}$ $8.50_{0.20}^{0.33}$	0 200.17	0 450.10	0 510.05	$8.79_{0.37}^{0.37}$ $8.85_{0.07}^{0.06}$	$8.50_{0.09}^{0.08}$	
DES16C2ggt DES16C3axz	$9.80_{0.29}$ $9.85_{0.33}^{0.33}$	0.200.78	0.650.44	$8.50_{0.20}$ $8.59_{0.37}^{0.37}$	8.28 <sup>0.17</sup> 8.520.05	$8.45^{0.10}_{0.11}$	8.51 <sup>0.05</sup> 8.50 <sup>0.02</sup>	8.070.03	$8.61_{0.03}^{0.09}$	
	9.05	$0.20_{0.56}^{0.56}$ $0.45_{0.68}^{1.00}$	$-9.65_{0.37}^{0.47}$ $-9.41_{0.41}^{0.68}$	$8.84_{0.23}^{0.06}$	$8.53_{0.05}^{0.05}$	$8.56_{0.03}^{0.03}$	$8.59_{0.02}^{0.02}$	$8.97_{0.04}^{0.03}$ $8.84_{0.23}^{0.12}$ $8.93_{0.34}^{0.11}$	$8.49_{0.22}^{0.14}$	
DES16C3gin	$9.85_{0.27}^{0.32}$ $9.57_{0.12}^{0.43}$	0.43 <sub>0.68</sub>	$-9.41_{0.41}$ $-9.50_{0.26}^{0.55}$	$8.93_{0.24}^{0.11}$	-	-	-	0.04 <sub>0.23</sub>	$8.49_{0.22}^{0.22}$ $8.59_{0.22}^{0.13}$	
DES16E2pv		$0.07_{0.39}^{0.98} \\ 0.30_{0.86}^{0.79}$	-9.50 <sub>0.26</sub>	$8.93_{0.34}^{0.11}$ $8.25_{0.17}^{0.10}$	0 110.25	9.250.07	0.210.05	8.93 <sub>0.34</sub>	8.59 <sub>0.33</sub>	
DES16S1bbp	$9.06_{0.39}^{0.30}$	$-0.49_{0.71}^{1.70}$	$-9.30_{0.26}$ $-8.75_{0.48}^{0.49}$ $-9.09_{0.43}^{1.29}$	8.25	$8.11_{0.33}^{0.25} \\ 8.33_{0.45}^{0.38}$	$8.25_{0.09}^{0.07}$	$8.31_{0.08}^{0.05} \\ 8.27_{0.13}^{0.07}$	$8.24_{0.12}^{0.13} \\ 8.27_{0.10}^{0.10}$	8.24 <sup>0.04</sup> 9.220.07	
DES16S1dxu	$8.60_{0.28}^{0.41}$		$-9.09_{0.43}^{1.23}$ $-10.69_{0.26}^{0.91}$	$8.26_{0.15}^{0.12}$	8.330.45	$8.24_{0.14}^{0.11}$	8.270.13	8.270.10	$8.22_{0.12}^{0.07}$	
DES16X1eho	10.700.42	$0.01_{0.50}^{1.33}$	-10.69 <sub>0.26</sub>	- 700.16	-	-	-	- 700.16	- 420.18	
DES16X3cxn	9.330.08	$-0.32_{0.19}^{1.00}$	$-9.88_{0.11}^{0.26}$	$8.79_{0.36}^{0.16}$ $8.58_{0.18}^{0.18}$	- 0.440.17	- 470.08	- 0.610.07	$8.79_{0.36}^{0.16}$	$8.43_{0.27}^{0.18}$	
DES16X3ega	$9.94_{0.06}^{0.18}$	$0.30_{0.21}^{0.26}$	$-9.64_{0.15}^{0.08}$	8.580.18	$8.44_{0.20}^{0.17}$	$8.47^{0.08}_{0.09}$	$8.61_{0.06}^{0.07}$	8.770.06	8.51 <sup>0.08</sup> 0.08	
DES16X3erw	$9.89_{0.25}^{0.26} \\ 9.56_{0.21}^{0.16} \\ 10.07_{0.14}^{0.23}$	$0.94_{0.75}^{0.62}$ $-0.02_{0.44}^{0.33}$ $0.84_{0.78}^{0.50}$	$-8.94_{0.51}^{0.36}$ $-9.58_{0.24}^{0.17}$ $-9.22_{0.64}^{0.27}$	$8.84_{0.13}^{0.09}$ $8.78_{0.31}^{0.17}$ $8.34_{0.14}^{0.12}$	-	-	-	8.84 <sup>0.09</sup> <sub>0.13</sub> 8.78 <sup>0.17</sup> <sub>0.31</sub> 8.29 <sup>0.11</sup> <sub>0.100</sub>	$8.49_{0.13}^{0.10} \\ 8.43_{0.25}^{0.19}$	
DES17C2hno	$9.56_{0.21}^{0.10}$	$-0.02_{0.44}^{0.53}$	-9.58 <sub>0.24</sub>	8. /80.17	-	- 0.050.18	- 200.08	8./80.1/	8.430.15	
DES17C3fwd	10.070.23	0.840.30	$-9.22_{0.64}^{0.27}$	8.340.12	-	$8.35_{0.17}^{0.18}$	$8.38_{0.12}^{0.08}$	8.290.11	8.300.09	
DES17C3gop	$9.90^{0.38}_{0.09}$	$0.19_{0.38}^{0.86} \\ 0.28_{1.91}^{0.70}$	-9./10.20	$8.79_{0.48}^{0.21}$ $8.75_{0.47}^{0.52}$	0.40	- 0.50 47	-	$8.79_{0.48}^{0.21}$	$8.43_{0.28}^{0.24}$	
DES17S2fee	$11.20_{0.14}^{0.22}$	$0.28_{1.91}^{0.70}$	$-10.92_{1.77}^{0.28}$	8.750.32	$8.55^{0.42}_{0.48}$	$8.95^{0.47}_{0.45}$	-		0 28	
DES17X3cds	$9.36_{0.12}^{0.31}$	$-0.18^{0.83}_{0.25}$	$-9.53_{0.13}^{0.52}$	$8.69_{0.51}^{0.26}$	-	-	-	$8.69_{0.51}^{0.26}$	$8.33_{0.28}^{0.28}$	
DES17X3dxu	$10.49_{0.32}^{0.22} \\ 8.60_{0.11}^{0.00}$	$1.02_{0.69}^{0.63} \\ -0.37_{0.11}^{0.10}$	$-9.47_{0.37}^{0.40}$ $-8.96_{0.00}^{0.10}$	0.26	-	-	-	0.26	0 18	
DES17X3hxi	$8.60^{0.00}_{0.11}$	$-0.37^{0.10}_{0.11}$	$-8.96^{0.10}_{0.00}$	$8.45_{0.45}^{0.26}$	-	-	-	$8.45_{0.45}^{0.26}$	$8.18_{0.18}^{0.18}$	

<sup>&</sup>lt;sup>a</sup> Linear combination of the likelihoods for D16, PP04 N2, PP04 O3N2, KK04 R23.

(DES16C1cbd) lies among passive galaxies. The spectrum of this object is red in colour, but does exhibit [OII] emission indicative of recent star-formation activity consistent with the upper end of the sSFR error bar. The hosts of individual rapid transients Cow, Koala and SN2018gep are lower in mass than the majority of RETs. SN2018gep and Cow lie along the SFMS, while Koala sits in the starburst regime, which is not heavily populated by DES RETs. SNLS04D4ec is consistent with the peak of the DES RET mass and SFR distributions.

Figs. 8 and 9 are similar to Figs. 6 and 9, except that here SFR has been normalised by stellar mass, and thus shows the specific star-formation rate (sSFR), which is a more representative measure of the star-forming efficiency.

<sup>&</sup>lt;sup>b</sup> Weighted average of PP04 N2, PP04 O3N2, and KK04 R23, where N2 and R23 were converted to PP04 O3N2 via Kewley & Ellison (2008).

