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HOW GAS HYDRATE SATURATION AND MORPHOLOGY CONTROL SEISMIC ATTENUATION: A CASE STUDY FROM THE SOUTH HYDRATE RIDGE

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1 HOW GAS HYDRATE SATURATION AND MORPHOLOGY CONTROL SEISMIC
2 ATTENUATION: A CASE STUDY FROM THE SOUTH HYDRATE RIDGE

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ABSTRACT

14 Prior studies have shown an ambiguous relationship between gas hydrate
15 saturation and seismic attenuation in different regions, but the effect of gas hydrate
16 morphology on seismic attenuation of hydrate-bearing sediments was often overlooked.
17 Here we combine seismic data with rock physics modeling to elucidate how gas hydrate
18 saturation and morphology may control seismic attenuation. To extract P-wave
19 attenuation, we process both the vertical seismic profile (VSP) data within a frequency
20 range of 30 – 150 Hz and sonic logging data within 10 – 15 kHz from three wells in the
21 south Hydrate Ridge, offshore of Oregon (USA), collected during Ocean Drilling
22 Program (ODP) Leg 204 in 2000. We calculate P-wave attenuation using spectral
23 matching and centroid frequency shift methods, and use Archie's relationship to derive
24 gas hydrate saturation from the resistivity data above the bottom simulating reflection
25 (BSR) at the same wells. To interpret observed seismic attenuation in terms of the effects
26 of both gas hydrate saturation and morphology, we employ the Hydrate-Bearing Effective
27 Sediment (HBES) rock physics model. By comparing the observed and model-predicted
28 attenuation values, we infer that: (1) seismic attenuation appears to not be dominated by
29 any single factor, instead, its variation is likely governed by both gas hydrate saturation
30 and morphology; (2) the relationship between seismic attenuation and gas hydrate
31 saturation varies with different hydrate morphologies; (3) the squirt flow, occurring at
32 different compliances of adjacent pores driven by pressure gradients, may be responsible
33 for the significantly large or small attenuation over a broad frequency range.

34 INTRODUCTION

35 Gas hydrates are ice-like crystalline compounds of water and gas (mainly
36 methane) molecules, existing naturally in sediments within a limited depth range on the
37 continental margins and permafrost environments at high pressure (>0.6 MPa) and
38 moderately low temperature (<300 K) conditions when methane exceeds its solubility
39 (Brooks et al., 1986; Kvenvolden et al., 1993; Sloan, 1998). Due to its widespread
40 occurrence in the continental margins and permafrost settings, numerous studies, in
41 particular, seismic surveys, have been carried out to investigate the effects of gas hydrate
42 on climate and seafloor stability as well as its important role as a potential energy
43 resource (Dickens, 2003; Bohannon, 2008; Collett et al., 2014a). In recent decades,
44 measurements of seismic velocity, attenuation and gas hydrate saturation have been
45 carried out in the Blake Ridge site off the southeast coast of the United States (Holbrook
46 et al., 1996; Guerin et al., 1999; Hornbach et al., 2008), at the Mallik site in the
47 Mackenzie Delta, Canada (Dallimore and Collett, 2002; Guerin and Goldberg, 2002;
48 Guerin et al., 2005; Bellefleur et al., 2007), in the Nankai Trough, offshore central Japan
49 (Matsushima, 2006), on the western Svalbard continental margin (Carcione et al., 2005;
50 Madrussani et al., 2010), in the Krishna-Godavari (KG) Basin off the eastern coast of
51 India (Jaiswal et al., 2012; Collett et al., 2014b; Shankar, 2016; Jyothi et al., 2017) and in
52 the Gulf of Mexico (Brooks et al., 1986; Phrampus and Hornbach, 2012; Wang et al.,
53 2017).

54 Generally, gas hydrate growth in the pore space alters the elastic properties of
55 host sediments, allowing us to characterize hydrate-bearing sediments with seismic
56 methods. Previous studies showed that gas hydrate growing in the pore space can stiffen

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3 57 the matrix of the bearing sediments by increasing their bulk and shear moduli, thereby
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58 leading to higher P-wave and S-wave velocities (Yuan et al., 1996; Chand et al., 2004;
59 Liu et al., 2020). However, the degree at which the elastic properties are altered depends
60 on both hydrate saturation and morphology (Riedel et al., 2010; Liu and Liu, 2018; Liu et
61 al., 2020). Hydrate morphology can be divided into two main types depending, mainly,
62 on the grain size of the host sediment (Ren et al., 2020). Hydrate growing in fine-grained
63 sediments tends to create particle displacing morphologies such as lenses, nodules, and
64 chunks or veins. In coarse-grained sediments hydrate growth tends to be pore invasive
65 (does not displace the grains) and can be subdivided into (i) pore-floating or pore-filling
66 and (ii) matrix-supporting, including load-bearing, contact-cementing and grain-coating
67 hydrate (Waite et al., 2009; Pan et al., 2019, 2020). For gas hydrate floating in the pore
68 space without any grain contact, the P-wave velocity increases while the S-wave velocity
69 remains almost unchanged due to unaffected shear modulus and minor decrease in bulk
70 density, as the density of gas hydrate is only slightly smaller than that of pore water.
71 When gas hydrate cements grains or acts as part of the load-bearing frame, both the P-
72 wave velocity and the S-wave velocity increase (Chand et al., 2004; Sava and Hardage,
73 2006; Sahoo et al., 2018, 2019).

74 Seismic attenuation provides complementary information to velocity on
75 constraining gas hydrate saturation and morphology. However, until now, the
76 investigation of seismic attenuation of hydrate-bearing sediments is scarce partly because
77 good field measurements are limited. As a result, how seismic attenuation varies with gas
78 hydrate properties is still poorly understood. In the field, the effect of gas hydrate
79 saturation on attenuation varies from region to region. For example, in the Mallik field in

80 Canada, where gas hydrate occurs in subpermafrost occupying up to 60% - 80% of the
81 pore space, the high gas hydrate saturation is associated with high attenuation with Q
82 values of less than 20 for P-waves (Guerin and Goldberg, 2002). A similar observation of
83 high attenuation has been shown in the Nankai Trough in Japan, where gas hydrate
84 locates on the continental margin occupying 20% - 30% of the pore space (Matsushima,
85 2006). In contrast, the high attenuation is associated with the low gas hydrate saturation
86 observed in the west Svalbard in Norway (Madrussani et al., 2010) and the KG basin in
87 India (Nittala et al., 2017). The measurements of these studies are summarized in Table 1.
88 We find that these measurements exhibit an ambiguous relationship between the
89 attenuation and gas hydrate saturation.

90 While previous studies attempted to derive the relationship between gas hydrate
91 saturation and seismic attenuation in hydrate-bearing sediments, gas hydrate morphology,
92 which may greatly impact the attenuation in hydrate-bearing sediments, was mostly
93 ignored due to the difficulties of conducting morphology measurements from limited
94 field samples. Therefore, it is challenging to systematically establish a database of gas
95 hydrate morphology from field studies. To explain this gas hydrate morphology-seismic
96 puzzle, rock physics modeling and laboratory measurements have been used. For
97 example, an early study by Lee and Collett (2000) used rock physics modeling to find the
98 relationship between morphology and seismic velocity; Choi et al. (2014) synthesized a
99 non-cementing form of gas hydrate in sandy sediments in the laboratory, measuring the
100 P-wave velocity of the sample during the synthesizing process, and trying to determine
101 gas hydrate morphology at each transition through rock physics modeling. Liu et al.
102 (2020) proposed a joint analysis of P-wave velocity and resistivity to identify hydrate

103 morphology and estimate hydrate saturation in a continuous depth profile. They
104 successfully used the cross plot of P-wave velocity and resistivity to identify fracture-
105 filling gas hydrate-bearing sediments (GHBS) from pore-filling GHBS in the South
106 China Sea based on the observations that fracture-filling gas hydrate-bearing sediments
107 exhibit higher resistivity but lower P-wave velocity than those of pore-filling GHBS in
108 the case of identical hydrate concentration. Recently, Zhan and Matsushima (2018) used
109 the Marin-Moreno's et al. (2017) Hydrate-Bearing Effective Sediment (HBES) model
110 and the Guerin and Goldberg's (2005) model to quantify the attenuation due to a single
111 morphology and multiple morphologies in the Nankai Trough, Japan. Their results
112 confirmed that the occurrence of gas hydrate in different morphologies can better explain
113 the seismic attenuation measurements. However, some questions remain: does the
114 attenuation mechanism behave the same at different hydrate sites and at different hydrate
115 saturation? Our study provides insights into this aspect at a relatively lower hydrate
116 saturation site compared to that of the Nankai Trough in Japan.

117 In this study, we aim to extend the understanding of how gas hydrate saturation
118 and morphology can alter seismic attenuation and finally elucidate the possible
119 attenuation mechanism in the south Hydrate Ridge. In this work we assume that hydrate
120 growth is pore invasive. We choose high quality vertical seismic profiling (VSP) and
121 sonic logging data to extract P-wave attenuation. We process the VSP data within a
122 frequency range of 30 – 150 Hz and sonic logging data within 10 – 15 kHz from three
123 wells in the south Hydrate Ridge, offshore of Oregon, collected during Ocean Drilling
124 Program (ODP) Leg 204 in 2000 (Tréhu et al., 2004). P-wave attenuation is estimated
125 using spectral matching and centroid frequency shift methods. Different models have

126 been proposed to interpret attenuation measurements on hydrate bearing sediments
127 (Chand and Minshull, 2004; Guerin et al., 2005; Best et al., 2013; Marín-Moreno et al.,
128 2017). To interpret the estimated attenuation from VSP and sonic logging data in terms of
129 the effects of both gas hydrate saturation and morphology, here we employ a frequency
130 dependent HBES rock physics model (Marín-Moreno et al., 2017). The HBES has been
131 recently proven successful in capturing VSP and sonic logging attenuation measurements
132 from natural hydrate bearing sediments in the eastern Nankai Trough (Zhan and
133 Matsushima, 2018). Then we discuss the possible attenuation mechanisms on the gas
134 hydrate sites.

135 Overall, our study helps elucidate the interaction between gas hydrate saturation
136 and morphology in the field, and bridge the gap between gas hydrate and seismic
137 attenuation, which underpins the pivotal role of combining rock physical modeling and
138 field observations in future gas hydrate studies.

