

A galaxy-driven model of type Ia supernova luminosity variations

P. Wiseman,^{1*} M. Vincenzi,^{2,1} M. Sullivan,¹ L. Kelsey,^{3,1} B. Popovic,² B. Rose,² D. Brout,^{4,5} T. M. Davis,⁶ C. Frohmaier,^{3,1} L. Galbany,^{7,8} C. Lidman,^{9,10} A. Möller,¹¹ D. Scolnic,² M. Smith,¹² M. Agüena,¹³ S. Allam,¹⁴ F. Andrade-Oliveira,¹⁵ J. Annis,¹⁴ E. Bertin,^{16,17} S. Bocquet,¹⁸ D. Brooks,¹⁹ D. L. Burke,^{20,21} A. Carnero Rosell,^{22,13,23} M. Carrasco Kind,^{24,25} J. Carretero,²⁶ F. J. Castander,^{7,8} M. Costanzi,^{27,28,29} M. E. S. Pereira,³⁰ S. Desai,³¹ H. T. Diehl,¹⁴ P. Doel,¹⁹ S. Everett,³² I. Ferrero,³³ D. Friedel,²⁴ J. Frieman,^{14,34} J. García-Bellido,³⁵ M. Gatti,³⁶ E. Gaztanaga,^{7,8} D. Gruen,¹⁸ J. Gschwend,^{13,37} G. Gutierrez,¹⁴ S. R. Hinton,⁶ D. L. Hollowood,³⁸ K. Honscheid,^{39,40} D. J. James,⁴ M. March,³⁶ F. Menanteau,^{24,25} R. Miquel,^{41,26} R. Morgan,⁴² A. Palmese,⁴³ F. Paz-Chinchón,^{24,44} A. Pieres,^{13,37} A. A. Plazas Malagón,⁴⁵ A. K. Romer,⁴⁶ E. Sanchez,⁴⁷ V. Scarpine,¹⁴ I. Sevilla-Noarbe,⁴⁷ M. Soares-Santos,¹⁵ E. Suchyta,⁴⁸ G. Tarle,¹⁵ C. To,³⁹ and T. N. Varga^{49,50,51}

(DES Collaboration)

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ABSTRACT

Type Ia supernovae (SNe Ia) are used as standardisable candles to measure cosmological distances, but differences remain in their corrected luminosities which display a magnitude step as a function of host galaxy properties such as stellar mass and rest-frame $U - R$ colour. Identifying the cause of these steps is key to cosmological analyses and provides insight into SN physics. Here we investigate the effects of SN progenitor ages on their light curve properties using a galaxy-based forward model that we compare to the Dark Energy Survey 5-year SN Ia sample. We trace SN Ia progenitors through time and draw their light-curve width parameters from a bimodal distribution according to their age. We find that an intrinsic luminosity difference between SNe of different ages cannot explain the observed trend between step size and SN colour. The data split by stellar mass are better reproduced by following recent work implementing a step in total-to-selective dust extinction ratio (R_V) between low- and high-mass hosts, although an additional intrinsic luminosity step is still required to explain the data split by host galaxy $U - R$. Modelling the R_V step as a function of galaxy age provides a better match overall. Additional age vs. luminosity steps marginally improve the match to the data, although most of the step is absorbed by the width vs. luminosity coefficient α . Furthermore, we find no evidence that α varies with SN age.

Key words: supernovae: general – galaxies: evolution – cosmology: observations – dust, extinction

1 INTRODUCTION

Empirical relationships among light-curve properties of type Ia supernovae (SNe Ia) are routinely used to standardize their peak brightnesses (Pskovskii 1977; Phillips 1993; Riess et al. 1996; Tripp 1998; Mandel et al. 2017) and facilitate their use as distance indicators in cosmological measurements (Riess et al. 1998; Perlmutter et al. 1999; Sullivan et al. 2011; Betoule et al. 2014; Scolnic et al. 2018; DES Collaboration et al. 2018; Brout et al. 2022). This standardisation typically exploits the ‘faster–fainter’ relation between peak brightness and SN light curve width, and the ‘bluer–brighter’ relation between peak brightness and SN optical colour. A detailed astrophysical understanding of these relationships is not yet in place, but they are assumed to derive from physical processes intrinsic to the SN progenitors and their explosions (e.g., SN explosion mechanism

and physics, or progenitor white dwarf mass, age or metallicity), or in the circumstellar medium (CSM) surrounding the explosion, or in the interstellar medium (ISM) between the SN and the observer (e.g., dust). Established correlations between the observed SN light-curve width and the properties of the SN host galaxy, with faster-declining SNe more readily located in more massive, less actively star-forming galaxies, presumably reflect some of this underlying astrophysics.

After standardising SN luminosities for SN light-curve width and SN colour, any remaining difference between inferred distances and those predicted by the cosmological model are termed ‘Hubble residuals’. These Hubble residuals are also observed to correlate with host galaxy properties, both those local to the SN and the integrated global properties. The most commonly used of these correlations is the between Hubble residual and host galaxy stellar mass (Kelly et al. 2010; Lampeitl et al. 2010; Sullivan et al. 2010; Childress et al. 2013), which is typically represented by a step function in SN peak brightness of order 0.1 mag at a host galaxy stellar mass

* E-mail: P.S.Wiseman@soton.ac.uk

of around $\log(M_*/M_\odot) = 10$. A similar step in peak brightness is also observed when considering other host galaxy properties, such as specific star-formation rate (sSFR; Sullivan et al. 2010; Rigault et al. 2020), gas phase metallicity (Gallagher et al. 2005; Childress et al. 2013), rest-frame galaxy colour (Roman et al. 2018; Kelsey et al. 2021), and emission line equivalent widths (Dixon et al. 2022). By assuming that the step is caused by two populations of SNe Ia that can be (imperfectly) probed by one of these environmental tracers, Briday et al. (2022) showed that the size of the step is correlated with the level of population contamination across the step location of the given host galaxy tracer, with spectroscopically observed local sSFR and global stellar mass displaying the least contamination and largest steps. One interpretation of local sSFR being the best step tracer is that SNe Ia of different ages have different standardised luminosities, perhaps representing two (or more) SN Ia populations. Indeed, the case for multiple populations is strengthened by Nicolas et al. (2021) who show that the distribution of light-curve widths can be modelled as a combination of a young, slow-declining population and an old, fast-declining population.

Recent analyses have revealed a dichotomy in sizes of host galaxy – Hubble residual steps when the SN samples are divided based on the colour of the SNe. In the 3-year spectroscopically-confirmed sample of SNe Ia from the Dark Energy Survey (DES3YR; Brout et al. 2019a), redder SNe have large steps as a function of host stellar mass, while bluer SNe have negligible steps (Kelsey et al. 2021). This effect is also present in the larger, photometrically-classified DES5YR sample Kelsey et al. (2022, hereafter K22), who also showed that the effect is present when considering host galaxy rest-frame $U - R$ colour in place of stellar mass.

To explain the complex relationships between host galaxy Hubble residual steps and SN colour BS21 introduced a framework where the slope of the dust extinction law along the SN line-of-sight, as governed by the total-to-selective extinction parameter R_V , is different in low mass (mean $R_V = 2.75$) and high mass (mean $R_V = 1.5$) SN host galaxies. The effect of this model is that dustier SNe are redder and fainter, but the level of dimming is different for a SN with the same colour depending on the stellar mass of its host galaxy. Recent analysis by Popovic et al. (2021a) of the larger Pantheon+ SN Ia compilation (Scolnic et al. 2021) confirms the results from BS21, and shows that correcting SN Ia luminosities using this model may improve their use as cosmological probes. However, although the BS21 model reproduces the divergence of the stellar mass step as a function of SN colour, K22 show that after correcting for the trend of Hubble residual versus SN colour when split by stellar mass, there remains a small residual step in host galaxy $U - R$, with SNe in redder galaxies appearing brighter after correction. As this $U - R$ colour is related to stellar population age, this suggests there may also be an age-based effect in addition to dust.

In this work, we build upon the model of BS21 to qualitatively investigate the relationships between host galaxy mass, R_V , galaxy colour, and age, with the aim of explaining the residual $U - R$ step observed by K22 by expanding upon the multiple SN Ia age populations modelled by Nicolas et al. (2021).

Most modern simulations of SN Ia data sets use the SuperNova Analysis (SNANA; Kessler et al. 2009) software which constructs an accurate representation of how a large sample of SNe would appear in a survey, at the image level, by passing synthesised light curves through telescope responses and survey cadences. In this analysis we build a simplified toy model of the DES-SN survey by approximating many of the complex steps involved in SNANA, such as the light curve generation and fitting. This streamlining reduces computation time, allowing us to explore a broad parameter space while tracing

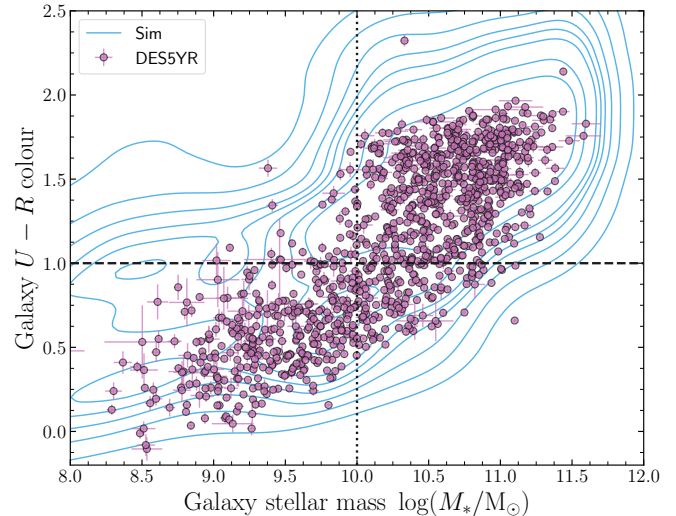


Figure 1. Rest frame galaxy $U - R$ colour versus stellar mass for the full simulated galaxy catalogue (blue contours) using the star-formation history models from W21, updated to include better quenching prescription. Magenta points are the observed stellar masses and colours from the DES5YR sample of SN Ia host galaxies (K22), and indicate where SN Ia hosts lie in this parameter space. The vertical dotted line indicates the $\log(M_*/M_\odot) = 10$ location assumed for mass-related transitions throughout this work, while the horizontal dashed line is located at the corresponding point used to separate galaxy colours, $U - R = 1$.

individual stellar populations through a large simulation of galaxies. Results of this qualitative analysis will then be used to inform full scale simulations using SNANA in the future.

We base our simulation on a forward model of host galaxy properties, and test the effects of galaxy and SN age on the light curves, dust extinctions and HRs of SNe Ia. While most SN simulations use empirical relationships to assign SNe to host galaxies that have minimal information (usually stellar mass and SFR), here we use the stellar mass assembly models of Wiseman et al. (2021, hereafter W21) in order to create a library of host galaxies for which the stellar age distribution of the galaxy is known at any given look-back time. We then derive a SN progenitor age distribution for each galaxy in the library, allowing us to trace the age of each simulated SN. In Section 2 we describe the galaxy simulation process and how we use stellar population synthesis models to translate star-formation histories into observables used in K22: stellar mass and rest-frame colour. In Section 3 we describe various models that we use to sample light-curve parameters for the SNe, paying notable attention to the SN progenitor age and host stellar ages. Section 4 presents a validation that the model approximates the observed SN parameter distributions and light-curve parameter – host galaxy relations. In Section 5 we evaluate the models based on their ability to reproduce the Hubble residual trends with SN colour, and in Section 6 we discuss the interpretations and implications of the results.

Where appropriate we assume a flat Λ CDM cosmological model with $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$ and $\Omega_M = 0.3$. Uncertainties are given as 1σ confidence intervals.

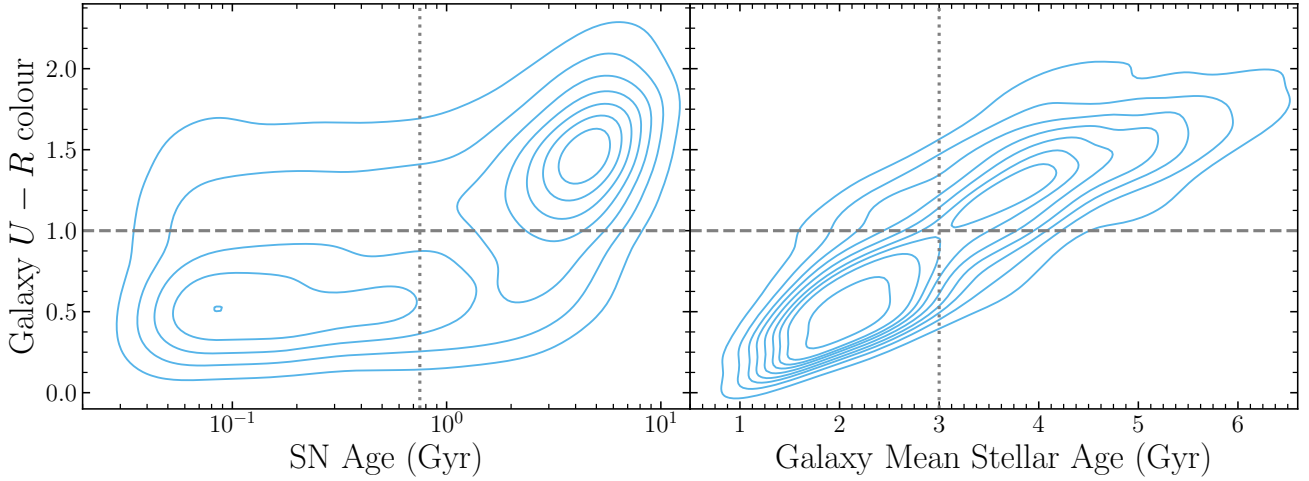


Figure 2. *Left:* Galaxy rest frame $U - R$ colour vs SN progenitor age at the time of explosion for SNe in the simulation. Contours represent lines of equal relative density. The horizontal line represents the colour $U - R = 1$ used to split observed populations, while the vertical line is the age used to split the young and old SN populations. *Right:* As left, but galaxy rest frame $U - R$ colour vs galaxy mass weighted mean stellar age. The vertical line shows the 3 Gyr split used for distinguishing young and old galaxies.

2 GALAXY SIMULATIONS

The main aim of this work is to create a realistic simulation of a population of SNe Ia with an accurate representation of the links between the light-curve properties of the SNe and the global properties of their host galaxies. This library is then used to link the host galaxy properties to Hubble residuals. In this section, we outline the method used to simulate the SN Ia host galaxies and their properties. In Section 3 we describe how we use the host galaxy library to simulate the SNe themselves.

We base our host galaxy simulation on the stellar mass assembly models of W21, which themselves are based on those of Childress et al. (2014). These simple analytical models of galaxy evolution are derived empirically from observed relationships between stellar mass, SFR and redshift, and take into account mass loss and quenching.

The ability of the models to reproduce observed stellar mass and SFR distributions is presented in Childress et al. (2014) and W21.