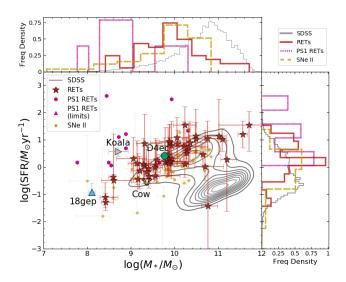
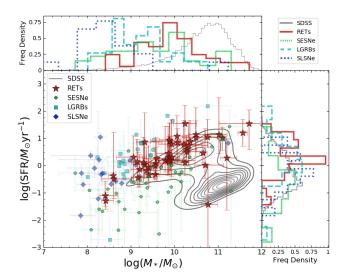


Figure 6. The  $M_*$ – SFR sequence of RET hosts, hosts of other rapid transients and SNe II, and SDSS field galaxies (grey contours). SFRs have not been corrected for redshift evolution. RET hosts lie slightly above the low-z star-formation main sequence upper half of the SDSS contours, and systematically avoid passive galaxies (dense SDSS contours in the lower right).



**Figure 7.** Same as Fig. 6, but comparing RET hosts with hosts of SESNe, LGRBs and SLSNe, and SDSS field galaxies (grey contours).

It is once again clear that RET hosts lie systematically above the majority of SDSS star-forming galaxies in terms of sSFR. Normalised by mass, it is here perhaps clearer to see that RET hosts lie at higher sSFR than CCSNe hosts, but not in the extremely star-forming environments of LGRBs or SLSNe.

We show the cumulative distribution of sSFR in Fig. 10. The RET hosts are clearly shifted to higher sSFRs than CCSNe. To statistically compare the host sSFR distribution of RETs with the other samples, we employ the method of W20. For each pair of samples, we model the PDFs as skewed normal distributions described by the parameters 'loc' (loca-

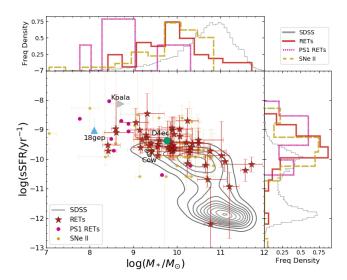
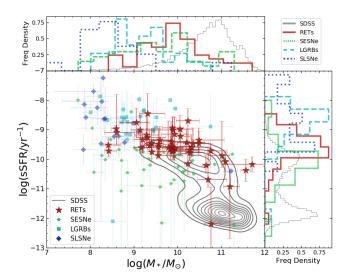


Figure 8. The  $M_*$ - sSFR sequence of RET hosts as well as local rapid transients and SNe II, along with the SDSS field galaxies.

tion, identical to the mean for zero skewness), 'scale' (spread, identical to the standard deviation for zero skewness)<sup>2</sup>, and 'skewness'. To impose priors on loc and scale, we combine the two samples and use normal priors centred on the combined mean and twice the combined standard deviation respectively, while for skewness the prior is a broad normal distribution centred on 0. We note that the loc parameter describes the location of the distribution (its relative position on the x-axis) and is not a mean, median, or mode. A highly-skewed distribution may have a loc that lies above almost the entire sample. A worked example as well as the results from the simultaneous fitting are displayed in Appendix B.

The comparison shows RET hosts to be shifted to higher sSFRs than CCSNe. In 98% of the posterior samples the RET sSFR distribution loc was at a higher value than SNe II, while the same was true 95% of the time for RETs when comparing with SESNe. The mean difference is 0.6 dex. To test whether some of this could be attributed to the difference in redshift between the samples we apply an approximate redshift correction based on the parameterisation of the SFMS at different redshifts by Salim et al. (2007) and Noeske et al. (2007), as has been done in other comparisons such as Taggart & Perley (2019). Transforming all host SFRs to their values at z = 0 would result in the CCSN SFRs dropping by an average of 0.05 dex, while the RETs would decrease 0.35 dex, i.e. a difference of 0.30 dex, or half of the observed difference. The remaining 0.30 dex is thus consistent with being an intrinsic difference. RET host galaxies are significantly lower in sSFR than LGRBs. While the distributions are similar in shape, with a mean difference in scale of 0.09, the mean difference in loc is -1.16, with no overlap between the posterior distributions. The sSFR distribution of SLSNe hosts is much broader than the RETs, with a scale of 1.12, twice that of the RETs. They are also

<sup>&</sup>lt;sup>2</sup> See W20 for a detailed description of the fitting procedure and the parameters describing the skewed normal distributions.



**Figure 9.** Same as Fig. 8 but comparing RET hosts with hosts of SESN, LGRB, and SLSN hosts, along with the SDSS field galaxies.

shifted to higher sSFRs than RETs, with the loc of the distribution on average 0.76 dex greater than RETs. The strong high-sSFR tail shows SLSNe occur in a different galaxy population to RETs.

#### 5.2 Metallicity

In the Section 5.1 we demonstrate that RETs occur in galaxies with systematically higher sSFR than CCSNe, to which one explanation is that they are related to more massive stars. A further property that could directly impact the composition of stellar populations harbouring potential RET progenitors is the metallicity. Using the gas-phase oxygen abundances calculated in Section 4.3 as a proxy for metallicity, we compare the chemical state of RET host galaxies with CCSNe and star-forming field galaxies. The cumulative distributions of metallicity are displayed in Fig. 11, and show RET hosts to be inconsistent with SNe II and field galaxies. The RET curve lies at lower metallicity than those galaxies, and appears visually similar to the curves for SESNe. The metallicity distribution of SESNe is, however, quite broad (e.g. Anderson et al. 2010), with different subclasses showing different trends (with SNe Ic host environments exhibiting higher metallicity than Ib, and IIb much lower). RETs occur, on average, in slightly more metal-rich environments than LGRBs and SLSNe.

We compare the metallicity distributions in the same way as the sSFRs, with the distribution fits shown in Appendix C. The RET host metallicity distribution shows a broad peak, leading to two families of skewed-Gaussians that fit it well, one with a low-valued centre  $(12 + \log{(O/H)} \sim 8.1)$  and a positive skew, and the other with a higher-valued centre  $(12 + \log{(O/H)} \sim 8.6)$ . Comparing the DES RETs to the Stoll et al. (2013) SNe II shows the latter to be centred around 8.8, with the centre being greater than the RETs in 94% of samples. We determine that the RET host metallicities are derived from a different population to the SNe II. On the other hand, simultaneous fits with SESNe show very

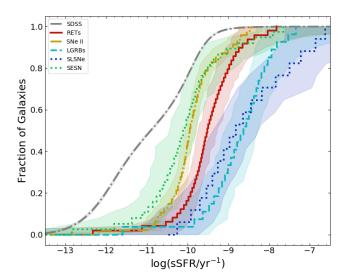
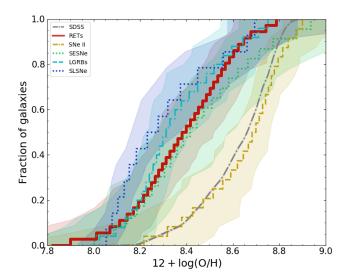


Figure 10. Cumulative distributions of the sSFR of RET hosts, compared to CCSNe and the low-z SDSS sample. Uncertainties have been estimated via a bootstrap Monte-Carlo technique and include limits.

similar distributions, including a smaller higher-metallicity peak, such that they are indistinguishable statistically.

The median CDFs of LGRBs and SLSNe show divergence from the RETs, particularly at low metallicity. As a result, the locs of their fits are shifted compared to the RETs. In 67% of samples, the RET sample had a higher loc than the LGRBS, with RETs also showing a broader distribution 67% of the time. While these effects are not as significant as with the RETs - SNe II comparison, there is mild evidence that RETs are located in galaxies with higher metal content than LGRBs. The effect is more pronounced for SLSNe, where the RETs have a higher metallicity for the distribution peak 92% of the time. The SLSN distribution is also more strongly skewed, with 89% of the posterior distribution being more strongly skewed than the RETs. There is thus mild-to-strong evidence that RETs occur in more metal-rich environments than SLSNe.

