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On the effects of walking speed, crowd density and human-to-source distance on pollutant dispersion in indoor spaces

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Abstract

The effects of walking speed, crowd density and human-to-source distance on pollutant dispersion in two scaled room models are investigated using simultaneous planar laser-induced fluorescence and particle-image velocimetry techniques. For a small 3 m high room, where the length-scales of the people and room are comparable, the walking motions significantly influenced the macro room mean flow patterns. This has a strong effect on the scalar dispersion properties as the magnitudes of the advective scalar fluxes are often comparable or larger than the turbulent scalar fluxes. As such, the scalar dispersion properties are case specific. For a large 9 m high room, the walking motion influenced only the local mean flow field. The increase in walking speed and crowd density improves the efficiency in which the scalar is transported and mixed with fresh ambient fluid out of the measurement plane (i.e. along the direction of the motion), leading to scalar-free zones observed on the opposite side of the room from the ventilation outlet. The area of the scalar-free zone increases with an increase in the walking speed and crowd density. The advective scalar fluxes are more sensitive to the human motion than the turbulent components, and as the mixing efficiency improves, the advective fluxes show a greater weakening with increased distance from source. The concentration PDFs in the near-source region can be described by the exponential function where the expected value at the 99% percentile can be derived as $C_{99}/c'_{rms}=4.61$, which agreed well with the experimental measurements of 4.1 to 5.9.

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Keywords: Indoor airflow, crowd effects, scalar dispersion, indoor air quality, risk assessment

1. Introduction

Understanding scalar dispersion in indoor spaces is an essential aspect of improving public health and well-being, as there are many infectious viral and bacterial diseases that can be transmitted through the airborne route. This includes the SARS-CoV-2 virus which is responsible for the COVID-19 pandemic, the common cold and influenza, tuberculosis and chickenpox. A recent multidisciplinary review has presented strong evidence of the association between indoor airflow and the transmission of airborne diseases [1] which highlights the importance of, and the need to, better understand indoor airflows.

Indoor airflow is complex due to the lack of clearly defined air flow paths and its sensitivity to a wide range of different parameters. In the absence of strong pressure variations (i.e. naturally ventilated empty rooms), thermal stratification tends to determine the flow patterns, with the three main forms of stratification being: ‘stable stratification’, ‘unstable stratification’ and the ‘gravity current’ [2]. For mechanically ventilated rooms, the air changes per

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hour (ACH) is a particularly important parameter, and its strong influence on scalar mixing and indoor air quality is widely recognised [3, 4, 5]. There are many other parameters that can influence indoor airflow and air quality. Recent studies have examined the effects of the relative position of inlet and outlet vents in convective flows [6], the effects of temperature and relative humidity on exposures to airborne pathogens [7], the seasonal variations in carbon dioxide levels in 36 naturally ventilated classrooms in the UK monitored throughout 2021 [8], and the impact of human movements on particle dispersion and resuspension [9].

In scenarios where hazardous air pollutants are released in crowded indoor spaces, which can include airborne pathogens or toxic chemical gases, our understanding of scalar dispersion is crucial to support accurate assessments of exposures and pre-planned decisions that can mitigate the consequences and save lives. These decisions can be in the form of official instructions for people to evacuate from the incident site or to seek shelter [10]. The effects of human activities may have a significant influence on the dispersion and concentrations of the pollutants, and is therefore, an important factor that needs to be considered to better inform decisions and to avoid unintended effects.

There have been several studies focusing on the effects of walking on scalar dispersion. Wang & Chow [11] numerically studied the effects of walking motion on expiratory droplets, and observed that higher walking speeds reduce the overall amount of suspended droplets by increasing the removal rate of pollutants through the ventilation outlet. Tan et al. [12] observed higher walking speeds of medical staff led to more significant secondary airflow and therefore, lower number of particles settling on a patient in an isolation ward. Mingotti et al. [13] experimentally investigated the repeated passage of a person along a corridor and focused on the effects of the wake-driven mixing on the turbulent diffusivity, which is a parameter that describes the rate of turbulent scalar mixing. Empirical observations suggest the turbulent diffusivity coefficient is a function of the characteristic length-scales of the corridor and the person, as well as the average frequency of the person moving along the corridor.

The impact of walking on scalar dispersion are often considered in the context of other key parameters, such as the room airflow patterns or the proximity to the inlet/outlet/region of interest. Hang et al. [14] investigated the effects of ACH, ventilation design, and walking motion of realistic vs idealised (rectangular block) body geometries. Their results showed the effects of ventilation design can be more significant in influencing gaseous pollutant spread than the human motion. Lv et al. [15] observed that different human walking patterns can lead to complex interactions between the thermal plume from a kitchen hob and the wake flow generated by the walking motion. Al Assaad et al. [16] simulated a static and a moving thermal manikin model, and showed the turbulence produced by a moving model can create local mixing effects, which enhanced the entrainment of contaminants from the room into a personalised ventilation jet. Liu et al. [17] performed Reynolds-averaged Navier-Stokes (RANS) simulation on a circulating nurse walking in an operating room to study its effect on the airflow and pollutant concentration. Although higher pollutant concentration was observed near the instrument tables, the airflow above the operating table which was slightly further from the circulating nurse was not affected, thus indicating that the walking-induced turbulence and mixing effects are localised.

The airflow is much more complex in crowd scenarios but has many important applications particularly in public spaces such as transport platforms. Zhao et al. [18] performed simulations to investigate the effects of human movements on droplet and aerosol transport in an airport terminal. When the walking direction is perpendicular to the room flow, the effect of the crosswind led to increased droplet dispersion perpendicular to the walking direction, as compared to the perpendicular dispersion when walking into or against the room flow. Marlow et al. [19] used a multidisciplinary simulation tool based on the lattice-Boltzmann method, a wake model and a human agent behavioral model to simulate an evacuation scenario. They observed higher concentrations of air pollutants near the room exit, due to pollutant entrainment by the wake flow of people leaving the room. Wu et al. [20] observed a much more extensive pollutant dispersion range in the multi-person walking configuration (as compared to the single-person), for aerosols and droplets exhaled by a moving person in the middle of the room.

Besides walking, there are several studies that have looked into the effects of other human-related factors on indoor scalar dispersion. For example, Hathway et al. [21] investigated the effects of door motion on mass exchange between two rooms and observed large-scale vortex and jet-like motions governing the transport and mixing processes, and a linear relationship between door hold-open time and volume flux between the two rooms. Kamar et al. [22] investigated turning movements of medical staff in an operating room. A bent forearm was observed to change the flow field more significantly than a straight forearm and increased the concentrations of particle settling on the patient. Abouelhamd et al. [23] investigated the effect of crowd density on the human breathing zone for social distancing applications. As the crowd density decreases from 3 people occupying one square meter (ppl m^{-2}) to 0.25 ppl m^{-2} ,

64 the horizontal propagation of the breathing zone can increase by up to three-fold.

65 To better understand the current trends of indoor airflow research, here, we discuss key takeaways from selected
66 review articles [24, 25, 9]. The advancements in the fundamental understanding of indoor airflow in the past few
67 decades have mainly focused on relatively idealistic room geometries [24, 25]. For realistic room flows, the occupied
68 regions of the room tend to have weak forced advection [26], which can lead to Reynolds number (Re) dependent flow
69 in some regions of the room despite the flow at the ventilation inlet having already achieved Re independence [27].
70 The complexity of indoor room flow is well documented and different flow dynamics (such as jet flows, recirculation
71 zones, laminar or fully turbulent flows) can exist in different regions of the room simultaneously [26, 28, 25]. As a
72 result, it can be challenging for Computational Fluid Dynamics (CFD) methods to model the problem. Kek et al. [9]
73 recently reviewed human movements in indoor airflow literature, and found RANS and URANS simulations are still
74 the most popular methods in indoor airflow studies involving either static or dynamic people. Many studies apply the
75 dynamic meshing technique to account for human movements, and treat the humans as the pollutant source (i.e. either
76 as a direct source of aerosols/droplets and/or an indirect source due to particle resuspension). Despite the popularity
77 of CFD approaches, experimental results are still considered vital as a validation dataset due to the complexity of
78 the airflow in a room. Nonetheless, accurate experimental results from full-scaled chamber could be challenging to
79 collect in dynamic airflow measurements, as human intrusion during the measurements would introduce uncertainties
80 [9].

81 The research objective of this study is to investigate the effects of walking speed, crowd density and human-to-
82 source distance on the full-field concentration statistics and scalar transport mechanisms of pollutants from a point-
83 source in two scaled indoor spaces (i.e. a small 3 m ceiling room and a large 9 m ceiling room). Simultaneous particle-
84 image velocimetry (PIV) and planar laser-induced fluorescence (PLIF) experiments on two scaled room models were
85 performed, which enabled access to completely non-intrusive and full-field concentration and velocity measurements.
86 To the best of the authors' knowledge, there are currently very limited studies with a similar research objective, and
87 this may be one of the very first experimental studies that offers a unique perspective on the full-field concentration
88 statistics (which include the peak-to-mean concentration and concentration probability density functions (PDFs)) and
89 the covariance of the concentration and velocity (i.e. turbulent scalar fluxes). As high-fidelity experimental dataset
90 are scarce, we hope that our open access dataset will be beneficial to the indoor airflow community.

91 In section 2, we present the description of our experimental setup and test parameters. Section 3 discusses the
92 effects of walking speed, crowd density and human-to-source distance on the concentration statistics. In section
93 4, we discuss the concentration-velocity statistics and the scalar transport mechanisms. Section 5 presents the key
94 conclusions of our study.