2.1 Updated galaxy models

In the simulations of W21, galaxies evolve following a deterministic set of equations, such that a simulated galaxy represents the average galaxy of a given stellar mass and redshift. While such a prescription is adequate for an overall modelling of SN rates as a function of mass and redshift, it does not result in an accurate *distribution* of host galaxy star formation rates, which is important when considering host galaxy colours. To mitigate this issue, we update our prescription of galaxy quenching to be more stochastic. We also add bursts of star formation to the models. These two additions are detailed in turn below.

2.1.1 Quenching

Each galaxy track in the simulation is represented by 100 implementations. At each time step of the simulation, each galaxy implementation is designated as either ‘non-quenched’ or ‘quenched’.

Quenched galaxies remain quenched, while non-quenched galaxies are subject to being quenched according to a quenching probability. This probability is a function of the stellar mass, similar to the average quenching penalty in the Childress et al. (2014) and W21 implementations of the model:

$$p_Q(M_*, z) = \frac{1}{2} \left[1 - \operatorname{erf} \left(\frac{\log(M_*) - \log(M_Q(z))}{\sigma_Q} \right) \right] - f_Q, \quad (1)$$

where f_Q is the fraction of the 100 implementations that have already quenched. In order to accurately reproduce the star formation main sequence and quenched fraction, we slightly adjust the reference quenching mass from W21 by 0.2 dex to:

$$\log(M_Q(z)/M_\odot) = 10.377 + 0.636z, \quad (2)$$

and the transition scale σ_Q from 1.1 to 1.5.

2.1.2 Bursts of star formation

A second update to the models of W21 is the addition of random bursts of star formation to the models. We add instantaneous, delta-function bursts that form stars at a rate of half of the main sequence value for the mass and redshift of the galaxy. Bursts occur at a constant probability of 0.05 for every time step in the simulation.

2.2 Simulating observable galaxy properties

We build upon the model SFHs by ascribing observable properties to each of the simulated galaxies. In particular, we are interested in the observed and rest frame spectral energy distributions (SEDs), which are the observed data used to infer global galaxy properties such as stellar mass, SFR, and stellar population age. For the latter, we focus on the rest frame absolute magnitudes in Bessell following Roman et al. (2018), Kelsey et al. (2021), and K22. In the simulation, we link the model SFHs (with known stellar mass, SFR, and age distributions) to synthetic broad-band photometry using a similar

method to that widely used in the inference of those observables: we convolve the SFHs with a grid of spectral templates of simple stellar populations (SSPs, Section 2.2.1) of different ages, summing up the individual SSP spectra according to the relative weights given by the SFH to produce a final composite spectrum S_{gal} . The luminosities of the combined spectra are scaled by the stellar mass of the galaxy, resulting in an accurate representation of the intrinsic rest-frame galaxy spectra. We add nebular emission from ionised gas (Section 2.2.2), and apply a reddening due to interstellar dust via attenuation (Section 2.2.3). This grid of host galaxies acts as a host galaxy library from which we sample a realistic SN Ia host population in Section 3.

2.2.1 Template libraries

The choice of spectral templates affects the relationship between galaxy properties and observables. We use the common [Bruzual & Charlot \(2003, hereafter BC03\)](#) library, and for consistency with the SN host galaxy literature we generate SSPs with a [Chabrier \(2003\)](#) initial mass function (IMF). The BC03 code draws upon the Padova 1994 evolutionary tracks ([Bertelli et al. 1994](#)). For simplicity, all SSPs are generated at solar metallicity (Z_{\odot}).

2.2.2 Nebular emission

In addition to the stellar continuum, our spectral templates incorporate nebular emission lines from H II regions. The BC03 models do not incorporate nebular emission, and thus we add it separately following [Boquien et al. \(2019\)](#). The line strengths are set relative to the ionizing photon flux, and are determined by the ratios at solar metallicity presented in [Inoue \(2011\)](#). The ionization parameter U is set at $\log(U) = -2$.

2.2.3 Dust attenuation

To arrive at realistic simulation of galaxy spectra, we apply the effects of dust attenuation to the final scaled spectral templates. We add dust according to the attenuation law of [Cardelli, Clayton & Mathis \(1989\)](#). For each galaxy there are 15 realisations with dust added in steps of equal size in V -band extinction A_V in the range $0 \leq A_V \leq 1.5$ mag; the final value of A_V for each host is chosen at a later stage along with the SN parameters, such that host has the same reddening $E(B-V)$ as the SN. To calculate the host extinction A_V we follow the results of [Salim et al. \(2018\)](#) and use a total-to-selective extinction R_V of 2.61 for star-forming galaxies with $\log(M/M_{\odot}) < 9.5$, 2.99 for $9.5 \leq \log(M/M_{\odot}) < 10.5$, and 3.47 for $\log(M/M_{\odot}) \geq 10.5$. For passive galaxies with $\log(\text{sSFR}) < -11$ we follow [Salim et al. \(2018\)](#) and assume the quiescent galaxy R_V of 2.61.

2.2.4 Synthetic galaxy photometry

We convert our simulated galaxy spectra into photometric observables by multiplying them with the transmission functions of desired passbands. This is possible with any filter at any wavelength – the galaxy templates extend from the far ultraviolet to the far infrared – but here we focus on optical wavelengths. For rest-frame galaxy photometry we calculate absolute magnitudes in the $UBVRI$ filters of [Bessell \(1990\)](#) following [K22](#), and in the observer frame we calculate apparent magnitudes in the DES *griz* filters (e.g., m_g^{host} , m_r^{host} ; [Flaugher et al. 2015](#)) for use in SN survey simulations (e.g., [Vincenzi et al. 2021](#)).

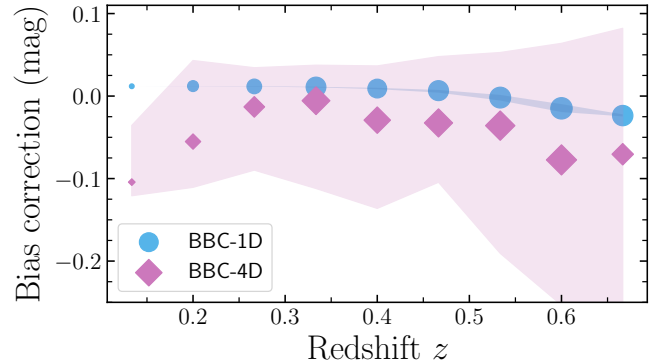


Figure 3. Average magnitude of a one-dimensional (BBC-1D; blue circles) and four-dimensional (BBC-BS20; magenta diamonds) bias correction as a function of redshift for a representative simulation. Points show the median in redshift bins; point sizes are proportional to the number of SNe in each redshift bin, while shaded regions show the 16th to 84th percentiles.

The resulting library of simulated galaxy stellar masses and rest-frame colours are shown in Fig. 1, with DES SN Ia hosts shown to illustrate how the simulation covers the necessary stellar-mass–colour parameter space from which to accurately draw SN hosts.

3 SUPERNOVAE

We next describe how we simulate a sample of SNe Ia. This includes ensuring a realistic association with their host galaxies, accurate distributions of their light-curve parameters, and the principle effects introduced by observational noise. We discuss each in turn.

3.1 SN host galaxy association

We start with our simulated galaxy models from Section 2, and take a flat distribution in stellar mass over the range $7 < \log(M_*/M_{\odot}) < 12$. We use a SN Ia delay-time distribution (DTD), which describes the rate at which SNe Ia occur as a function of the delay time τ_A since an episode of star formation, and convolve this DTD with each galaxy’s SFH. In our simulations we use the DTD of [W21](#), which takes a power-law form with index -1.13 and takes effect 40 Myr after star formation. The result is a SN Ia rate per year in each model galaxy (R_G). We then calculate a volumetric rate of SNe by multiplying R_G by a stellar-mass function ($\phi(M_*, z)$) at the redshift of the simulated galaxy. The combination $R_G \times \phi(M_*, z)$ is then used as a weighting to select a representative sample of SN Ia host galaxies in which SNe can be simulated. We use the redshift-dependent stellar-mass function from [ZFOURGE \(Tomczak et al. 2014\)](#).

The relationship between observed galaxy colour and SN and galaxy ages is shown in Fig. 2. There is a distinction between the host $U-R$ populations for young and old SNe, although there is a significant fraction of old SNe that exist in blue hosts. The relationship between galaxy colour and galaxy age is even clearer, with very little crossover between the old-and-red, and young-and-blue populations.

3.2 Light curve properties

For each simulated SN in a given host galaxy, we then assign ‘observed’ light-curve properties by drawing randomly from statistical

distributions. We use the standard SALT2 light-curve fitter (Guy et al. 2007) framework of x_1 (light-curve width, or ‘stretch’) and c (SN peak rest-frame optical colour), together with a Tripp (1998)-like standardisation relation that relates the inferred distance modulus μ_{obs} to SN Ia light-curve parameters. In most current SN Ia analyses, this takes the form

$$\mu_{\text{obs}} = m_B - M_B + \alpha x_1 - \beta c + \mu_{\text{bias}}, \quad (3)$$

where M_B is the peak SN absolute magnitude in the rest-frame B band, m_B is the peak rest-frame apparent magnitude, the α and β coefficients parametrize the stretch–luminosity and colour–luminosity relationships, and the μ_{bias} term is a correction for biases introduced by survey selection effects. Hubble residuals, μ_{res} , can then be calculated using

$$\mu_{\text{res}} = \mu_{\text{obs}} - \mu_{\text{model}}, \quad (4)$$

where μ_{model} depends on the cosmological parameters, which in this analysis we hold fixed at our fiducial model.

This simple framework assumes all colour variation in SNe Ia can be captured in a simple ‘ βc ’ term. However, we are explicitly interested in this analysis in the effect of dust extinction and reddening on SNe Ia, and thus we also assign each simulated SN a colour excess (i.e., $E(B-V)$) due to dust reddening that we denote E_{SN} , and a total-to-selective extinction parameter R_V . We next outline the methods used to draw the samples for each light-curve and dust parameter.

3.2.1 SN Ia absolute magnitude

We simulate SNe using a fixed peak B -band absolute magnitude of $M_B = -19.325$ mag. We allow for the possibility of host galaxy ‘steps’ of size γ , where SNe on either side of a threshold in a given host galaxy property X (such as stellar mass) have a factor of $\Gamma \times \gamma$ added to their simulated absolute magnitude, where $\Gamma = +0.5$ for hosts below the threshold and $\Gamma = -0.5$ for hosts above the threshold X_{split} . Thus,

$$M_B = -19.325 + \Gamma_X \times \gamma_X \text{ mag}, \quad (5)$$

with

$$\Gamma_X = \begin{cases} 0.5, & X < X_{\text{split}} \\ -0.5, & X \geq X_{\text{split}} \end{cases}, \quad (6)$$

where X_{step} is the location of the step in the given host parameter.

3.2.2 SN x_1

We include two prescriptions for the SALT2 x_1 parameter. The first parametrises the x_1 distribution as an asymmetric Gaussian, while the second assumes two Gaussians dependent on the age of the SN progenitor’s stellar population.

The first x_1 population model is used in BS21, which in turn follows Scolnic & Kessler (2016), where x_1 is drawn from an asymmetric Gaussian model defined by parameters $\sigma_-, \sigma_+, \bar{x}_1^{-1}$, giving a probability distribution

$$P(x_1) = \begin{cases} e^{[-(x_1 - \bar{x}_1)^2 / 2\sigma_-^2]}, & x_1 < \bar{x}_1 \\ e^{[-(x_1 - \bar{x}_1)^2 / 2\sigma_+^2]}, & x_1 > \bar{x}_1 \end{cases}. \quad (7)$$

The second x_1 model builds upon the evolving double-Gaussian of

¹ We use the notation \bar{x}_1 instead of the conventional μ_{x_1} to avoid confusion with the distance modulus μ , and maintain this notation for all parameter means.

Table 1. Outline of the different SN Ia models used in this analysis.

Model name	SN x_1 distribution	SN c distribution	R_V distribution
BS21	SK16	Gauss + exp (mass)	Step (stellar mass)
Age R_V	N21	Gauss + exp (age)	Step (galaxy age)
Age R_V Linear	N21	Gauss + exp (age)	Linear (WD age)

Nicolas et al. (2021) with the x_1 probability distribution represented as a Gaussian mixture model – a combination of two Gaussian distributions, one for a ‘young’ SN population and one an ‘old’ population. In the Nicolas et al. (2021) model, environmental age is traced by the local specific SFR (lsSFR). SNe in young (high lsSFR) environments are assumed to belong solely the old x_1 population, while SNe in old (low lsSFR) environments can be drawn from either the old or young x_1 modes, with the relative weight of each determined by a mixing probability a_{x_1} . Since the fraction of old and young SNe varies with redshift, so does the x_1 distribution (see Fig. 8 in Nicolas et al. 2021).

Instead of using lsSFR as an observational indicator of progenitor age, our simulations allow us to track the progenitors directly. For each galaxy, the likelihood that a SN exploding at time t_0 comes from a progenitor of age τ_A is

$$P(t_0, \tau_A) = \psi(t_0 - \tau_A) \Phi(\tau_A), \quad (8)$$

where $\psi(t_0 - \tau_A)$ is the SFH and $\Phi(\tau_A)$ the DTD. We sample SN age τ_A from this distribution, and subsequently derive x_1 from the age by adapting the Nicolas et al. (2021) prescription

$$x_1 \sim \begin{cases} \mathcal{N}(\bar{x}_{1,1}, \sigma_{x_1,1}), & \tau_A < \tau_{A,\text{thresh}} \\ a_{x_1} \times \mathcal{N}(\bar{x}_{1,1}, \sigma_{x_1,1}) + \\ (1 - a_{x_1}) \times \mathcal{N}(\bar{x}_{1,2}, \sigma_{x_1,2}), & \tau_A > \tau_{A,\text{thresh}} \end{cases} \quad (9)$$

where $\bar{x}_{1,[1,2]}$ and $\sigma_{x_1,[1,2]}$ are the mean and standard deviation of the young and old Gaussians respectively, a_{x_1} is the mixing coefficient which corresponds to the probability that the progenitor belongs to the young x_1 mode, and $\tau_{A,\text{thresh}}$ is the age which separates young and old timescales.

3.2.3 SN colour

We follow the method of BS21, which itself builds on Jha et al. (2007) and Mandel et al. (2011, 2017). These model the observed SN Ia colour distribution as a combination of an intrinsic colour c_{int} and the colour excess due to dust reddening E_{SN} . c_{int} is normally distributed

$$c_{\text{int}} \sim \mathcal{N}(\bar{c}_{\text{int}}, \sigma_{c_{\text{int}}}), \quad (10)$$

and E_{SN} is drawn from an exponential distribution following BS21

$$P(E_{\text{SN}}) = \frac{1}{\tau_E} e^{-E_{\text{SN}}/\tau_E}, \quad (11)$$

where τ_E is the mean reddening of the SN population.

