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ABSTRACT

We report on X-ray (NICER/NuSTAR/MAXI/Swift) and radio (MeerKAT) timing and spectroscopic analysis from a three-month monitoring campaign in 2022 of a high-intensity outburst of the dipping neutron star low-mass X-ray binary 1A 1744–361. The 0.5–6.8 keV NICER X-ray hardness-intensity and color-color diagrams of the observations throughout the outburst suggests that 1A 1744–361 spent most of its outburst in an atoll-state, but we show that the source exhibited Z-state-like properties at the peak of the outburst, similar to a small sample of other atoll-state sources. A timing analysis with NICER data revealed several instances of an ≈ 8 Hz quasi-periodic oscillation (QPO; fractional rms amplitudes of $\sim 5\%$) around the peak of the outburst, the first from this source, which we connect to the normal branch QPOs (NBOs) seen in the Z-state. Our observations of 1A 1744–361 are fully consistent with the idea of the mass accretion rate being the main distinguishing parameter between atoll- and Z-states. Radio monitoring data by MeerKAT suggests that the source was at its radio-brightest during the outburst peak, and that the source transitioned from the ‘island’

spectral state to the ‘banana’ state within ~ 3 days of the outburst onset, launching transient jet ejecta. The observations present the strongest evidence for radio flaring, including jet ejecta, during the island-to-banana spectral state transition at low accretion rates (atoll-state). The source also exhibited Fe XXV, Fe XXVI $K\alpha$, and $K\beta$ X-ray absorption lines, whose origins likely lie in an accretion disk atmosphere.

Keywords: stars: neutron – stars: oscillations (pulsations) – X-rays: binaries – X-rays: individual (1A 1744–361)

1. INTRODUCTION

1.1. Neutron Star Low-Mass X-ray Binaries

Low-mass X-ray binaries (LMXBs) are binary systems comprising a compact object (i.e., a black hole or neutron star) accreting matter from a low-mass companion star ($M_c < 1M_\odot$; see Bahramian & Degenaar 2022, for a review). Neutron star (NS) LMXBs have traditionally been broken down into two classes based on their correlated spectral and timing properties and on the patterns they trace in the X-ray color-color diagram (CCD): atoll sources trace out atoll-shaped regions in the CCD, tend to have lower luminosities ($\sim 0.001 - 0.5 L_{\text{Edd}}$), and often exhibit sporadic transient outbursts followed by extended periods of quiescence. Z sources trace a ‘Z’ shape in the CCD and tend to be highly luminous and persistent, accreting at close to the Eddington limit for a $1.4M_\odot$ NS ($> 0.5 L_{\text{Edd}}$; Hasinger & van der Klis 1989; van der Klis 1995; Belloni et al. 2002; van der Klis 2004).

Atoll sources have been observed in multiple spectral states, namely the ‘island’ and ‘banana branch’ states, historically defined by their positions in the CCD (Hasinger & van der Klis 1989; Wijnands et al. 2017). Typically, the island states have low accretion rates, and therefore low X-ray luminosities. Furthermore, the power spectrum is characterized by strong band-limited noise ($\gtrsim 10 - 20\%$ in fractional rms amplitude). At high(er) accretion rates, an atoll source occupies the banana state, and the power spectrum shows a weak ($\sim 3\%$) red noise component (Wijnands et al. 1999). The island X-ray spectra are hard (i.e., dominated by high-energy X-ray photons), and the spectra usually include a significant nonthermal component often parameterized by a power law (or broken power law). In contrast, the banana state X-ray spectra are dominated by soft, thermal X-ray photons from a blackbody (of the NS surface), multi-colored disk blackbody (from the accretion disk; Lin et al. 2007), or the boundary layer between the NS surface and accretion disk. Some atolls have not been observed to show state transitions, existing either only in the island state (e.g., 4U 1812–12: Tarana et al. 2006; Gladstone et al. 2007; XTE J1814–338: van Straaten et al. 2005) or banana state (e.g., GX 9+9:

Hasinger & van der Klis 1989; Kong et al. 2006; Fridriksson 2011).