95 2. Experimental methodology

96 2.1. Setup and data acquisition

97 The experiments were performed in the University of Southampton's Boldrewood Campus Recirculating Water
98 Tunnel (flume) facility. A schematic of the experimental setup is shown in figure 1(a). The room model was installed
99 in an upside-down configuration to enable the mechanical traverse system to actuate the cylinder models, which
100 represent people, from the top of the flume. The room model was fully submerged, spanned across the entire width of
101 the flume ($W = 1200$ mm), and flow was driven by the flume's pumps through a channel of height $h_{channel} = 40$ mm.
102 After the inlet channel, the flow entered the room model through 6×2 $h_{inlet} = 60$ mm diameter round ventilation
103 inlets with 80 mm diameter round deflector plates. The ventilation inlets were spaced at 180 mm from each other as
104 shown in figure 1(b). The flow exited the room through a one-sided ground level opening (with height h_{exit}), which is
105 representation of an open doorway as shown in figure 1(c). The motivation and design of the room geometry can be
106 found in an earlier study [29] which readers can refer to for more details.

107 Figure 1(c) shows the key dimensions of the room model, which was machined entirely out of clear acrylic to allow
108 optical access, and covered with matt black vinyl tape to reduce laser reflections. The design and key dimensions of
109 the ventilation inlet are shown in figure 1(b), which was also machined using clear acrylic. An image of the partially
110 assembled room is shown in figure 2(a), where the top plate (i.e., room floor) was lifted to show the ventilation inlet.
111 We investigated two different room designs in the experiments. Room A is a 20:1 full-to-model scale room that has
112 $h_{exit,A} = 100$ mm and represents a 3 m ceiling room that is typical of small offices or houses. Cylinders in room A

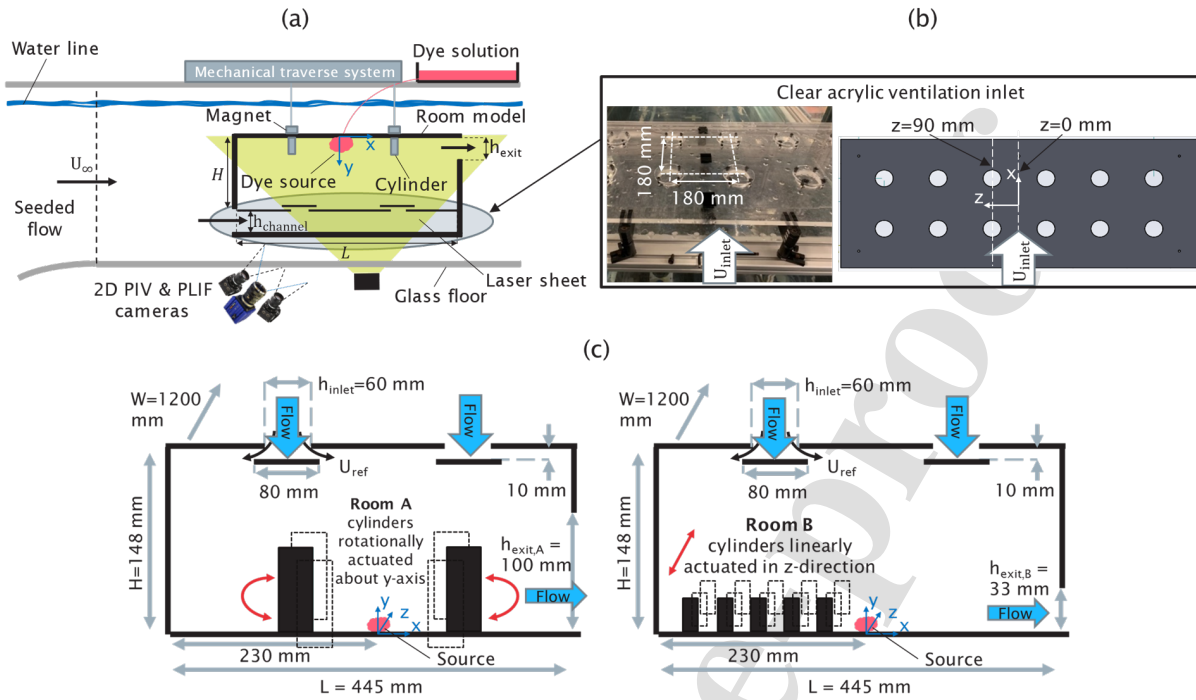


Figure 1. Schematic of (a) the experimental setup with room model in an upside down configuration, (b) the ventilation inlet grid and (c) key dimensions of room A (left) and room B (right).

113 which were used as a proxy for people have a height of 75 mm and diameter of $D_A=25$ mm. Room B is a 60:1
 114 full-to-model scale room that has $h_{exit,B}=33$ mm and represents a 9 m ceiling room that is typical of bigger rooms such
 115 as an auditorium or a lecture hall. Cylinders in room B have a height of 25 mm and diameter of $D_B=8$ mm.

116 The patterns of people's movement in indoor spaces are linked to behavioral patterns and would therefore be dep-
 117 endent on specific applications. Valid motions can range from static, if people are resting, to random directional
 118 motion, or the convergence of people at a common location in an evacuation scenario. Complex behaviour can also
 119 emerge, such as lane formation in the flows of people in opposing directions [30]. In this study, complex behaviour
 120 patterns are not taken into account, due to the complexity of the experimental setup, with people's movement simpli-
 121 fied into two main types of motion.

122 Room B cylinder arrays (figures 2(d,f)) were linearly actuated using the flume's linear z-direction mechanical
 123 traverse system as shown in figures 1(c) and 2(b). Their movement was prescribed with a triangular displacem-
 124 ent-time waveform with peak-to-peak amplitude of 160 mm (i.e. $20D_B$) for all test cases in this study. Room B cylinder
 125 arrays occupy only one side of the room, on the opposite side of the room to the outlet, as shown in the schematic in
 126 figure 1(c). This is of interest as it represents a realistic dilemma where a decision has to be made on whether people
 127 should evacuate through the scalar source to the exit on the other side of the room, or to seek shelter on their side of
 128 the room. Note that the side of the room closer to the exit is empty as we have assumed people close to the room
 129 outlet would always choose to evacuate.

130 Room A cylinder pairs (figures 2(d,e)) were actuated with a stepper motor to achieve rotational motion about the y-
 131 axis of the point source as shown in figures 1(a) and 2(c). This was prescribed a triangular angular displacem-
 132 ent-time waveform with peak-to-peak amplitudes of 60° , 20° and 19.2° at constant cylinder-to-source distances of $2D_A$, $4D_A$
 133 and $6D_A$ respectively. All of these correspond to a peak displacement of 50 mm in the z-direction. In the rotational
 134 motion case, the positions of the two cylinders are always equidistant from the scalar source in the x-direction.

135 The cylinders in the room model cannot be actuated by machining holes in the top plate (room floor) as it would
 136 lead to tracer dye escaping the room. To circumvent this problem, two strong round neodymium magnets were

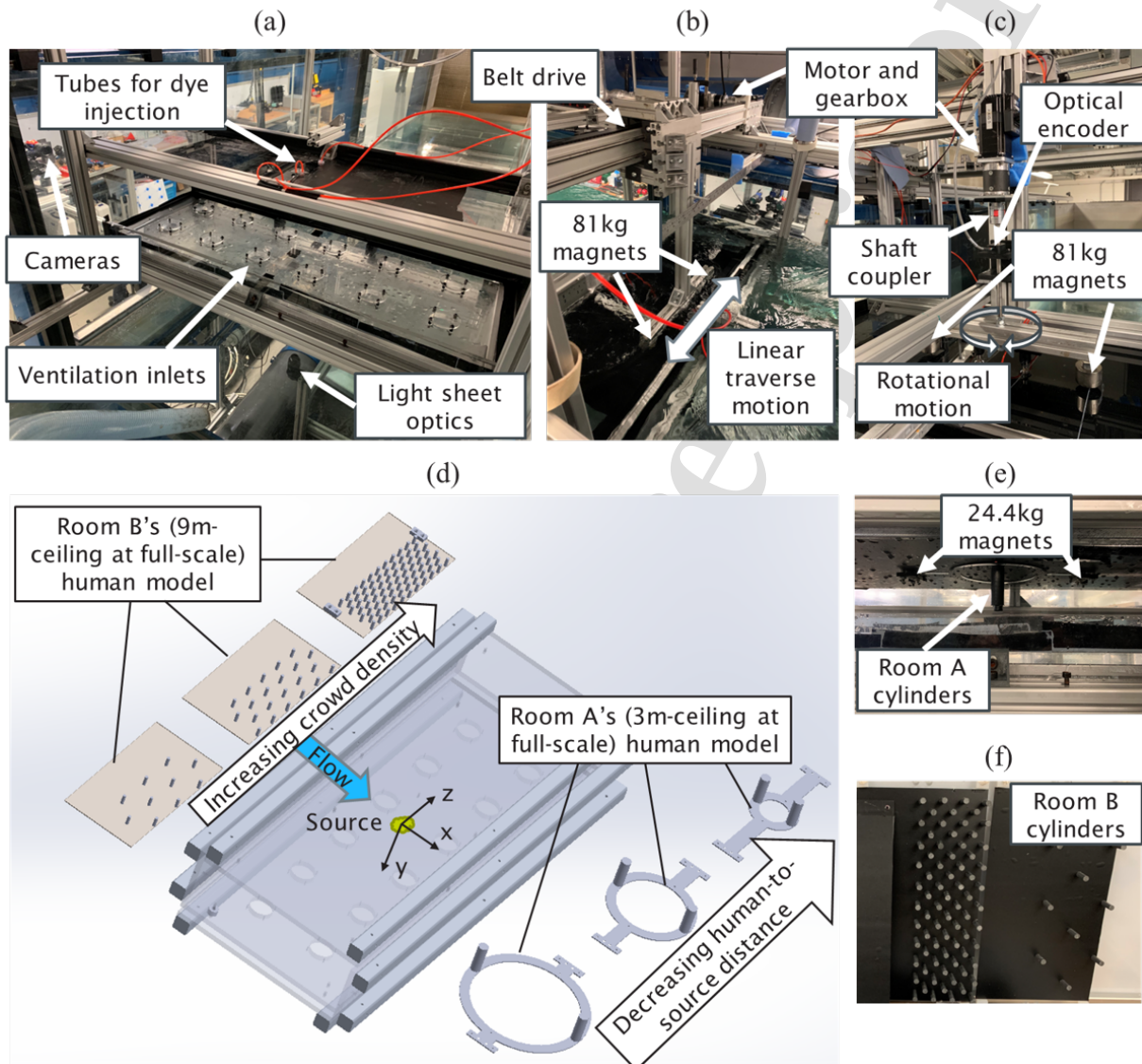


Figure 2. (a) Partial assembly of the room setup in the upside down configuration, (b) linear traverse for room B cylinder motions, (c) rotational setup for room A cylinder motions, (d) schematic of test models, (e) room A cylinders in upside down configuration held by magnets and (f) the construction of room B cylinder array.