139 GEOLOGY AND DATA DESCRIPTION OF THE HYDRATE RIDGE

140 The Hydrate Ridge is located on the Cascadia continental margin which overlies
141 the subduction zone where the Juan de Fuca plates is subducting beneath the North
142 American plate (Figure 1a). Along this margin, two main basins (Cascadia Basin and
143 Gorda Basin), several fans and an accretionary complex have developed. A wide variety
144 of Pleistocene and Holocene turbidites generated most of the deposits in which bottom-
145 simulating reflections (BSRs) are widely recognized. To investigate gas hydrates, a suite
146 of well log data including caliper, gamma ray, resistivity, sonic and VSP was acquired
147 during ODP Leg 204 in the south Hydrate Ridge. The ODP Leg 204 scientific report

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3 148 (Tréhu et al., 2006c, 2006b) indicated the presence of gas hydrate and revealed its
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5 149 occurrence within turbiditic silty layers. The strong BSR suggests the presence of gas
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7 150 hydrate bearing sediments overlying sediments hosting free gas. An average gas hydrate
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9 151 saturation of ~10% has been estimated in the targeted formations of Leg 204. However, it
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1 152 can reach 20% - 30% in specific locations (e.g. southern summit of the ridge) (Milkov et
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3 153 al., 2003; Tréhu et al., 2004). In this study, we process the VSP (30 – 150 Hz) and sonic
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5 154 logging data (10 – 15 kHz) from three selected sites (Site 1244, 1247 and 1250) (Figure
6
7 155 1b) due to their high data quality to calculate the P wave attenuation over different
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9 156 frequency bands. We obtain hydrate saturation from the resistivity logging data in those
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1 157 wells (Tréhu et al., 2006b) using Archie's relationship (Archie, 1942).

158 SEISMIC DATA AND ATTENUATION CALCULATION

159 **Field data**

160 Figure 2 shows a two-dimensional (2-D) seismic profile across the location of one
161 of our selected wells. The seismic profile clearly reveals a BSR at 100-150 m depth
162 beneath seafloor (seafloor is at 905 m depth in 1244E well). In this study, we choose high
163 quality near-offset (55 m) VSP data (Figure 3) and monopole sonic data (Figure 4) for
164 attenuation calculation in two frequency bands. Taking 1244E well as an example, VSP
165 data was collected from 85-245 m depth beneath seafloor with a receiver interval of 5 m.
166 In order to make a more reliable calculation, we first apply a median filter to separate the
167 upgoing and downgoing waves, then used a bandpass filter with a frequency of (20-40-
168 80-100 Hz) to filter out noises of downgoing waves. Filtered downgoing waves are

169 windowed with a cosine window (length = 0.025 s) centered at the first peak for later
170 attenuation calculation (Figure 5a). Regarding to sonic data, the transmitter is 108 inches
171 (2.74 m) away from the first receiver and the interval between each receiver is 6 inches
172 (0.15 m). Due to the significant energy concentrating in the relatively lower frequencies
173 (Figure 5b), we apply a bandpass filter (5-6-7-8 kHz) and then a cosine taper window
174 (length = 0.45 ms). Figures 5c and 5d show amplitude spectra of VSP data (Figure 5c)
175 and sonic data (Figure 5d), respectively. The above-mentioned data processing was done
176 using open-source MATLAB software CREWES.

177 The fact that high frequencies are attenuated more than lower frequencies
178 motivates us to explore the spectra-based methods to calculate seismic attenuation. In this
179 study, we use the spectral matching method proposed by Blias (2012) to estimate VSP
180 attenuation and the centroid frequency method (Quan and Harris, 1997) to estimate sonic
181 attenuation following the workflow as stated below (Figure 6).

182 **Attenuation (1/Q) estimation**

183 Seismic attenuation is often referred to 1/Q, including the scattering and intrinsic
184 attenuation. When the heterogeneities and wavelength are comparable, scattering which
185 is the reflection of the wave in directions other than its original propagation direction will
186 occur. The intrinsic attenuation is the absorbed wave energy converted to heat, often
187 quantified with the inverse of quality factor of the media (Q_{intr}^{-1}). The estimated
188 attenuation which we call effective attenuation (Q_{eff}^{-1}) is the combination of the scattering
189 (Q_{scat}^{-1}) and intrinsic attenuation (Q_{intr}^{-1}), i.e., $Q_{eff}^{-1} = Q_{scat}^{-1} + Q_{intr}^{-1}$. In practice, we are most

190 interested in the intrinsic attenuation which can be compared by calculating its inverse –
 191 the quality factor Q for mathematical simplification.

192 *Spectral Matching Method*

193 The spectral matching method is one of the most popular methods for estimating
 194 attenuation where data from two receiver depths are selected. Let us assume that a
 195 seismic wavelet with amplitude spectrum $S_1(f)$ has an amplitude spectrum $S_2(f)$ after
 196 travelling in an attenuating medium for an interval time t . Then, we can describe the
 197 seismic attenuation process as:

$$|S_2(f)| = G|S_1(f)|e^{-\frac{\pi f t}{Q}}, \quad (1)$$

198 where $S_1(f)$ and $S_2(f)$ are amplitude spectra of downgoing waves at the depth z_1 and z_2 ,
 199 f is the frequency and G represents all the frequency independent amplitude loss in total,
 200 including spherical divergence, reflection and transmission loss.

201 Equation 1 is based on the following assumptions: (1) the source and geophone
 202 coupling does no change between the two levels; (2) there is no interference from
 203 reflected waves; (3) Q is frequency independent (Tonn, 1991; Harris et al., 1997; Blias,
 204 2012). This formula can be treated as the foundation for most spectral methods to
 205 estimate Q . To process VSP data, we use the spectra matching method proposed by Blias
 206 (2012). $S_1(f)$ is modified by varying Q until an optimum approximation to $|S_2(f)|$ is
 207 obtained:

$$Q_{test} = \min_Q \left\| |S_2(f)| - G|S_1(f)|e^{-\frac{\pi f(t_2 - t_1)}{Q}} \right\|^2, \quad (2)$$

208 where G addresses the frequency independent energy loss and can be calculated as:

$$G = \frac{\int_{-\infty}^{\infty} |S_2(f)| |S_1(f)| e^{\frac{-\pi f(t_2 - t_1)}{Q}} df}{\int_{-\infty}^{\infty} |S_1(f)|^2 e^{\frac{-2\pi f(t_2 - t_1)}{Q}} df}. \quad (3)$$

209 Equation 2 was applied to the windowed VSP data described above. The downgoing
 210 arrivals at receivers spaced by different intervals (5, 10, 20, 30 m) were tested for Q
 211 estimation. We find that the 10 m interval works the best even though they all contain
 212 some negative values largely due to scattering or equipment coupling.

213 *Centroid Frequency Shift Method*

214 With observations of significant frequency down-shift in sonic data, we select the
 215 centroid frequency shift method (Quan and Harris, 1997) to estimate sonic data
 216 attenuation.

217 With the assumption of Gaussian wavelet, the centroid frequency shift can be
 218 linked to wave attenuation as:

$$\int \frac{\pi}{Qv} dl = \frac{f_1 - f_2}{\sigma_1^2}. \quad (4)$$

219 Where f_1 is the centroid frequency of reference seismogram, f_2 is the centroid frequency
 220 of the seismogram at target receivers, and σ_1^2 is the variance of the seismogram at
 221 receiver. Therefore, following Quan and Harris (1997), Q between the i_{th} and $(i + 1)_{th}$
 222 receiver is defined as:

$$Q = \frac{\pi \sigma_i^2 \Delta t_i}{\Delta f_i}. \quad (5)$$

223 In equation 5, $\Delta f_i = (f_i - f_{i+1})$ is the centroid frequency difference between the i_{th} and
224 $(i + 1)_{th}$ layers, $\Delta t_i = t_i - t_{i+1}$ is the traveltime difference between the two layers, and
225 σ_i^2 is the variance at the i_{th} receiver. Equation 5 is applied to the windowed sonic data
226 described above with the centroid frequencies between 6 kHz to 7 kHz.

227 VSP AND SONIC ATTENUATION RESULTS

228 The logging data, including sonic attenuation, of three wells are shown in Figure
229 7, Figure 8 and Figure 9 for 1244E, 1247B and 1250F, respectively. The gas hydrate
230 zones are identified above the BSR (grey dashed lines) at three wells, which are around
231 1025 m in 1244E, 976 m in 1247B and 917 m in 1250F below sea level (Tréhu et al.,
232 2006b). In these wells, changes in lithology with depth are likely small, based on the
233 small depth variation of gamma ray (green lines) between 50-70 API. Therefore, we
234 assume that the variation of seismic attenuation with depth in this study area is mainly
235 caused by gas hydrate properties varying with depth.

236 The gas hydrate saturation is derived from resistivity data using Archie's equation
237 (details in Appendix A). The calculated VSP attenuations range between 0.004 – 0.013 at
238 site 1244E, and 0.004 at site 1247B and 1250F (red dots). They are calculated at certain
239 layers and are generally smaller than the sonic attenuation (black lines in Figure 7, Figure
240 8 and Figure 9). The uncertainty analysis is implemented using a Monte Carlo method
241 (Riedel et al., 2013; Wang et al., 2017) and errors are illustrated in terms of the standard
242 deviation and 95% confidence interval of average attenuation (details in Appendix B).

243 ROCK PHYSICS MODELING

244 To further quantify the effect of gas hydrate saturation and morphology on
245 seismic attenuation, we employ the frequency dependent HBES model (Marín-Moreno et
246 al., 2017) that builds on Best et al.'s (2013) model. The HBES formulation considers that
247 fluid inclusions in the hydrate, which are inclusions of water and/or methane in the
248 hydrate connected to the intergranular pores, can modify the elastic moduli of the hydrate
249 and sediments; fluid inclusions in the hydrate are ellipsoidal, one-sided connected to the
250 intergranular pores, homogeneously distributed in the hydrate, and independent on
251 hydrate morphology; contact-cementing, grain-coating and pore-filling/floating hydrate
252 morphologies as defined by Ecker et al. (1998, 2002). Regarding the application of the
253 HBES to this work, we consider that the pore space above the BSR contains only gas
254 hydrate and water in the pore space and that the hydrate is homogeneously distributed in
255 the pore space.

256 The two categories of gas hydrate morphology were defined based on whether the
257 hydrate exists adhering to host grains or floating in the pore space and initially deduced
258 from the effect of gas hydrate morphology on the elastic wave velocity (Ecker et al.,
259 1998). The idealized conceptual model of the microstructure of hydrate-bearing
260 sediments is shown in Figure 10.

261 This model uses the Biot-Stoll poro-elastic theory (Biot, 1956a, 1956b), gas
262 hydrate morphologies as defined by Ecker et al. (1998) and the formulation for squirt
263 flow given by Leurer (1997) to predict seismic attenuation of hydrate bearing sediments.
264 Let us consider the gas hydrate as a compliant composite porous material with inclusions
265 like gas or water, rather than as a solid. The inclusions are a consequence of isolated

266 pockets of gas or water trapped during gas hydrate formation. In this case, the gas hydrate
267 can behave like other microporous, compliant minerals like clay assemblages (Leurer,
268 1997). When an elastic wave passes through it, the compliant porous host (e.g., host grain
269 framework) and the porous gas hydrate grains create local fluid pressure gradients
270 between the gas hydrate inclusions and the sand frame pores, leading to viscous fluid
271 flow (squirt flow) of water/gas and corresponded wave energy loss. The HBES model
272 allows seismic attenuation to be estimated as a function of frequency, gas hydrate
273 saturation and different combinations of gas hydrate morphologies. A simplified version
274 of the HBES model workflow is shown in Figure 11. The input parameters and symbols
275 of this model are shown in Table C1 and C2 in Appendix C.