In Fig. 12 we show the MZR for the RET and comparison samples. The contours show the MZR for low-redshift  $(\hat{z} = 0.08)$  star-forming galaixes from SDSS, adjusted to the PP04 O3N2 diagnostic. We use the MZR parameterisation Zahid et al. (2014) to show the best fit to the MZR for starforming galaxies. The blue dashed line shows the fit to the low-z data, while the green dashed line corresponds to the MZR at z = 0.45, the mean redshift of the RET host sample. The RET hosts lie systematically below the galaxy MZR fits as well as the bulk of the SDSS galaxies, meaning that for a given stellar mass they have a lower metallicity. They populate similar regions to SESNe, LGRBs and SLSNe but are clearly offset from the SNe II. SN2018gep, the Koala, and SNLS04D4ec appear in line with the RETs, while the Cow lies quite distinctly above the MZR, and is even outside the bulk of the local field galaxies.



**Figure 11.** Cumulative distributions of the gas-phase oxygen abundances of RET hosts, hosts of the comparison samples, and the SDSS sample. Uncertainties have been estimated via a bootstrap Monte-Carlo technique and include limits.

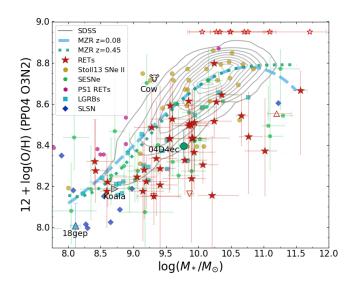


Figure 12. The mass-metallicity relation (MZR) for RET host galaxies and comparison samples. Upward- and downward-pointing triangles reflect lower and upper limits respectively. The DES RETs with no metallicity measurement have been placed at the top of the figure for completeness. The dashed lines represent MZR parameterisations from Zahid et al. (2014).

# 6 DISCUSSION

# 6.1 Selection Biases

The properties presented in Section 5 are derived from a subset of the total sample of RETs. Of 106 objects, under half (52/106) have secure host galaxy redshifts. Three of these were obtained from transient spectra, for which we are unable to disentangle the host and transient contributions, and four were obtained by programmes for which we do not have access to the spectra. Of the remaining 45, it was possible to derive a metallicity or at least a limit for

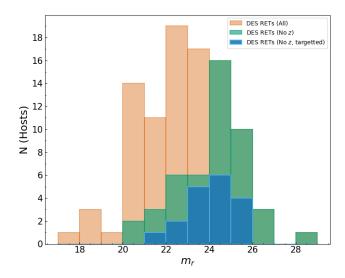


Figure 13. Observer-frame r-band magnitude distribution for the host galaxies of RETs in DES. The orange histogram represents the 96/106 DES RETs for which a host was detected. The green histogram shows those that did not have a successful redshift measurement, while blue shows those with no redshift despite being targetted by OzDES.

**Table 3.** Summary statistics for the samples compared in this work. Here we show the mean and standard deviation of the stellar mass, sSFR, and oxygen abundance, but note that the values for these parameters are not necessarily normally distributed for each sample.

Sample	$\langle \log{(M_*)} \rangle$	$\langle \log (sSFR) \rangle$	$\langle 12 + \log{(O/H)} \rangle$
	$(M_{\odot})$	$\left( yr^{-1}\right)$	
DES RETs	$9.9 \pm 0.8$	$-9.6 \pm 0.6$	$8.4 \pm 0.2$
PS1 RETs	$9.0 \pm 0.8$	$-9.2 \pm 0.8$	$8.5 \pm 0.2$
SNe II	$9.8 \pm 0.8$	$-9.8 \pm 0.7$	$8.6 \pm 0.2$
SESNe	$9.5 \pm 0.9$	$-10.2 \pm 1.0$	$8.4 \pm 0.2$
LGRBs	$9.2 \pm 0.6$	$-8.9 \pm 0.8$	$8.4 \pm 0.2$
SLSNe	$8.5 \pm 0.9$	$-8.7\pm0.9$	$8.3 \pm 0.2$

40 host galaxies, while five exceeded the redshift range for the necessary lines to fall within the wavelength coverage of AAOmega. The observed metallicity distribution could have arisen if the galaxies without redshifts (and metallicities) are systematically higher in metallicity than those for which measurements were possible. For low SNR objects, redshifts are typically obtained from only two of the strongest lines (e.g.  ${\rm H}\alpha$ ,  ${\rm H}\beta$ , [OIII], and [OII]). It is likely that the redshifts were not obtained because the galaxies are physically smaller or are at higher redshift. However, galaxies with high metallicity have weaker [OIII] lines, meaning they are less likely to have a redshift detection compared to less enriched galaxies with the same mass and redshift. Future, deeper spectral observation programmes as well as large, complete low-redshift samples are necessary to reduce this possible bias.

Another possibility is that the hosts without a redshift are mostly non star-forming, passive galaxies, for which a redshift is typically harder to obtain than for emission-line galaxies (Yuan et al. 2015; Childress et al. 2017; Lidman et al. 2020). To test this possibility, we examined the RETs

Cut	Number of remaining objects
All RETs	106
No redshift	57
Has host in SN Deep	47
Has host in SVA1	40
Targetted by OzDES	18

**Table 4.** Numbers of RETs passing various cuts relating to redshift targetting and completeness. Each row is a subset of the row above

that do not have a host galaxy redshift. Table 4 shows the numbers of RETs that failed various stages of the redshifting process, and is summarised in Fig. 13. Of the 57 objects without a redshift, 47 of them have host galaxies detected in the SN Deep coadds of W20. Of more significance is that only 40 have host galaxies in the SVA1 catalogues which were used for targeting during the OzDES campaign. The other, 'hostless', objects are either transients that are located remotely from a galaxy that was detected, or are hosted by a galaxy that was not detected. Non-detected hosts are either intrinsically faint and thus low in mass, situated at high redshift, or both. Neither are expected to be systematically higher in metallicity than the detected hosts. Similarly, a further 22 hosts were detected but not targeted by OzDES, due to being too faint to pass the selection criteria ( $m_r < 24.5$ ), leaving 18 that were targetted but no redshift was found. The resulting redshift completeness of targeted objects is 71% (83% for objects brighter than  $m_r = 24$  mag), which is in line with the average for OzDES (Lidman et al. 2020). In Fig. 14 we show the observer-frame r-band magnitudes and g-i colours for all RET hosts that were detected. The 47 objects with detected hosts but no redshift lie at fainter magnitudes, and appear to extend to bluer colours than those with secure redshifts. This is contrary to the hypothesis that they are high-redshift and/or passive hosts, but instead are low-mass, star-forming galaxies whose line fluxes were not strong enough to be detected. We thus conclude that the results presented from the subset of hosts with measured redshifts are at the very least representative of the star-forming nature of the population of RET hosts.