137 mounted on the traverse system (this is visible in figure 2(c), each magnet is capable of pulling 81 kg) which attracted
138 two smaller rectangular (40 mm × 20 mm × 10 mm) neodymium magnets (each capable of pulling 24.4 kg) mounted
139 onto thin (up to 3 mm) sheets on which the cylinders were attached. The use of very strong magnets is necessary
140 due to the spatial separation (i.e. air gap) between the magnets. To avoid clamping the cylinder model to the room
141 floor, two ball transfer units were used to support the round magnets on the traverse system, and four miniature ball
142 transfer units were used to support the plates on which the cylinders were mounted. This also helped to reduce friction
143 between the cylinder plate and the room floor. Figure 2(e) shows room A cylinder model held in place by the magnets
144 and figure 2(f) shows the construction of room B cylinder arrays, where the cylinders were machined out of nylon
145 rods and screwed onto the 3 mm thin, 400 mm × 200 mm acrylic sheets (i.e. cylinder plate). Room A cylinders were
146 constructed in a similar manner, but with a different cylinder plate design, as shown in figures 2(d,e).

147 The concentration and velocity statistics were measured simultaneously in the xy -plane at $z=0$ (see figure 1) using
148 PLIF and PIV techniques. For PLIF, Rhodamine 6G fluorescent dye which has a Schmidt number of $Sc=2500\pm300$
149 [31] was used as a passive scalar tracer and released at the origin via a 2.5 mm inner diameter tube. The source
150 concentration was $C_s=10 \text{ mg L}^{-1}$ for all test cases. The dye flow rate of $Q_{dye}=10 \text{ mL min}^{-1}$ was maintained using
151 a needle valve and a Mariotte's bottle reservoir. As the dye flow rate is very low, the dye equilibrates rapidly with
152 the room flow and would therefore not have any influence on the scalar dispersion properties [32]. The choice of
153 our scalar tracer meant that our findings are applicable to gaseous pollutants and very fine particulate matter where
154 the particle velocity response time (which is a function of the particle density and diameter) is low. The effects of
155 pollutant resuspension or dry deposition are not considered in this study. For PIV, 50 μm polyamide seeding particles
156 were recirculated in the flume until uniform seeding was achieved. The water quality was closely monitored and
157 treated to ensure any changes to the background levels of Rhodamine 6G dye, agglomeration of seeding particles or
158 bio-fouling did not contribute to experimental uncertainties. A Nd:YAG double-pulsed laser (emission wavelength of
159 532 nm) and laser optics were used to form a laser sheet to illuminate the seeding particles and excite the fluorescent
160 dye which fluoresce at 554 nm. To separate the signals from the seeding particles and fluorescent dye, a 540 nm
161 long-pass filter was fitted to the PLIF camera (5.4 MP 16-bit depth Imager sCMOS) and 532 nm laser-line bandpass
162 filters were fitted to the PIV cameras (two 4 MP Imager MX CMOS cameras in side-by-side configuration). For all
163 test cases, 1000 realisations of the flow field were captured at an acquisition frequency of either 2 Hz or 2.5 Hz. This
164 is dependent on the exact test cases as the dataset presented in this study came from multiple experimental campaigns
165 and has negligible effects on the results and discussions in this study.

166 The image post-processing procedures are similar to several of our recent studies which readers can refer to for
167 more details [32]. The particle images were post-processed using LaVision DaVis 10 software. Image pre-processing
168 was performed to remove background noise before performing multi-pass, multi-grid cross-correlation analysis with
169 an overlap ratio of 50%. Spurious vectors were removed and replaced using local neighbourhood standard deviation
170 and median filters. The final vector resolution is 1.56 mm. A conservative estimate of the velocity bias is 0.5% and the
171 standard error of the mean statistics estimated based on the standard bootstrap with replacement statistical procedure
172 is 3% at 95% confidence interval. The PLIF post-processing was implemented using an inhouse code where PLIF
173 images were calibrated against the background, 0.03 mg L^{-1} and 0.05 mg L^{-1} concentrations of dye in two calibration
174 tanks, to obtain the linear relationship between the fluorescence intensity and dye concentration. Light attenuation in
175 the calibration tanks was accounted for using the Beer-Lambert's law and temporal variations in the laser pulses were
176 accounted for using an energy monitor device attached to the laser. The PLIF resolution is 0.174 mm. A conservative
177 estimate of the bias error is 10.2% at 95% confidence interval and the standard error is estimated to be 1% based on
178 the bootstrap with replacement statistical procedure. An image masking procedure was performed for regions where
179 the field-of-view was obstructed, and consequently, these regions are not considered in the discussions.

180 2.2. Test parameters

181 The test parameters investigated in this study are presented in table 1. The first group of eight rows presents the
182 test parameters used to investigate the effects of walking speed for rooms A and B, characterised by the Reynolds
183 number based on the diameter and mean velocity of the cylinders. The dynamic similarity of the walking motion was
184 matched using Re_{crowd} , which was calculated based on the diameters of the cylinders and their walking speed (U_{crowd}).
185 The Re_{crowd} investigated in this study are lower than that of a full-scale human walking which is approximately
186 $Re \approx 20,000$, but the cylinder flow regimes are the same (subcritical regime with fully turbulent vortex streets [33])
187 hence results and trends are expected to be representative. The next group of four rows (separated by the black line)

Table 1. Test parameters for rooms A (20:1) and B (60:1). The subscript ‘FS’ denotes full-scale quantities where applicable. Test cases with repeated entries are marked by superscripts * and #.

Test case	U_{ref} [m s ⁻¹]	$U_{ref,FS}$ [m s ⁻¹]	Re_{inlet}	ACH_{FS} [h ⁻¹]	U_{crowd} [m s ⁻¹]	$U_{crowd,FS}$ [m s ⁻¹]	Re_{crowd}	Crowd density [ppl m ⁻²]	Human-to-source distance (D)
B-WS0	0.29	0.07	17,400	1.9	0	0	0	0.125	4.3
B-WS1	0.29	0.07	17,400	1.9	0.10	0.03	768	0.125	4.3
B-WS2	0.29	0.07	17,400	1.9	0.19	0.05	1536	0.125	4.3
B-WS3*	0.29	0.07	17,400	1.9	0.32	0.08	2560	0.125	4.3
A-WS0	0.19	0.14	11,200	10.8	0	0	0	-	4
A-WS1	0.19	0.14	11,200	10.8	0.08	0.06	2000	-	4
A-WS2#	0.19	0.14	11,200	10.8	0.16	0.12	4000	-	4
A-WS3	0.19	0.14	11,200	10.8	0.22	0.17	5500	-	4
B-CD0	0.29	0.07	17,400	1.9	0.32	0.08	2560	0	4.3
B-CD1	0.29	0.07	17,400	1.9	0.32	0.08	2560	0.05	4.3
B-CD2*	0.29	0.07	17,400	1.9	0.32	0.08	2560	0.125	4.3
B-CD3	0.29	0.07	17,400	1.9	0.32	0.08	2560	0.5	4.3
A-HS1	0.19	0.14	11,200	10.8	0.16	0.12	4000	-	2
A-HS2#	0.19	0.14	11,200	10.8	0.16	0.12	4000	-	4
A-HS3	0.19	0.14	11,200	10.8	0.16	0.12	4000	-	6

188 presents the test parameters used to investigate the effects of crowd density, and is characterised by the number of
 189 people occupying one square meter (ppl/m²) at full-scale. The last group of three rows presents the test parameters
 190 used to investigate the effects of human-to-source distance. The human-to-source distance is measured from the center
 191 of the cylinders in the row that is closest to the source. Some of the test cases are repeated in different groups, and are
 192 organised in this manner to allow readers to easily compare the test parameters of specific crowd effects.

193 The dynamic similarity of the room flow to the full-scale problem is achieved by matching the Reynolds number
 194 (Re_{inlet}) based on U_{ref} and h_{inlet} [34], which represents the characteristic velocity and length scales associated with
 195 the ventilation inlet respectively. Note that U_{ref} is in effect the spatially-averaged flow velocity at the inlet, as it
 196 was calculated using the room flow rate and the inlet area, where the room flow rate was estimated using velocity
 197 profiles measured at the outlet for $z=0$ and $z=90$ mm planes. Our Re_{inlet} is similar to some [35, 29] and higher than
 198 other studies [34, 26]. The ventilation time scale was calculated using the room volume and room flow rate (i.e.
 199 $1/ACH \times 3600$), which led to full-to-model ventilation time scale ratios of $t_{ratio}=26.7$ for room A (geometric scale
 200 ratio, $L_{ratio}=20$) and $t_{ratio}=240$ for room B (geometric scale ratio, $L_{ratio}=60$). Using these ratios, full-scale parameters
 201 can be calculated (denoted by subscript ‘FS’ in this study) [29]. An example equation to convert the model-scale U_{ref}
 202 to full-scale is given by:

$$U_{ref,FS} = U_{ref} \times \frac{L_{ratio}}{t_{ratio}}. \quad (1)$$

203 The same methodology was used to calculate the ACH_{FS} , which in this study (table 1), are comparable to several past
 204 studies and representative of indoor spaces [36, 3].