276 COMPARISON AND DISCUSSION

277 **Relation between attenuation and the hydrate properties**

278 Figure 12 shows calculated attenuation from sonic logging data and gas hydrate
279 saturation with depth. Since gas hydrate exists only above the BSR indicated by the grey
280 dashed line in Figure 12 and the VSP data is limited, we only focus on the sonic
281 attenuation above the BSR. Some negative values of attenuation may be resulted from
282 scattering of background or equipment coupling (Matsushima, 2006). The trend between
283 saturation and attenuation is not obvious in Figure 12. We quantify the trend by
284 calculating the correlation coefficients between seismic attenuation and gas hydrate
285 saturation. Small coefficients (-0.34, -0.09 and 0.32 for 1244E, 1247B and 1250F,
286 respectively) imply that there may not be a strong linear relation between seismic

287 attenuation and gas hydrate saturation. In order to see if there is a regional pattern
288 between the seismic attenuation and saturation, we also make a scatter plot of the seismic
289 attenuation and saturation from 1244E (blue circles), 1247B (red circles), and 1250F
290 (yellow circles) in Figure 13. Not surprisingly, there is no clear correlation between gas
291 hydrate saturation and seismic attenuation though we do observe that the attenuation
292 values are more concentrated at lower hydrate saturation and relatively scattered at higher
293 hydrate saturation.

294 To combine the effect of gas hydrate morphology on the attenuation, we overlay
295 rock physics modeling results on the scatter plot representing different combinations of
296 gas hydrate morphologies between two end-members: cementing and pore-filling hydrate
297 (Marín-Moreno et al., 2017). Then, the modeled attenuations from the HBES model
298 (solid lines) are compared with our calculated attenuations from three wells (Figure 13).

299 For the modeled results shown as solid lines in Figure 13, we clarify that the
300 fraction attached with pore-filling or cementing means a fraction of the total hydrate
301 saturation, i.e., percentage \times hydrate saturation. The attenuation-versus-hydrate saturation
302 curves predicted from the HBES model based on different hydrate morphologies clearly
303 indicate that P-wave attenuation is strongly dependent on gas hydrate morphology. We
304 note that the modeled P-wave attenuation curve with 100% cementing hydrate (red line)
305 appears to capture the upper boundary of calculated sonic attenuations with the increasing
306 of the hydrate saturation, while the modeled P-wave curve with 60% cementing hydrate
307 (yellow line) seems to define the lower boundary of calculated sonic attenuations, in
308 particular at lower saturation. For both the pure pore-filling hydrate (green line) and the
309 multiple morphologies (containing both the pore-filling hydrate and the cementing

310 hydrate) with more pore-filling hydrate (purple and blue lines), the predicted attenuation
311 has a monotonic increase with the saturation. For the multiple morphologies with more
312 cementing hydrate (yellow and black lines), the predicted attenuation can be separated
313 into three stages: (1) at very low saturation (below ~2%), the attenuation increases with
314 the saturation and reaches the peak at around 2%; (2) with saturation increasing from 2%
315 to 10 – 15%, the predicted attenuation decreases gradually; (3) at higher saturations, the
316 predicted attenuation slightly increases. For the pure cementing hydrate model, the
317 predicted attenuation shows the same trend, but the attenuation peak occurs around 6%,
318 then it shows a very smooth decrease at a relatively wider saturation range.

319 To consider possible frequency effects on attenuation, we compare VSP and sonic
320 attenuations with those modeled at certain layers having gas hydrate saturations of 30%,
321 24%, 16% and 4% as a function of frequency (Figure 14). Since the VSP attenuations are
322 much smaller and closer for all the three sites, we use the same yellow symbol to display
323 all of them (Figure 14). In general, the VSP attenuation is smaller than the sonic
324 attenuation. Considering the scattering of calculated attenuation from the field data, the
325 modeled results could be treated as a good fit with our measured sonic attenuation. The
326 possible gas hydrate morphologies at some depth in the south Hydrate Ridge in the
327 Cascadia Margin could be around 85% cementing hydrates, which is consistent with the
328 observation of most attenuation data being well represented by hydrate with 60% to
329 100% cementing morphology (Figure 13).

330 Discussion

331 Our results show that the effect of gas hydrate saturation on seismic attenuation
332 can be different: (1) when the percentage of pore-filling hydrates is larger than that of
333 cementing hydrates (see 100% pore-filling, and 20% and 40% cementing curves in Figure
334 13), in general the attenuation increases with the increasing of saturation; (2) when the
335 percentage of cementing hydrates is larger than pore-filling hydrates (see 60%, 80% and
336 100% cementing curves in Figure 13), the increase of saturation can lead to a decrease of
337 the attenuation at a certain saturation range controlled by the percentage of cementing
338 hydrate (i.e. more cementing hydrate leads to a decrease of attenuation at a wider
339 saturation range), and then the attenuation increases again as the saturation keeps
340 increasing.

341 Priest et al. (2006) proposed a reasonable theory for the case with more cementing
342 hydrates: (1) at the first stage (Figure 15a), let us consider a dry sediment specimen with
343 a small volume of adsorbed water, which tends to condense at grain contacts due to
344 surface tension and capillary pressure. When a seismic wave passes through, the
345 generated pressure leads to squirt flow at grain contacts, resulting in energy loss.
346 However, since the amount of adsorbed water is not large enough to fill many grain
347 contact micropores, the energy loss is very small and so minimal attenuation occurs; (2)
348 at the second stage (Figure 15b), hydrate starts to grow in the pore space at grain contacts
349 as a cement, which causes the increase of effective area of the grain contacts followed by
350 more squirt flow. Therefore, more energy could be lost and the attenuation increases
351 rapidly until a critical hydrate saturation; the first two stages are consistent with our
352 results at low saturation (below about 5%); (3) when reaching the third stage (Figure

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3 353 15c), beyond the critical hydrate saturation, increasing growth of hydrates starts to
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5 354 envelop the host grains and impede the movement of squirt flow, which in turn leads to a
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7 355 reduction in attenuation. However, gas hydrate also has a porous structure (Staykova et
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9 356 al., 2003; Stern et al., 2004), which may still contribute to attenuation caused by squirt
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11 357 flow. Therefore, the attenuation could be close to constant during a certain range of
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13 358 hydrate saturation. In this case, this may correspond to our results at saturations from 5%
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15 359 to 20%; (4) finally, at the fourth stage (Figure 15d), the hydrate starts to create a well-
16
17 360 formed interconnected grain-hydrate network resulting in a larger increase in the
18
19 361 effective area of the grain contacts, and consequently more energy loss caused by squirt
20
21 362 flow, so the attenuation increases again. Our results shown in Figure 13 seems to be
22
23 363 consistent with Priest et al. (2006) in that when gas hydrate saturation is pretty low (<
24
25 364 5%), gas hydrate exhibits more likely cementing behavior; once gas hydrate saturation
26
27 365 exceeds ~ 5%, the effects of pore-filling and cementing hydrate on seismic attenuation
28
29 366 cannot be distinguished. That is why we cannot observe a linear relationship between
30
31 367 hydrate saturation and seismic attenuation: effects of hydrate morphology and saturation
32
33 368 on seismic attenuation overlap with each other. However, several studies (Tohidi et al.,
34
35 369 2001; Yun, 2005; Lee et al., 2010) contradict this, in that cementing behavior cannot be
36
37 370 observed at low hydrate saturation (ca. <40%) based on laboratory work. The
38
39 371 disagreement may result from the difference between sand specimens most used in the
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41 372 laboratory and the more complicate lithologies in the field, which in our case are clays
42
43 373 and silty turbidites. In addition, our study site on the south Hydrate Ridge is located in the
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45 374 subduction zone which could cause more uncertainties than in the laboratory due to the
46
47 375 heterogeneity and complexity of the subduction zone setting. Thus, further investigation

376 on the different scales between laboratory and field work combined with rock physics
377 modeling need to be conducted in order to understand the contribution of hydrate to the
378 elastic properties of sediments, particularly seismic attenuation.

379 Regarding the limitation of the HBES model, as we mention in the previous
380 section, it only assumes two idealized hydrate morphologies, contact-cementing and
381 pore-filling and weights their individual contribution based on their concentration.
382 Though, it could be the case that their average contribution to changes in elastic response
383 of sediments which is not volumetric. In addition, the model assumes initial elliptical
384 macropores and only one aspect ratio of ellipsoidal pores when hydrate forms, whereas in
385 natural environments, it is more likely that a range of aspect ratios are generated.

386 Fundamentally, estimating seismic attenuation of hydrate-bearing sediments
387 remains challenging. Seismic attenuation comes not only from the intrinsic attenuation of
388 the hydrate-bearing sediments, but also from the scattering due to the heterogeneity of
389 natural sediments. In different hydrate sites, the estimated attenuation values exhibit
390 different ranges even if they were estimated at the same frequency range, as shown in
391 Table 1. Our estimated VSP attenuation values for the Cascadia Margin and the
392 attenuation values at the Nankai Trough (Matsushima, 2006) are lower than those at the
393 Mallik site in Canada (Dvorkin and Uden, 2004) but larger than those in the Krishna-
394 Godavari Basin (KG) in India (Nittala et al., 2017) and offshore Svalbard (Rossi et al.,
395 2007). At different frequency bands, the VSP and sonic attenuation are very different, as
396 shown in Figure 14. However, they are similar to those at the Mallik site in Canada
397 (Dvorkin and Uden, 2004; Pratt et al., 2005). Because of different lithologies in different
398 sites, the dominant attenuation mechanisms can be also different. For example, at the

399 Mallik site, hydrate saturations up to 60%-80% are inferred in individual sand layers of
400 up to 40 m thick (Bellefleur et al., 2007), and macroscopic squirt flow from elastic
401 heterogeneity in the rock frame elastic moduli may be responsible for the significant
402 attenuation (Dvorkin and Uden, 2004). In the KG Basin in India, there are mainly shale
403 layers, and Nittala et al. (2017) explained the attenuation by accounting for horizontal
404 transverse isotropy. In the Nankai Trough, Zhan and Matsushima (2018) suggested squirt
405 flow due to fluid inclusions in a microporous hydrate and the Biot-squirt (BISQ)
406 mechanism in pore spaces between hydrates and host grains might be the dominant
407 attenuation mechanism in the sonic frequency range, while squirt flow might be the
408 dominant attenuation mechanism in the seismic frequency range. The different thickness
409 of hydrate-bearing sediments may be a factor that influences the attenuations at different
410 hydrate sites. Moreover, gas hydrate can be stored either in marine sediments or
411 permafrost regions, which could result in various hydrate morphologies. We should also
412 note that the free gas coming out from the dissociation of gas hydrate may be captured
413 into a new gas hydrate on its way upwards given proper conditions. Therefore, different
414 lithologies and tectonic settings in different natural hydrate sites can provide significant
415 insights for more comprehensive studies about hydrate morphologies which always show
416 localized characteristics.