# 6.2 Origin of RETs

The sample of DES RETs shows a preference for lowmetallicity, strongly star-forming host environments. The PDF of their metallicities displays a strong similarity to the hosts of SESNe, as well as LGRBs. There is a clear difference to the PDF of SNe II, which follow SDSS field galaxies. The preference for low-metallicity systems is not as strong as for LGRBs or SLSNe, but the highest metallicities found in all three samples are very similar at around solar metallicity. This result is suggestive of a stripped-envelope, massivestar origin for RETs. The population of RET hosts lies, on average, between CCSNe and LGRBs/SLSNe in terms of both star formation and metallicity. A loose correlation exists between the luminosity and rarity of events, and the host galaxy conditions required for their formation - on average, rarer events occur in more extreme environments. The approximate rate of RETs ( $\geq 10^{-6} \text{Mpc}^{-3} \text{yr}^{-1}$ ), Drout et al. 2014, P18, Tampo et al. 2020) is  $\sim 1\%$  of the CCSN rate (Li et al. 2011; Horiuchi et al. 2011; Strolger et al. 2015),

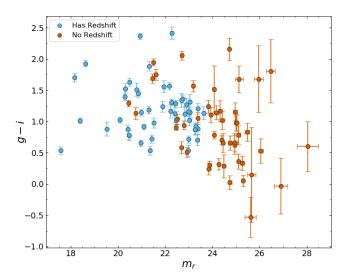


Figure 14. The colour-magnitude distribution of RET hosts with (cyan) and without (orange) redshifts. There is an excess of objects with blue colours that do not have redshift measurements.

which itself is divided into the more common SNe II and subdominant SESNe (Kelly & Kirshner 2012; Frohmaier et al. 2020). At  $\sim 1\%$  of the CCSN rate, RETs are more common than SLSNe ( $\sim 0.01 - 0.05\%$  of CCSNe; McCrum et al. 2015; Prajs et al. 2017; Frohmaier et al. 2020) and LGRBs (intrinsically ~ 0.08% when accounting for beaming; Graham & Schady 2016). These figures place the rate of DES RETs between extreme objects (SLSNe, LGRBs) and more common SNe (SNe II, SESNe) in terms of rate, matching the location of RET hosts in the various host galaxy parameter spaces presented in Section 5. While stressing rates are uncertain and host galaxy parameters span wide ranges for all transients, they are both linked to the respective transients' progenitor channels. Based upon both indicators, it is reasonable to infer that RETs are linked to very massive stars, potentially stripped of their envelopes, and possibly sharing some of the extreme properties of SLSNe or LGRBs. This hypothesis also suggests that RETs are an intermediate and/or precursory step, whereby the initial collapse of the star occurs leading to shock breakout and subsequent cooling, but conditions are not highly tuned enough for a LGRB or SLSN and the respective central engine does not form.

# 6.3 Correlations between lightcurve and host galaxy properties

Many classes of transients show trends between properties intrinsic to the objects themselves and their host galaxies. For example, SNe Ia lightcurves appear to be broader in less massive galaxies with higher sSFRs (Sullivan et al. 2006; Neill et al. 2009; Howell et al. 2009; Sullivan et al. 2010; Roman et al. 2018; Kelsey 2020), while SLSN lightcurves that have been fit with a magnetar model show a tentative relationship between the magnetar spin period and host galaxy metallicity (Chen et al. 2016). In Fig. 15 we show the RET peak magnitude (upper panels) and lightcurve width param-

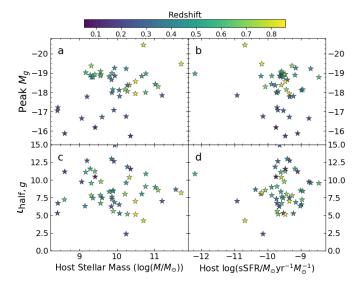


Figure 15. RET lightcurve properties as a function of host galaxy measurements.

eterised as  $t_{half}$ , the time the lightcurve is above half the peak brightness (lower panels), and their correlation with host galaxy stellar mass (left-hand panels) and sSFR (right-hand panels). The decline rates have been converted to the restframe of the transients, while the peak magnitudes have been k-corrected assuming a blackbody SED. There is no correlation between decline rate and either stellar mass or sSFR, while there are hints of a trend between peak magnitude and both mass and sSFR. These apparent trends are driven by the more extreme hosts (the three with  $\log (M_*/M_{\odot}) < 9$  and one with very high mass/low sSFR). Assuming that these points are not outliers, the trends are still likely driven by selection effects. At higher redshifts, only the brighter transients are recovered by the survey and our selection method. while at those high redshifts only the more massive galaxies are detected. This effect can be seen in Fig. 15a, with redshift increasing from the lower left to the upper right, while the same is true from the upper left to lower right in Fig. 15b. It is hoped that a more complete, volume-limited sample of RETs will be obtained by The Rubin Observatory Legacy Survey of Space and Time (LSST) in order to reveal any underlying relationships.

#### 6.4 Comparison with individual RETs

The nearby transient AT2018cow has drawn many comparisons to the cosmological RETs from DES and PS1 (e.g. Perley et al. 2019; Margutti et al. 2019; Fox & Smith 2019; Mohan et al. 2020) due to its rapid evolution and blue colour. AT2018cow displayed a contracting photosphere as well as evidence for central-engine power alongside an unusual spectrum that showed similarities to broad-lined SNe Ic (SN Icbl) at early stages (e.g. Xu et al. 2018; Izzo et al. 2018), developing to something entirely different at later epochs (Perley et al. 2019) with hints of similarities to interacting SNe Ibn (Fox & Smith 2019). There have been several suggestions that AT2018cow is indeed an analogue of the high-z RETs. The host galaxy of AT2018cow is to be moderately star forming and lies very close to the centre of the SFMS

(Lyman et al. 2020; Figs. 6,8), along with many of the DES RET hosts. However, the host lies somewhat above the fiducial MZR in Fig. 12, suggesting that it has an unusually high metallicity for its stellar mass. While consistent with SNe II, this is in contrast to the DES RET hosts which are systematically less enriched for a given stellar mass.

Other local rapid transients include SN2018gep (Ho et al. 2019), a spectroscopically classified SN Ic-bl with a rapid rise. The host of SN2018gep appears more similar to the DES RET sample, lying in the same  $M_*$ -SFR and  $M_*$ -SFR plane, as well as lying below the MZR. While the SN2018gep host is lower in stellar mass than any DES RET (log  $(M/M_{\odot})$  = 8.11), galaxies of that mass are unlikely to have been detected at the redshifts of the DES RETs (Wiseman et al. 2020). The authors' conclusion that SN2018gep is related to a shock-breakout of a massive, stripped-envelope star is similar to that posited in Section 6.2.

The rapidly evolving lightcurve of ZTF18abvkwla ("the Koala") has been attributed to shock interaction, while radio emission can be explained by a collimated jet. The host of the Koala is a low metallicity starburst more typical of LGRBs and SLSNe, and places this transient at the very extreme end of the DES RET host population. While we note that the Ho et al. (2020) study made multiple non-detections of radio emission from the DES RETs, these were taken at very late epochs ( $\geq$  1 year), so the presence of jets in the early evolution is not ruled out. Similarly, we cannot rule out that the Koala comes from the same population of transients as the DES RETs.

SN2018kzr (McBrien et al. 2019) is one of the most rapidly declining transients ever discovered, with spectral signatures similar to SNe Ic. While host galaxy properties are not derived, the authors of that paper refer to narrow emission from the host galaxy, along with an apparently small, blue, star-forming host and is thus consistent with the DES RETs.

#### 7 CONCLUSIONS

By analysing the host galaxies of 49 rapidly evolving transients (RETs) discovered in the Dark Energy Survey, we have been able to place constraints on the nature of these as-yet unexplained phenomena. We conclude that RETs are strongly linked to massive stars, due to their hosts all exhibiting signatures of star formation. They likely originate from stars more massive, on average, than than those that cause SNe II, and perhaps all SESNe, as they occur in galaxies with higher sSFR. RET hosts are significantly lower in metallicity than SN II hosts, and marginally lower than SESN hosts, suggesting some reliance on rotational energy or other metallicity-dependent effects. Of the RET analogues discovered in modern large-area, high-cadence surveys, ZTF18abvkwla shares the most similar host galaxy characteristics with the DES RET population. SN2018gep appears in a galaxy too faint to have been detected by DES-SN at the redshift of most of the DES RETs, while the host of AT2018cow is higher in metallicity.

While current surveys such as ZTF (Bellm et al. 2019), GOTO (Dyer et al. 2018), and BlackGEM (Bloemen et al. 2016) are well equipped to find low-redshift RETs, a sample similar to that presented here will likely not be collected

until LSST comes online. With several hundreds of objects, detailed studies of RETs and their hosts will be possible in a systematic and more complete manner as has been achieved with LGRBs and SLSNe.