The X-ray spectral states are correlated with the properties of the relativistic jets fueled by the accretion flow (e.g., Migliari & Fender 2006; Muñoz-Darias et al. 2014). In the island state most dominated by hard X-ray photons (i.e., the ‘extreme island’ state; see Méndez et al. 1997; van der Klis 2006; Wijnands et al. 2017), the accretion flow fuels a continuous, compact jet that exhibits a flat or inverted radio spectral index ($F_\nu \propto \nu^\alpha$; $\alpha \gtrsim 0$) up to a break frequency (often at sub-mm wavelengths) where the radio emission is optically thin ($\alpha \sim -0.7$; e.g., Migliari et al. 2010; Russell et al. 2013; Díaz Trigo et al. 2018). In some atolls, the compact jet is quenched in the banana states, becoming undetectable after the transition (e.g., Aql X–1; Miller-Jones et al. 2010). This behavior is analogous to the jet quenching observed in all black hole (BH) LMXBs during the transition from the hard to soft state (e.g., Coriat et al. 2011) and motivated the typical comparisons between the extreme island and banana states with the hard and soft states of BH LMXBs. However, some atolls have exhibited banana state radio detections (e.g., 4U 1820–30, Ser X–1; Migliari et al. 2004). Many BH LMXBs exhibit bright radio flaring during a hard to soft state transition (Fender 2006). In contrast to compact jet emission, radio flares show temporally evolving spectra that rapidly transition to optically thin (i.e., steep) spectral indices at radio frequencies. The flaring is thought to result from ejection events, and thus are de-coupled from the accretion flow and do not follow the L_R-L_X relation. Although a similar ejection process has been suggested for atolls (Migliari & Fender 2006), there are no observations of bright flaring or spatially resolved jet ejecta during the ‘island-to-banana’ transition. Therefore, there may be a weak quenching mechanism or possibly a different jet launching process in atoll NS LMXBs (see, van den Eijnden et al. 2021, for a summary of the differences between jet quenching in NS and BH LMXBs).

For Z sources, the CCD spectral states alternate between the ‘horizontal branch’ (HB; top of the Z), the ‘normal branch’ (NB; the diagonal), and the ‘flaring branch’ (FB; bottom). The X-ray spectra of Z sources

tend to be softer than atoll sources (Muno et al. 2002b). Like some atolls, Z sources experience quenching of the compact jet in the FB and flaring during NB–HB transitions (Migliari & Fender 2006). The Z source radio flaring is significantly more common than observed in atolls, thought to be the result of frequent transitions between the NB and HB. These flares have been spatially resolved as discrete ejection events in a number of Z source LMXBs (e.g., Cyg X–2: Spencer et al. 2013; Sco X–1: Fomalont et al. 2001; Motta & Fender 2019).

1.2. Quasi-Periodic Oscillations

Quasi-periodic oscillations (QPOs) are broad, ephemeral features in the power spectra of NS and BH LMXBs usually described with Lorentzians (van der Klis 2000). QPOs are observed throughout the Z-track in NS LMXBs, where the HB QPOs (HBOs) are roughly between 15–60 Hz, the NB QPOs (NBOs) are between 5–7 Hz, and the FB QPOs (FBOs) are between 7–20 Hz (e.g., Middleditch & Priedhorsky 1986; van der Klis 1989; Motta et al. 2017). The HBOs are thought to be linked to the disk truncation radius (e.g., Stella & Vietri 1998; Ingram & Done 2010; Motta & Fender 2019), while NBOs are thought to be connected to subsonic density (radial) oscillations in a thick disk (Fortner et al. 1989; Alpar et al. 1992) or relativistic and transient ejections of plasma (e.g., Fomalont et al. 2001; Migliari & Fender 2006).