205 3. Concentration statistics

206 In this section, we examine the effects of the walking speed, crowd density and human-to-source distance on the
 207 concentration statistics, which may be useful to inform guidelines for incidents where hazardous air pollutants are
 208 accidentally released in indoor spaces. As the movement of the cylinder models for room A and B were not synchronised
 209 with the PLIF-PIV imaging system, our analysis will only focus on the statistics of the xy-plane at $z=0$ and
 210 the effect of the full period of motion in our analysis (instead of examining the flow field at different phases of the

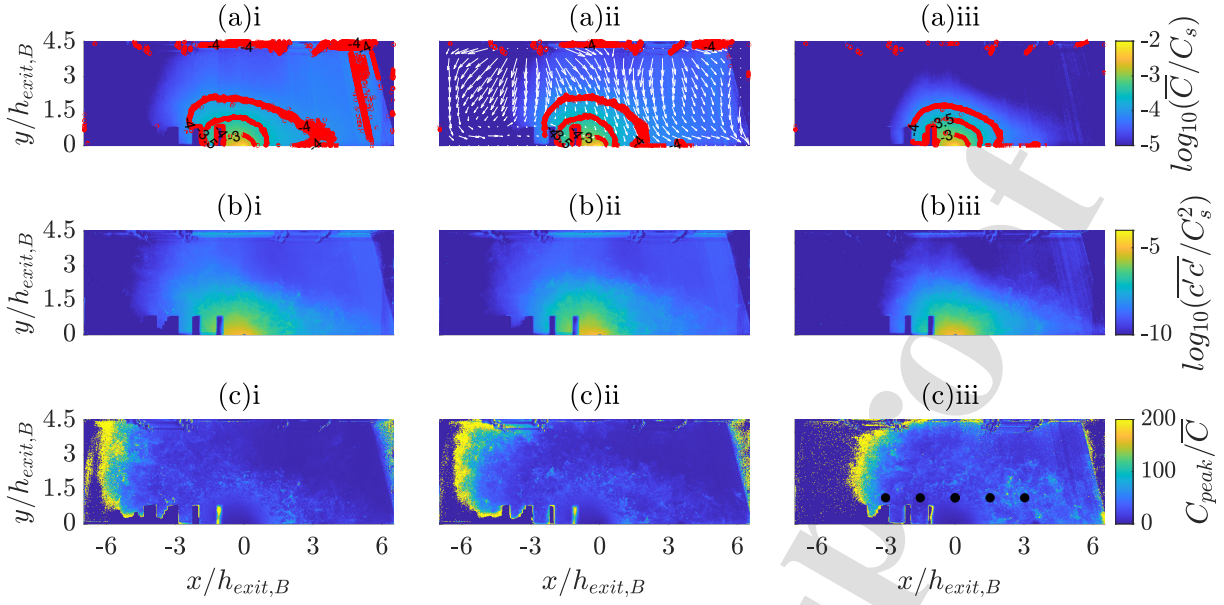


Figure 3. Effects of walking speed, room B. (a) Mean concentration, (b) concentration variance and (c) peak-to-mean concentrations for test cases (i) B-WS0, (ii) B-WS1 and (iii) B-WS3 (i \rightarrow iii: increasing walking speed). The white arrows represent mean velocity vectors, red represents isocontour lines of the mean concentration and black dots represent the locations where the concentration PDF is extracted.

211 cylinder motion). In this study, \bar{C} and c' represent the time-averaged (mean) concentration and instantaneous concentration fluctuation respectively. C_{peak} represents the maximum concentration measured over the entire experimental acquisition time.

214 3.1. Effects of walking speed

215 Figure 3 shows the concentration statistics for room B. The mean velocity vectors are overlaid for only the B-WS1 test case in figure 3a(ii) as they are similar across the three test cases. The similarity of the flow field across the test cases can be attributed to the relatively small length-scales of the cylinders compared to the room, hence differences in their walking speed did not significantly affect the room flow at the macro scale.

219 With an increase in walking speed (i.e. figure 3i \rightarrow 3iii), lower magnitudes of mean concentrations are observed at regions far from the source. The isocontour lines at the near-source region take on a more uniform hemispherical shape (figure 3a(iii)) and in particular, closer isocontour lines to the right of the source show an increase in the magnitude of the mean concentration gradient particularly towards the room outlet. The magnitudes of the concentration variance are highest at the source and decay with distance from source. The variance magnitudes are observed to decay more quickly as the walking speed increases (i.e. figure 3(b)i \rightarrow 3(b)iii), particularly in the left side of the room, which is occupied by the cylinders.

226 The trends observed here offer insights to the scalar mixing in the out-of-plane z-direction and the in-plane xy-direction, as they have competing effects on the mean concentration and concentration variance of the xy-plane at $z=0$ as presented in figure 3. As an example, if there is infinitely efficient scalar mixing in the out-of-plane z-direction, and negligible scalar mixing in the in-plane xy-direction, then we can expect negligible mean concentration and concentration variance maps in the xy-plane at $z=0$. This is because the scalar is transported out of the xy-plane much quicker than it is allowed to mix in the in-plane direction. Similarly, with an infinitely efficient scalar mixing in the in-plane xy-direction, and negligible scalar mixing in the out-of-plane z-direction, we would expect the mean concentration maps for xy-plane at $z=0$ to show uniform magnitudes in the entire room.

234 Since the increase in walking speed led to lower magnitudes of the mean concentration and concentration variance as shown in figure 3, this is indicative that there is much more significant improvement of the scalar mixing in the out-of-plane z-direction, which is aligned with the direction of the human motion, than in the in-plane xy-direction (i.e.

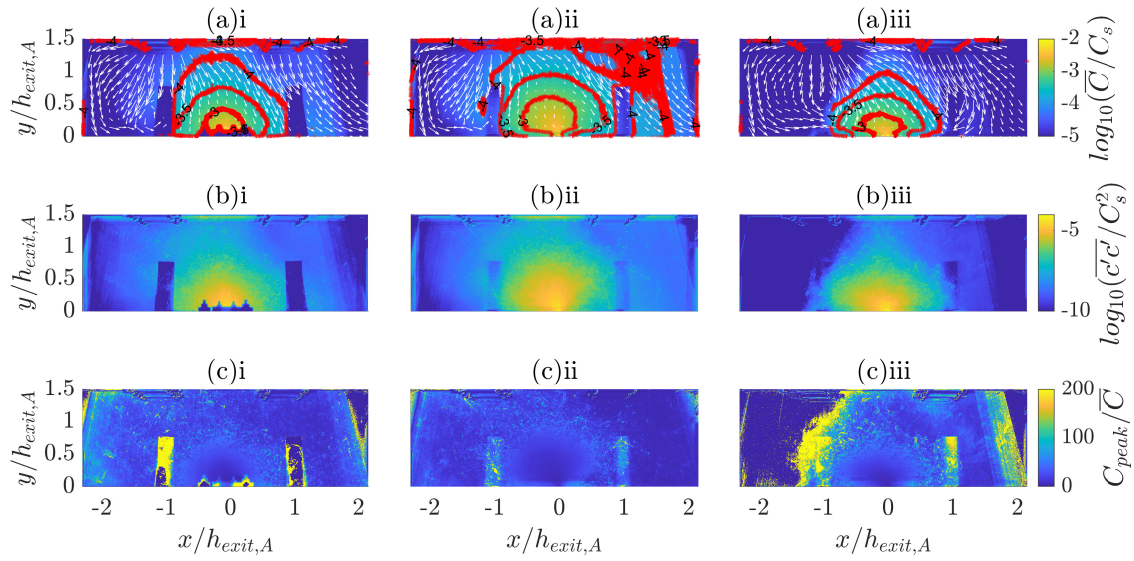


Figure 4. Effects of walking speed, room A. (a) Mean concentration, (b) concentration variance and (c) peak-to-mean concentrations for test cases (i) A-WS0 baseline, (ii) A-WS1 and (iii) A-WS3 (i → iii: increasing walking speed). The white arrows represent mean velocity vectors and red represents isocontour lines of the mean concentration.

237 scalar mixing is non-uniform and anisotropic). This observation is consistent with the literature, where Marlow et al.
 238 [19], Mingotti et al. [13] and Hang et al. [14] have all observed pollutant entrainment due to the wake flow of people,
 239 and in particular, Marlow et al. [19] observed higher concentrations of air pollutants close to where people cluster
 240 together. Therefore, the improvement in scalar mixing in the z-direction is likely due to the z-direction movements of
 241 the cylinder arrays in room B.

242 In the context of assessing exposures to toxic chemicals with high toxic load exponents [37, 38], insight into
 243 short-duration events is needed. The peak-to-mean concentration ratio (C_{peak}/\bar{C}) is a useful parameter that represents
 244 the maximum concentration relative to the time-averaged (i.e. mean) concentration measured in the entire acquisition
 245 time. The C_{peak}/\bar{C} maps are presented in figure 3(c), where large magnitudes are observed in the left side of the
 246 room occupied by the cylinders. Intriguingly, the magnitude peaks are always connected to the cylinder array, and
 247 as the walking speed increases, they shift towards rows of cylinders that are closer to the source. Even further to the
 248 left, the row of cylinders furthest from the source appear to be in a scalar-free zone, which is defined as any region
 249 where negligible magnitudes of mean concentrations and C_{peak}/\bar{C} are observed. We will return to this interesting flow
 250 feature and discuss it in greater detail in section 3.4.