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This work makes extensive use of Astropy,<sup>3</sup> a community-developed core Python package for Astronomy (Astropy Collaboration et al. 2013, 2018), Pandas (Mckinney 2010), and matplotlib (Hunter 2007).

#### REFERENCES

Abazajian K. N., et al., 2009, ApJS, 182, 543 Ahumada R., et al., 2019, arXiv:1912.02905

Anderson J. P., Covarrubias R. A., James P. A., Hamuy M., Habergham S. M., 2010, MNRAS, 407, 2660

Anderson J. P., Habergham S. M., James P. A., Hamuy M., 2012, MNRAS, 424, 1372

Angus C. R., Levan A. J., Perley D. A., Tanvir N. R., Lyman J. D., Stanway E. R., Fruchter A. S., 2016, MNRAS, 458, 84
Angus C. R., et al., 2019, MNRAS, 487, 2215

Arcavi I., et al., 2016, ApJ, 819, 35

Astropy Collaboration et al., 2013, A&A, 558

Astropy Collaboration et al., 2018, ApJ, 156, 123

Baldry I. K., et al., 2018, MNRAS, 474, 3875

Baldwin J. A., Phillips M. M., Terlevich R., 1981, PASP, 93, 5

Bellm E. C., et al., 2019, PASP, 131, 018002

Bernstein J. P., et al., 2012, ApJ, 753, 152

Bloemen S., et al., 2016, in Ground-based and Airborne Telescopes VI. SPIE, p. 990664, doi:10.1117/12.2232522

Brinchmann J., Charlot S., White S. D., Tremonti C., Kauffmann G., Heckman T., Brinkmann J., 2004, MNRAS, 351, 1151

Bruzual G., Charlot S., 2003, MNRAS, 344, 1000

Cappellari M., 2017, MNRAS, 466, 798

Cappellari M., Emsellem E., 2004, PASP, 116, 138

Cappellari M., Michele 2012, ascl, p. ascl:1210.002

Cardelli J. A., Clayton G. C., Mathis J. S., 1989, ApJ, 345, 245Chabrier G., 2003, PASP, 115, 763

Chen T. W., Smartt S. J., Yates R. M., Nicholl M., Krühler T., Schady P., Dennefeld M., Inserra C., 2016, MNRAS, 470, 3566 Childress M., et al., 2013, ApJ, 770, 107

Childress M. J., et al., 2017, MNRAS, 472, 273

<sup>&</sup>lt;sup>3</sup> http://www.astropy.org

```
Coil A. L., et al., 2011, \ensuremath{\mathsf{ApJ}}\xspace, 741, \, 8
Cool R. J., et al., 2013, ApJ, 767, 118
Cooper M. C., et al., 2012, MNRAS, 425, 2116
Croom S., Saunders W., Heald R., 2004, AAONw, 106, 12
D'Andrea C. B., et al., 2018, arxiv: 1811.09565
De Cia A., et al., 2018, ApJ, 860, 100
Dopita M. A., Kewley L. J., Sutherland R. S., Nicholls D. C.,
   2016, Ap&SS, 361, 61
Driver S. P., et al., 2009, Astronomy and Geophysics, 50, 12
Drout M. R., et al., 2014, ApJ, 794, 23
Dyer M. J., et al., 2018, in SPIE. SPIE-Intl Soc Optical Eng, p. 14
   (arXiv:1807.01614), doi:10.1117/12.2311865
Filippenko A. V., 1988, ApJ, 96, 1941
Filippenko A. V., 2002, Annual Review of Astronomy and Astro-
   physics, 35, 309
Flaugher B., et al., 2015, ApJ, 150, 150
Foreman-Mackey D., 2016, The Journal of Open Source Software,
   1, 24
Fox O. D., Smith N., 2019, MNRAS, 488, 3772
Frohmaier C., Angus C. R., Sullivan M., 2020, MNRAS, in prep
Fruchter A. S., et al., 2006, Nature, 441, 463
Gal-Yam A., 2012, Science, 337, 927
Galama T. J., et al., 1998, Nature, 395, 670
Galbany L., et al., 2018, ApJ, 855, 107
Gallazzi A., Bell E. F., 2009, Astrophysical Journal, Supplement
   Series, 185, 253
Gallazzi A., Charlot S., Brinchmann J., White S. D., Tremonti
   C. A., 2005, MNRAS, 362, 41
Goldstein D. A., et al., 2015, ApJ, 150, 82
Graham J. F., Schady P., 2016, ApJ, 823, 154
Hjorth J., et al., 2003, Nature, 423, 847
Ho A. Y. Q., et al., 2019, ApJ, 871, 73
Ho A. Y. Q., et al., 2020, arXiv:2003.01222
Hoffman M. D., Gelman A., 2011, Journal of Machine Learning
   Research, 15, 1593
Horiuchi S., Beacom J. F., Kochanek C. S., Prieto J. L., Stanek
   K. Z., Thompson T. A., 2011, ApJ, 738, 154
Howell D. A., et al., 2009, ApJ, 691, 661
Hunter J. D., 2007, Computing in Science and Engineering, 9, 99
Iglesias-Páramo J., et al., 2013, A&A, 553, L7
Iglesias-Páramo J., et al., 2016, ApJ, 826, 71
Izzo L., et al., 2018, ATel, 11753, 1
James P. A., Anderson J. P., 2006, A&A, 453, 57
Japelj J., et al., 2016, A&A, 590, A129
Kauffmann G., et al., 2003, MNRAS, 341, 33
Kelly P. L., Kirshner R. P., 2012, ApJ, 759, 107
Kelly P. L., Kirshner R. P., Pahre M., 2008, ApJ, 687, 1201
Kelsey L., 2020, MNRAS, in prep
Kessler R., et al., 2015, ApJ, 150, 172
Kewley L. J., Ellison S. L., 2008, ApJ, 681, 1183
Kewley L. J., Dopita M. A., Leitherer C., Davé R., Yuan T., Allen
   M., Groves B., Sutherland R., 2013, ApJ, 774, 100
Klebesadel R. W., Strong I. B., Olson R. A., 1973, ApJ, 182, L85
Kobulnicky H. A., Kewley L. J., 2004, ApJ, 617, 240
Krühler T., et al., 2015, A&A, 581, A125
Kuin N. P. M., et al., 2019, MNRAS, 487, 2505
Le Floc'h E., Charmandaris V., Forrest W. J., Mirabel I. F., Ar-
   mus L., Devost D., 2006, ApJ, 642, 636
Leloudas G., et al., 2015, MNRAS, 449, 917
Levesque E. M., Kewley L. J., Graham J. F., Fruchter A. S., 2010,
   ApJ, 712
Li W., Chornock R., Leaman J., Filippenko A. V., Poznanski D.,
   Wang X., Ganeshalingam M., Mannucci F., 2011, MNRAS,
   412, 1473
Lidman C., et al., 2020, MNRAS, accepted
Lunnan R., et al., 2014, ApJ, 787, 138
```

```
Lyman J. D., Galbany L., Sanchez S. F., Anderson J. P., Kun-
   carayakti H., 2020, arXiv:2005.02412
Lyutikov M., Toonen S., 2019, MNRAS, 487, 5618
Margutti R., et al., 2019, ApJ, 872, 18
McBrien O. R., et al., 2019, ApJ, 885, L23
McCrum M., et al., 2015, MNRAS, 448, 1206
Mckinney W., 2010, in PROC. OF THE 9th PYTHON IN SCI-
   ENCE CONF. p. 51
Minkowski R., 1941, PASP, 53, 224
Modjaz M., et al., 2008, ApJ, 135, 1136
Mohan P., An T., Yang J., 2020, ApJ, 888, L24
Morokuma-Matsui K., et al., 2019, ApJ, 879, L13
Neill J. D., et al., 2009, ApJ, 707, 1449
Neill J. D., et al., 2011, ApJ, 727, 15
Noeske K. G., et al., 2007, ApJ, 660, L43
Oke J. B., Gunn J. E., 1983, ApJ, 266, 713
Palmerio J. T., et al., 2019, A&A, 623, A26
Perley D. A., et al., 2016a, ApJ, 817, 8
Perley D. A., et al., 2016b, ApJ, 830, 13
Perley D. A., et al., 2019, MNRAS, 484, 1031
Pettini M., Pagel B. E. J., 2004, MNRAS, 348, L59
Prajs S., et al., 2017, MNRAS, 464, 3568
Prentice S. J., et al., 2018, ApJ, 865, L3
Pursiainen M., et al., 2018, MNRAS, 481, 894
Pursiainen M., et al., 2020, MNRAS, accepted
Quimby R. M., et al., 2011, Nature, 474, 487
Rest A., et al., 2018, Nat. Astron., 2, 307
Roman M., et al., 2018, A&A, 615, A68
Sako M., et al., 2008, ApJ, 135, 348
Salim S., et al., 2007, ApJS, 173, 267
Salpeter E. E., 1955, ApJ, 121, 161
Sanders N. E., et al., 2012, ApJ, 758, 132
Schlegel E., 1990, MNRAS, 244, 269
Schlegel D. J., Finkbeiner D. P., Davis M., 1998, ApJ, 500, 525
Schulze S., et al., 2018, MNRAS, 473, 1258
Soker N., Grichener A., Gilkis A., 2019, MNRAS, 484, 4972
Speagle J. S., Steinhardt C. L., Capak P. L., Silverman J. D.,
   2014, ApJS, 214, 15
Stanek K. Z., et al., 2006, Acta Astronomica, 56, 333
Stoll R., Prieto J. L., Stanek K. Z., Pogge R. W., 2013, ApJ, 773,
   12
Strolger L.-G., et al., 2015, ApJ, 813, 93
Sullivan M., et al., 2006, ApJ, 648, 868
Sullivan M., et al., 2010, MNRAS, 406, 782
Swann E., 2020, MNRAS, in prep
Taggart K., Perley D., 2019, arXiv, 1911.09112
Tampo Y., et al., 2020, ApJ, accepted
Tremonti C. A., et al., 2004, ApJ, 613, 898
Uno K., Maeda K., 2020, ApJ, submitted, arXiv:2003.05975
Vazdekis A., Sánchez-Blázquez P., Falcón-Barroso J., Cenarro
   A. J., Beasley M. A., Cardiel N., Gorgas J., Peletier R. F.,
   2010, MNRAS, 404, 1639
Vergani S. D., et al., 2015, A&A, 581, A102
Vergani S. D., et al., 2017, A&A, 599, A120
Wiseman P., et al., 2020, MNRAS, accepted, arXiv:2001.02640
Woosley S. E., 1993, ApJ, 405, 273
Woosley S. E., Bloom J. S., 2006, Annu. Rev. Astron. Astrophys,
   44, 507
Woosley S. E., Heger A., 2006, ApJ, 637, 914
Xu D., et al., 2018, ATel, 11740, 1
York D. G., et al., 2000, ApJ, 120, 1579
Yuan F., et al., 2015, MNRAS, 452, 3047
Zahid H. J., Dima G. I., Kudritzki R.-P., Kewley L. J., Geller
    M. J., Hwang H. S., Silverman J. D., Kashino D., 2014, ApJ,
```