3.2.4 Extinction law R_V

Related to the SN dust reddening E_{SN} is the slope of the reddening law R_V , such that the absolute V -band extinction A_V is

$$A_V = R_V E_{\text{SN}}, \quad (12)$$

Table 2. SN population parameters for the age-based models. Explanation of parameters can be found in Section 3.2.

$\overline{x_{1,1}}$	$\sigma_{x_{1,1}}$	$\overline{x_{1,2}}$	$\sigma_{x_{1,2}}$	a_{x_1}	τ_A , thresh	$\overline{c_{\text{int}}}$	$\sigma_{c_{\text{int}}}$	$\tau_{E,1}$	$\tau_{E,2}$
0.22	0.61	-1.22	0.56	0.38	0.75 Gyr	-0.075	0.042	0.135	0.135

and for the B -band magnitudes typically used in SN Ia analyses,

$$A_B = R_B E_{\text{SN}} \equiv (R_V + 1) E_{\text{SN}}. \quad (13)$$

We investigate three models for how R_V varies among SNe. These are inspired by the model of BS21, where R_V follows a step with stellar mass, and motivated by the evidence that the x_1 distribution is related to SN or host galaxy age. Our R_V that depend on some host galaxy property Y are:

- (i) BS21 ($Y = M_*$): R_V is drawn from a Gaussian distribution:

$$R_V \sim \mathcal{N}(\overline{R_V}, \sigma_{R_V}), \quad (14)$$

with $\overline{R_V}$ and σ_{R_V} taking on different values for host galaxies on either side of Y_{split} , a split point in stellar mass;

(ii) Age R_V ($Y = \tau_G$): The same model as BS21, but using mass-weighted mean stellar age of the host galaxy, τ_G , instead of host stellar mass, as the parameter to split R_V . τ_G is linked to both the observed stellar mass and the $U - R$ colour;

(iii) Age R_V Linear ($Y = \tau_A$): $\overline{R_V}$ follows a linear relationship with progenitor age τ_A . We parametrise the relationship by fixing the value of $\overline{R_V}$ for very young (t_1) and very old (t_2) progenitors as

$$\overline{R_V} = R_V^{t_1} + (\tau_A - t_1) \frac{R_V^{t_2} - R_V^{t_1}}{t_2 - t_1}, \quad (15)$$

and the standard deviation σ_{R_V} following a similar relationship in order to encapsulate having more diverse dust in younger systems:

$$\sigma_{R_V} = \sigma_{R_V}^{t_1} + (\tau_A - t_1) \frac{\sigma_{R_V}^{t_2} - \sigma_{R_V}^{t_1}}{t_2 - t_1}. \quad (16)$$

3.2.5 Mean reddening τ_E

For setting the mean reddening τ_E we follow BS21, fixing the value separately in low- and high-mass galaxies. We also implement a slight variation where the split is defined on the mass-weighted mean stellar age of the galaxy.

3.2.6 Colour-luminosity coefficient β

As per BS21 we do not assume a universal relationship between SN peak brightness and c_{int} . Instead we allow the coefficient relating them, denoted β_{sim} , to vary between SNe, and draw it from a Gaussian distribution:

$$\beta_{\text{sim}} \sim \mathcal{N}(\overline{\beta}, \sigma_{\beta}), \quad (17)$$

with mean $\overline{\beta}$ and standard deviation σ_{β} . In this model β_{sim} is independent of any host galaxy, progenitor, or dust property.

3.2.7 Stretch-luminosity coefficient α

For the baseline model we adopt a value of 0.15 for the stretch-luminosity coefficient α_{sim} . In cases where a mass or age luminosity step are introduced, the measured value of α can be different to the input value α_{sim} . In such cases we vary α_{sim} such that the measured value matches that found by K22. We further investigate the effects of α and its importance in age-based models in Section 6.3.1.

3.3 Observables

Given a set of simulated SN Ia light curve parameters, the synthetic m_B is then calculated by analogy to equation 3, modified to explicitly distinguish dust from SN colour variations, i.e.,

$$m_B = M_B + \mu(z)_{\text{sim}} - \alpha_{\text{sim}} x_1 + \beta_{\text{sim}} c_{\text{int}} + E_{\text{SN}}(R_V + 1) + \Delta M_B, \quad (18)$$

where $\mu(z)_{\text{sim}}$ is the distance modulus at the simulated redshift z given by the reference cosmology.

Even following the standardisation of SNe Ia there remains additional scatter in their observed M_B . We describe this in our simulated SNe using the parameter σ_{int} . We draw an offset (ΔM_B) to M_B from this scatter according to the distribution:

$$\Delta M_B \sim \mathcal{N}(0, \sigma_{\text{int}}). \quad (19)$$

The value of σ_{int} is tuned so that in combination with the rest of the model the overall unexplained scatter in the Hubble residuals matches that of the data, which is of the order ~ 0.1 mag.

3.4 Application to DES5YR data

The models introduced above are generic and intended to approximate the intrinsic universal properties of SNe Ia. Here we describe the data set against which we compare the simulations, and the survey-specific additions to the simulation that are necessary in order to replicate the selection effects of the data set.

We use an identical sample to that presented in K22: namely a photometrically classified SN Ia sample using five years of DES observations (DES5YR), as described in Vincenzi et al. (2021) and Möller et al. (2022), with a redshift limit of $z \leq 0.7$. Despite the redshift cut, the sample is magnitude-limited, and thus there are selection effects which we must approximate, which we discuss below.

3.4.1 Host galaxy spectroscopic efficiency

To be included in the sample, DES SNe Ia require a spectroscopic redshift of the host galaxy, and the efficiency of obtaining one in the spectroscopic follow-up campaigns is a strong function of host galaxy apparent magnitude, m_r^{host} . Vincenzi et al. (2021) presented the DES SN spectroscopic-redshift efficiency function $\epsilon_{z_{\text{spec}}}(m_r^{\text{host}})$, which we apply to our simulated SN hosts. In the simulation, we assign a probability of successfully obtaining a spectroscopic redshift for each galaxy at random from a Bernoulli distribution with success probability $p = \epsilon_{z_{\text{spec}}}(m_r^{\text{host}})$. In this work we do not model further survey efficiencies (e.g., SN detection efficiency; Brout et al. 2019b; Kessler et al. 2019). Therefore, instead of simulating a population of SNe that follows the true intrinsic volumetric rate as a function of redshift, we sample redshift from an ad-hoc probability distribution that matches the observed DES5YR redshift distribution. We find that sampling with a probability that increases directly proportional to redshift (up until $z = 0.7$) matches the data well. We validate this z distribution in Section 4.

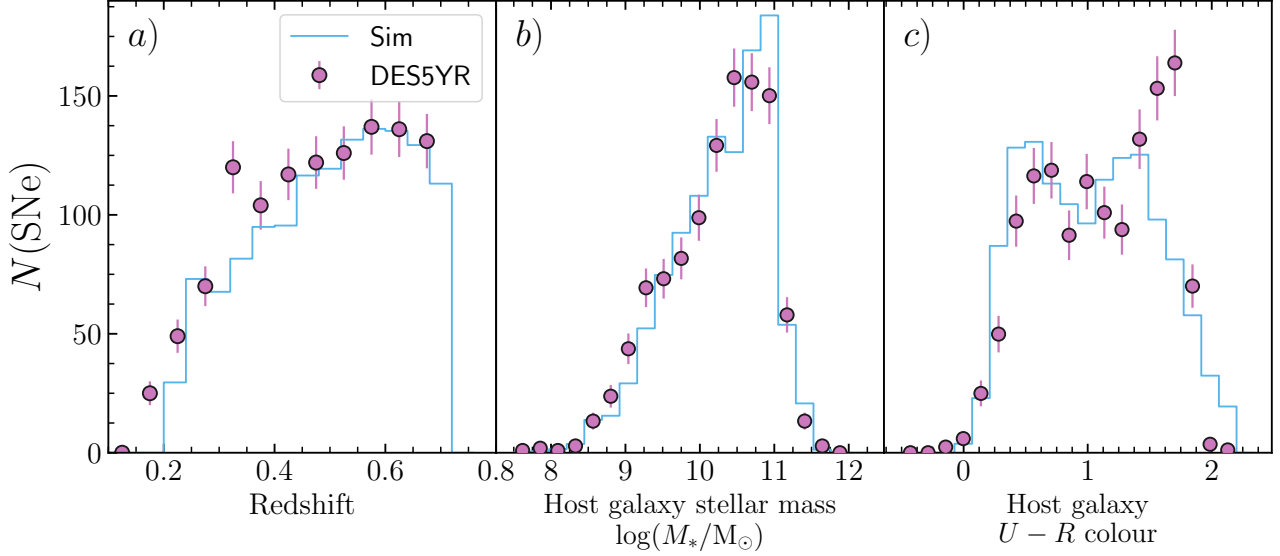


Figure 4. Population distributions of observed DES5YR SNe Ia compared to SNe Ia samples simulated in this work, in redshift, host galaxy stellar mass, and host galaxy rest-frame $U-R$ colour. Counts from the simulation have been scaled to match the total number of DES5YR SNe.

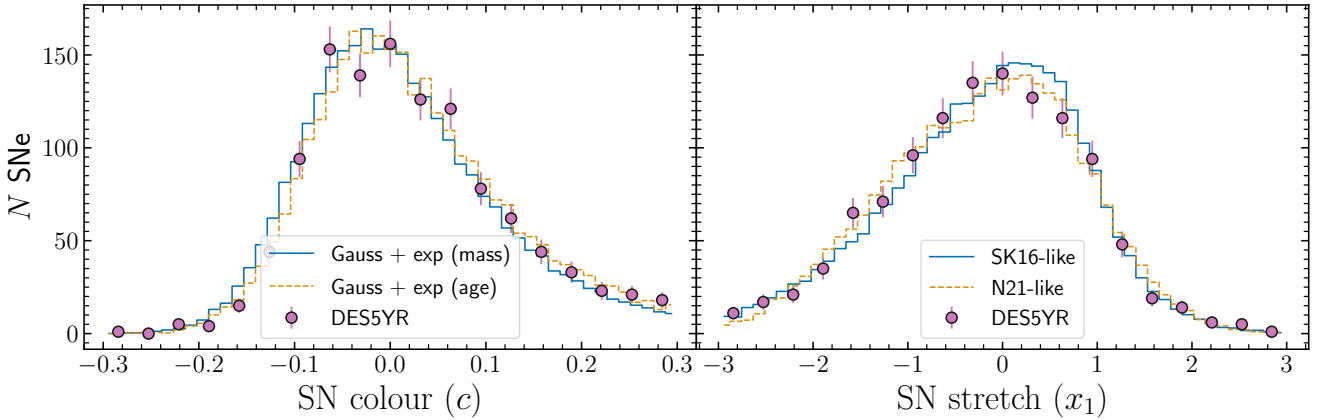


Figure 5. *Left:* distribution of SN SALT2 colour c for DES5YR and simulations drawing c from an intrinsic Gaussian and an exponential dust distribution. *Right:* distribution of SN SALT2 stretch x_1 for DES5YR and simulations drawing from an asymmetric Gaussian (yellow dashed) as per Scolnic & Kessler (2016) and a double Gaussian based on SN progenitor age (red solid) as per N21.

3.4.2 Supernova uncertainties

Without fully simulating SN light curves and their subsequent observation by the survey, we must approximate the uncertainties on our SN light curve parameters. We simulate uncertainties σ_{m_B} on m_B by approximating the relationship between $\sigma_{m_B, \text{obs}}$ and $m_{B, \text{obs}}$ in the DES5YR data. The uncertainty on x_1 (σ_{x_1}) is then simulated according to σ_{m_B} following a linear least-squares fit to the observed DES5YR σ_{m_B} and σ_{x_1} . The same procedure is followed for the colour uncertainty, σ_c . A full description of the uncertainties is presented in Appendix A.

3.4.3 Bias corrections

With simulated brightnesses and light curve parameters, their respective uncertainties, as well as a redshift distribution that all match the data, the simulation mimics the selection effects and observational biases inherent in the data. We therefore calculate the distance modulus to each simulated SN Ia using an identical method to K22 using the BEAMS with Bias Corrections (BBC; Kessler & Scolnic 2017) framework. BBC follows the method of Marriner et al. (2011), a χ^2 minimization of Hubble residuals binned in redshift. The BEAMS part of BBC is not used, as we do not deal with contamination from core-collapse SNe and all simulated SNe are assigned as SNe Ia with probability 1.0. In BBC, we fit for Eq. 3. We do not fit for the additional host step terms γ_X , as we investigate how these terms evolve as a function of SN colour in this work.

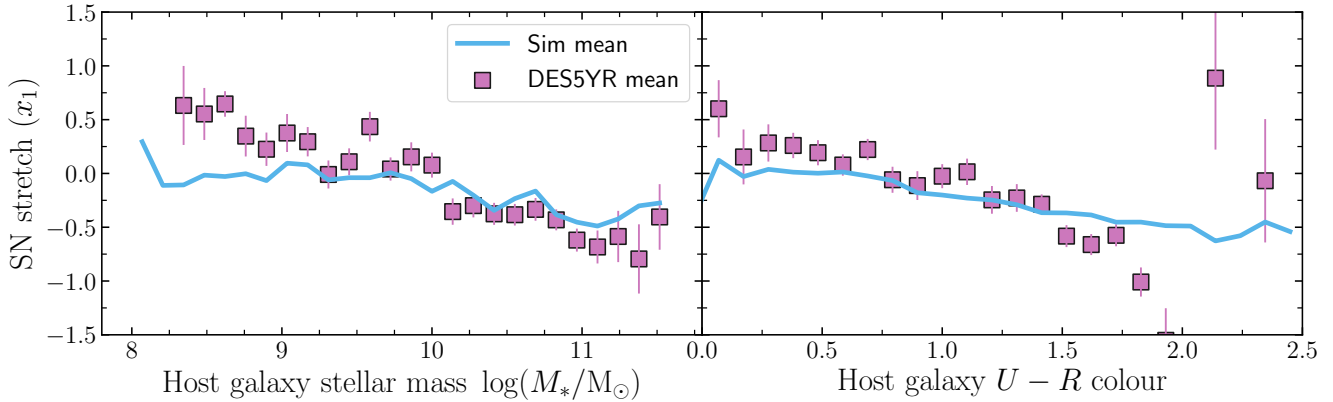


Figure 6. *Left:* distribution of SN x_1 as a function of host galaxy stellar mass for DES5YR data and simulations using the N21 double Gaussian x_1 model based on SN progenitor age; *Right:* as left, but for host galaxy $U - R$ colour.