251 The effects of walking speed for room A are examined with the mean concentration statistics presented in figure
 252 4. With an increase in walking speed, the mean velocity vector maps show slightly different flow fields for each of
 253 the test cases, indicating the motion of the cylinders has an influence on the room flow. This is likely due to the
 254 comparable length-scales of the cylinders and the room, which was not the case in the results shown for room B. The
 255 mean concentration isocontour lines are relatively symmetrical about the source as compared to the previous room B
 256 configuration, likely due to the symmetric cylinder positions. Similar to room B, the increase in walking speed led
 257 to lower magnitudes of mean concentration and concentration variances in regions far from the source, particularly
 258 in the left side of the room, as shown in figures 4(a)iii and 4(b)iii. The C_{peak}/\bar{C} ratio is presented in figure 4(c).
 259 High magnitudes of the C_{peak}/\bar{C} ratio are observed only for the A-WS3 test case. This is likely because the rotational
 260 motion of the left cylinder (about the vertical y-axis) has improved scalar mixing and transport of the scalar out of the
 261 xy-plane much more significantly than the scalar transport and mixing within the xy-plane. As such, the magnitudes
 262 of the mean concentration as shown in figure 4(a)iii are low, which is a necessary condition for high magnitudes of
 263 the C_{peak}/\bar{C} ratio. We will return to discuss the C_{peak}/\bar{C} parameter in greater detail by examining the concentration

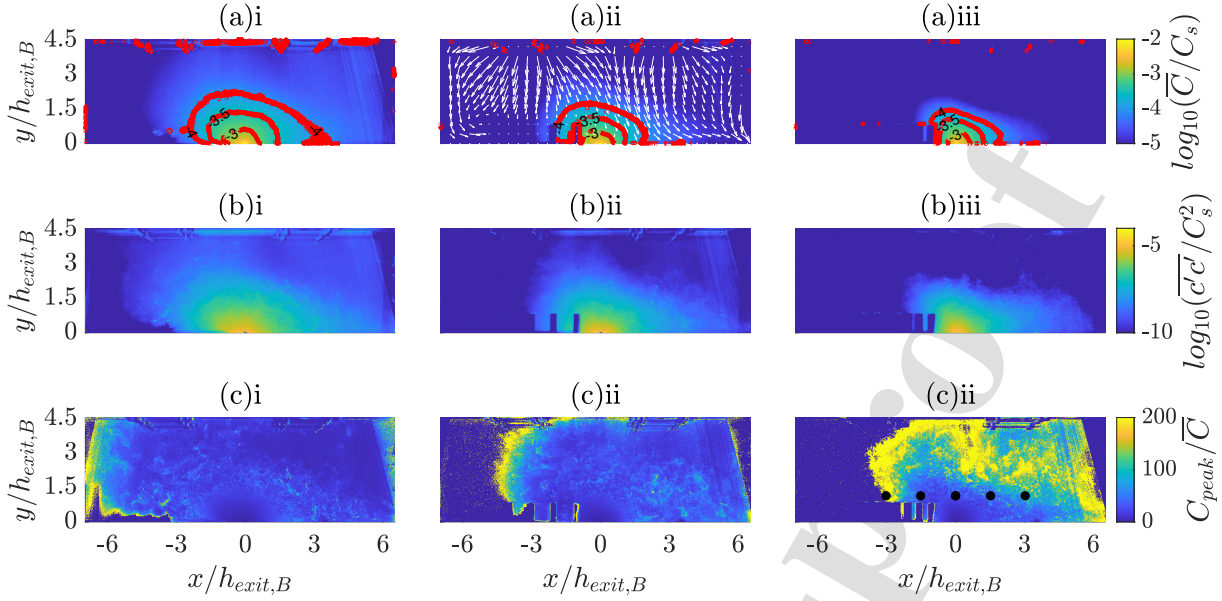


Figure 5. Effects of crowd density, room B. (a) Mean concentration, (b) concentration variance and (c) peak-to-mean concentrations for test cases (i) B-CD0, (ii) B-CD2 and (iii) B-CD3 (i \rightarrow iii: increasing crowd density). The white arrows represent mean velocity vectors, red represents isocontour lines of the mean concentration and black dots represent the locations where the concentration PDF is extracted.

264 PDFs in section 3.4.

265 3.2. Effects of crowd density

266 The effects of crowd density for room B are presented in figure 5, where mean velocity vectors are shown in
 267 only one test case as they are similar across all three test cases. With an increase in the crowd density, the shape of
 268 the mean concentration isocontour lines remain similar, but the magnitudes of the mean concentration decay more
 269 rapidly with distance from source, thus indicating more efficient out-of-plane scalar mixing in the near-source region.
 270 The concentration variances follow the same trends and are observed to decay more rapidly, particularly in regions
 271 close to the ceiling and at the left side of the room (occupied by the cylinder array) where negligible magnitudes
 272 were measured. The peak magnitudes of the C_{peak}/\bar{C} ratio has shifted to originate from rows of cylinders closer to
 273 the source, with the region beyond $x/h_{exit,B} < -3$ a scalar-free zone in which the scalar dye does not penetrate into,
 274 evidenced by negligible magnitudes of both \bar{C} and C_{peak}/\bar{C} . These observations are similar to the effects of increasing
 275 the walking speed for room B previously observed in figure 3(c).

276 3.3. Effects of human-to-source distance

277 The effects of human-to-source distance for room A are presented in figure 6. As discussed in section 3.1, room
 278 A cylinders have an influence on the macro room flow while room B cylinders were only able to influence the local
 279 mean flow field due to the relative length-scales of the cylinders to the room. The mean velocity vectors presented in
 280 figure 6 show that the human-to-source distance has a very significant effect on the macro room flow. In fact, the mean
 281 flow patterns in the left side of the room can be observed to be completely different for all three test cases. Smaller
 282 but non-negligible differences can be observed at the right side of the room. This can be attributed to the relatively
 283 closer proximity to the outlet, which resulted in the mean flow being less affected by the motion of the cylinders.
 284 The A-HS1 test case, which has the cylinders closest to the source, appears to have the best scalar mixing properties
 285 along the direction of the motion (out-of-plane). This results in low magnitudes of mean concentrations observed
 286 in most regions of the room (figure 6(a)). The observation of non-uniform and anisotropic scalar mixing, yet again,
 287 indicate the need to consider more complex human motion for subsequent studies. Closer proximity to source does
 288 not necessarily always lead to better scalar mixing properties, as exemplified by the A-HS2 test case, which bucks this

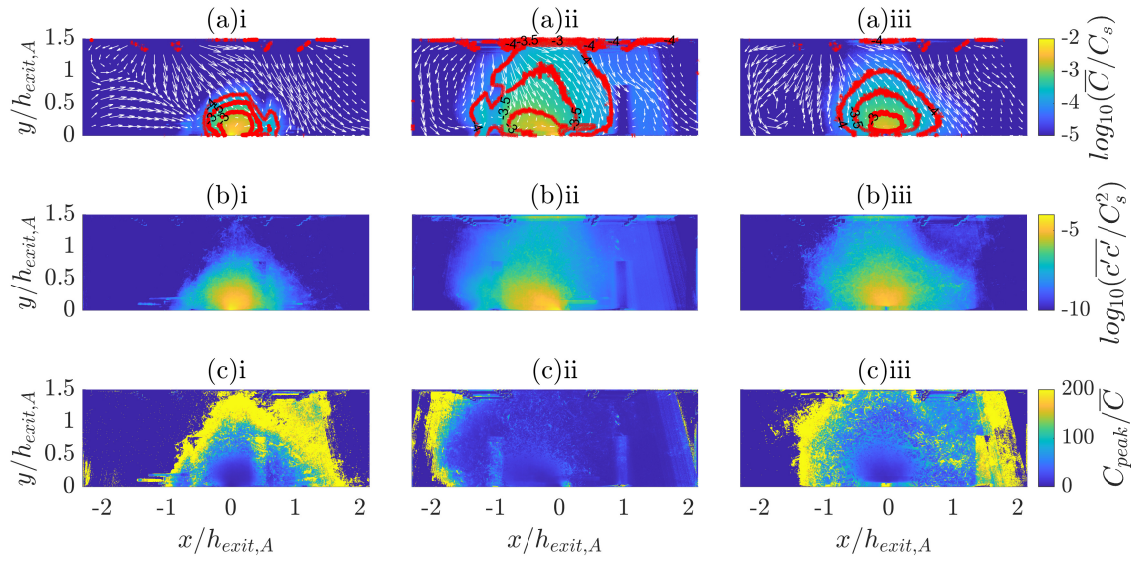


Figure 6. Effects of human-to-source distance, room A. (a) Mean concentration, (b) concentration variance and (c) peak-to-mean concentrations for test cases (i) A-HS1, (ii) A-HS2 and (iii) A-HS3 (i → iii: increasing human-to-source distance). The white arrows represent mean velocity vectors and red represents isocontour lines of the mean concentration.

289 trend (in comparison with A-HS1 and A-HS3), and has relatively high magnitudes of mean concentrations in most
 290 regions of the room. This is likely because of the non-linear interaction between the room flow and wake flow of the
 291 cylinders' which affects the scalar dispersion properties of all three test cases differently.

292 The concentration variance maps (figure 6(b)) follow similar patterns to the mean concentration, with peak magni-
 293 tudes observed at the source and magnitudes decay with distance from source. These observations are consistent with
 294 the earlier results on the effects of walking speed and crowd density, and appears to be universal across the range of
 295 test cases investigated in this study. For the peak-to-mean concentration ratio (figure 6(c)), the A-HS2 test case shows
 296 completely different C_{peak}/\bar{C} patterns to the other two test cases. Since this parameter is dependent on the mean flow
 297 field and intermittent scalar structures transporting high concentrations of dye, it is expected that the patterns will be
 298 very different if there are significant changes to the mean flow field of the room. We will discuss this in greater detail
 299 in the subsequent section.