791, 130

Lunnan R., et al., 2018, ApJ, 852, 81

# 18 P. S. Wiseman

Table A1. Emission line fluxes for DES RET host galaxies. Values are given in units of erg s<sup>-1</sup> cm<sup>-2</sup> Å<sup>-1</sup>, and have been corrected for Milky Way reddening using Schlegel et al. (1998) assuming a Cardelli et al. (1989) reddening law with  $R_V = 3.1$ , but have not been corrected for intrinsic host galaxy reddening.

	[OII]3727	[OIII]4960	[OIII]5007	[NII]6549	[NII]6585	[SII]6717	[SII]6731	$\mathrm{H}\delta$	Нγ	$H\beta$	$_{ m H}\alpha$
${\rm DES13X3gms}$	$1.2 \pm 14.8$	$0.3 \pm 1.8$	$0.9 \pm 1.8$	-	-	-	-	$0.0\pm1.9$	$15.1 \pm 2.0$	$0.7 \pm 1.5$	-
DES13C1tgd	$1.6 \pm 1.5$	$0.0 \pm 0.5$	$0.0 \pm 0.5$	$0.7 \pm 0.6$	$2.2 \pm 0.6$	$1.5\pm0.3$	$1.0\pm0.4$		$1.1\pm0.5$	$0.3 \pm 0.5$	$5.8 \pm 0.8$
DES13S2wxf	$32.7 \pm 10.5$	$1.3 \pm 1.1$	$4.1 \pm 1.1$	-	-	-	-	$1.9 \pm 1.2$	$2.1\pm1.0$	$1.8 \pm 1.1$	-
DES13X1hav	$3.0 \pm 1.2$	$1.0 \pm 0.6$	$3.0 \pm 0.6$	-	-	-	-	$0.0 \pm 0.3$	$0.5 \pm 0.5$	$0.9 \pm 0.3$	-
DES13X3nyg	$2.6 \pm 0.6$	$0.1 \pm 0.4$	$0.3 \pm 0.4$	-	-	-	-	$0.1 \pm 0.2$	$0.0 \pm 0.2$	$0.1 \pm 0.1$	-
DES13X3gmd	$2.7 \pm 0.7$	-	-	-	-	-	-	$0.1 \pm 0.4$	$0.3 \pm 0.4$	$1.3 \pm 1.0$	-
DES13C3bcok	$3.6 \pm 3.8$	$0.7 \pm 1.2$	$2.2 \pm 1.2$	-	-	-	-	$0.0 \pm 1.2$	$2.0\pm1.4$	$2.3\pm1.0$	$10.8 \pm 3.0$
DES13X2wvv	$16.4 \pm 2.7$	$2.9 \pm 1.4$	$8.8 \pm 1.4$	-	-	-	-	$1.1 \pm 1.3$	$0.6 \pm 1.2$	$4.3 \pm 1.1$	-
DES14X1bnh	$8.3 \pm 0.8$	-	-	-	-	-	-	$0.1\pm0.3$	$1.0\pm0.3$	-	-
DES15S1fli	$15.0 \pm 1.6$	$1.3 \pm 1.0$	$3.9 \pm 1.0$	-	-	-	-		$2.2 \pm 0.6$	$4.3 \pm 0.5$	-
DES15S1fll	$10.3 \pm 2.5$	$5.0\pm2.0$	$15.2 \pm 2.0$	$0.5 \pm 1.1$	$1.4 \pm 1.1$	$0.8 \pm 0.8$	$1.5\pm0.8$		$2.5\pm1.3$	$4.2 \pm 1.5$	$15.4 \pm 1.4$
DES14X3pkl	$2.5 \pm 1.2$	$0.8 \pm 0.4$	$2.3 \pm 0.4$	$0.2 \pm 0.2$	$0.5 \pm 0.2$	$0.8 \pm 0.2$	$0.4 \pm 0.2$	$0.4 \pm 0.5$	$0.6 \pm 0.6$	$0.0 \pm 0.7$	$2.3 \pm 0.3$
DES13X3npb	$15.9 \pm 6.6$	$0.8 \pm 0.8$	$2.4 \pm 0.8$	-	-	-	-	$2.0 \pm 0.7$	$2.6 \pm 0.4$	$3.3 \pm 0.8$	-
DES15X2ead	$31.4 \pm 23.5$	$1.6 \pm 3.0$	$4.8 \pm 3.0$	$1.1 \pm 1.6$	$3.5 \pm 1.6$	$5.3 \pm 1.7$	$3.2 \pm 1.9$	$9.6 \pm 7.9$	$4.5 \pm 6.0$	$5.4 \pm 3.1$	$16.7 \pm 1.9$
DES14S2plb	$54.5 \pm 7.2$	$4.5 \pm 2.1$	$13.5 \pm 2.1$	$5.6 \pm 1.8$	$16.9 \pm 1.8$	$9.5 \pm 0.9$	$7.1 \pm 0.7$	$6.5 \pm 2.0$	$7.9 \pm 1.7$	$19.3\pm1.5$	$53.4 \pm 1.7$
DES14S2pli	$7.7 \pm 1.1$	$0.4 \pm 0.4$	$1.2 \pm 0.4$	-	-	-	-	$0.0\pm0.6$	$1.1\pm0.5$	$2.4 \pm 0.3$	-
DES14C3tnz	$4.2 \pm 0.9$	$0.4 \pm 0.3$	$1.3 \pm 0.3$	-	-	-	-	$0.0 \pm 0.3$	$0.4 \pm 0.3$	$0.4 \pm 0.5$	-
DES15X3mxf	$4.8 \pm 0.7$	$0.4 \pm 0.3$	$1.2 \pm 0.3$	-	-	-	-	$0.6 \pm 0.5$	$0.7 \pm 0.4$	$1.1\pm0.3$	-
DES15C3lpq	$3.9 \pm 0.8$	$0.5 \pm 0.3$	$1.6 \pm 0.3$	-	-	-	-	$0.5 \pm 0.3$	$0.2\pm0.2$	$0.6 \pm 0.2$	-
DES15C3nat	$3.9 \pm 0.6$	-	-	-	-	-	-	$0.2 \pm 0.1$	$0.2 \pm 0.1$	-	-
DES15C3mgq	$4.1 \pm 2.5$	$1.1 \pm 0.5$	$3.2 \pm 0.5$	$0.1 \pm 0.2$	$0.2 \pm 0.2$	$0.2 \pm 0.2$	$0.0 \pm 0.4$	$0.6 \pm 0.8$	$0.4 \pm 0.7$	$1.2 \pm 0.5$	$1.4 \pm 0.2$