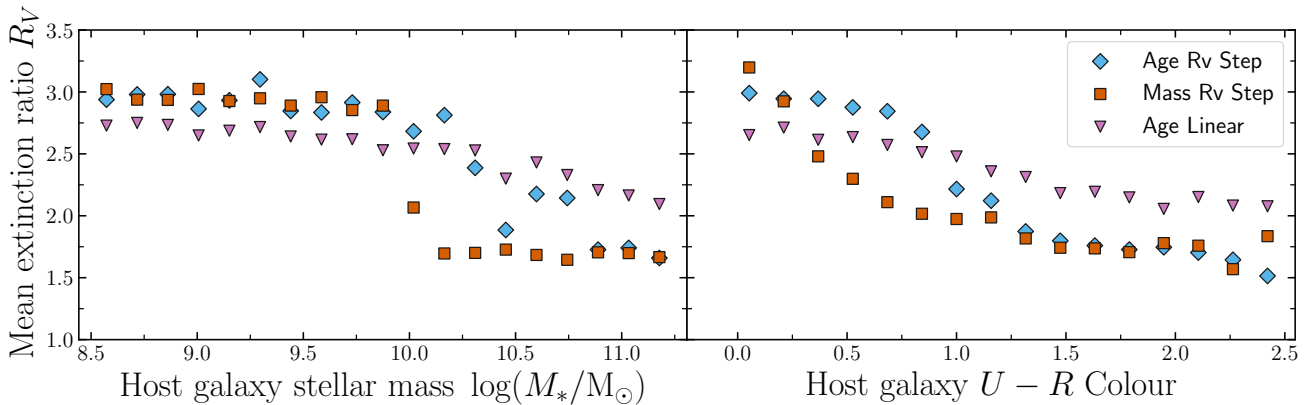


Figure 7. *Left:* Relationship between total-to-selective dust extinction R_V and host galaxy stellar mass for the three models used in this work. *Right:* As left, but shown against host galaxy $U - R$ colour.

The simplest form of bias correction (μ_{bias} ; equation 3) implemented in BBC is dependent on redshift only: a one-dimensional correction (‘BBC-1D’). However, [Scolnic & Kessler \(2016\)](#) showed that the asymmetric nature of the c and x_1 distributions leads to distance biases as a function of those parameters, which must be corrected as a function of α , x_1 , β , and c in addition to z (a ‘BBC-5D’ bias correction). In implementations of such 5D corrections, the treatment of relationships between x_1 and host stellar mass must also be carefully accounted for to avoid introducing subtle biases ([Smith et al. 2020](#)). [Popovic et al. \(2021b\)](#) introduced a further two dimensions to account for such correlations between light-curves and host galaxies (‘BBC-7D’), as well as a method that assumes the BS21 colour model (‘BBC-BS20’).

In this work we use the 1D (redshift only) bias correction to validate and compare models, and also compare our best models to the data with the BBC-BS20 corrections. We include a x_1 -host mass correlation in the bias correction simulation as per [Vincenzi et al. \(2021\)](#), and for BBC-BS20 we use the best-fit R_V values from BS21 in the bias correction simulation. Both approaches require that the simulation adequately includes the same observational biases and selection

effects that have been introduced onto the data, which although not guaranteed has been well validated in the following section. The average and range of bias correction magnitudes for both cases are shown in Fig. 3 and show that, in particular for the BBC-1D case, the corrections are small compared to the step size in the redshift range considered.

3.5 Baseline model

For the validation of the model (Section 4), we introduce baseline simulations on top of which the dust and luminosity step parameters are subsequently varied (Section 5). For the BS21 model we fix the x_1 and c population parameters similarly to those used in BS21. We use the $x_{1,\text{sim}}$ parameters from [Scolnic & Kessler \(2016\)](#), while for c_{sim} we use the values in Table 2. The parameters for both $x_{1,\text{sim}}$ and c_{sim} used in our baseline age-based model are also presented in Table 2. Dust parameters R_V are held constant at 2.75, although R_V is not important for the validation presented below. The width-luminosity coefficient is fixed at $\alpha_{\text{sim}} = 0.15$ and the colour-luminosity coefficient has mean $\bar{\beta} = 1.98$ and standard deviation $\sigma_{\beta} = 0.35$

4 VALIDATION

Before presenting the results, we describe the metrics used to validate the simulations introduced in Sections 2 and 3. We validate our model by comparing it to the DES5YR SN Ia data presented in K22. A set of qualitative and quantitative validations are outlined as follows. In the following sections, we define the chi-squared statistic on a histogram with both the data and simulation in equivalent bins:

$$\chi^2 = \sum_i \frac{(N_i^{\text{data}} - N_i^{\text{sim}})^2}{e_i^2}, \quad (20)$$

where N_i^{data} is the count of the observed data in the i th bin, N_i^{sim} is the equivalent for the simulation, and e_i^2 the error on the data approximated by $\sqrt{N_i^{\text{data}}}$. While this estimate of the error does not take into account uncertainty on the measurements themselves, doing so would require a further level of resampling that would render the procedure computationally prohibitive. For goodness of fit and model comparison, we also employ the reduced chi-squared:

$$\chi_v^2 = \chi^2 / N_{\text{bins}}. \quad (21)$$

4.1 Host galaxies

We validate the host galaxy simulations by investigating the one-dimensional distributions of redshift, stellar mass (M_*), and rest-frame $U - R$ (Fig. 4). The overall trend of the redshift distribution is recovered by construction, as outlined in Section 3.4.1. The distribution of observed M_* is well-replicated. The observed distribution of $U - R$ is bimodal, and this general behaviour is recovered by the model, although the peak of red galaxies in the simulation lies around $U - R \sim 1.5$, compared to the data which peak closer to 1.8. This difference could be due to a number of reasons, such as the simplistic treatment of metallicity and an incomplete implementation of galaxy quenching. While not perfect, we proceed satisfied that both the observed and simulated distributions can be separated at $U - R = 1$. Since the Hubble residual steps tested using this split point, differing shape of the red peak of the distribution should not affect the results significantly.

We further validate this argument by comparing the observed and simulated correlations between M_* and $U - R$ (Fig. 1): the simulations accurately replicate the positive trend between $U - R$ and M_* , as well as the characteristic step from majority blue to majority red galaxies as their mass increases beyond $\log(M_*/M_\odot) = 10$. This qualitative validation reinforces trust that our galaxy evolution model reproduces, on a general level, the complex relationships between galaxy stellar mass, age, and colour, and that the DTD–SFH convolution selects a representative sample of SN Ia hosts from the global galaxy population.

4.2 Light curve parameters

We next validate the simulated distributions of SN x_1 and c , and their relationships with host galaxy attributes M_* and $U - R$. The simulated distributions are shown in Fig. 5, and the parameters used for the baseline age-based model are presented in Table 2.

The c distribution is well-matched by the ‘Gaussian plus dust reddening’ model. We find the data are described well ($\chi_v^2 = 0.95$) using similar parameters to BS21 for both the intrinsic Gaussian and exponential tail, with the mean of the colour distribution shifting marginally bluer. We find a single value of $\tau_E = 0.135$ describes the

data well, which implies there is no strong difference in the mean reddening for SNe in different galaxies, whether they are split by their stellar mass or their age.

Both the asymmetric Gaussian model and Gaussian mixture model provide good fits to the x_1 distribution, with the mixture model offering a slightly improved fit around the peak. The best-matching parameters for the mixture model are similar to those from N21, with the only significant difference being the mixture coefficient a_{x_1} , for which N21 report ~ 0.5 while we find 0.38. In our model, the value of a_{x_1} is degenerate with the choice of age threshold τ_A , an investigation into which is left for future work. There is a degree of overfitting, with the best fit model having $\chi_v^2 = 0.65$, due to the large uncertainties on measured x_1 in the SN Ia data.

The relationship between x_1 and host galaxy properties is shown in Fig. 6. The N21 model implemented in our simulations reproduces the general trends in the data, but does not achieve the same strength of relation, despite the one-dimensional x_1 distribution being adequately replicated (Fig. 5). The likely cause is the inconsistency between the simulated and observed $U - R$ distribution, due to the under-representation of passive/red galaxies in the simulation. The underprediction of the slope of the x_1 versus M_* relation has consequences in the recovery of the α parameter (Dixon 2021; Rose et al. 2021).

5 RESULTS

The principle aim of this work is to explain the observed trends between Hubble residual and SN c , with the sample divided by both host galaxy M_* and host galaxy $U - R$ colour. In this section, we present the results of simulating a sample of SN Ia host galaxies (Section 2), light curve properties (Section 3.2), and fitting their Hubble residuals (Section 3.4.3) using different assumptions of the generative models for x_1 and R_V and various forms of intrinsic luminosity step.

For each of our model variations, all of the parameters defining the population distributions are held fixed apart from three parameters of interest, for which we simulate a coarse grid. These parameters are the two extinction law R_V parameters and the luminosity step size. We vary $\overline{R_{V,1}}$ between [1.5 – 2.75] in steps of 0.25, $\overline{R_{V,2}}$ in the range [2.5 – 3.75] with steps of 0.25, and γ_X between [0 – 0.2] with steps of 0.05. Future analyses will employ inference techniques to precisely constrain the best fit parameters of these models for their use in cosmological analyses, as has been performed for the BS21 model in Popovic et al. (2021a). The relationship between simulated SN R_V as a function of host galaxy stellar mass and $U - R$ colour is shown in Fig. 7, for a representative simulation with $\overline{R_{V,1}} = 1.75$, $\overline{R_{V,2}} = 3.0$. It is clear that an R_V step in age is better recovered by $U - R$ colour than stellar mass, and that stellar mass steps are not particularly well recovered by splitting colour at $U - R = 1$, highlighting the need for both diagnostics.

Each model is compared to the Hubble residuals measured in the DES5YR data of K22 split by M_* and $U - R$. A summary of the results of our simulations and comparison to the data is presented in Tables 3 & 4.

Since SNe are simulated stochastically by sampling from a number of interlinked distributions and a numerical host galaxy library, an analytical evaluation of the relationship between Hubble residual and c in the model is not possible. We thus compare the simulated and observed trends in a similar way to the one-dimensional parameter distributions (Section 4.2).

For each bin i in a host galaxy property X , we group simulated

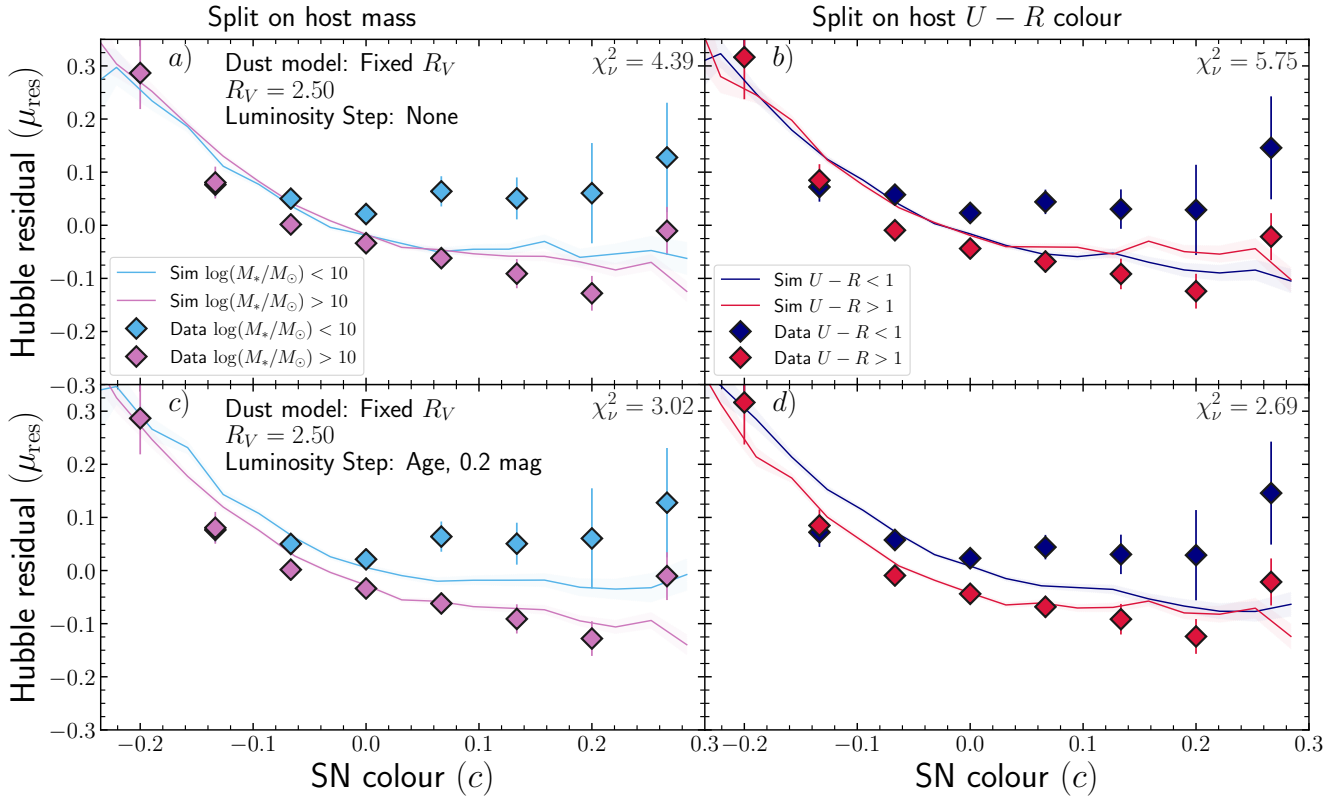


Figure 8. SN Ia Hubble residual versus SN colour, with the DES5YR SN Ia sample shown as the data points and the simulations shown as the solid lines. Left hand panels show SNe split by their host stellar mass at split at $\log(M/M_*) = 10$, and right hand panels show SNe split by their host colour at $U - R = 1$. Panel *a*) shows simulations from the baseline model with R_V fixed across all hosts and no intrinsic luminosity step. Hubble residuals for data and simulations have been measured with BBC-1D. *b*): as the left panel, but for the data and simulations split at a host galaxy rest-frame $U - R$ colour of $U - R = 1$. The χ^2_{ν} statistic for both comparisons is shown. *c*): as *a*) but for simulations with an additional 0.2 mag intrinsic luminosity step at a SN progenitor age of 0.75 Gyr. *d*): as *c*) but with SNe split by their host galaxy colour.

Table 3. Dust population parameters for the models presented in this work when compared to DES5YR data split at host stellar mass $\log(M_*/M_\odot) = 10$ (left hand panels of Figs. 9 & 10.) Y represents the host galaxy or SN property used to determine R_V , while X is the property on which an intrinsic luminosity step (of magnitude γ_X ; Eq. 6) is placed.

Model Name	BiasCor	$\overline{R_{V,1}}$	$\sigma_{R_{V,1}}$	$\overline{R_{V,2}}$	$\sigma_{R_{V,2}}$	Y	Y_{split}	Y_1	Y_2	X	X_{split}	γ_X	$\chi^2_{\nu, \mu_{\text{res}}}(M_*)$
BS21	1D	1.75	1.0	3.0	1.0	M_*	$10^{10} M_\odot$	-	-	M_*	$10^{10} M_\odot$	0.0	1.63
BS21	1D	1.75	1.0	3.0	1.0	M_*	$10^{10} M_\odot$	-	-	τ_A	0.75 Gyr	0.0	2.03
Age R_V	1D	1.5	1.0	3.0	1.0	τ_G	3 Gyr	-	-	τ_A	0.75 Gyr	0.0	1.94
Age R_V	1D	1.5	1.0	2.75	1.0	τ_G	3 Gyr	-	-	M_*	$10^{10} M_\odot$	0.0	1.86
Age R_V Linear	1D	1.5	1.0	2.5	1.0	τ_A	-	0.1 Gyr	10 Gyr	τ_A	0.75 Gyr	0.15	2.27

Table 4. Dust population parameters for the models presented in this work when compared to DES5YR data split at host $U - R = 1$ (right hand panels of Figs. 9 & 10.)