300 3.4. Discussions on concentration PDFs

301 The concentration statistics previously presented in figures 3(c), 4(c), 5(c) and 6(c) showed interesting C_{peak}/\bar{C}
 302 patterns. In particular, high magnitudes are observed to originate from rows of cylinders closer to the source with an
 303 increase in walking speed and crowd density for room B. To better understand the C_{peak}/\bar{C} parameter, we consider the
 304 conditions where high magnitudes of C_{peak}/\bar{C} are likely to occur. Due to the definition of C_{peak}/\bar{C} , low magnitudes
 305 of mean concentration is a necessary condition for high magnitudes of C_{peak}/\bar{C} as shown in figures 3(c), 4(c), 5(c)
 306 and 6(c). However, high values of the concentration variance do not necessarily lead to high C_{peak} as shown in the
 307 results. As such, high magnitudes of C_{peak}/\bar{C} should be expected only when the sampled location typically has very
 308 low concentrations, with a few rare instances where highly intermittent scalar structures transport high concentrations
 309 of dye to it.

310 Since the intermittency of the scalar structures affects the C_{peak}/\bar{C} ratio, it would be useful to examine the PDF
 311 of the concentration. The concentration PDFs of B-WS3 test case sampled at 5 selected locations (see figure 3(c)iii)
 312 are presented in figure 7(a). Far from the source, the concentration values are almost always close to zero with only a
 313 couple of events out of the 1000 PLIF realization contributing to large concentration values (not visible in the plots).
 314 Close to the source, wider PDFs are observed which is consistent with the non-negligible mean concentration values

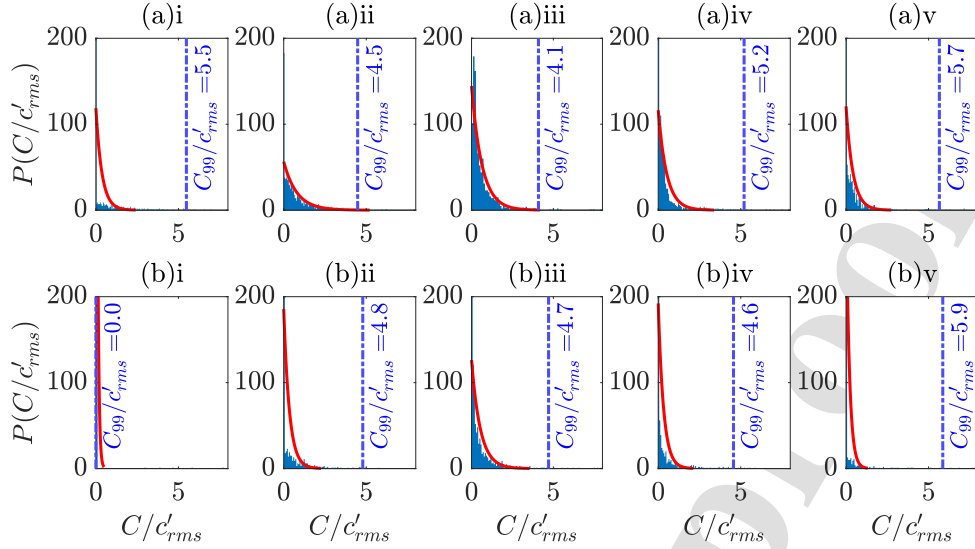


Figure 7. PDF for test case (a) B-WS3 (also equivalent to test case B-CD2) and (b) B-CD3 extracted at $y/h_{exit,C}=1.14$ and $x/h_{exit,C}=(i) -3.03$, (ii) -1.52 , (iii) 0 , (iv) 1.52 and (v) 3.03 . Extracted locations are plotted as black dots on figures 3(c)iii and 5(c)iii. The red line shows an exponential fit.

315 observed in figure 3(a). The exponential function is a good fit to the PDF, which is consistent with the results observed
 316 for an empty room [29]. In the same plot, we have also appended the C_{99}/c'_{rms} ratio which is another parameter that
 317 can be used to determine exposures. This is defined as the ratio in which the concentration value is exceeded 1%
 318 of the time to the root-mean-square concentration fluctuations (i.e., exposure to a concentration of at least C_{99}/c'_{rms}
 319 for 1% of the time). This ratio generally ranges between 4.1 to 5.7 and does not appear to have a clear correlation
 320 with distance from source. The expected value for an exponential distribution at the 99% percentile can be derived as
 321 $C_{99}/c'_{rms}=4.61$ [29] which agrees well with the values here.

322 The PDFs associated with test case B-CD3 are presented in figure 7(b) (see figure 7(a) for PDFs of test case B-
 323 CD2). The general observations are similar to that of test case B-CD2. Wider PDFs are observed in the near-source
 324 region which indicates frequent scalar events that carry varying concentrations of dye to those locations, while far
 325 from the source where negligible mean concentration were measured, the PDFs showed concentrations values are
 326 almost always zero with only a few events contributing to high C_{peak}/\bar{C} values. For most of the sampled locations,
 327 C_{99}/c'_{rms} ratio ranges between 4.6 to 5.9 but there is no clear correlation with the C_{peak}/\bar{C} values or distance from
 328 source. Figure 7(b)i is a particularly interesting example that has high C_{peak}/\bar{C} (see figure 5(c)iii) but zero C_{99}/c'_{rms}
 329 values, and can be attributed to the intermittent scalar events with high concentrations of dye occurring at less than
 330 1% of the time in the 1000 PLIF realisations. This result shows that the use of both C_{peak}/\bar{C} and C_{99}/c'_{rms} parameters
 331 are useful and provide different information for the purpose of determining exposures.

332 In figures 3(c), 4(c), 5(c) and 6(c), peak magnitudes of C_{peak}/\bar{C} are observed to link the cylinders to the room
 333 ceiling. In particular, the increase in walking speed and crowd density for room B led to peak magnitudes shifting
 334 towards rows of cylinders closer to the source. This can be attributed to cylinder-generated turbulence that is more
 335 effective at expelling high concentrations of dye vertically out of the cylinder array and into the region of the room
 336 which typically has low dye concentrations. The dye is then entrained by large-scale structures of the length-scale of
 337 the room and advected to the room ceiling.

338 4. Concentration-velocity statistics and the scalar transport mechanisms

339 4.1. Effects of walking speed

340 To better understand the mean concentration patterns and underlying scalar transport mechanism, we examine the
 341 maps of the advective and turbulent scalar fluxes. The effects of walking speed for room B are presented in figure

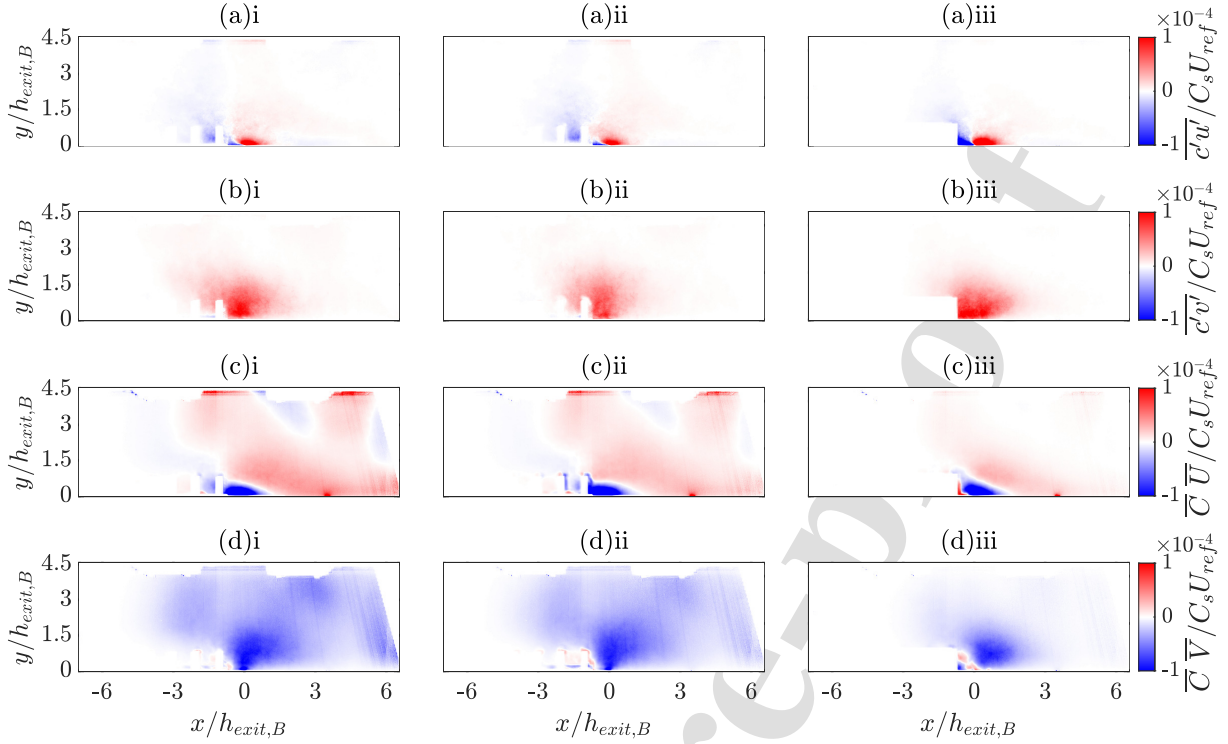


Figure 8. Effects of walking speed, room B. Mean (a) horizontal and (b) vertical turbulent scalar fluxes. Mean (c) horizontal and (d) vertical advective scalar fluxes. Test cases (i) B-WS0, (ii) B-WS1 and (iii) B-WS3 (i \rightarrow iii: increasing walking speed).