417 Additionally, the laboratory experiments by Priest et al. (2006) indicate that
418 significant seismic attenuation can be caused by the squirt flow due to the adsorbed thin
419 water film between host grains and cementing hydrates. Combining the modeled
420 attenuation using the HBES model with our calculated attenuation from field data, we
421 suggest that squirt flow in the microporous hydrate could play a significant role in

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3 422 seismic attenuation over a broad frequency range in the south Hydrate Ridge. In this
423 study, the VSP attenuation is pretty low compared with the sonic attenuation, which may
424 result from its lower frequency. There are other possible causes, e.g., scattering energy
425 mixing with first arrivals due to strong heterogeneity in the gas hydrate zone, even over
426 shot distance (Tréhu et al., 2004). Moreover, inspired by previous studies mentioned
427 above (Dvorkin and Uden, 2004; Matsushima, 2006; Rossi et al., 2007; Nittala et al.,
428 2017), source-coupling of VSP and sonic equipment could also be a factor. However,
429 since the VSP data is limited and hydrate saturation is only up to about 40%. Hence, for a
430 more comprehensive understanding on the potential role that different seismic
431 frequencies play in regulating the seismic attenuation at gas hydrate sites, complementary
432 studies elsewhere that consider low frequency surface seismic (10 Hz) and/or higher gas
433 hydrate saturations are required.

434 Compared to the frequent investigation of P-wave attenuation of gas hydrate,
435 lesser quantifications were conducted with S-wave attenuation because the frame rigidity
436 and shear modulus are unaffected if gas hydrate simply fills the space with little grain
437 contact. Future work is needed to fill the gap through comparing different characteristics
438 of P and S-wave in different gas hydrate sites.

439 CONCLUSION

440 We have presented a case study using field seismic attenuation measurements and
441 rock physics modeling to investigate the effect of gas hydrate saturation and morphology
442 on seismic attenuation in the south Hydrate Ridge in the Cascadia Margin. We derived
443 seismic VSP and sonic attenuation from field data and interpreted the interrelation

444 between attenuation and gas hydrate saturation and morphology using rock physics
445 modeling, namely HBES hydrate model. The limited data on VSP attenuation shows
446 relatively small values compared to the sonic attenuations. The sonic attenuation shows a
447 scattered distribution with hydrate saturation, suggesting that attenuation is not only
448 controlled by gas hydrate saturation but also likely by morphology. The joint analysis of
449 seismic attenuation measurements and the HBES modeling results demonstrates the
450 theoretical possibility of multiple morphologies coexisting in the pores and the effect of
451 both gas hydrate saturation and morphology in the hydrate-bearing sediments. We
452 propose that in the south Hydrate Ridge (1) cementing hydrate may be predominant at low
453 hydrate saturation (<5%), whereas the effects of cementing and pore-filling hydrate
454 cannot be distinguished at relatively higher hydrate saturation; (2) gas hydrate
455 morphology may change with the gas hydrate saturation; (3) squirt flow is responsible for
456 the attenuation changes in the hydrate-bearing sediments at sonic frequencies. Overall,
457 this study provides insights into interpreting the seismic attenuation in the hydrate-
458 bearing sediments using theoretical rock physics models.

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460 All well log data in this study are available to download in the public data
461 repository of the Ocean Drilling Program
462 (<https://mlp.ldeo.columbia.edu/data/odp/leg204/1244E/>,
463 <https://mlp.ldeo.columbia.edu/data/odp/leg204/1247B/> and
464 <https://mlp.ldeo.columbia.edu/data/odp/leg204/1250F/>). VSP and sonic data processing
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473 APPENDICES

474 **Appendix A**

475 **Archie's relationship**

476 Archie's equation is used to calculate the hydrate saturation according to
 477 resistivity logging profile.

$$S = 1 - (R_t/R) \quad (A1)$$

$$R_t = aR_w\phi^{-j} \quad (A2)$$

478 where a , j are Archie parameters, ϕ is the porosity, R_t is the resistivity of formation, R_w
 479 is the resistivity of connate water and R is the recorded resistivity with depth. We choose
 480 $aR_w = 0.55 \Omega m$ and $j = 1.3$ based on Leg204 gas report (Tréhu et al., 2006a). Note that
 481 better saturation calculation could be obtained from recent work (Pan et al., 2019).

482 **Appendix B**

483 **Uncertainty analysis of results**

484 Attenuation can be calculated with centroid frequency shift method: $Q = \frac{\pi\sigma_i^2\Delta t_i}{\Delta f_i}$,
485 changes of the centroid frequency Δf_i , velocity, distance ($t=\text{distance}/\text{velocity}$) between
486 two receivers and the variance of centroid frequency (σ_i^2) can cause the uncertainty on the
487 attenuation calculations. Since variance is also associated with the frequency, and
488 distance can be measured very accurately, we only consider the uncertainty caused by the
489 frequency and the velocity. To this end, we apply the Monte Carlo method for calculating
490 the attenuation with three receiver pairs (receiver1-6, receiver1-7 and receiver1-8). We
491 randomly select the reasonable frequency range 1000 times, i.e., low frequency and high
492 frequency, which is the same for seismic traces at two receivers. Low frequency
493 randomly changes between 5 – 7 kHz while high frequency randomly changes between 9
494 – 10 kHz. Arrival time is determined by distance/velocity and velocity is assumed to have
495 an uncertainty of 7.5% (Baron and Holliger, 2010). Therefore, for each interval at each
496 well (1244E, 1247B and 1250F), we can get 3000 Q values (i.e., $1000 \times$ three receiver
497 pairs). To make sure there are enough data samples (>300) for analysis, we filter the
498 interval with more than 300 infinite Q values which are invalid.

499 The next step is to combine the selected Q values with hydrate saturation at the
500 same depth. In order to quantify the reliability of the results, we plot all the results with
501 their standard deviations shown as different colors (Figure B1), from which we can see
502 most of the results have the standard deviation lower than 0.1. In addition, higher
503 attenuation shows a higher standard deviation which means lower attenuation is much

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3 504 more reliable. However, the standard deviation cannot provide a reliable range of our
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505 calculated attenuations. We therefore apply the 95% confidence interval of the mean
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506 attenuation to quantify the range. For this purpose, we employ two approaches: (1) Since
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507 my sample size is large enough (>100) to use the z-distribution directly, we can simply
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508 use the formula: $\bar{x} \pm z \frac{\sigma}{\sqrt{n}}$, where z equals 1.96 after looking for 95% confidence interval;
9
509 \bar{x} is the sample mean, σ is the sample standard deviation and n is sample size; (2) To
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510 verify the 95% confidence interval from approach (1), we use 'bootstrap' to calculate the
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511 95% confidence interval again: at each depth, there are n (300<n<3000) valid Q values,
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512 we randomly choose 1000 Q values for 3000 times and calculate mean Q values each
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513 time, then we can get 3000 mean Q.

514 The mean Q are normally distributed, and the 95% confidence interval can be
515 easily calculated, which is the same as approach (1) and the results are shown in Figure
516 B2. The black dots in Figure B2 are mean attenuations after resampling selected
517 attenuations at each interval 1000 times, the error bar shows the 95% confidence interval
518 for each mean.

519 To see the mean attenuation confidence interval in each well and if the resampled
520 results show a similar pattern with Figure 13, we plot these dots in different colors in
521 Figure B3, which still show the same trend and distribution as Figure 13. The modeled
522 results are also plotted in Figure B4, the upper boundary (cementing hydrate) and lower
523 boundary (60% cementing hydrate) restrict the scattering results from the field, the same
524 as Figure 13.

525 **Appendix C**

526 **Parameters in the HBES model**

527 The description of input parameters in the HBES model are shown in Tables C1
528 and C2.

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- 728

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 781 respectively.

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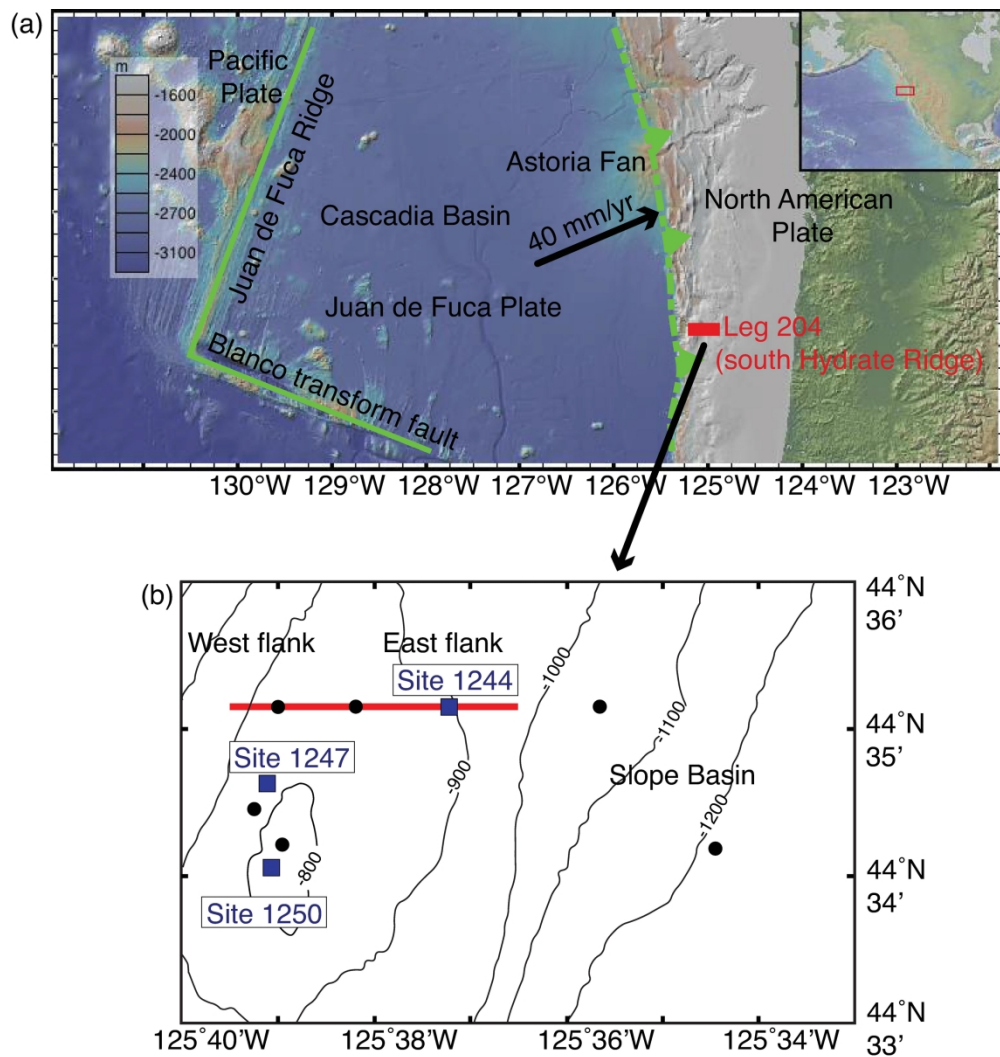


Figure 1. Location of Leg 204 on the South Hydrate Ridge on the Cascadia continental margin which overlies the subduction zone of Juan de Fuca plates thrusting beneath the North American plate. Seismic data is from Sites 1244, 1247 and 1250; black circles represent other drillings; red line shows the surface seismic line crossing site 1244 (Modified after Leg 204 Report, 2006).