Model Name	BiasCor	$\overline{R_{V,1}}$	$\sigma_{R_{V,1}}$	$\overline{R_{V,2}}$	$\sigma_{R_{V,2}}$	Y	Y_{split}	Y_1	Y_2	X	X_{split}	γ_X	$\chi^2_{\nu, \mu_{\text{res}}}(U - R)$
BS21	1D	1.5	1.0	3.0	1.0	M_*	$10^{10} M_\odot$	-	-	M_*	$10^{10} M_\odot$	0.10	1.70
BS21	1D	1.5	1.0	3.0	1.0	M_*	$10^{10} M_\odot$	-	-	τ_A	0.75 Gyr	0.10	1.59
Age R_V	1D	1.75	1.0	2.5	1.0	τ_G	3 Gyr	-	-	τ_A	0.75 Gyr	0.15	1.70
Age R_V	1D	1.5	1.0	2.5	1.0	τ_G	3 Gyr	-	-	M_*	$10^{10} M_\odot$	0.1	1.65
Age R_V Linear	1D	1.5	1.0	2.5	1.0	τ_A	-	0.1 Gyr	10 Gyr	τ_A	0.75 Gyr	0.20	1.99

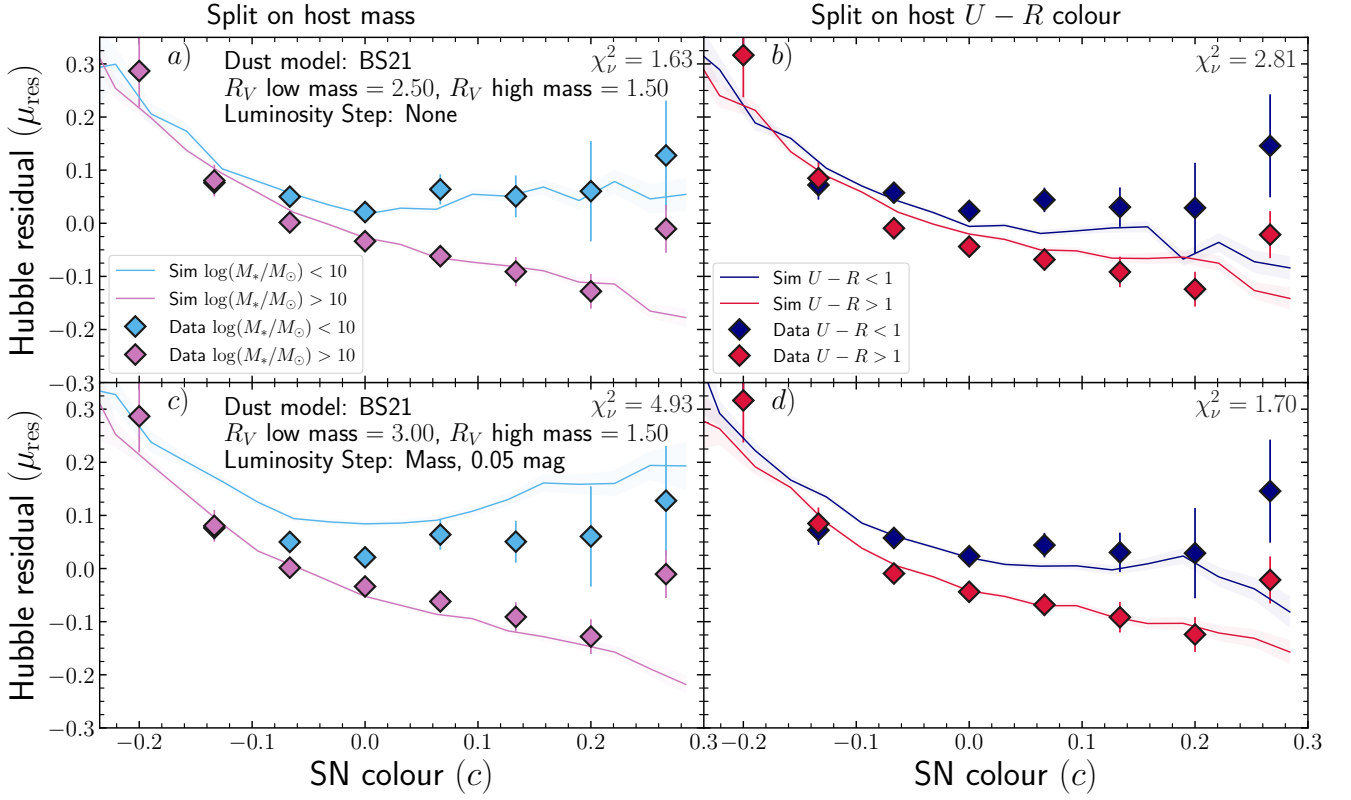


Figure 9. SN Ia Hubble residual versus SN colour, as in Fig. 8, but for models with R_V that varies with host galaxy stellar mass. *a*): Simulations from the BS21 model with $\overline{R_V} = 2.5$ in low mass host galaxies and $\overline{R_V} = 1.5$ in high mass host galaxies. Hubble residuals for data and simulations have been measured with BBC-1D; *b*): as *a*, but for the data and simulations split host galaxy $U - R = 1$. The model that describes the mass-split data well does not reproduce the full difference between Hubble residuals split on $U - R$. *c*) & *d*): as *a* & *b* but for the model parameters that best describe the $U - R$ data (right hand panel): $\overline{R_V} = 3.0$ in low mass host galaxies and $\overline{R_V} = 1.5$ in high mass host galaxies, and an intrinsic luminosity step at $\log(M/M_\odot) = 10$ of size 0.1 mag.

SNe in the same bins c_j as used for the data. The reduced chi-squared is then

$$\chi^2_{\nu, \mu_{\text{res}}} = \frac{1}{N_{\text{bins}}} \sum_i^X \sum_j^{N_{\text{bins}}} \frac{(\overline{\mu_{\text{res, data, } i, j}} - \overline{\mu_{\text{res, sim, } i, j}})^2}{e_{i, j}^2}, \quad (22)$$

where $\overline{\mu_{\text{res, data, } i, j}}$ and $\overline{\mu_{\text{res, sim, } i, j}}$ are the weighted mean Hubble residuals in each host bin i and colour bin j for data and simulations, respectively, and $e_{i, j}$ is the standard error on the means of the data.

Each of the following sections present the results for one model of how R_V varies with host or SN properties. For each model we present the data and simulations (Figs. 8-10) split by their stellar mass (left-hand panels) and $U - R$ (right-hand panels). The values from our grid search that result in the smallest $\chi^2_{\nu, \mu_{\text{res}}}$ for each model when split by stellar mass and $U - R$ are presented in Tables 3 and 4 respectively. Unless otherwise stated, parameters reported below are determined using a 1D (BBC-1D) bias correction. Corresponding figures and tables for BBC-BS20 can be found in Appendix B. We stress that the R_V values here are indicative only, since the generative models for the simulated data and bias corrections are not the same. A thorough fitting via the method of Popovic et al. (2021a) is left for future work.

5.1 Fixed R_V , no luminosity step

We begin by simulating SNe from a baseline model with no luminosity step and $\overline{R_V}$ fixed at 2.5 across all host galaxies. The resulting Hubble residuals shown in Fig. 8 (panels *a*, *b*) display a trend with SN colour c , as expected due to the effects of parameter migration (Scolnic & Kessler 2016), but they do not reproduce the differences between low and high stellar mass, or between blue and red host galaxies. To test whether a simple luminosity step as a function of SN age can explain the diverging Hubble residuals, we add a luminosity age step $\gamma_{\tau_A} = 0.2$ mag to the simulation (Fig. 8, panels *c*, *d*). Again, the diverging Hubble residuals are not reproduced. Notably the simulated 0.2 mag luminosity step is not recovered in the Hubble residuals, with a step of only around 0.05 mag appearing. The reasons behind this are discussed in Section 6.3.

5.2 R_V split on stellar mass (BS21)

Motivated by the clear signal in the data of diverging Hubble residuals as a function of SN c , we continue by implementing the model of BS21. Our BS21 simulation is shown in Fig. 9 (upper panels), and qualitatively replicates the trend in the data when split by stellar mass. Simulations with $\overline{R_V} = 3.0$ in low mass galaxies and 1.75 in high mass hosts match the data well ($\chi^2_{\nu} = 1.63$), values that are consistent with those found by BS21 and Popovic et al. (2021a).

However, this model that describes the mass-split data adequately

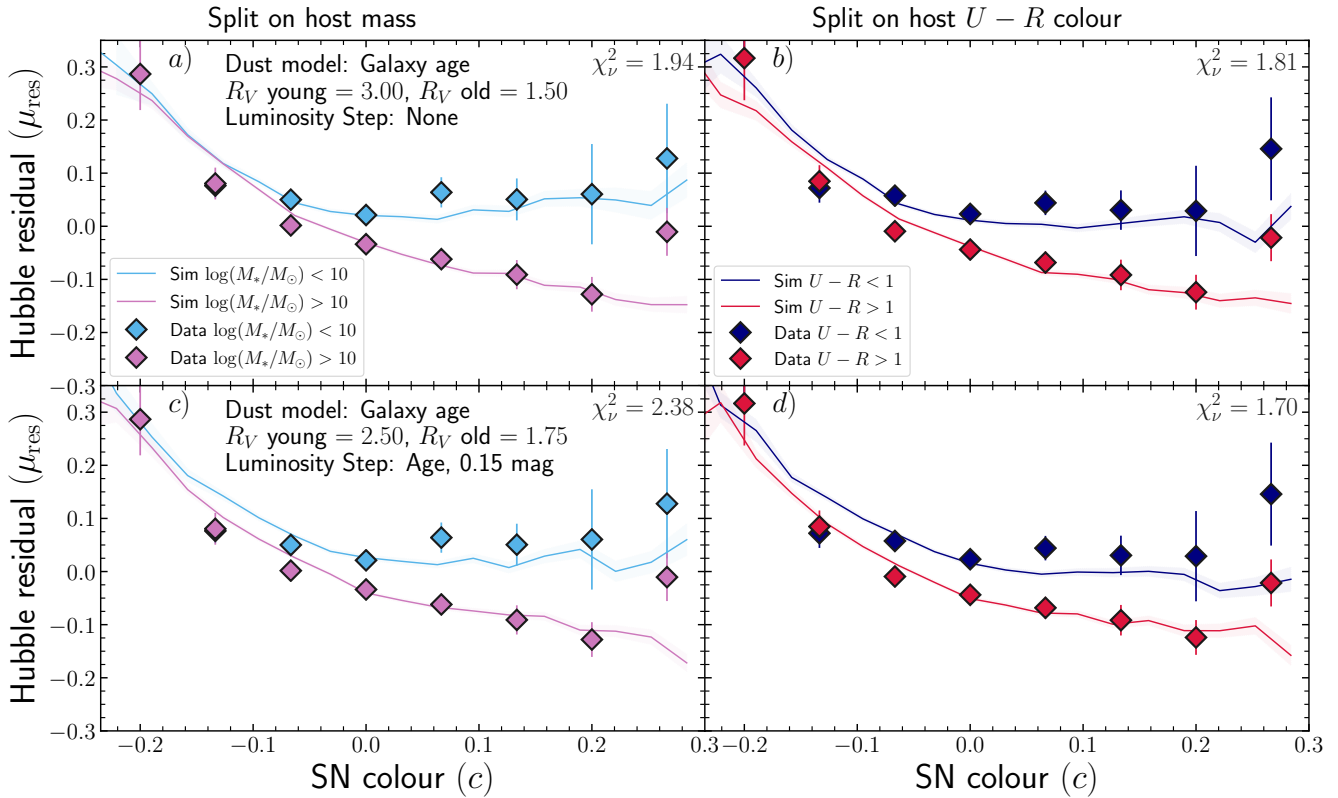


Figure 10. SN Ia Hubble residual versus SN colour, as in Fig. 9 but for the model with R_V changing with galaxy age. *a)*: The model that best fits the SN data when split by stellar mass includes $\overline{R_V} = 3.0$ in young hosts, $\overline{R_V} = 1.5$ in old hosts and no intrinsic luminosity step. *b)*: as *a)*, but showing SNe split at host galaxy $U - R = 1$. *c)* & *d)*: as for *a)* & *b)* but for the best fitting model to the data split by host $U - R$ colour: $\overline{R_V} = 2.5$ in young hosts, $\overline{R_V} = 1.75$ in old hosts, and a 0.15 mag intrinsic luminosity step at a progenitor age of 0.75 Gyr.

performs less well when the data are split by host $U - R$ colour ($\chi^2_\nu = 2.81$), and neither the difference between Hubble residuals in blue nor red SNe are matched as well. The best match ($\chi^2_\nu = 1.70$), shown in the lower panels of Fig. 9, has a steeper extinction law in high mass galaxies with $\overline{R_V} = 1.5$, as well as the addition of an intrinsic 0.1 mag luminosity step on stellar mass. However, the addition of this step has the effect of drastically reducing the quality of the match when the data are split by stellar mass ($\chi^2_\nu = 4.93$).

5.3 R_V split on galaxy age

The BS21 model is able to reproduce the observed trends between Hubble residual and SN c when split by stellar mass but the model requires different parameters and an intrinsic luminosity step when matching the data split by host $U - R$ colour. Motivated by a hypothesis that dust parameters could be driven by stellar age rather than galaxy mass, we implement the ‘Age R_V ’ model, which is identical to the BS21 model but with R_V described by a step function at a mean stellar age of 3 Gyr rather than $\log(M/M_\odot) = 10$. This model also incorporates the age-based Nicolas et al. (2021) model of x_1 .

As with the mass-based model, the better matching simulations are those with differing R_V in young and old galaxies, with larger differences in $\overline{R_V}$ required compared to the BS21 model. When splitting Hubble residuals on stellar mass (Fig. 10 upper panels, Table 3), the favoured models ($\chi^2_\nu = 1.94$) include a similar difference in R_V between old (1.5) and young (2.75) hosts and no intrinsic

luminosity step is required. When splitting the data on $U - R$ this model provides a significantly better match than the mass- R_V -step model of BS21 (when implemented with no further luminosity step, i.e. Fig. 10b has $\chi^2_\nu = 1.70$ compared to Fig. 9b with $\chi^2_\nu = 2.44$). The best match to the $U - R$ split data has a smaller difference between R_V values (1.5 and 2.5) plus additional 0.1 mag intrinsic luminosity step on stellar mass, or an even smaller R_V difference (1.75 and 2.5) and a 0.15 mag step on stellar age (Fig. 10, lower panels). As described in the results of the previous models, a much smaller step is evident in the Hubble residuals than the age step input into the SNe. We address this effect in Section 6.3.