Atoll sources show a variety of QPO phenomena, including mHz QPOs (e.g., Revnivtsev et al. 2001; Heger et al. 2007; Mancuso et al. 2023), ~ 1 Hz QPOs seen in some atolls exhibiting absorption dips (e.g. Jonker et al. 1999; Homan et al. 1999), and hecto-Hz QPOs (e.g., Ford & van der Klis 1998; Altamirano et al. 2008; Bult & van der Klis 2015). A class of atoll sources also show NBO-like QPOs near the highest source luminosities; they include XTE J1806–246, the ultracompact X-ray binary 4U 1820–30, and the transient source Aql X–1 (Wijnands et al. 1999; Wijnands & van der Klis 1999; Reig et al. 2004). These QPOs only lasted for hundreds of seconds, and were found in the tip of the upper banana branch in the CCD; they resembled the FBOs and NBOs seen in Z sources when the source accretion rate was close to the Eddington limit (van der Klis 1989). Both atolls and Z sources also exhibit kHz QPOs (e.g., van der Klis et al. 1996; Motta et al. 2017). The power spectra of the NS and BH LMXBs are additionally characterized by several broadband and higher frequency noise components (van der Klis 2000).

1.3. Z and Atoll States

Some NS LMXBs have transitioned between atoll and Z-like behavior at high mass accretion rates (Lin

et al. 2009). Detailed X-ray timing and spectral studies with the archetypal transitional Z/atoll source, XTE J1701–462, suggested that the transitions between the two spectral states are driven by changes in the mass accretion rate (Homan et al. 2007a,b; Lin et al. 2009; Homan et al. 2010; Fridriksson et al. 2015). The transient NS LMXB IGR J17480–2446 had also been observed to show transitions between Z-source behavior and atoll source behavior (Altamirano et al. 2010; Chakraborty et al. 2011). While the transitions themselves were not directly observed, Cir X–1 (Oosterbroek et al. 1995; Shirey et al. 1998, 1999) also showed Z-source-like properties at the higher luminosities ($> 0.5 L_{\text{Edd}}$) and atoll source-like properties at lower luminosities ($\sim 0.001 - 0.5 L_{\text{Edd}}$). Finally, the NBO-like QPOs observed from atolls XTE J1806–246, 4U 1820–30, and Aql X–1 (mentioned above; Wijnands et al. 1999; Wijnands & van der Klis 1999; Reig et al. 2004) at the highest source luminosities suggest a luminosity threshold at which the NBO production mechanism is activated. *In this work, we will discuss our results in the framework of Lin et al. (2009) and Homan et al. (2010), i.e., that atoll sources and Z sources are not distinct subpopulations of NS LMXBs but, instead, represent two states on a continuum defined by the mass accretion rate (atoll-state at low luminosities, Z-state at high luminosities). We will refer to them as atoll-state sources and Z-state sources from here on.*

1.4. Geometric Properties of NS LMXBs

Some LMXBs show absorption dips in their X-ray light curves (e.g., 4U 1624–490: Smale et al. 2001; 4U 1254–690: Smale et al. 2002). Given the expected viewing geometry of these phenomena, the binary inclination angle of non-eclipsing dippers is expected to be $i \approx 60 - 75^\circ$ (Frank et al. 1987). The absorption dips are thought to be due to intervening cold and/or partially ionized absorbers, around the interaction point where the accretion stream hits the accretion disk, along the line-of-sight (Díaz Trigo et al. 2006), or the partial covering by an accretion disk corona from an extended absorber (Church et al. 1998, 2005; Bałucińska-Church et al. 2011). These absorption dips are observed to be energy-dependent, and are accompanied by an increase in spectral hardness due to scattering of softer photons by the gas cloud. Some sources exhibit both eclipses and absorption dips (e.g., Swift J1858.6–0814: Buisson et al. 2021; EXO 0748–676: Parmar et al. 1986), and so the binary inclination angle of such systems is expected to be $i \approx 75 - 80^\circ$ (Frank et al. 1987).