8, where negative and positive values represent the directions of the fluxes relative to the tracer source at the origin. The magnitudes of the turbulent fluxes are highest at the source, decay rapidly with distance from source, and are negligible in regions far from the source. This can be attributed to the rapid decay of the concentration variance with distance from source (see figure 3(b)). The general patterns of the advective fluxes are similar, but they decay more slowly with increasing distance from source than the turbulent fluxes.

In the near-source region, the advective fluxes are of the same order of magnitude as the turbulent fluxes. As such, the scalar transport in the near-source region is dependent on the relative magnitudes and directions of the advective and turbulent fluxes. In particular, the vertical advective (downward) and turbulent (upward) fluxes are opposite in directions, which lead to a competing effect on whether the dye remains close to the ground. However, since the turbulent fluxes decay more rapidly with distance from source, and the vertical advective flux is downward in most regions of the room, the dye generally remains close to the ground. In the far-field, the advective fluxes dominate over the turbulent fluxes, and the scalar transport predominantly depends on the mean flow field of the room.

With an increase in walking speed, the magnitudes of the horizontal turbulent scalar flux at the near-source region increase, which suggests more rapid near-source turbulent scalar dispersion and greater turbulent entrainment of dye into the cylinder array. Dye is removed from the $z=0$ measurement plane due to the direction of the motion of the cylinder array, which is consistent with the observations of Marlow et al. [19] where they observed pollutant entrainment by the wake flow of the people. This leads to less dye being transported to the far-field and hence, lower magnitudes of the advective fluxes are observed in regions far from the source. Notably, low magnitudes of both advective and turbulent fluxes are observed in the left side of the room at higher walking speeds. This is consistent with low mean concentrations observed in this same region.

The scalar fluxes corresponding to the effects of walking speed for room A are presented in figure 9. Similar trends to room B are observed, with the turbulent scalar fluxes decaying more rapidly with distance from source and the advective fluxes dominant in regions far from the source. The increase in walking speed generally led to slightly higher magnitudes of the turbulent scalar fluxes at the near-source region, with the exception of the vertical turbulent

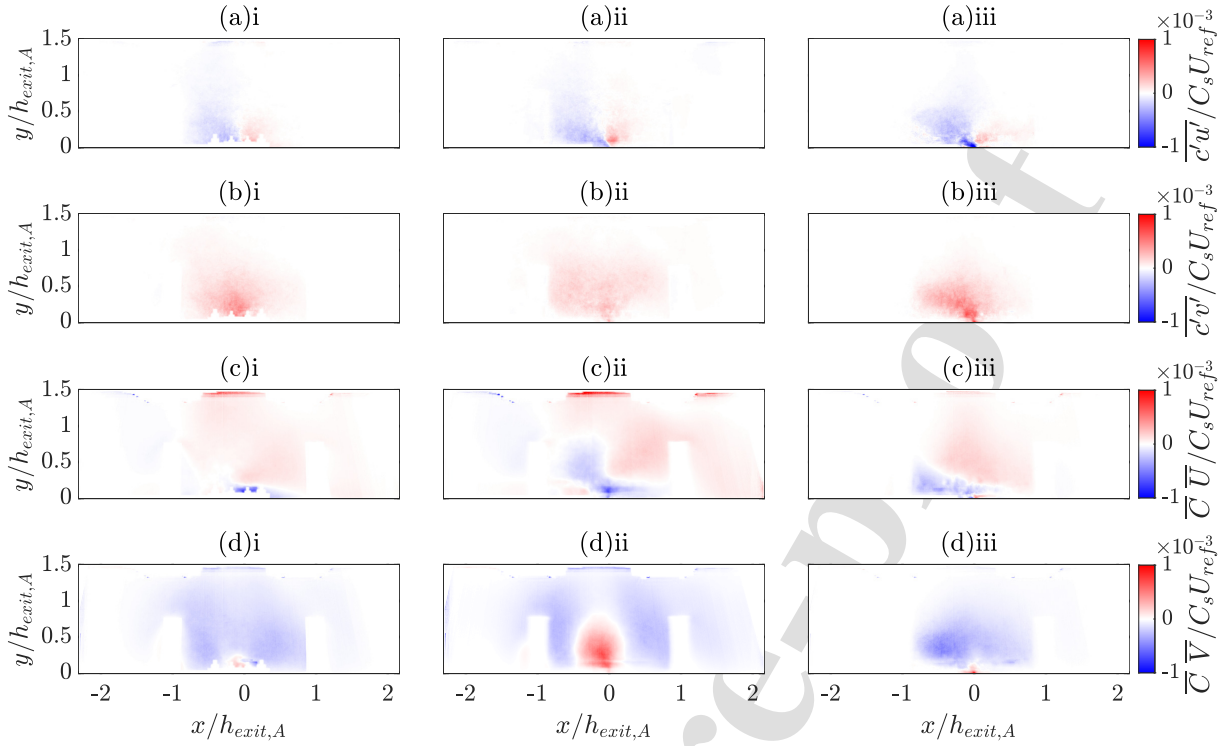


Figure 9. Effects of walking speed, room A. Mean (a) horizontal and (b) vertical turbulent scalar fluxes. Mean (c) horizontal and (d) vertical advective scalar fluxes. Test cases (i) A-WS0 baseline, (ii) A-WS1, (iii) A-WS3 (i \rightarrow iii: increasing walking speed).

366 scalar flux for test case A-WS1, which has slightly lower magnitudes than the other two test cases (see figure 9b(ii)).
 367 This can be attributed to significant differences in the mean flow field, with test case A-WS1 observed to have mean
 368 upward flow close to the source (see velocity vectors in figure 4a(ii)), which is also captured as positive values of
 369 vertical advective scalar flux (see figure 9d(ii)).

370 For the advective scalar fluxes, the effect of higher walking speed generally led to more dye being advected towards
 371 the left cylinder, evidenced by greater magnitudes of the negative horizontal advective scalar flux for both A-WS1 and
 372 A-WS3 test cases relative to the baseline A-WS0 test case as shown in figure 9(c). At the highest walking speed
 373 (test case A-WS3), negligible magnitudes of the advective fluxes are observed beyond the cylinder positions (i.e.,
 374 $x/h_{exit,B} < -1$ and $x/h_{exit,B} > 1$). Possibly, the increase in walking speed has led to more efficient dye entrainment by
 375 the cylinders' wake flow [19], with the dye subsequently being transported out of the measurement plane due to the
 376 direction of the cylinders' motion.

377 4.2. Effects of crowd density

378 To illustrate the effects of crowd density, the turbulent scalar fluxes for room B are presented in figures 10(a-b).
 379 Similar to the walking speed (room B) test cases previously presented in figure 8, the magnitudes of the turbulent
 380 fluxes here decay more rapidly with distance from source than the advective fluxes, resulting in the advective fluxes
 381 dominating scalar transport in regions far from the source. With an increase in the crowd density, the magnitudes of
 382 the turbulent scalar flux remain similar at the near-source region and decreases slightly for regions farther from the
 383 source.

384 The effect of increasing the crowd density is more significant for the advective fluxes as shown in figures 10(c-d).
 385 In the baseline B-CD0 test case, regions far from source show non-negligible magnitudes of the advective fluxes,
 386 which decrease to negligible values as the crowd density increases. Interestingly for the B-CD3 test case, the near-
 387 source advective flux pattern has changed significantly compared to the other two test cases, with positive magnitudes
 388 of the horizontal and vertical advective scalar fluxes observed close to the source. Although the length-scales of the

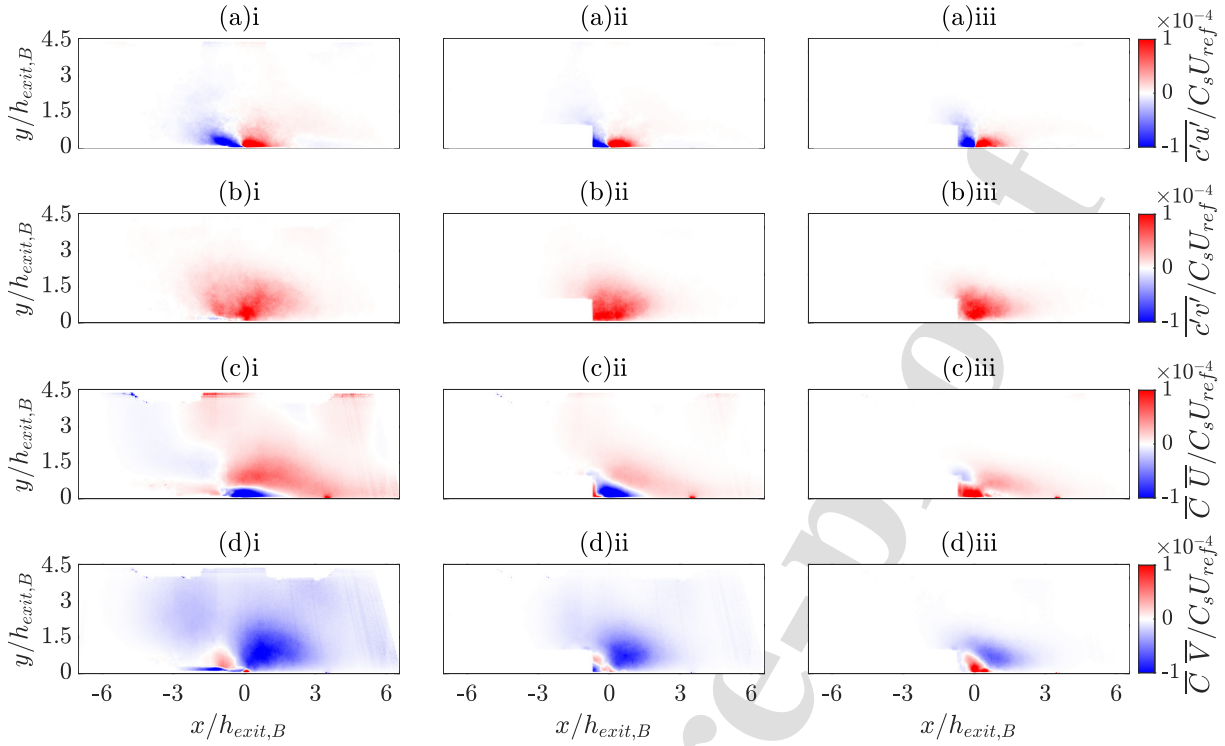


Figure 10. Effects of crowd density, room B. Mean (a) horizontal and (b) vertical turbulent scalar fluxes. Mean (c) horizontal and (d) vertical advective scalar fluxes. Test cases (i) B-CD0, (ii) B-CD2 and (iii) B-CD3 (i \rightarrow iii: increasing crowd density).