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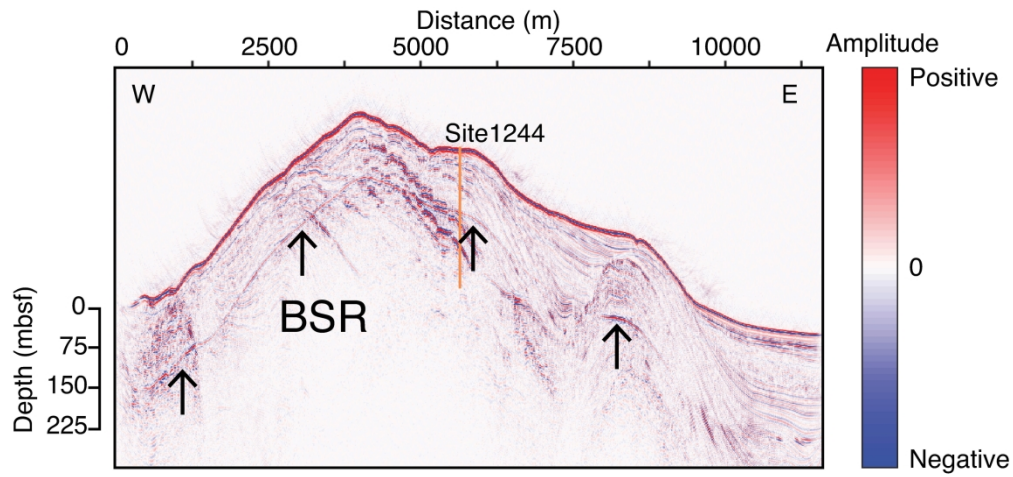


Figure 2. Seismic section across Site 1244 where BSR (indicated by arrows) can be recognized at 100 – 150 m depth.

289x137mm (300 x 300 DPI)

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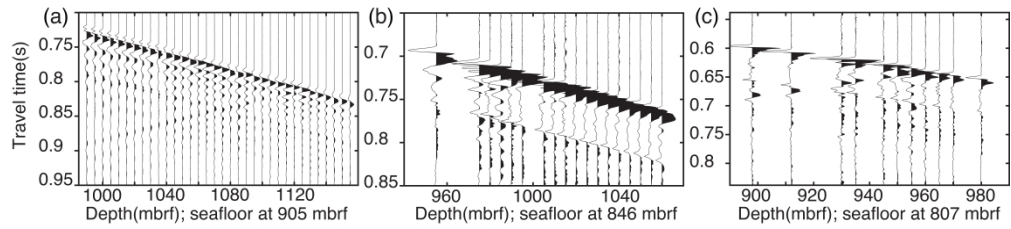


Figure 3. Raw z-component near-offset (55 m) VSP seismic data in (a) 1244E; (b) 1247B; (c) 1250F.

1192x266mm (600 x 600 DPI)

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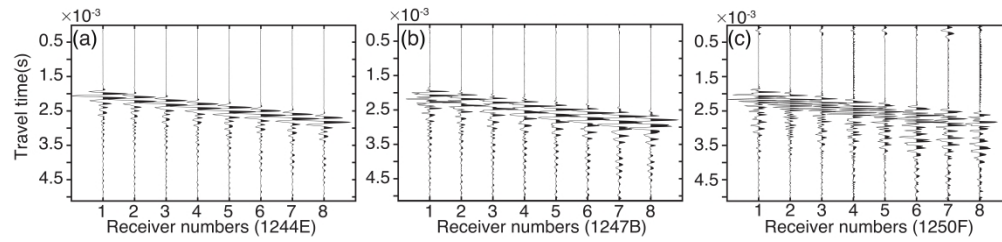


Figure 4. Raw sonic logging data in (a) 1244E; (b) 1247B; (c) 1250F.

1202x290mm (600 x 600 DPI)

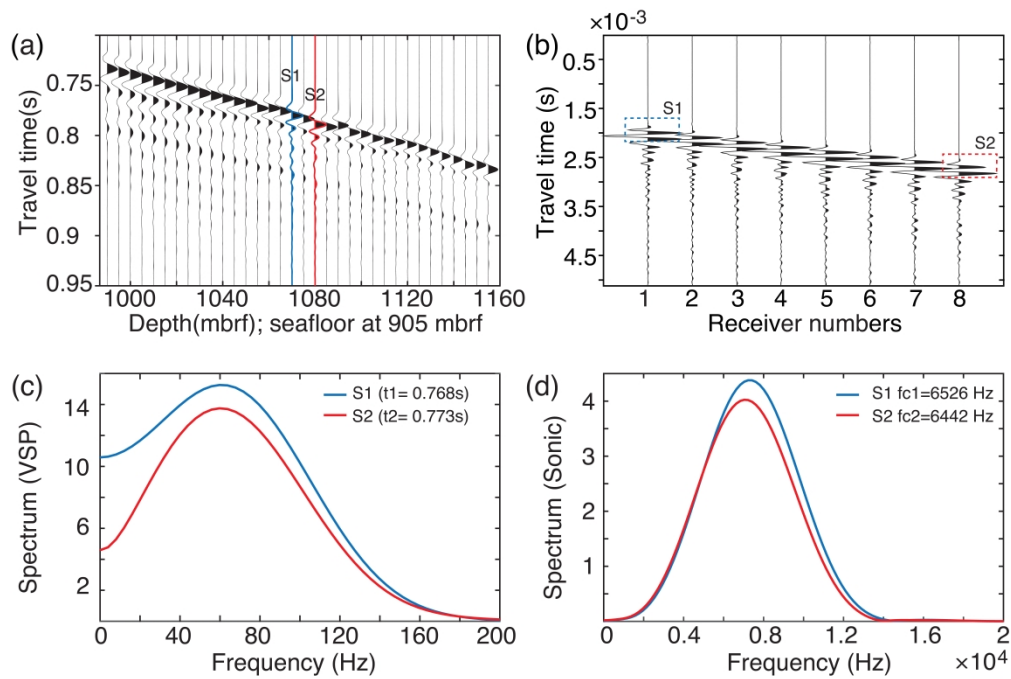


Figure 5. (a) Z component of near offset (55 m) VSP downgoing recording after pre-processing; (b) Monopole seismic traces obtained from sonic logging; (c) Spectrum of VSP traces for shallower trace (blue) and deeper trace (red); (d) Spectrum of sonic traces at receiver 1 (blue) and receiver 8 (red).

1242x839mm (600 x 600 DPI)

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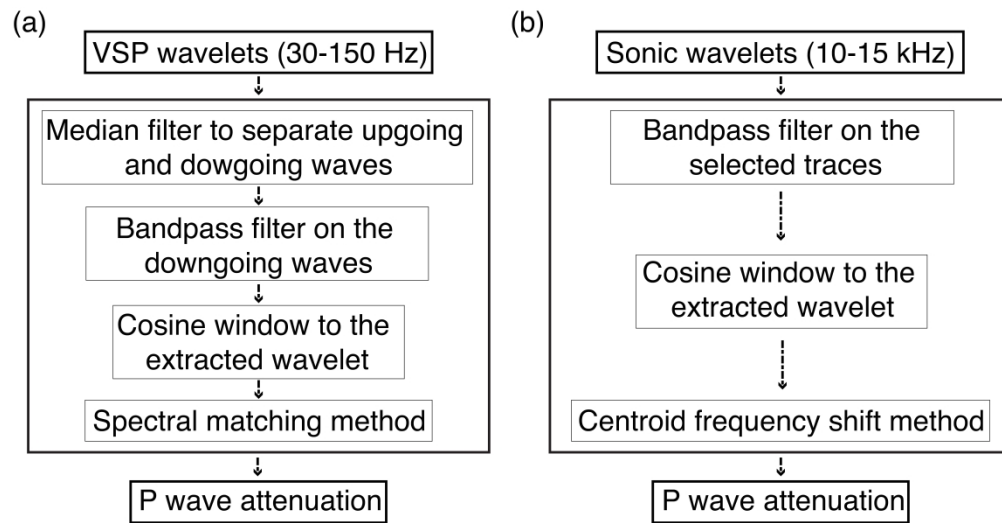


Figure 6. The data processing procedure for (a) the VSP and (b) the sonic logging to estimate the seismic attenuation.

737x380mm (600 x 600 DPI)

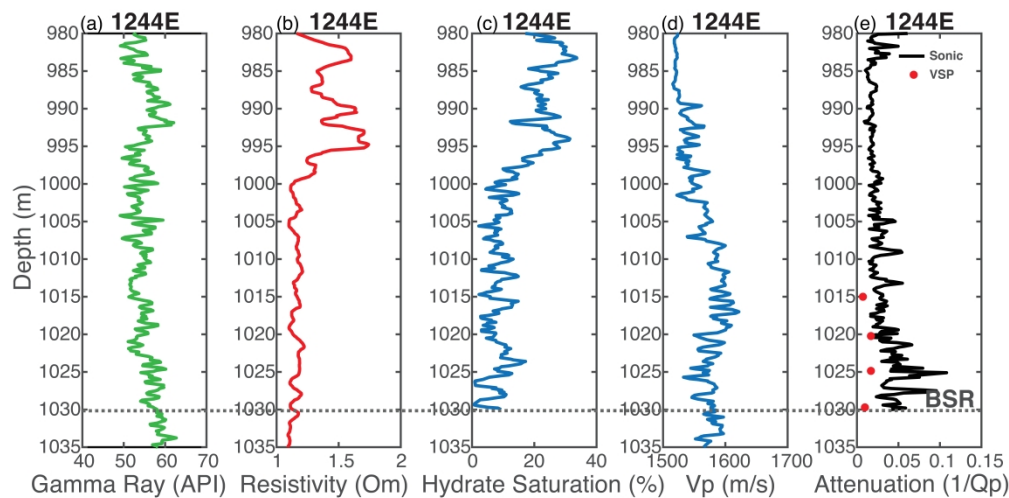


Figure 7. 1244E well logging data: (a) gamma ray, (b) resistivity, (c) gas hydrate saturation derived from the log of resistivity using Archie's equation, (d) P-wave velocity and (e) attenuation. The dashed line represents BSR.