1.5. 1A 1744–361

Tobrej et al. (2023) recently reported on joint NICER/NuSTAR spectroscopic analysis of 1A 1744–361, showing that a phenomenological model comprising a power law with a high energy cutoff described the spectrum well, and hinted at existence of a reflection component. They also showed prominent absorption features at 6.92 keV and 7.98 keV, which they attributed to H-like iron and a blend of Fe XXV and Ni XXVII lines, respectively. Mondal et al. (2023) also recently reported spectral analysis with NuSTAR data which found that the observations showed relativistic disk reflection and similar absorption features. The spectral analysis in both work identified the source to be in the banana branch state.

In this work, we present the X-ray (NICER, NuSTAR, MAXI, and Swift) and radio (MeerKAT) monitoring campaign of the dipping NS LMXB 1A 1744–361, which was first discovered by the Ariel 5 satellite in 1976 due to an outburst (Davison et al. 1976; Carpenter et al. 1977). 1A 1744–361 has exhibited several known outbursts between 1989 and the most recent outburst in 2022. Observations of the high-intensity 2003 outburst reported in Bhattacharyya et al. (2006a) revealed two sets of energy-dependent absorption dips in two successive RXTE orbits, which led the authors to infer an orbital period of $P_{\text{orb}} = 97 \pm 22$ min, though they did not rule out periods twice or half as long. Observations of the low-intensity 2004 outburst revealed the presence of a significant (4.3σ) $\nu \approx 3.5$ Hz QPO with a quality factor $Q \approx 2$ and a fractional rms amplitude of $\approx 5\%$ (Bhattacharyya et al. 2006a). They also reported a low significance (2.3σ) lower kHz QPO with $\nu \approx 800$ Hz, high Q value of $Q \approx 62.5$, and a fractional rms amplitude of $\sim 6\%$ (Bhattacharyya et al. 2006a). The QPOs were connected to the island state (Bhattacharyya et al. 2006b), similar to 4U 1746–37 and EXO 0748–676 (Homan et al. 1999; Jonker et al. 2000; Homan 2012), while the energy-dependent absorption dips were only observed in the high-intensity outbursts (Bhattacharyya et al. 2006b). Bhattacharyya et al. (2006a) have also reported the NS in 1A 1744–361 to have a spin frequency of $\nu \approx 530$ Hz, based on the discovery of thermonuclear burst oscillations in a type I X-ray burst seen in the high-intensity 2005 outburst.

1A 1744–361 went into outburst in June 2022 for about three months in its longest known outburst yet. Radio observations taken by MeerKAT detected bright radio emission during the early stages of the outburst (Hughes et al. 2022). In this paper, we present X-ray and radio observations of 1A 1744–361 taken over the span of almost three months to characterize the (aperiodic and periodic) timing and spectral evolution of the

source. In §2, we outline the observations and data analysis procedures. In §3, we present the results from the timing and spectral analyses. We will discuss the results in §4 and finally conclude in §5. All uncertainties reported represent 1σ confidence intervals unless otherwise noted.

2. OBSERVATIONS AND DATA ANALYSIS

2.1. Neutron star Interior Composition Explorer (NICER)

NICER, an external payload on the International Space Station, consists of 52 operating co-aligned X-ray concentrator optics and silicon drift detectors in focal plane modules (FPMs). NICER has fast-timing capabilities in the 0.2–12.0 keV energy range, allowing for a GPS time-tagging accuracy of 100 ns (Gendreau et al. 2016; LaMarr et al. 2016; Prigozhin et al. 2016).

NICER observed 1A 1744–361 starting from 2022 June 3 through 2022 August 31, with ObsIDs starting with 5202 and 5406. We processed the NICER observations with HEASOFT version 6.31.1 and the NICER Data Analysis Software (NICERDAS) version 10a (2022-12-16-V010a) with calibration version xti20221001. Our data processing criteria included the following: a source angular offset of $\text{ANG_DIST} < 0^\circ 015$; elevation angle from the Earth limb $\text{ELV} > 20^\circ$; NICER being outside the South Atlantic anomaly; bright Earth limb angle $\text{BR_EARTH} > 30^\circ$; undershoot rate (dark current; per FPM) range of $\text{underonly_range} = 0\text{--}500$; overshoot rate (charged particle saturation; per FPM) range of $\text{overonly_range} = 0\text{--}1.5$. We also applied COR_SAX (magnetic cut-off rigidity in GeV/c) filtering of $\text{COR_SAX} > 1.5$ to filter out background flares. This resulted in 53.5 ks of filtered exposure for scientific analysis out of 95.3 ks of unfiltered exposure.