389 cylinders are much smaller than that of the room and that their motion does not affect the macro room flow, this result
 390 shows it is still possible to induce changes to the near-source region which leads to changes in the advective fluxes.
 391 This is likely because of the relatively close proximity of the cylinders to the source. If the cylinders were farther from
 392 the source, the near-source advective flux patterns are not expected to be affected by the crowd density.

393 In addition, the more rapid decay of scalar fluxes with distance from source and negligible magnitudes of the
 394 mean concentration in the far field are similar to those on the effects of walking speed in room B (figures 5 and 8).
 395 This signifies more effective scalar mixing in the z-direction, which can be attributed to more rows of cylinders in
 396 the near-source region that are responsible for entraining the scalar dye out of the measurement plane and bringing in
 397 fresh ambient fluid.

398 4.3. Effects of human-to-source distance

399 The scalar fluxes corresponding to the effects of human-to-source position (in room A) are presented in figure 11.
 400 Although the variation of the human-to-source distance do not produce any significant trends, much larger differences
 401 are observed for the advective scalar fluxes compared to the turbulent scalar fluxes across the three test cases. This
 402 indicates that the advective scalar transport is the primary mechanism responsible for the changes in the mean con-
 403 centration fields observed in figure 6(a). The fact that the cylinders can significantly influence the macro-scale flow
 404 and alter the flow field to the extent where trends in the scalar dispersion properties can no longer be observed, shows
 405 the human-to-source results discussed here are likely to be case specific. As such, if the length-scales of the people
 406 and the room are comparable (i.e. people in a room with 3 m ceiling), due care must be exercised when attempting to
 407 generalise the results to other applications.

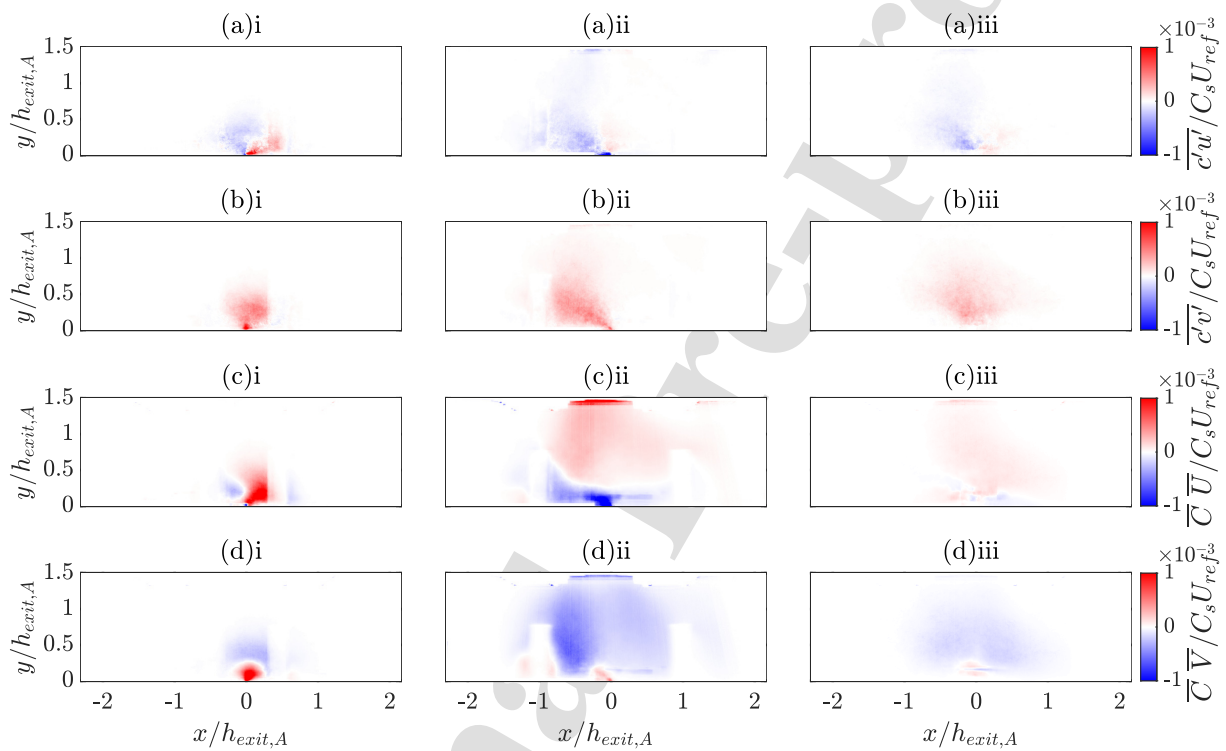


Figure 11. Effects of human-to-source distance, room A. Mean (a) horizontal and (b) vertical turbulent scalar fluxes. Mean (c) horizontal and (d) vertical advective scalar fluxes. Test cases (i) A-HS1, (ii) A-HS2 and (iii) A-HS3 (i \rightarrow iii: increasing human-to-source distance).

5. Conclusions

In this study, the effects of walking speed, crowd density and human-to-source distance on the scalar dispersion of a passive point source in two scaled room models were experimentally investigated using simultaneous PIV and PLIF techniques in a water flume facility. For all test cases, the magnitudes of the mean concentration and concentration variance were observed to peak at the source and decay with distance from source. Low magnitudes of mean concentration were observed to be a necessary condition to achieve high magnitudes of peak-to-mean concentration ratios, together with instances of highly intermittent scalar structures transporting high concentrations of dye.

Room A, a 20:1 full-to-model scale model representing a 3 m ceiling room, showed that the macro flow field of the room is affected by the people's motion due to the comparable length-scales of the cylinders and the room. Varying the walking speed and human-to-source distance showed that the advective scalar fluxes undergo more significant differences than the turbulent scalar fluxes due to changes in the mean flow field. However, the lack of generic trends for the advective fluxes, and that their magnitudes are often comparable or even larger than the turbulent fluxes, indicate the scalar dispersion patterns for room A are likely to be case specific and cannot be generalised to other applications.

Room B, a 60:1 full-to-model scale model representing a 9 m ceiling room, showed the motion of the people only had a local influence on the flow field due to the relatively larger difference in length-scales of the cylinders and the room. The increase in walking speed and crowd density led to improved scalar mixing with the fresh ambient fluid particularly in the z-direction, evidenced by low mean concentrations in regions far from the source and higher magnitudes of the concentration gradient at the near-source region. The peak magnitudes of the peak-to-mean concentration ratio were observed to connect the cylinder arrays to the ceiling. With an increase in walking speed or crowd density, this region shifts to rows of cylinders closer to the source. This suggests high concentrations of dye that were swept into the cylinder array from the source are removed quicker due to the higher cylinder-generated turbulence, which were then entrained and removed by the large-scale structures in the room. This results in a larger scalar-free zone with less rows of cylinders exposed to high magnitudes of mean concentration and peak-to-mean concentration. Nonetheless, if the ventilation design, room geometry or source location changes, the flow field at the near-source region may change drastically (e.g. shifting the source location right beside the outlet can result in the dye being advected directly out of the room [29]), and hence these results may no longer be valid.

There are a few limitations of this study which we hope to address in the future. Firstly, the location and design of the ventilation inlets and outlets are expected to have a significant influence in the mean flow patterns of the room but were not considered in this study. Secondly, we have not considered the effects of buoyancy due to the thermal plume of the people, temperature differences at the room boundaries or other heat sources in the room. This is expected to be more significant in determining the room flow characteristics when there are low pressure variations within the room, and include scenarios where the crowd is stationary or when the ACH is low. Thirdly, the Reynolds number of the crowd motion was limited to $O(10^3)$ in this study. Although the turbulence regime of cylinder wake flows are expected to be similar up to $O(10^5)$, the results in this study have shown that the motion of cylinders can affect the macro room flow. In the context of indoor scalar dispersion, Re_{crowd} must be considered relative to the macro room flow characteristics, and a systematic study on the ratio of the Reynolds number of the room flow/inlet to the crowd motion would be helpful. Fourthly, the crowd motion was limited to two types of simplified movement in this study. This may not be representative of realistic scenarios and more complex crowd motion should be considered. Finally, the measurements of the turbulent diffusivity is an important aspect of this study that has many practical applications but has not been discussed. We hope to address these limitations in future studies.

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The authors report no conflict of interest.

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455 All the quantitative maps are made available in the University of Southampton data repository to enable comparison
456 and validation with future research.

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Highlights:

- Advective and turbulent scalar flux maps reveal underlying transport mechanism
- Crowd effects can induce non-uniform and anisotropic scalar mixing
- Scalar-free zone observed on the opposite side of the room from ventilation outlet
- Dispersion patterns in small rooms are case specific due to non-linear interactions

Declaration of interests

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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