882x435mm (600 x 600 DPI)

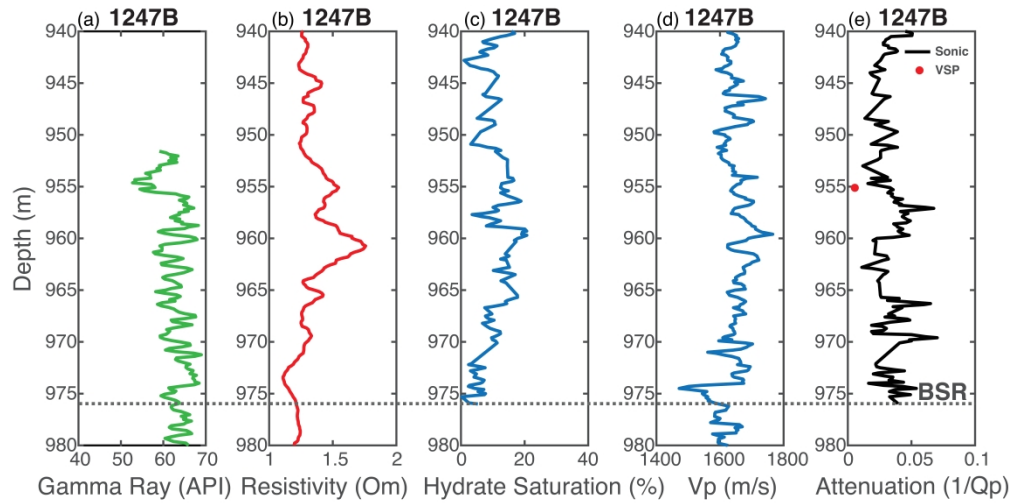


Figure 8. 1247B well logging data: (a) gamma ray, (b) resistivity, (c) gas hydrate saturation derived from the log of resistivity using Archie's equation, (d) P-wave velocity and (e) attenuation. The dashed line represents BSR.

888x443mm (600 x 600 DPI)

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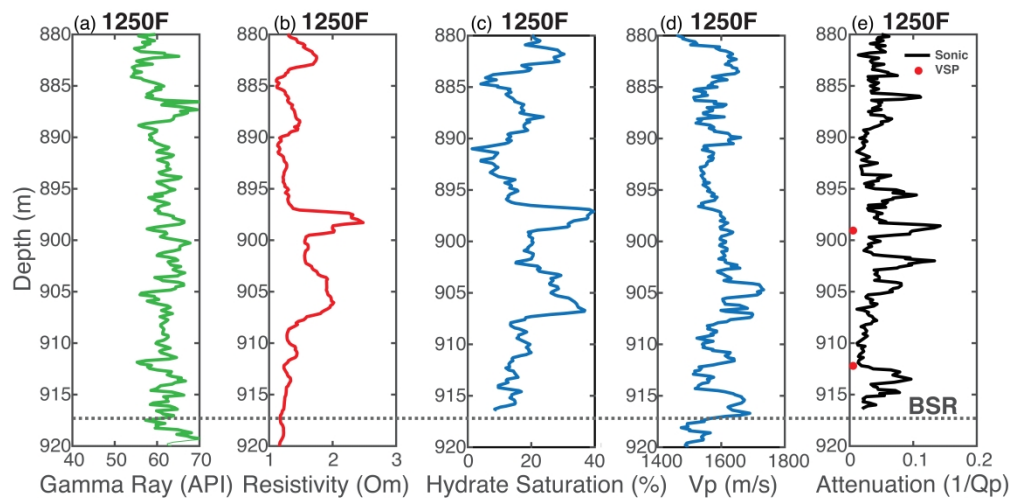


Figure 9. 1250F well logging data: (a) gamma ray, (b) resistivity, (c) gas hydrate saturation derived from the log of resistivity using Archie's equation, (d) P-wave velocity and (e) attenuation. The dashed line represents BSR.

890x441mm (600 x 600 DPI)

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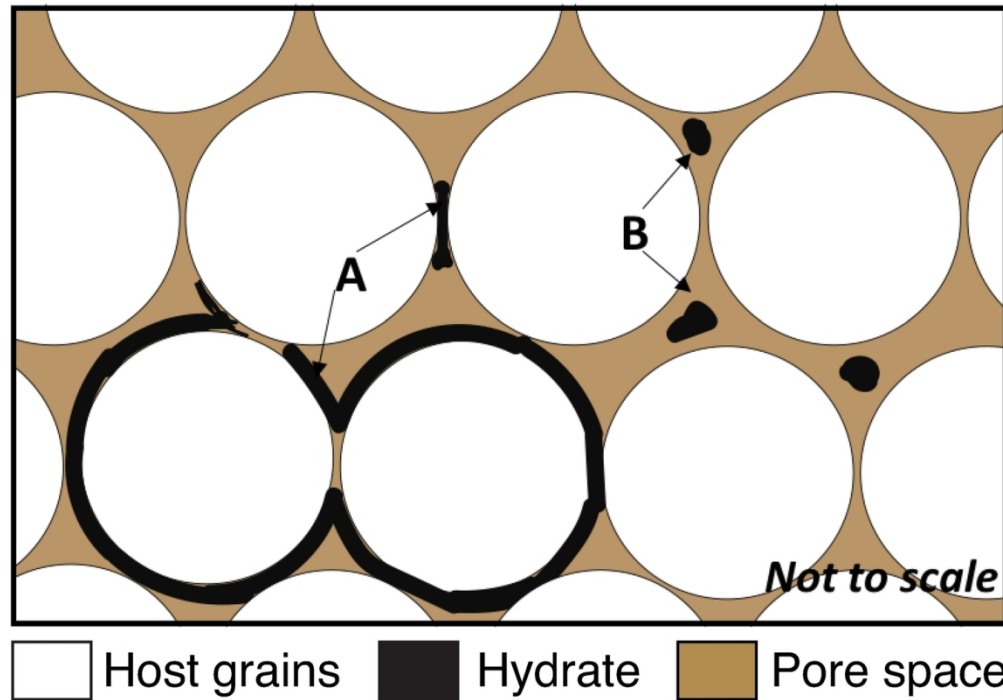


Figure 10. Idealized conceptual illustration of the microstructure of hydrate-bearing sediments. A = Cementing at grains contact; B = Pore-filling hydrates framework.

134x95mm (300 x 300 DPI)

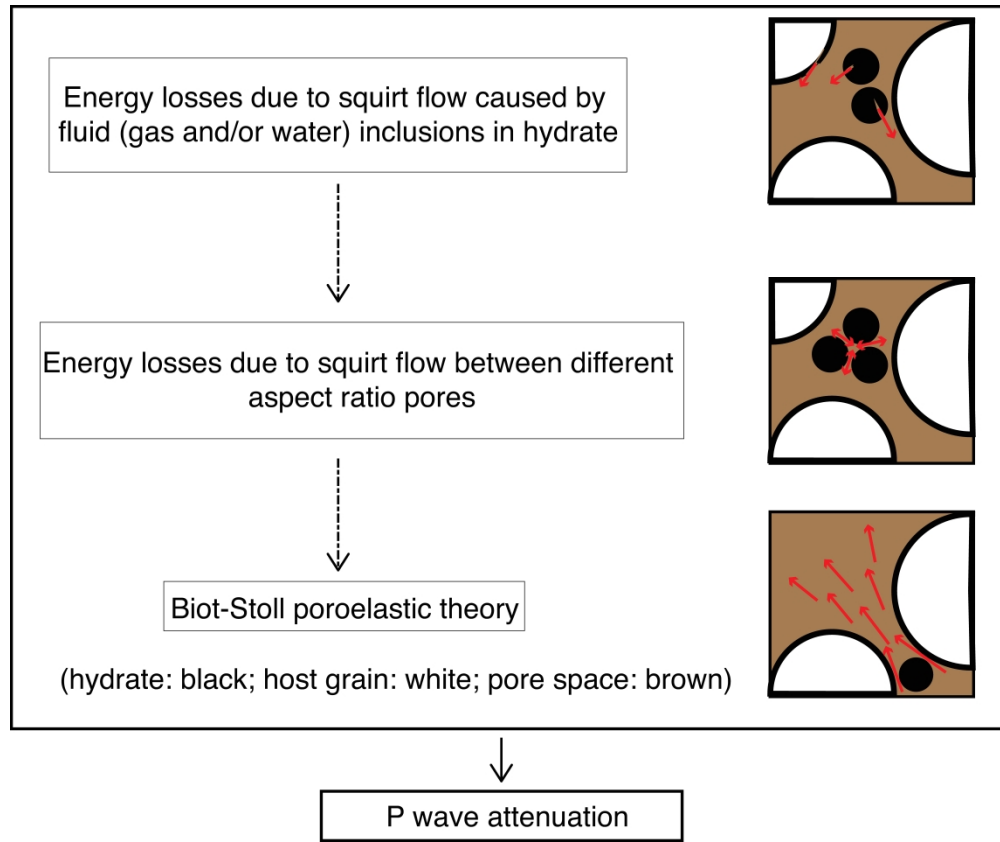


Figure 11. Simplified workflow of the Hydrate-Bearing Effective Sediment (HBES) model. More detailed descriptions and procedures can be referred to Marin-Moreno et al. (2017).