5.4 R_V linear with SN age

Modelling R_V as a linear function of SN progenitor age is not as successful as the previous step-based models, with consistently higher χ^2_ν (Tables 3 & 4). This R_V model is the only one that also requires an age-luminosity step to explain both the mass-split and $U - R$ split data, with $\gamma_X \geq 0.15$ mag.

6 DISCUSSION

The results of our comparisons of simulations to data support the notion that the SN colour–luminosity relation in SNe Ia is not linear with SN colour, consistent with the model of BS21 where the dust

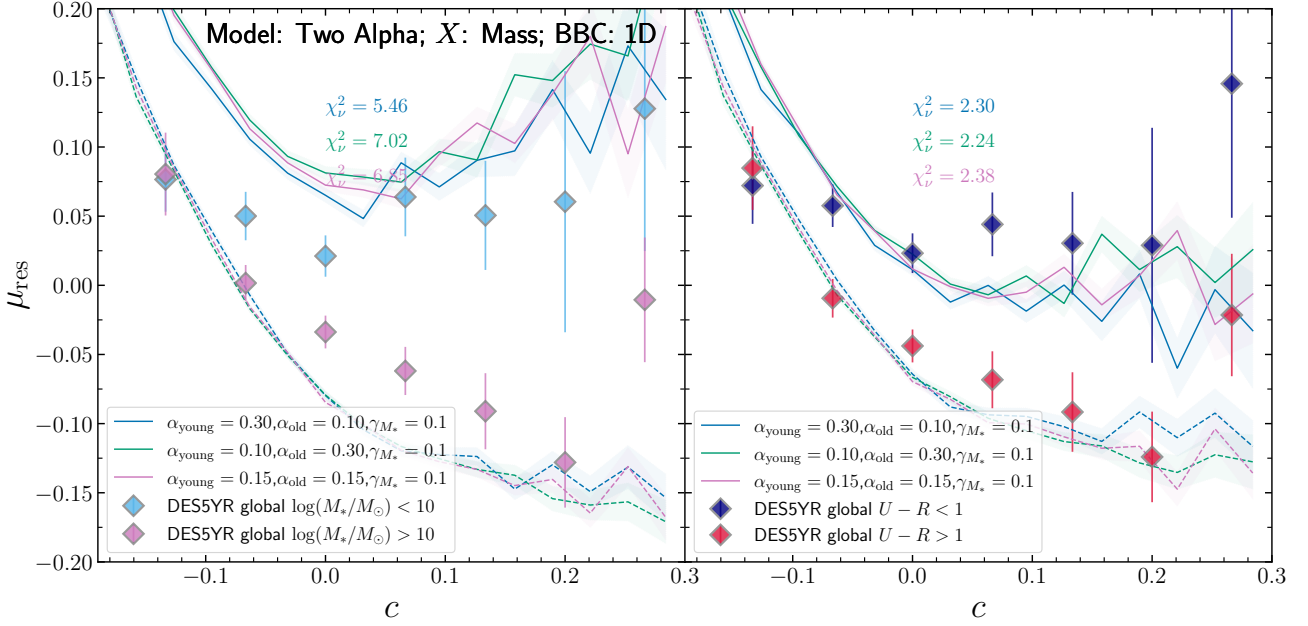


Figure 11. *Left:* As Fig. 10 with fixed R_V ($R_{V,1} = 2.0$; $R_{V,2} = 3.75$, $Y = \tau_G$), an intrinsic mass step $\gamma_{M_*} = 0.1$ mag, and varying α_{sim} for young and old populations of SNe. There is no significant difference in the trend of HRs with c , nor the size of the measured mass step, despite large variations in α_{sim} ; *Right:* as left, but for host galaxy $U - R$ colour. As with mass, there is no significant difference in the HRs despite large variations in α_{sim} and the recovered $U - R$ step is not affected.

extinction law along the line-of-sight to SNe Ia correlates with the global properties of host galaxies. The reason for, and best tracer of, the difference in R_V is less clear and is discussed in the following sections. As with the previous section, the discussion here focuses on BBC-1D bias corrections, but the results are consistent when using BBC-BS20 bias corrections.

6.1 R_V in young, low mass galaxies

We find no significant difference in the fit quality when the R_V changes with stellar mass or stellar age, the two of which are themselves strongly correlated. When implementing the BS21 model we find $\overline{R_{V,2}}$ is broadly consistent with the results from BS21 and Popovic et al. (2021a) when splitting on stellar mass. When R_V changes with stellar age the simulations are more consistent with the data across both stellar mass and $U - R$, but $2.5 \leq \overline{R_{V,2}} \leq 3.0$ are generally smaller than the 3.0 in Popovic et al. (2021a). These values for R_V consistent with those typically measured in SN Ia light curves and spectra for similar hosts, with Cikota et al. (2016) finding R_V values of 2.71 ± 1.58 in spiral galaxies.

6.2 R_V in old, high mass galaxies

In all of the models trialled the R_V value in high mass/old galaxies ($\overline{R_{V,1}}$) is required to be significantly lower than that for low mass/young galaxies in order to replicate the trends of Hubble residual against SN colour. Generally, lower values are preferred when the data are split on $U - R$. Values in the range $1.5 \leq \overline{R_{V,1}} \leq 1.75$ indicate the dust in older, more massive SN hosts is composed mainly of small grains leading to a steep extinction curve. How this observation relates to studies of the general galaxy population is complicated

by different relationships seen between dust and different galaxy descriptors. While Reddy et al. (2018) and Salim et al. (2018) show that R_V increases with stellar mass for star-forming galaxies, Salim et al. (2018) also show that R_V is lower in passive (and starburst) galaxies than those on the star forming main sequence, and that the R_V of star-forming galaxies is more closely related to its overall extinction A_V than its stellar mass. However, the quiescent galaxies in Salim et al. (2018) still have an average R_V of 2.61, far greater than that found here and indicating that the SNe do not trace the average dust composition in the galaxy. We note however that the SN R_V measurements are somewhat dependent on the assumed SALT2 colour law, rendering direct comparisons less informative.

6.3 The effect of intrinsic luminosity steps

While models of an R_V step in stellar mass or galaxy age provide reasonable matches to the Hubble residuals split by either stellar mass or host $U - R$ colour, they do not fully match both host properties simultaneously – different R_V steps are required to match mass and to match $U - R$. Furthermore, K22 show that there are residual steps in the data in one parameter (e.g. stellar mass) if the effect of another (e.g. $U - R$) is subtracted. We find that in general the step seen in our simulated HRs is smaller in $U - R$ than in M_* , but that step is rarely fully removed using R_V only and that adding an intrinsic luminosity step is required. However, such a step introduces a number of its own issues.

Firstly we consider adding a step based on the SN age. We find that adding a 0.2 mag step has only a marginal effect on the HRs – instead, the best fit value α is much smaller than the simulated value α_{sim} . This effect is caused in this model by the dependence on age of both the intrinsic step and x_1 , which act on the SN luminosity in the

opposite direction to each other. Such an affect has been observed before (e.g. [Rose et al. 2021](#)), who suggest that steps γ_X should be fit simultaneously with the other nuisance parameters. However, with the step size evolving with SN colour, this is not trivial.

An intrinsic mass step is affected much less absorbed by α , with a 0.1 mag step able to explain the $U - R$ data for most R_V model and bias correction combinations (e.g. [Fig. 9](#)). However, this means that the discrepancy between models with an intrinsic mass step and the data split by mass is much larger than when the intrinsic step is on age and no particularly good combinations are found.

6.3.1 Varying the width-luminosity coefficient α

Neither an intrinsic luminosity step based on SN age or host mass are able to simultaneously model the data split by stellar mass and by $U - R$. One logical explanation for this is that the relationships between SN age, x_1 , and α are linked with different strengths to the host parameters M_* and $U - R$. The $U - R$ colour traces age more directly than M_* , and with the bimodal [N21](#) model x_1 is strongly linked to SN age, meaning intrinsic steps in luminosity are washed out by being absorbed into the alpha parameter, and this occurs more strongly in the age-like $U - R$ than M_* which is less related to SN age. To counter this effect, we hypothesise that the two x_1 populations could follow different intrinsic values of α while simultaneously having different mean absolute magnitudes.

To test the two- α hypothesis we run a further set of simulations. We fix R_V to those from the best Age- R_V model as measured with BBC-ID and as compared to data split by mass: ($R_{V,1} = 1.5$, $R_{V,2} = 3.0$). We simulate over a new grid, varying α_{sim} from 0.05 to 0.30 for the young and old populations, as well as the intrinsic luminosity step γ_X . The combinations of α_{sim} and γ_X that we investigate are constrained by the fitted α . For X , we trial both mass and age steps between zero and 0.25 mag.

The results of varying α_{sim} are summarised in [Fig. 11](#), for which the models included a 0.1 mag step on M_* . Varying α_{sim} for the young and old SNe makes no difference to the inferred step in either M_* or $U - R$ and does not solve the discrepancy between the size of recovered steps – whatever the values of α_{sim} , the resulting HRs are split much more strongly by M_* than by $U - R$. The same results were found when including a 0.2 mag luminosity step on SN age. These simulations thus indicate a universal value for α even if the SNe are divided into two populations of x_1 . Such a lack of bimodality for α is consistent with previous results, e.g. [Sullivan et al. \(2011\)](#).

6.3.2 Reconciling mass and $U - R$ steps

With multiple α values improbable, we investigate whether the discrepancy is caused by limitations in our model. For example, the smooth SFHs lead to a single Gaussian distribution of $U - R$ while the data show bimodality, while the simulation also slightly underestimates the number of hosts in the stellar mass range $9 \leq \log(M/M_\odot)$. Although [Fig. 1](#) shows that the models do reproduce the $M_* - U - R$ relation in general, there may be subtle second order effects introducing inconsistencies for individual simulated SNe. Such effects include the treatment of metallicity which is fixed at solar in the simulation but evolves over both the mass and redshift range of the data (e.g. [Tremonti et al. 2004](#); [Zahid et al. 2014](#)), the lack of bursts of star formation in the simulation, the choice of BC03 templates for the simulation, and incorrect modelling of survey selection effects. These effects will be investigated in subsequent work.

A second possibility is that the inconsistencies between the model

and data are caused by physical effects that have not been fully incorporated or accurately implemented in the model. The relationships between the R_V along a SN sight line and its host galaxy properties are evidently more complicated than a simple step function of galaxy age or stellar mass – when dealing with integrated galaxy attenuation, R_V is inversely correlated with A_V , which we have not included in the model for SN R_V . Meanwhile, A_V correlates strongly with stellar mass ([Zahid et al. 2013](#)), meaning the R_V of galaxies also increases with stellar mass ([Salim et al. 2018](#)). On the contrary, the results of [BS21](#), [Popovic et al. \(2021a\)](#), and this work, indicate that R_V on SN sight lines decreases with stellar mass, suggesting that SN sight lines are systematically different than those of the integrated galaxy light. Meanwhile, [Salim et al. \(2018\)](#) also show that galaxies with old stellar populations have lower R_V than those with young populations, which is in line with the age – R_V steps found here. Because our models track individual stellar populations, it is possible to use our method to compare the predictions to measurements from integral field spectroscopy (e.g. [Galbany et al. 2014, 2016, 2018](#)) to measure dust properties in the very local environments of SNe, which we defer to future work.

Our model has no implementation of a relation between galaxy mass and metallicity. Metallicity is strongly correlated with stellar mass and more weakly with SFR (e.g. [Yates et al. 2012](#)), is tied to the SFH ([Bellstedt et al. 2020](#)) as well as affecting the strengths of nebular emission lines (e.g. [Kewley & Ellison 2008](#)), themselves affecting the $U - R$ colour. If there is an intrinsic correlation between metallicity and SN luminosity (e.g. [Höflich et al. 1998](#); [Kasen et al. 2009](#); [Moreno-Raya et al. 2016](#)) it may affect the observed mass and age steps in different amounts, but be hidden amongst the more dominant R_V effects.

7 CONCLUSIONS AND FUTURE WORK

This work presents a host-galaxy oriented framework for simulating populations of SNe Ia in a cosmological context. We trace stellar population ages throughout the build up of stellar mass of galaxies, and use this information combined with the SN Ia DTD from [W21](#) in order to associate SNe to hosts at realistic rates. The resulting host galaxy library is made available to the community for use in future cosmological SN Ia simulations². SNe are then simulated according to their host galaxy properties as highlighted below:

- light curve width x_1 is drawn from a two-population age-based model based on [N21](#), which accurately reproduces the x_1 distribution. The x_1 vs stellar mass and x_1 vs $U - R$ relations are also relatively well modelled, although the strength of the relation is stronger in the data than we recover in the model.
- we that 68% of SNe in old environments belong to the low-stretch mode. Further work is necessary in order to determine the cause of these two modes, whether they are related to the progenitor scenario, white dwarf composition, or explosion mechanism, and why the transition occurs around 0.75 Gyr;
- SN colour c is well described by a combination of an intrinsic Gaussian and a supplementary exponential distribution attributed to dust reddening. The DESSYR data is best modelled by very similar values to those for the Pantheon+ dataset ([Popovic et al. 2021a](#)).

By running simulations and data through the same BBC framework

² https://github.com/wisemanp/des_sn_hosts/tree/main/simulations/data

we obtain distance estimates and compare the evolution of the Hubble residuals with c . Our results support the findings of BS21 and Popovic et al. (2021a) that the extinction law slope R_V changes depending on host galaxy properties. Additionally we find that:

- if galaxy age is the driver of R_V change, then the evolution is different to if it is driven by stellar mass. However the change (with mass) is opposite to that observed in the galaxy population, while the change with age is consistent with the difference between star forming and passive galaxies;
- when Hubble residuals are split based on host stellar mass, the R_V models can fully explain the trends between Hubble residual and c ;
- when the data are split by host $U - R$, the addition of an intrinsic luminosity step to the model slightly improves the fit to the data.
- there is no preference for the intrinsic luminosity step to be based on stellar mass or SN age, but neither are able to simultaneously reproduce the data when split by host stellar mass and by host $U - R$.
- varying the width-luminosity coefficient α between populations does not solve the discrepancy.

The reason for the discrepancy is therefore either a shortcoming of the simulation or an unmodelled physical effect such as metallicity.

Future work will build upon the host galaxy models developed here. In order to test for effects such as metallicity, the galaxy evolution model needs to be treated in a more complex way by introducing metallicity evolution and bursts of star formation. Such models have been used in the modelling of SN Ia rates and DTDs (e.g. Gandhi et al. 2022) and can be extended to cosmological implementation using the methods outlined in this paper.

SOFTWARE

All software used in this publication are publicly available. The SN and galaxy evolution code can be found at https://github.com/wisemanp/des_sn_hosts. Additionally we made extensive use of numpy (Harris et al. 2020), Astropy (Astropy Collaboration et al. 2018), matplotlib (Hunter 2007), SciPy (Virtanen et al. 2020), and pandas (McKinney 2010).