For the timing analysis, we applied Solar system barycenter corrections in the ICRS reference frame (so that times are in Barycentric Dynamical Time; TDB), with source coordinates $\text{R.A.} = 267^\circ 0548$, $\text{Decl.} = -36^\circ 13251$, using barycorr in FTOOLS with the JPL DE421 solar system ephemeris (Folkner et al. 2009). The spectral analysis was conducted with XSPEC 12.13.0c (Arnaud 1996). The spectra were first grouped with the optimal binning scheme and rebinned such that each bin had at least 25 counts (Kaastra & Bleeker 2016). Background spectra were constructed with the nicerl3-spect model (Remillard et al. 2022) available through nicerl3-spect . The associated response matrices were generated with nicerarf and nicerrmf for each observation.

2.2. Nuclear Spectroscopic Telescope Array (NuSTAR)

NuSTAR is the first hard X-ray focusing telescope and consists of two co-aligned grazing incidence telescopes (Harrison et al. 2013). NuSTAR observed 1A 1744–361 between 2022 June 8 19:36 UTC (MJD 59738.81747; TT) and 2022 June 9 18:11 UTC (MJD 59739.75854; TT) for a total filtered exposure time of approximately 34.3 ks. The NuSTAR data reduction was performed with the standard pipeline tools (`nupipeline` and `nuproducts`) in NuSTARDAS v.1.9.7 and CALDB 20221229. We generated the final source and background event files with 100'' regions around the source, and in a source-free region near the source, respectively.

Initial analysis of 1A 1744–361 revealed relativistic reflection and strong emission and absorption features from the ~ 6.96 keV and ~ 8.0 keV Fe XXVI $K\alpha$ and $K\beta$ lines (Pike et al. 2022). In this work, we will focus on characterizing the timing properties of the source from the NuSTAR observation and defer detailed analysis of the spectroscopic properties to another publication (Pike et al. 2024, in prep.)

2.3. MeerKAT Radio Telescope

The MeerKAT radio telescope is an interferometric array located in South Africa. We observed 1A 1744–361 with MeerKAT as a part of the large survey project ThunderKAT (Fender et al. 2016). Our observing began with a single rapid response observation on 2022 May 31 (MJD 59730), ~ 2 days after the first MAXI/GSC X-ray detection (and ~ 8 hours after the outburst’s initial reporting; Kobayashi et al. 2022). Following this rapid response, we began a monitoring campaign on 2022 June 3 (MJD 59733), observing the source every ~ 7 days until 2022 August 27 (MJD 59818). Each observation consisted of a single scan with 15 minutes on-source flanked by two 2-minute scans of a nearby gain calibrator (J1830–3602). Each epoch also included a 5-minute scan of PKS B1934–638 (J1939–6342), for flux and bandpass calibration. Our observations used MeerKAT’s L-band receiver, with a central frequency of 1.28 GHz and a total (un-flagged) bandwidth of 856 MHz.

We performed flagging, calibration, and imaging using a modified version of the semi-automated routine OxKAT¹ (Heywood 2020), originally developed for the MeerKAT International GHz Tiered Extragalactic Exploration (MIGHTEE) Survey (see Heywood et al. 2022, for a detailed description of the routine). The first step (1GC) uses CASA (v5.6; CASA Team et al. 2022) to remove radio frequency interference (RFI) and apply the

standard flux density, bandpass, and complex gain calibrations. The second step (FLAG) applies a second round of flagging using TRICOLOUR (Hugo et al. 2022), followed by preliminary imaging of the source field using WSCLEAN (v2.9; Offringa et al. 2014). The preliminary image is then used to create an imaging mask. To maximize imaging sensitivity, we modified OxKAT, changing the Briggs’ robustness weight (Briggs 1995) from -0.3 (the default) to 0. We cannot increase the robustness further (i.e., > 0) as the MeerKAT synthesized beam becomes significantly non-Gaussian. The final step (2GC) begins with a masked deconvolution before using the model image for direction-independent (DI) self-calibration with CUBICAL (Kenyon et al. 2018). We then performed a second round of masked deconvolution using the DI self-calibrated visibilities. Given the high image fidelity of the 2GC images, we did not perform any direction-dependent self-calibration and adopted the 2GC images as our final data products.