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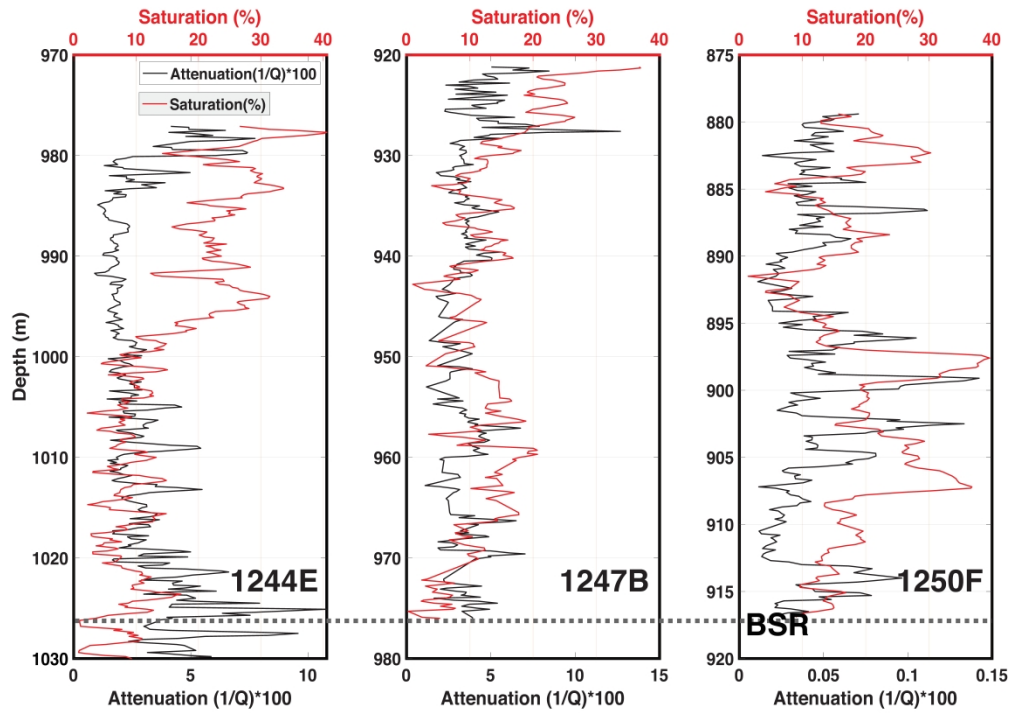


Figure 12. Comparison between the seismic attenuation and gas hydrate saturation above the BSR (grey dotted line) at (a) 1244E, (b) 1247B and (c) 1250F, respectively.

626x443mm (600 x 600 DPI)

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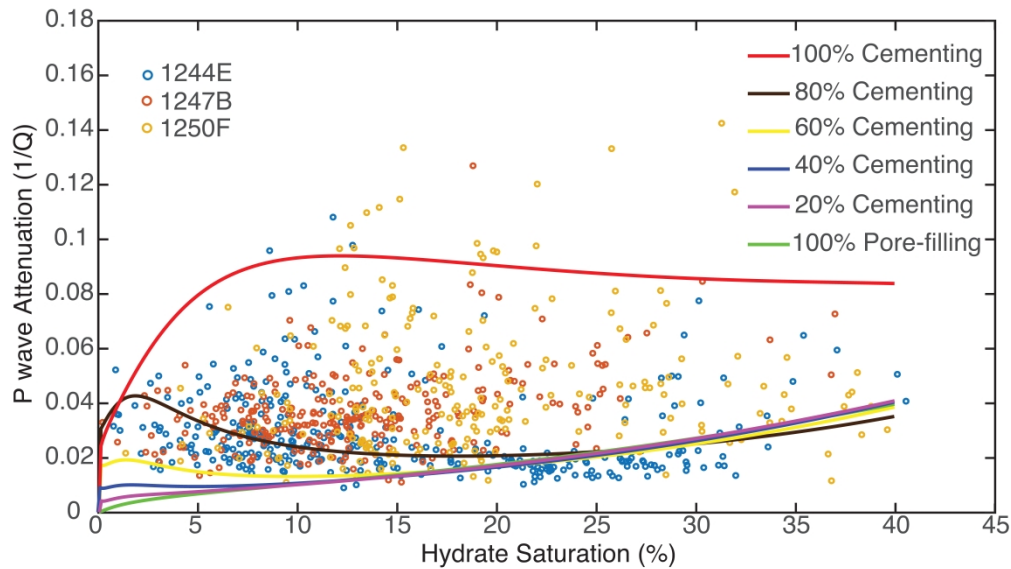


Figure 13. Measured attenuation (dots) and modeled attenuation (lines) as a function of gas hydrate saturation of 1244E (blue circles), 1247B (red circles) and 1250F (yellow circles), respectively. Note that the summation of pore-filling and cementing hydrate adds 100%.

841x479mm (600 x 600 DPI)

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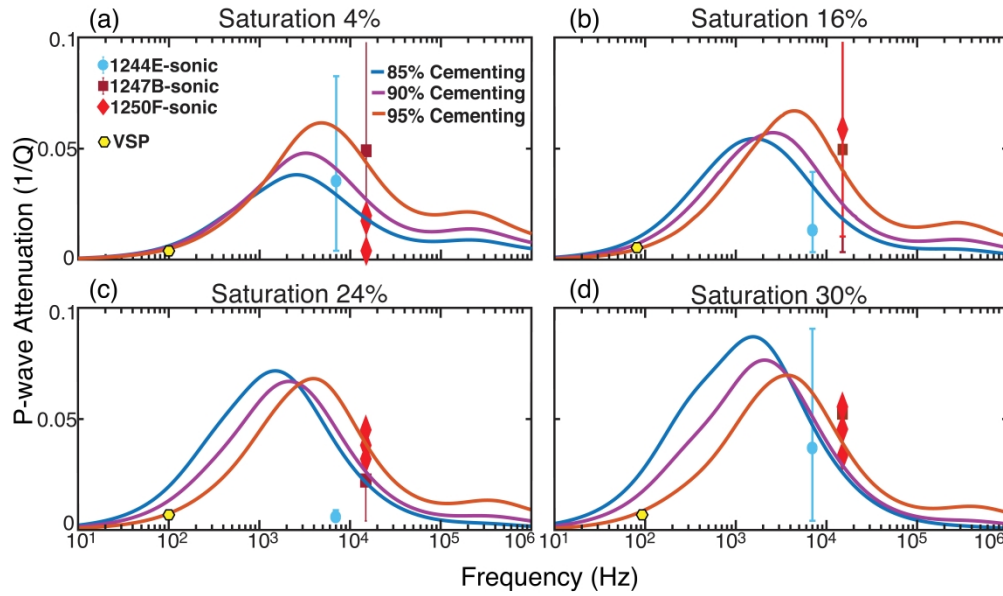


Figure 14. Measured attenuation (dots) and modeled attenuation (lines) of P wave as a function of frequency at (a) 4%, (b) 16%, (c) 24% and (d) 30% gas hydrate saturation. Note that 85%, 90% and 95% cementing refer to the fraction of cementing of the total hydrate saturation, i.e., 85% (or 90%/95%) \times hydrate saturation.

817x484mm (600 x 600 DPI)

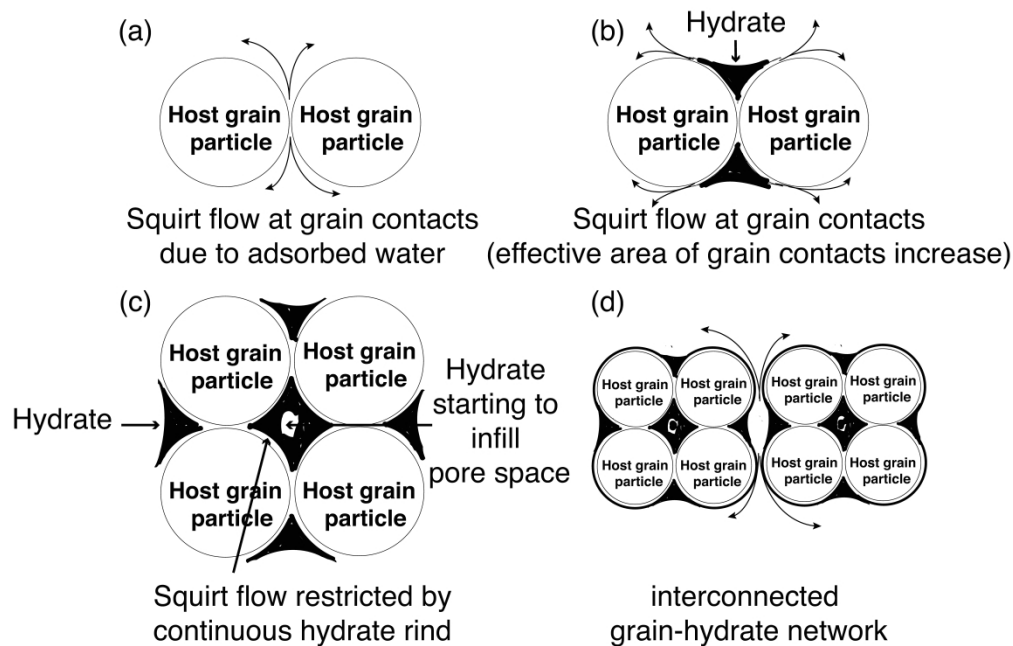


Figure 15. Conceptual model of cementing hydrates growing at grain contacts with increasing gas hydrate saturation. (a) Host grains without hydrates; (b) Hydrates growing at grain boundaries; (c) Enveloped grains with hydrates; (d) Interconnected grain-hydrate network (Modified after Priest et al., 2006).

737x477mm (600 x 600 DPI)

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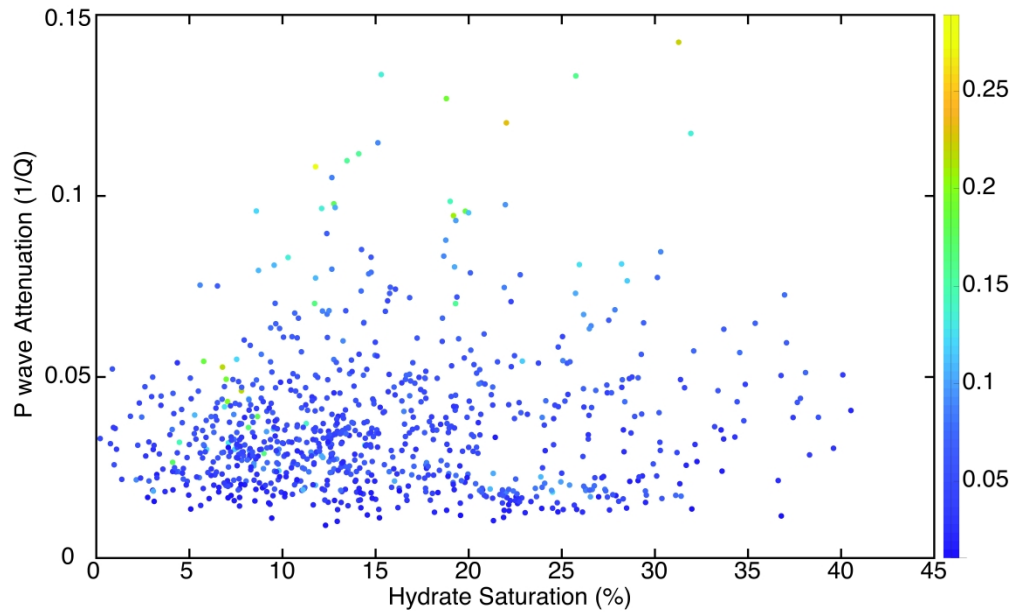


Figure B1. Scattering plot of the seismic attenuation and gas hydrate saturation above the BSR from 1244E, 1247B and 1250F. Different colors represent the standard deviation of P wave attenuation.