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DATA AVAILABILITY

The light curves for the DES-SN photometric SN Ia catalogue and associated host galaxy data will be made available as part of the DES5YR SN cosmology analysis at <https://des.ncsa.illinois.edu/releases/sn>.

REFERENCES

- Astropy Collaboration et al., 2018, *The Astronomical Journal*, 156, 123
- Bellstedt S., et al., 2020, *Monthly Notices of the Royal Astronomical Society*, 498, 5581
- Bertelli G., et al., 1994, *Astronomy and Astrophysics Supplement Series*, 106, 275
- Bessell M. S., 1990, *Publications of the Astronomical Society of the Pacific*, 102, 1181
- Betoule M., et al., 2014, *Astronomy & Astrophysics*, 568, A22
- Boquien M., Burgarella D., Roehlly Y., Buat V., Ciesla L., Corre D., Inoue A. K., Salas H., 2019, *Astronomy & Astrophysics*, 622, A103
- Briday M., et al., 2022, *Astronomy & Astrophysics*, 657, A22
- Brout D., Scolnic D., 2021, *The Astrophysical Journal*, 909, 26
- Brout D., et al., 2019a, *The Astrophysical Journal*, 874, 106
- Brout D., et al., 2019b, *The Astrophysical Journal*, 874, 150
- Brout D., et al., 2022, eprint arxiv:2202.04077
- Bruzual G., Charlot S., 2003, *Monthly Notices of the Royal Astronomical Society*, 344, 1000
- Cardelli J. A., Clayton G. C., Mathis J. S., 1989, *The Astrophysical Journal*, 345, 245
- Chabrier G., 2003, *Publications of the Astronomical Society of the Pacific*, 115, 763
- Childress M., et al., 2013, *Astrophysical Journal*, 770, 108
- Childress M. J., Wolf C., Zahid H. J., 2014, *Monthly Notices of the Royal Astronomical Society*, 445, 1898
- Cikota A., Deustua S., Marleau F., 2016, *The Astrophysical Journal*, 819, 152
- DES Collaboration et al., 2018, *Physical Review Letters*, 122, 171301
- Dixon S., 2021, *Publications of the Astronomical Society of the Pacific*, 133, 054501
- Dixon M., et al., 2022, eprint arXiv: 2206.12085
- Flaugher B., et al., 2015, *The Astronomical Journal*, 150, 150
- Galbany L., et al., 2014, *Astronomy & Astrophysics*, 572, 22
- Galbany L., et al., 2016, *Monthly Notices of the Royal Astronomical Society*, 455, 4087
- Galbany L., et al., 2018, *The Astrophysical Journal*, 855, 107
- Gallagher J. S., et al., 2005, *The Astrophysical Journal*, 634, 210
- Gandhi P. J., Wetzel A., Hopkins P. F., Shappee B. J., Wheeler C., Faucher-Giguère C.-A., 2022, eprint arxiv:2202.10477
- Guy J., et al., 2007, *Astronomy & Astrophysics*, 466, 11
- Harris C. R., et al., 2020, *Nature*, 585, 357
- Höfllich P., Wheeler J., Thielemann F., 1998, *The Astrophysical Journal*, 495, 617
- Hunter J. D., 2007, *Computing in Science and Engineering*, 9, 99
- Inoue A. K., 2011, *Monthly Notices of the Royal Astronomical Society*, 415, 2920
- Jha S., Riess A. G., Kirshner R. P., 2007, *The Astrophysical Journal*, 659, 122
- Kasen D., Röpke F. K., Woosley S. E., 2009, *Nature*, 460, 869
- Kelly P. L., Hicken M., Burke D. L., Mandel K. S., Kirshner R. P., 2010, *Astrophysical Journal*, 715, 743
- Kelsey L., et al., 2021, *Monthly Notices of the Royal Astronomical Society*, 501, 4861
- Kelsey L., Sullivan M., Wiseman P., Smith M., Vincenzi M., Scolnic D., Brout D., 2022, *MNRAS*, in prep
- Kessler R., Scolnic D., 2017, *The Astrophysical Journal*, 836, 56
- Kessler R., et al., 2009, *Astrophysical Journal Supplement Series*, 185, 32
- Kessler R., et al., 2019, *Monthly Notices of the Royal Astronomical Society*, 485, 1171
- Kewley L. J., Ellison S. L., 2008, *The Astrophysical Journal*, 681, 1183
- Lampeitl H., et al., 2010, *Monthly Notices of the Royal Astronomical Society*, 401, 2331
- Mandel K. S., Narayan G., Kirshner R. P., 2011, *The Astrophysical Journal*, 731, 120
- Mandel K. S., Scolnic D. M., Shariff H., Foley R. J., Kirshner R. P., 2017, *The Astrophysical Journal*, 842, 93
- Marriner J., et al., 2011, *ApJ*, 740, 72
- McKinney W., 2010, in *PROC. OF THE 9th PYTHON IN SCIENCE CONF.* pp 56–61, doi:10.25080/Majora-92bf1922-00a
- Möller A., et al., 2022, eprint arxiv:2201.11142
- Moreno-Raya M. E., Mollá M., López-Sánchez Á. R., Galbany L., Vílchez J. M., Rosell A. C., Domínguez I., 2016, *The Astrophysical Journal*, 818, L19
- Nicolas N., et al., 2021, *Astronomy & Astrophysics*, 649, A74
- Perlmutter S., et al., 1999, *The Astrophysical Journal*, 517, 565
- Phillips M. M., 1993, *The Astrophysical Journal*, 413, L105
- Popovic B., Brout D., Kessler R., Scolnic D., 2021a, eprint arxiv:2112.04456
- Popovic B., Brout D., Kessler R., Scolnic D., Lu L., 2021b, *The Astrophysical Journal*, 913, 49
- Pskovskii I. P., 1977, *SvA*, 21, 675
- Reddy N. A., et al., 2018, *The Astrophysical Journal*, 853, 56
- Riess A. G., Press W. H., Kirshner R. P., 1996, *The Astrophysical Journal*, 473, 88
- Riess A. G., et al., 1998, *The Astronomical Journal*, 116, 1009
- Rigault M., et al., 2020, *Astronomy & Astrophysics*, 644, A176
- Roman M., et al., 2018, *Astronomy & Astrophysics*, 615, A68
- Rose B. M., Rubin D., Strolger L., Garnavich P. M., 2021, *The Astrophysical Journal*, 909, 28
- Salim S., Boquien M., Lee J. C., 2018, *The Astrophysical Journal*, 859, 11
- Scolnic D., Kessler R., 2016, *The Astrophysical Journal*, 822, L35
- Scolnic D. M., et al., 2018, *The Astrophysical Journal*, 859, 101
- Scolnic D., et al., 2021, eprint arxiv:2112.03863
- Smith M., et al., 2020, *Monthly Notices of the Royal Astronomical Society*, 494, 4426
- Sullivan M., et al., 2010, *Monthly Notices of the Royal Astronomical Society*, 406, no
- Sullivan M., et al., 2011, *Astrophysical Journal*, 737, 102
- Tomczak A. R., et al., 2014, *The Astrophysical Journal*, 783, 85
- Tremonti C. A., et al., 2004, *The Astronomical Journal*, 613, 898
- Tripp R., 1998, *Astronomy and Astrophysics*, 331, 815
- Vincenzi M., et al., 2021, *Monthly Notices of the Royal Astronomical Society*, 505, 2819
- Virtanen P., et al., 2020, *Nature Methods*, 17, 261
- Wiseman P., et al., 2021, *Monthly Notices of the Royal Astronomical Society*, 506, 3330
- Yates R. M., Kauffmann G., Guo Q., 2012, *Monthly Notices of the Royal Astronomical Society*, 422, 215
- Zahid H. J., Yates R. M., Kewley L. J., Kudritzki R. P., 2013, *Astrophysical Journal*, 763, 92
- Zahid H. J., Dima G. I., Kudritzki R.-P., Kewley L. J., Geller M. J., Hwang H. S., Silverman J. D., Kashino D., 2014, *The Astrophysical Journal*, 791, 130

APPENDIX A: SIMULATING UNCERTAINTIES

We simulate uncertainties σ_{m_B} on m_B by approximating the relationship between $\sigma_{m_B, \text{obs}}$ and $m_{B, \text{obs}}$ in the observed data set of K22.

$$\sigma_{m_B} \sim \text{Max} \left\{ \frac{\mathcal{N}(\overline{\sigma_{m_B}}, \sigma_{\sigma_{m_B}})}{0.025} \right\}, \quad (\text{A1})$$

where

$$\overline{\sigma_{m_B}} = 10^{(0.395(m_B - 1.5) - 10)} + 0.03, \quad (\text{A2})$$

where $\overline{\sigma_{m_B}}$ is represented by the dashed line in Fig. A1. The scatter on the uncertainties $\sigma_{\sigma_{m_B}}$ also increases as a function of m_B :

$$\sigma_{\sigma_{m_B}} = \text{Max} \left\{ \frac{0.003(m_B - 20)}{0.003} \right\}. \quad (\text{A3})$$

The uncertainty σ_{m_B} is calculated after the intrinsic m_B has been adjusted for x_1 and c via Eq. 18.

Uncertainties on x_1 and c are estimated by linear least squares fits to the DES5YR $\sigma_{x_1} - \sigma_{m_B}$ and $\sigma_c - \sigma_{m_B}$ relations respectively. The

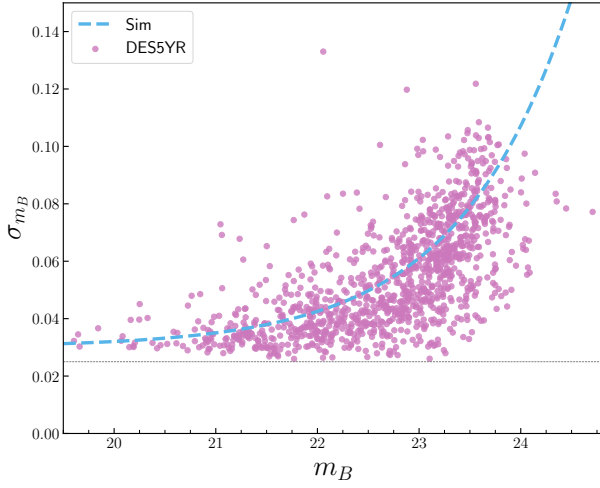


Figure A1. The relationship between peak B -band brightness m_B and its uncertainty, σ_{m_B} . Points are from the DES5YR data and the relationship used in Eq. A2.

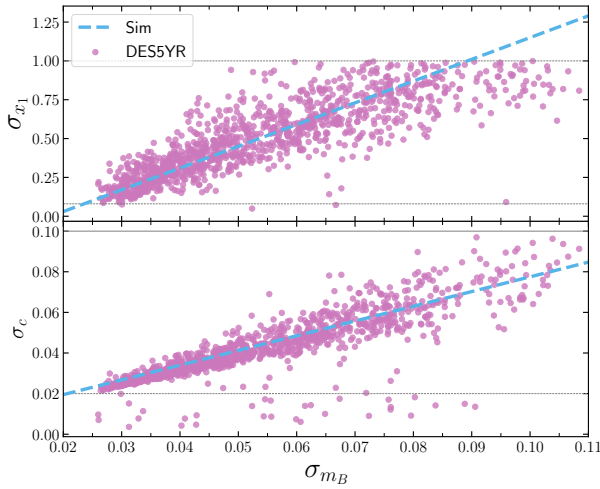


Figure A2. Relationships between σ_{m_B} and σ_{x_1} (upper), and σ_c lower. Points are from the DES5YR data and the relationship used in Eq. A2.

observed and fitted relationships between the uncertainties are shown in Figs. A1 & A2. We use these uncertainties to add noise to m_B , x_1 , and c by drawing them randomly from Gaussian distributions centred at 0 as per Eq. 19.

APPENDIX B: BBC-BS20 BIAS CORRECTIONS

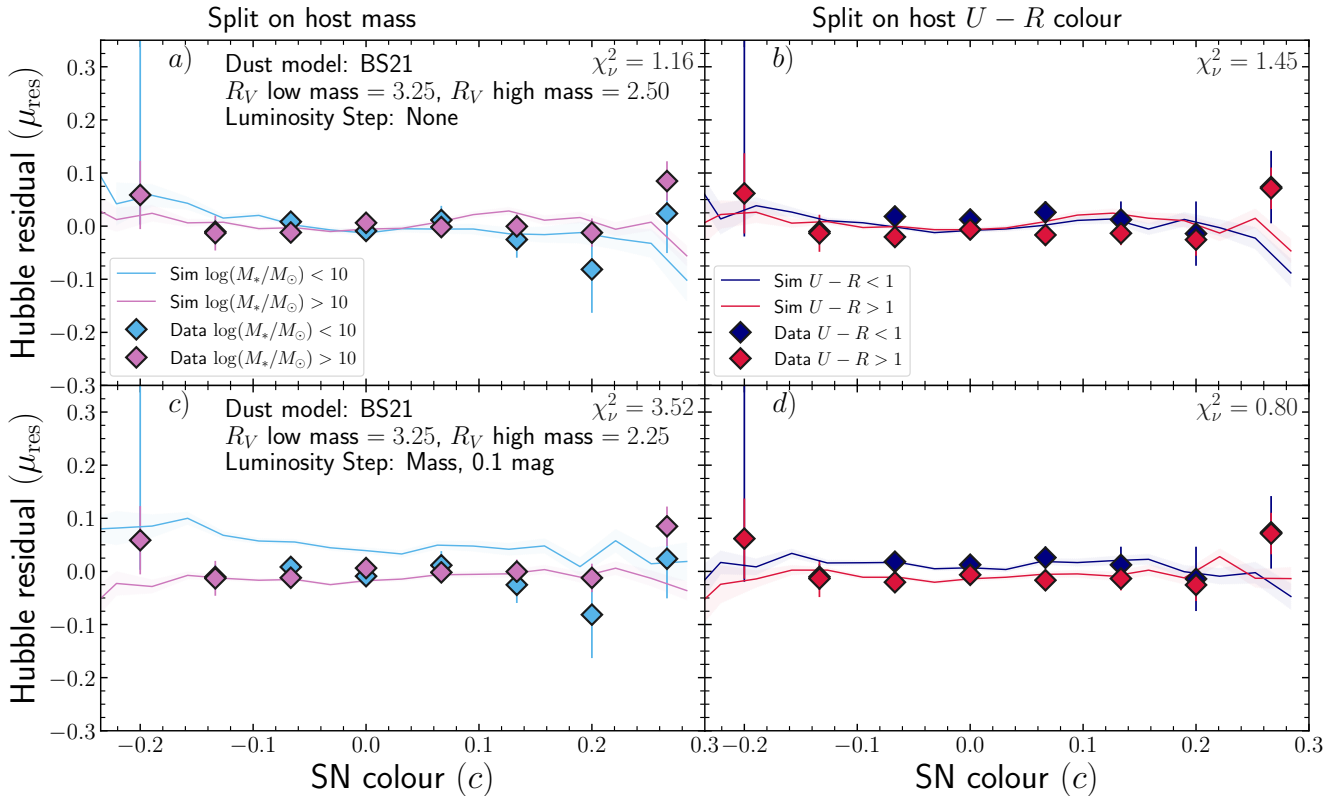
Here we present results in the same way as Section 5 but with Hubble residuals measured using a BBC-BS20 bias correction, which inherently assumes R_V changing with stellar mass. The models shown in Figs. B1-B3 correspond to the best match parameters when splitting on stellar mass (blue) and $U - R$ (green), which are presented in Tables B1 and B2 respectively.