DES15E2nqh	$0.7 \pm 1.0$	$0.9 \pm 0.5$	$2.8 \pm 0.5$	-	-	-	-	$0.4 \pm 0.4$	$0.9 \pm 0.5$	$0.8 \pm 0.5$	-
DES15C3opk	$1.6 \pm 0.8$	$0.5 \pm 0.3$	$1.4 \pm 0.3$	-	-	-	-	$0.1 \pm 0.3$	$0.3 \pm 0.2$	$0.4 \pm 0.3$	-
DES15C3opp	$2.5 \pm 0.7$	$0.2 \pm 0.2$	$0.5 \pm 0.2$	-	-	-	-	$0.2 \pm 0.4$	$0.5 \pm 0.3$	$0.4 \pm 0.1$	-
DES16E2pv	$2.5 \pm 0.6$	$0.7 \pm 1.0$	$2.2 \pm 1.0$	-	-	-	-	$0.0 \pm 0.4$	$0.7 \pm 0.5$	$1.3 \pm 0.6$	-
DES16S1bbp	$21.9 \pm 3.5$	$3.8 \pm 0.9$	$11.4 \pm 0.9$	$0.3 \pm 0.4$	$1.0\pm0.4$	$1.5 \pm 0.7$	$0.9 \pm 0.3$	$1.4 \pm 1.1$	$3.4 \pm 0.9$	$6.4 \pm 1.0$	$11.1 \pm 0.6$
DES16X3cxn	$2.3 \pm 0.5$	$0.3 \pm 0.4$	$0.9 \pm 0.4$	-	-	-	-	$0.0 \pm 0.3$	$0.2 \pm 0.2$	$0.6 \pm 0.2$	-
DES16C1cbd	$4.4 \pm 1.8$	$0.4 \pm 0.8$	$1.2 \pm 0.8$	-	-	-	-	$0.4 \pm 0.4$	$0.4 \pm 0.4$	$1.1 \pm 0.4$	-
DES16C3axz	$26.8 \pm 2.6$	$2.3 \pm 0.8$	$7.0 \pm 0.8$	$2.6 \pm 0.6$	$7.7 \pm 0.6$	$5.4 \pm 0.5$	$4.0 \pm 0.5$	$2.7 \pm 0.9$	$5.6 \pm 0.8$	$9.9 \pm 0.8$	$29.1 \pm 0.7$
DES16X3erw	$5.9 \pm 0.9$	$0.8 \pm 0.6$	$2.6 \pm 0.6$	-	-	-	-	$0.6 \pm 0.4$	$0.4 \pm 0.4$	$1.9 \pm 0.3$	-
DES16C3gin	$5.8 \pm 1.2$	$0.5 \pm 0.4$	$1.6 \pm 0.4$	-	-	-	-	$0.5 \pm 0.6$	$0.0 \pm 0.6$	$1.6 \pm 0.4$	$2.5 \pm 0.4$
DES16S1dxu	$29.2 \pm 3.6$	$6.4 \pm 1.9$	$19.5 \pm 1.9$	$0.4 \pm 1.0$	$1.2 \pm 1.0$	$0.5 \pm 1.0$	$1.4 \pm 0.7$	$1.7 \pm 1.1$	$2.9 \pm 0.9$	$8.6 \pm 1.1$	$16.5 \pm 0.8$
DES16X1eho	$2.5 \pm 0.9$	-	-	-	-	-	-	$0.0 \pm 0.3$	$0.2 \pm 0.4$	$0.7 \pm 0.5$	-
DES17C3gop	$2.0 \pm 0.8$	$0.2 \pm 0.3$	$0.6 \pm 0.3$	-	-	-	-	$0.2 \pm 0.3$	$0.6 \pm 0.4$	$0.5 \pm 0.2$	-
DES17S2fee	$0.8 \pm 5.1$	$0.0 \pm 1.2$	$0.0 \pm 1.2$	$0.4 \pm 1.0$	$1.2 \pm 1.0$	$0.4 \pm 1.1$	$1.4 \pm 1.0$	$0.0 \pm 1.2$	$0.6 \pm 1.2$	$0.0 \pm 1.0$	$2.3 \pm 0.7$
DES17X3dxu	$3.1 \pm 0.5$	-	-	-	-	-	-	$0.7 \pm 0.5$	$0.7 \pm 0.2$	-	-
DES17X3cds	$3.3 \pm 1.8$	$0.4 \pm 0.4$	$1.4 \pm 0.4$	-	-	-	-	$0.0 \pm 0.4$	$0.3\pm0.3$	$0.8 \pm 0.3$	-
DES17C3fwd	$14.7 \pm 1.5$	$2.0\pm0.6$	$6.2 \pm 0.6$	$0.3 \pm 0.6$	$0.8 \pm 0.6$	-	-	$1.1\pm0.7$	$0.5 \pm 0.7$	$3.8 \pm 0.5$	$6.2 \pm 1.5$
DES17X3hxi	$5.4 \pm 1.8$	$2.3 \pm 0.9$	$7.1 \pm 0.9$	-	-	-	-	$1.0\pm0.7$	$1.6\pm0.7$	$1.6\pm0.5$	-
DES13E2lpk	$8.9 \pm 1.5$	$0.4 \pm 0.5$	$1.2\pm0.5$	-	-	-	-	$0.4 \pm 0.6$	$0.8 \pm 0.5$	$2.4 \pm 0.4$	-
DES15C2eal	$1.9 \pm 0.6$	$0.2 \pm 0.2$	$0.6 \pm 0.2$	$0.1 \pm 0.2$	$0.2\pm0.2$	$0.4 \pm 0.2$	$0.5 \pm 0.2$	$0.0 \pm 0.2$	$0.3 \pm 0.2$	$0.4 \pm 0.3$	$2.1 \pm 0.3$
DES16C2ggt	$11.8 \pm 1.2$	$1.1 \pm 0.4$	$3.2 \pm 0.4$	$0.4 \pm 0.4$	$1.4 \pm 0.4$	$1.7\pm0.3$	$1.1\pm0.3$	$0.4 \pm 0.5$	$1.5\pm0.6$	$3.4 \pm 0.4$	$6.9 \pm 0.5$
DES17C2hno	$3.7 \pm 1.1$	$0.2 \pm 0.2$	$0.7 \pm 0.2$	-	-	-	-	$0.4 \pm 0.5$	$0.6 \pm 0.4$	$0.9 \pm 0.2$	-

#### APPENDIX A: SPECTRAL LINE FLUXES

Table A1 presents the line fluxes for all DES RET hosts for which spectra were available. Spectra are available from the public OzDES DR2 at https://docs.datacentral.org.au/ozdes/overview/dr2/.