881x538mm (600 x 600 DPI)

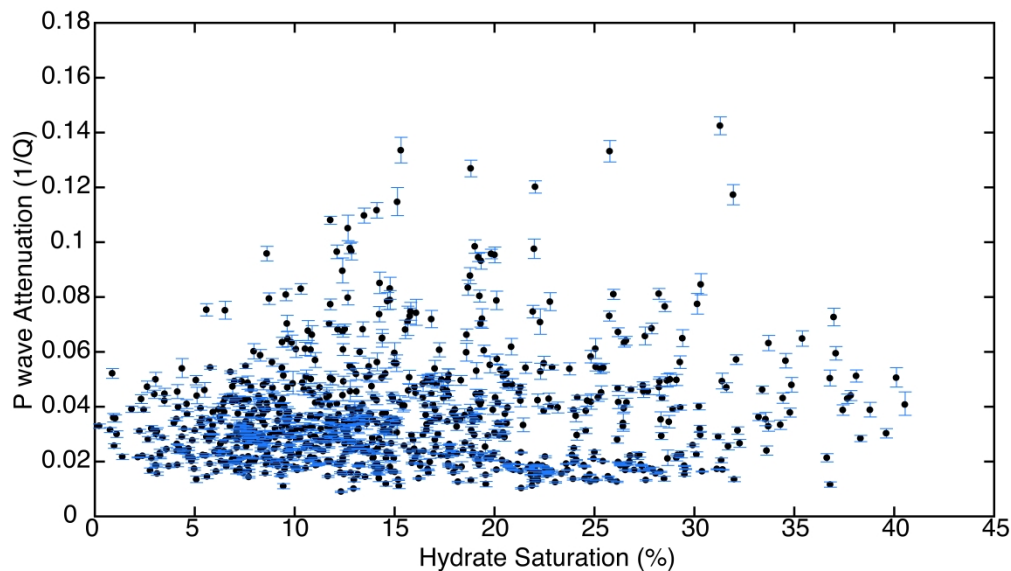


Figure B2. Measured mean attenuation (dots) with error bar showing the 95% confidence interval as a function of hydrate saturation.

838x479mm (600 x 600 DPI)

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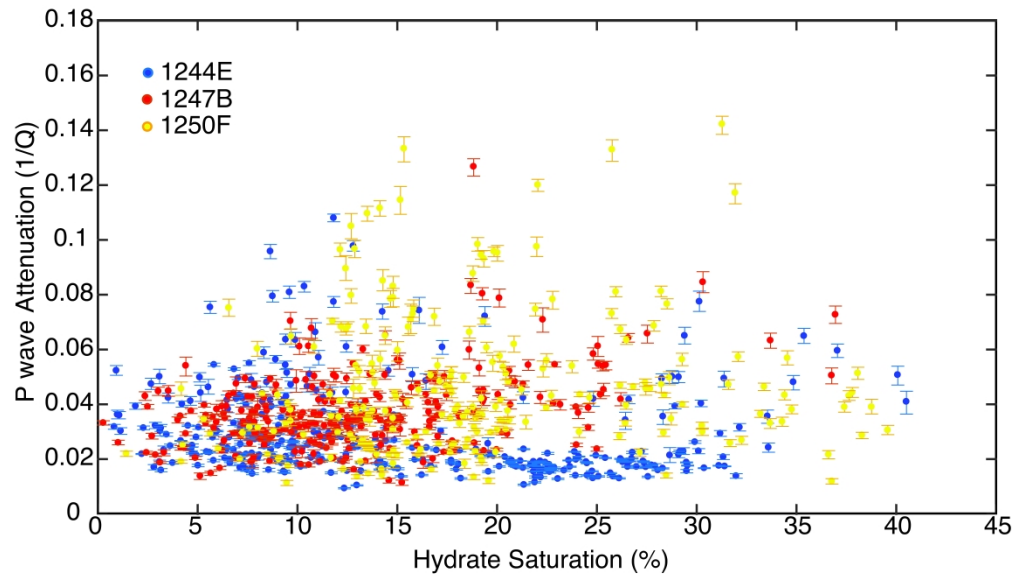


Figure B3. Measured mean attenuation (dots) with error bar showing the 95% confidence interval as a function of gas hydrate saturation for 1244E, 1247B and 1250F, respectively.

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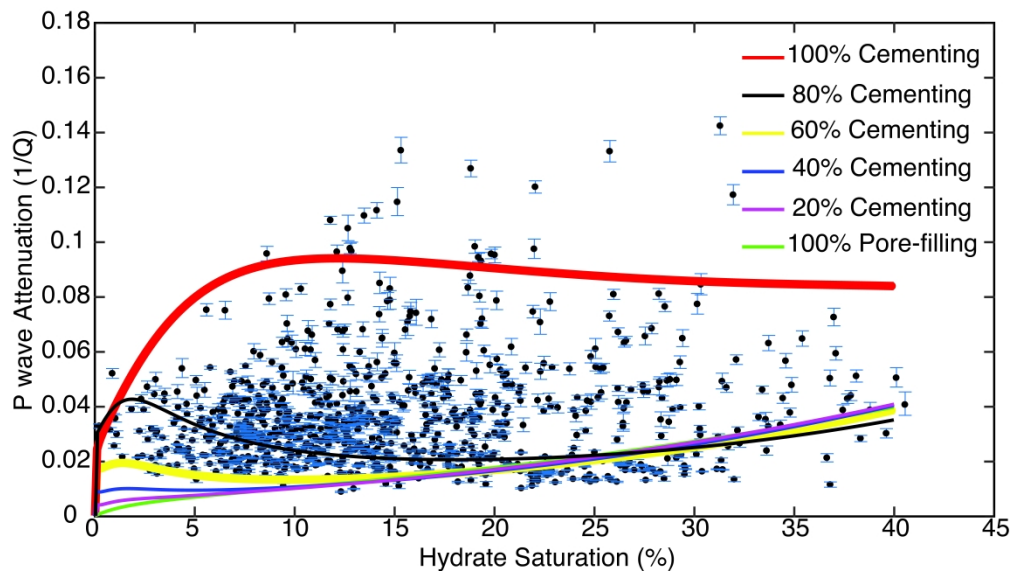


Figure B4. Measured mean attenuation (dots) with error bar showing the 95% confidence interval and modeled attenuation (lines) as a function of gas hydrate saturation.

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Table 1. Published and this study's seismic attenuation (Q value is the inverse of the attenuation) for hydrate-bearing sediments.

Location	Frequency (Hz)	Q value	Saturation
Blake Ridge, US (Wood et al., 2000)	20-150 (Single-channel)	200-300	3%-4%
	10-120 (VSP)	>20	
	10k-20k (Sonic logging)	No data	
Mallik field, Canada (Dvorkin & Uden 2004, Bauer et al., 2008)	10k-15k (Sonic logging)	6-20	60%-80%
	10-200 (VSP)	7-13	
	150-500 (Crosshole)	5-11	
Nankai Trough, Japan (Matsushima et al., 2006)	10k-20k (Sonic logging)	25-100	20%-30%
	30-110 (VSP)	>100	
Western Svalbard, Norway (Madrussani et al., 2010)	20-200 (3D seismic)	150-200	6%-10%
Krishna-Godavari, India (Nittala et al., 2017)	8k-15k (Sonic logging)	344.82	50%-80%
	5-50 (3D seismic)	80-81	
	5-120 (Multi-channel at 01-10 site)	160-320	
Gulf of Mexico (Wang et al., 2017)	8k-24k (Sonic logging)	50	40%-75%
Cascadia, Oregon (this study)	25-200 (3D seismic)	>50	8%-10%
	30-150 (VSP)	>80	
	3k-15k (Sonic logging)	>90	

Table C1. Fixed input parameters used in the HBES model.

Parameter	Value	Unit	Reference
Confining pressure (P_c)	1.14×10^7	Pa	
Pore fluid pressure (P_p)	0.92×10^7	Pa	
Temperature (T)	9	$^{\circ}C$	
Hydrate bulk modulus (K_H)	7.9×10^9	Pa	(Best et al., 2013)
Hydrate shear modulus (G_H)	3.3×10^9	Pa	(Best et al., 2013)
Hydrate Poisson's ratio (ν_H)	0.32		(Marín-Moreno et al., 2017)
Hydrate Density (ρ_H)	925	Kg/m^3	(Helgerud et al., 2009)
Methane bulk modulus (K_{CH_4})	$K_{CH_4}(P_p, T)$	Pa	(Millero et al., 1980)
Methane density (ρ_{CH_4})	$\rho_{CH_4}(P_p, T)$	Kg/m^3	(Millero et al., 1980)
Methane viscosity (μ_{CH_4})	$\mu_{CH_4}(P_p, T)$	$Pa s$	(Millero et al., 1980)
Methane irreducible saturation (S_{rCH_4})	0.02		(Reagan and Moridis, 2008)
Grain bulk modulus (K_s)	36×10^9	Pa	(Ecker et al., 2000)
Grain shear modulus (G_s)	45×10^9	Pa	(Ecker et al., 2000)
Grain Poisson's ratio (ν_s)	0.062		(Marín-Moreno et al., 2017)
Grain density (ρ_s)	2650	Kg/m^3	(Ecker et al., 2000)
Grain diameter (d_s)	1×10^{-4}	m	(Best et al., 2013)
Grain coordination number (n)	8.5		(Ecker et al., 2000)

Water bulk modulus (K_W)	$K_W(P_p, T)$	Pa	(Setzmann and Wagner, 1991)
Water density (ρ_W)	$\rho_W(P_p, T)$	Kg/m^3	(Setzmann and Wagner, 1991)
Water salinity (s)	3.5	$\% wt$	
Water irreducible saturation (S_{rW})	0.2		(Reagan and Moridis, 2008)
Porosity without hydrate (φ_0)	0.40		(Daigle et al., 2015)
Critical porosity (φ_c)	0.38		(Best et al., 2013)
Intrinsic permeability without hydrate (K_0)	10^{-13}	m^3	(Daigle et al., 2015)

Table C2. Case dependent input parameters used in the HBES model.

Parameter	Value	Unit	Reference
Aspect ratio of inclusions containing methane or water	$\alpha\{iCH_4, iW\}$		
Concentration of cementing hydrate (<i>C</i>) and pore-filling hydrate (<i>PF</i>)	$c\{C, PF\}$		
Concentration of inclusions in hydrate	c_i	<i>Pa</i>	
Frequency	<i>f</i>	<i>Hz</i>	
Porosity	φ		
Saturation of hydrate, methane, and water in the pore space	$S\{H, CH_4, W\}$		

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DATA AND MATERIALS AVAILABILITY

Data associated with this research are available and can be obtained by contacting the corresponding author.