Table B1. Dust population parameters for the models presented in this work when compared to DES5YR data split at host $\log(M_*/M_\odot) = 10$ and fit with BBC-BS20 (left hand panels of Figs. B1-B3.)

Model Name	BiasCor	$\overline{R_{V,1}}$	$\sigma_{R_{V,1}}$	$\overline{R_{V,2}}$	$\sigma_{R_{V,2}}$	Y	Y_{split}	Y_1	Y_2	X	X_{split}	γ_X	$\chi^2_{\nu, \mu_{\text{res}}}(U-R)$
BS21	BS20	2.5	1.0	3.25	1.0	M_*	$10^{10} M_\odot$	-	-	M_*	$10^{10} M_\odot$	0.00	1.16
BS21	BS20	2.25	1.0	3.0	1.0	M_*	$10^{10} M_\odot$	-	-	τ_A	0.75 Gyr	0.05	1.23
Age R_V	BS20	1.75	1.0	3.25	1.0	τ_G	3 Gyr	-	-	τ_A	0.75 Gyr	0.0	1.11
Age R_V	BS20	1.75	1.0	3.25	1.0	τ_G	3 Gyr	-	-	M_*	$10^{10} M_\odot$	0.0	1.11
Age R_V Linear	BS20	1.5	1.0	3.75	1.0	τ_A	-	0.1 Gyr	10 Gyr	τ_A	0.75 Gyr	0.20	1.40

Table B2. Dust population parameters for the models presented in this work when compared to DES5YR data split at host $U-R = 1$ and fit with BBC-BS20 (right hand panels of Figs. B1-B3.)

Model Name	BiasCor	$\overline{R_{V,1}}$	$\sigma_{R_{V,1}}$	$\overline{R_{V,2}}$	$\sigma_{R_{V,2}}$	Y	Y_{split}	Y_1	Y_2	X	X_{split}	γ_X	$\chi^2_{\nu, \mu_{\text{res}}}(U-R)$
BS21	BS20	2.25	1.0	3.25	1.0	M_*	$10^{10} M_\odot$	-	-	M_*	$10^{10} M_\odot$	0.1	0.80
BS21	BS20	1.75	1.0	3.75	1.0	M_*	$10^{10} M_\odot$	-	-	τ_A	0.75 Gyr	0.05	0.76
Age R_V	BS20	2.0	1.0	2.5	1.0	τ_G	3 Gyr	-	-	τ_A	0.75 Gyr	0.15	0.89
Age R_V	BS20	1.5	1.0	2.5	1.0	τ_G	3 Gyr	-	-	M_*	$10^{10} M_\odot$	0.05	0.82
Age R_V Linear	BS20	1.5	1.0	3.75	1.0	τ_A	-	0.1 Gyr	10 Gyr	τ_A	0.75 Gyr	0.20	1.22

**Figure B1.** BS21 models as Fig. 9 but where Hubble residuals for data and simulations have been measured with BBC-BS20. *a)*: Simulations from the BS21 model with $\overline{R_V} = 3.25$ in low mass host galaxies and $\overline{R_V} = 2.5$ in high mass host galaxies; *b)*: as *a*, but for the data and simulations split host galaxy $U-R = 1$. *c)* & *d)*: as *a* & *b* but for the model parameters that best describe the $U-R$ data (right hand panel): $\overline{R_V} = 3.25$ in low mass host galaxies and $\overline{R_V} = 2.25$ in low mass host galaxies, and an intrinsic luminosity step at $\log(M/M_\odot) = 10$ of size 0.1 mag.

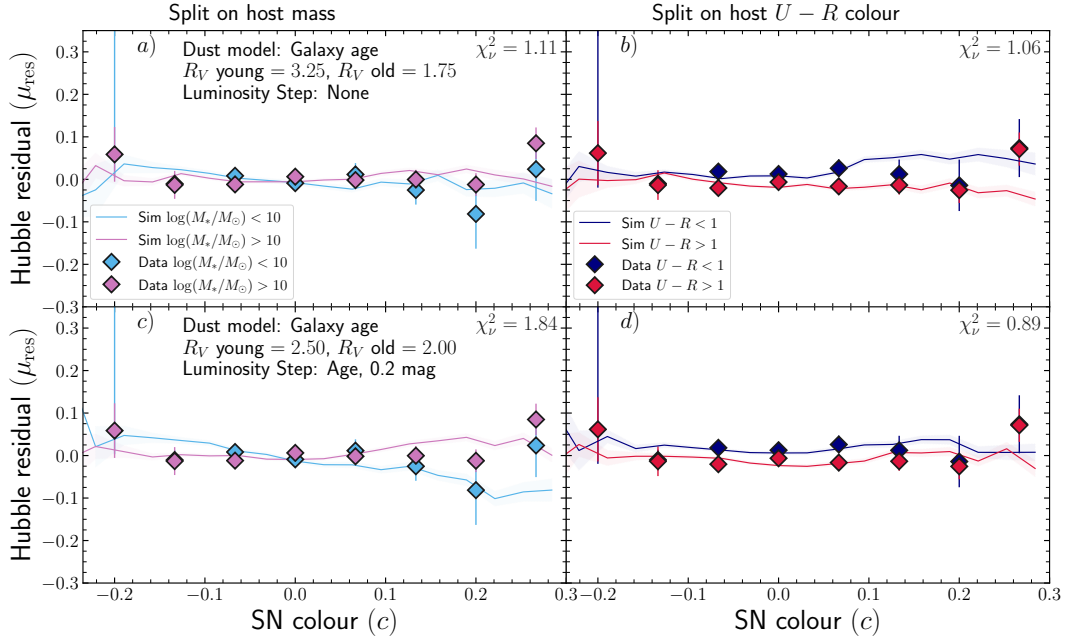


Figure B2. Age- R_V split models as Fig. 10, but where Hubble residuals for data and simulations have been measured with BBC-BS20.

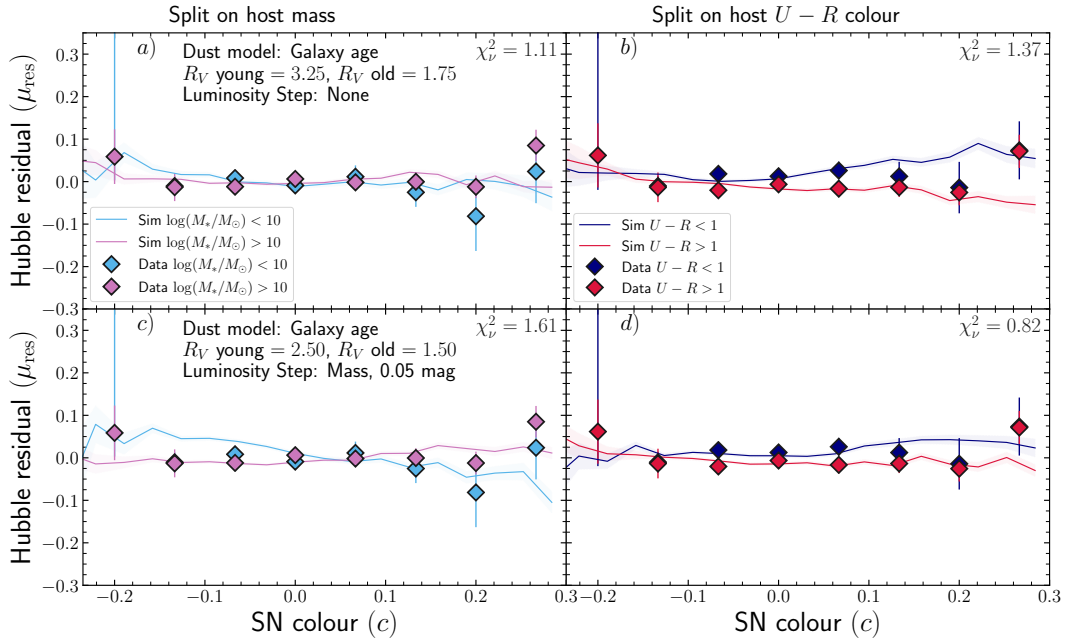


Figure B3. As Fig. B3 but for the models best fitting when introducing an intrinsic stellar mass step instead of a SN age step.

- ¹ School of Physics and Astronomy, University of Southampton, Southampton, SO17 1BJ, UK
- ² Department of Physics, Duke University Durham, NC 27708, USA
- ³ Institute of Cosmology and Gravitation, University of Portsmouth, Portsmouth, PO1 3FX, UK
- ⁴ Center for Astrophysics | Harvard & Smithsonian, 60 Garden Street, Cambridge, MA 02138, USA
- ⁵ Einstein Fellow
- ⁶ School of Mathematics and Physics, University of Queensland, Brisbane, QLD 4072, Australia
- ⁷ Institut d'Estudis Espacials de Catalunya (IEEC), 08034 Barcelona, Spain
- ⁸ Institute of Space Sciences (ICE, CSIC), Campus UAB, Carrer de Can Magrans, s/n, 08193 Barcelona, Spain
- ⁹ Centre for Gravitational Astrophysics, College of Science, The Australian National University, ACT 2601, Australia
- ¹⁰ The Research School of Astronomy and Astrophysics, Australian National University, ACT 2601, Australia
- ¹¹ Centre for Astrophysics & Supercomputing, Swinburne University of Technology, Victoria 3122, Australia
- ¹² Univ Lyon, Univ Claude Bernard Lyon 1, CNRS, IP2I Lyon / IN2P3, IMR 5822, F-69622 Villeurbanne, France
- ¹³ Laboratório Interinstitucional de e-Astronomia - LIneA, Rua Gal. José Cristino 77, Rio de Janeiro, RJ - 20921-400, Brazil
- ¹⁴ Fermi National Accelerator Laboratory, P. O. Box 500, Batavia, IL 60510, USA
- ¹⁵ Department of Physics, University of Michigan, Ann Arbor, MI 48109, USA
- ¹⁶ CNRS, UMR 7095, Institut d'Astrophysique de Paris, F-75014, Paris, France
- ¹⁷ Sorbonne Universités, UPMC Univ Paris 06, UMR 7095, Institut d'Astrophysique de Paris, F-75014, Paris, France
- ¹⁸ University Observatory, Faculty of Physics, Ludwig-Maximilians-Universität, Scheinerstr. 1, 81679 Munich, Germany
- ¹⁹ Department of Physics & Astronomy, University College London, Gower Street, London, WC1E 6BT, UK
- ²⁰ Kavli Institute for Particle Astrophysics & Cosmology, P. O. Box 2450, Stanford University, Stanford, CA 94305, USA
- ²¹ SLAC National Accelerator Laboratory, Menlo Park, CA 94025, USA
- ²² Instituto de Astrofísica de Canarias, E-38205 La Laguna, Tenerife, Spain
- ²³ Universidad de La Laguna, Dpto. Astrofísica, E-38206 La Laguna, Tenerife, Spain
- ²⁴ Center for Astrophysical Surveys, National Center for Supercomputing Applications, 1205 West Clark St., Urbana, IL 61801, USA
- ²⁵ Department of Astronomy, University of Illinois at Urbana-Champaign, 1002 W. Green Street, Urbana, IL 61801, USA
- ²⁶ Institut de Física d'Altes Energies (IFAE), The Barcelona Institute of Science and Technology, Campus UAB, 08193 Bellaterra (Barcelona) Spain
- ²⁷ Astronomy Unit, Department of Physics, University of Trieste, via Tiepolo 11, I-34131 Trieste, Italy
- ²⁸ INAF-Osservatorio Astronomico di Trieste, via G. B. Tiepolo 11, I-34143 Trieste, Italy
- ²⁹ Institute for Fundamental Physics of the Universe, Via Beirut 2, 34014 Trieste, Italy
- ³⁰ Hamburger Sternwarte, Universität Hamburg, Gojenbergsweg 112, 21029 Hamburg, Germany
- ³¹ Department of Physics, IIT Hyderabad, Kandi, Telangana 502285, India
- ³² Jet Propulsion Laboratory, California Institute of Technology, 4800 Oak Grove Dr., Pasadena, CA 91109, USA
- ³³ Institute of Theoretical Astrophysics, University of Oslo. P.O. Box 1029 Blindern, NO-0315 Oslo, Norway
- ³⁴ Kavli Institute for Cosmological Physics, University of Chicago, Chicago, IL 60637, USA
- ³⁵ Instituto de Física Teórica UAM/CSIC, Universidad Autónoma de Madrid, 28049 Madrid, Spain
- ³⁶ Department of Physics and Astronomy, University of Pennsylvania, Philadelphia, PA 19104, USA
- ³⁷ Observatório Nacional, Rua Gal. José Cristino 77, Rio de Janeiro, RJ - 20921-400, Brazil
- ³⁸ Santa Cruz Institute for Particle Physics, Santa Cruz, CA 95064, USA
- ³⁹ Center for Cosmology and Astro-Particle Physics, The Ohio State University, Columbus, OH 43210, USA
- ⁴⁰ Department of Physics, The Ohio State University, Columbus, OH 43210, USA
- ⁴¹ Institució Catalana de Recerca i Estudis Avançats, E-08010 Barcelona, Spain
- ⁴² Physics Department, 2320 Chamberlin Hall, University of Wisconsin-Madison, 1150 University Avenue Madison, WI 53706-1390
- ⁴³ Department of Astronomy, University of California, Berkeley, 501 Campbell Hall, Berkeley, CA 94720, USA
- ⁴⁴ Institute of Astronomy, University of Cambridge, Madingley Road, Cambridge CB3 0HA, UK
- ⁴⁵ Department of Astrophysical Sciences, Princeton University, Peyton Hall, Princeton, NJ 08544, USA
- ⁴⁶ Department of Physics and Astronomy, Pevensey Building, University of Sussex, Brighton, BN1 9QH, UK
- ⁴⁷ Centro de Investigaciones Energéticas, Medioambientales y Tecnológicas (CIEMAT), Madrid, Spain
- ⁴⁸ Computer Science and Mathematics Division, Oak Ridge National Laboratory, Oak Ridge, TN 37831
- ⁴⁹ Excellence Cluster Origins, Boltzmannstr. 2, 85748 Garching, Germany
- ⁵⁰ Max Planck Institute for Extraterrestrial Physics, Giessenbachstrasse, 85748 Garching, Germany
- ⁵¹ Universitäts-Sternwarte, Fakultät für Physik, Ludwig-Maximilians Universität München, Scheinerstr. 1, 81679 München, Germany

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