We measured the radio properties of the source in each image using the CASA task `imfit`. The task fits an elliptical Gaussian component in a small sub-region around the source, measuring the position and flux density. As the source was unresolved, we fixed the component shape to be the synthesized beam of each image. We measured the uncertainty on the flux measurement using the local root-mean-square (rms) noise. For each epoch, we extracted the rms from an annular region using the CASA task `imstat`. Each annulus was centered on the position of the Gaussian component. The inner radius was fixed at the major axis of the synthesized beam, and we scaled the outer radius such that the annular area was equal to 100 synthesized beams.

2.4. Swift X-ray Telescope

As a part of our ThunderKAT monitoring, for each radio epoch, we acquired quasi-simultaneous (within \sim days) X-ray observations using the X-ray telescope aboard the Neil Gehrels Swift Observatory (henceforth Swift/XRT; Gehrels et al. 2004; Burrows et al. 2005). The weekly-cadence monitoring began on 2022 June 1 (MJD 59731) and continued until 2022 August 8 (MJD 59798), totaling 10 individual epochs (target ID: 31222). We increased the cadence of our Swift/XRT monitoring (one observation every ≤ 3 days) after the source’s X-ray flux began decaying. We acquired an additional 12 epochs of this ‘high-cadence’ monitoring that began on 2022 August 10 (MJD 59801) and ended on 2022 August 24 (MJD 59815). There was no further Swift/XRT monitoring.

We extracted 0.5–10 keV light curves using the Python API version of the Swift/XRT pipeline, `swifttools`