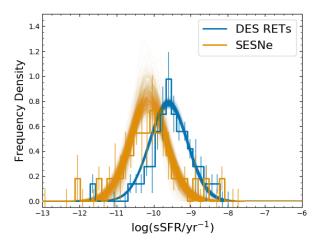


Figure B1. Results of the MCMC fits to the PDFs of sSFR for DES RETs and SESNe, accounting for uncertainties in each bin.

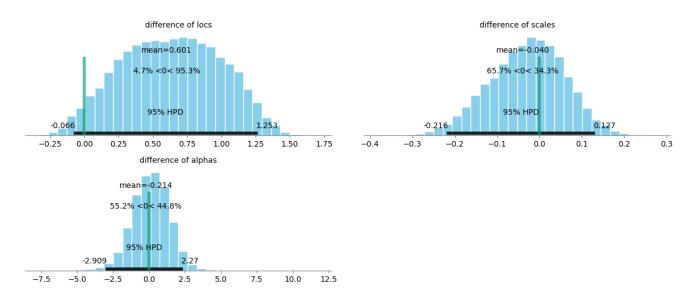


Figure B2. Histograms showing the differences between the fit parameters across the MCMC samples for the comparison between RETs and SESNe.

# APPENDIX B: BAYESIAN FITS - SSFR

To evaluate the likelihood that two independent distributions are from the same parent population we follow the method outlined in W20. We fit the PDFs, along with the uncertainty on the value in each bin, simultaneously with the same priors using the No U-Turn Sampler (NUTS; Hoffman & Gelman 2011) Hamiltonian Monte Carlo algorithm via the  $pymc3^4$  package to explore the posterior distribution. We utilise two chains, for a warm-up period of  $5 \times 10^3$  iterations per chain and a fit period of  $5 \times 10^3$  iterations per chain. Fig. B1 displays an example of the resulting fit where the DES RETs and Sanders et al. (2012) sSFR distributions are compared. Each resulting distribution is described by the 'loc' (location), 'scale' (spread), and 'alpha' (skewness). We then compare the differences in these parameters, as seen in Fig. B2, by reporting objectively the percentages of posterior samples that overlap, and subjectively what this means for the similarity of the distributions.

<sup>4</sup> https://docs.pymc.io/

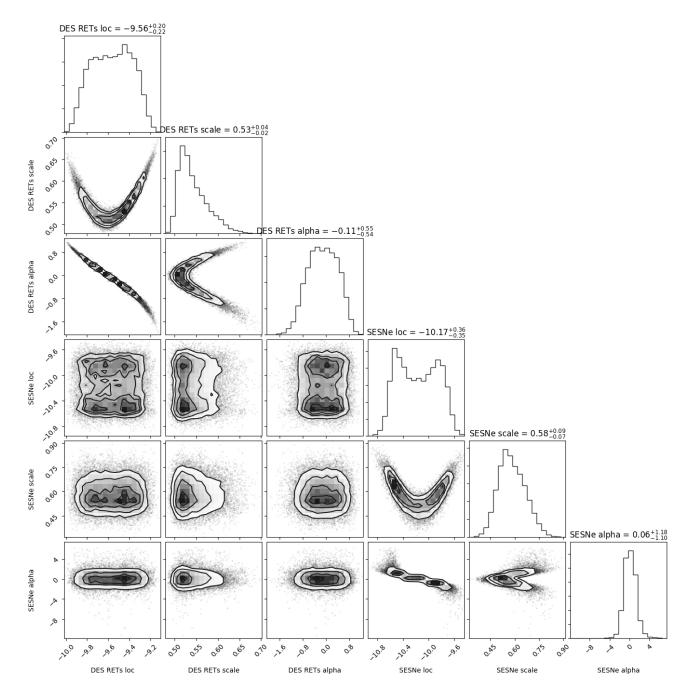


Figure B3. Corner plot showing the posterior samples from the MCMC fit to the DES RETs and Sanders et al. (2012) sSFRs. Notable features are: 1) the RET distribution is better constrained than the SESNe (S12); 2) the scale vs alpha and loc vs scale distributions are two-tailed due to alpha being centred close to 0; 3) there is a degeneracy between loc and alpha for the same reason. Loc and scale have units of yr<sup>-1</sup>, while alpha is dimensionless. Figure produced using the corner package (Foreman-Mackey 2016).

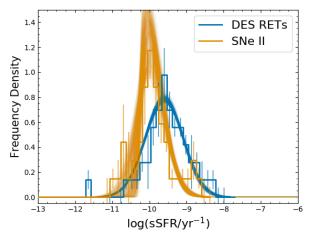
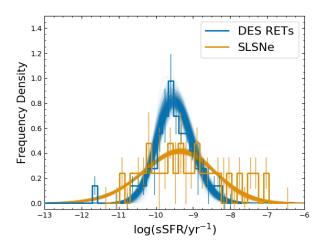


Figure B4. Results of the MCMC fits to the PDFs of sSFR for DES RETs and SNe II, accounting for uncertainties in each bin.



**Figure B6.** Results of the MCMC fits to the PDFs of sSFR for DES RETs and SLSNe, accounting for uncertainties in each bin.

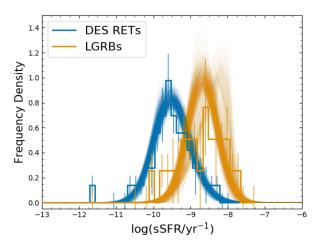
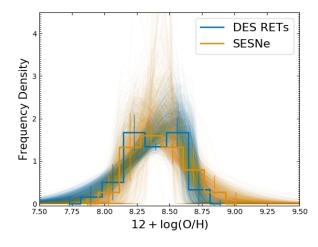


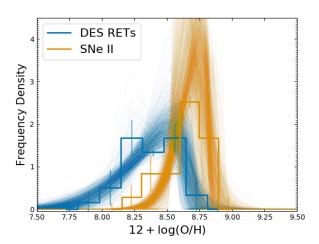
Figure B5. Results of the MCMC fits to the PDFs of sSFR for DES RETs and LGRBs, accounting for uncertainties in each bin.



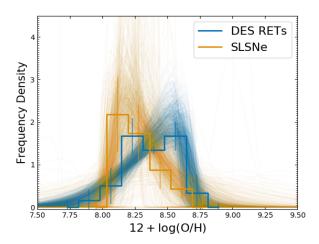
**Figure C1.** Results of the MCMC fits to the PDFs of metallicity for DES RETs and SESNe, accounting for uncertainties in each bin.

# APPENDIX C: BAYESIAN FITS - METALLICITY

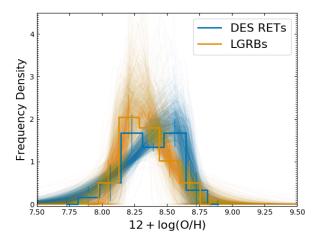
In this section, we present the Bayesian fits to the metallicity distributions of RETs and the comparison samples.



**Figure C2.** Results of the MCMC fits to the PDFs of metallicity for DES RETs and SNe II, accounting for uncertainties in each bin.



**Figure C4.** Results of the MCMC fits to the PDFs of metallicity for DES RETs and SLSNe, accounting for uncertainties in each bin



**Figure C3.** Results of the MCMC fits to the PDFs of metallicity for DES RETs and LGRBs, accounting for uncertainties in each bin.

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