¹ <https://github.com/IanHeywood/oxkat>

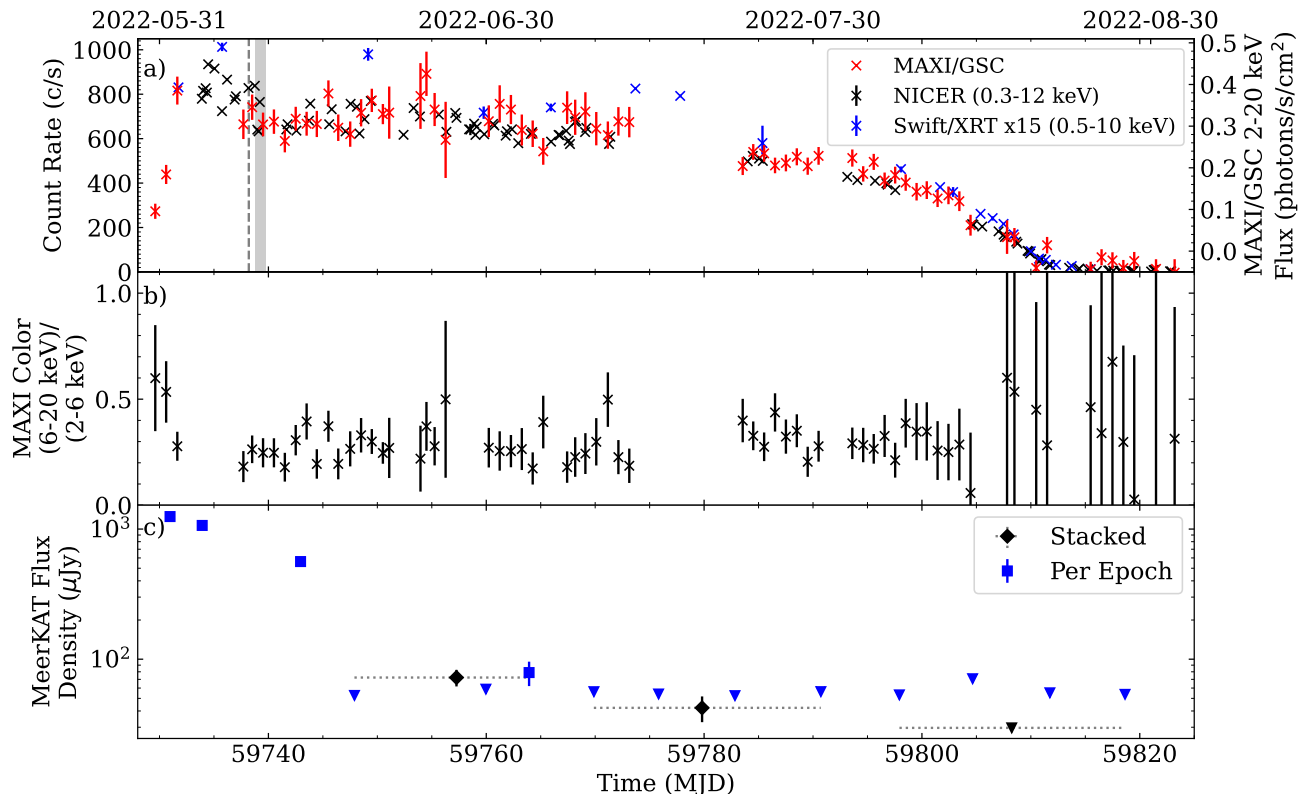


Figure 1. a) Count rate evolution of the rise and decay of the outburst of 1A 1744–361 from NICER (black; 10 ks bins; 0.3–12.0 keV), Swift/XRT (blue; rescaled by a factor of 15; 0.5–10 keV), and MAXI/GSC (red; fluxes in 24 h bins; 2–20 keV). MAXI/GSC data were extracted from the MAXI On-demand service (Matsuoka et al. 2009). The vertical dotted line corresponds to the absorption dip (started on MJD 59738.19985, lasted for ~ 100 s) that is shown in Figure 4. The shaded gray region corresponds to the NuSTAR observation interval. b) Temporal evolution of the X-ray color as defined by the ratio of 6–20 keV and 2–6 keV fluxes from MAXI/GSC; c) 1.3 GHz flux densities (μJy) from MeerKAT. Blue squares and filled triangles are the flux density per observation and the upper limits, respectively. Black diamonds denote stacked observations to increase the S/N of the detection (the corresponding observation epochs included are indicated by the dotted lines). The last data point for stacked observations (black filled triangle) is the 3σ upper limit at $30 \mu\text{Jy}$. The values for the flux densities are given in Table A1 in the Appendix.

(Evans et al. 2007, 2009), extracting a single data point for each observation, shown in panel a of Figure 1. Although our X-ray spectroscopy results presented in this work largely derive from the extensive and more sensitive NICER coverage, we also extracted the spectrum from the first Swift/XRT observation on 2022 June 1 (MJD 59731) with the online Swift/XRT product generator, since this observation was two days prior to the first NICER observation (Evans et al. 2009).

3. RESULTS

We first show the evolution of the 2022 outburst of 1A 1744–361 in panel a) of Figure 1 with NICER, Swift/XRT coverage, and MAXI/GSC monitoring. The NuSTAR observation interval is represented by the gray

shaded region. For context, we also show light curves from long-term monitoring of the field by RXTE/ASM and MAXI/GSC in Figure 2. The 2022 outburst of 1A 1744–361 had the longest known outburst duration from the source, and had a similar brightness level to the 2003 outburst.

To gain an understanding of the spectral variability of 1A 1744–361, using NICER data, we constructed the HID which is shown in Figure 3, where the soft color was defined as the ratio of background-subtracted count rates in the 1.8–3.5 keV and 0.5–1.8 keV bands, and the X-ray intensity is represented by background-subtracted count rates in the 0.5–6.8 keV band. We also generated the CCD, shown in Figure 3. We used the same definition for the soft color, and for the hard color we

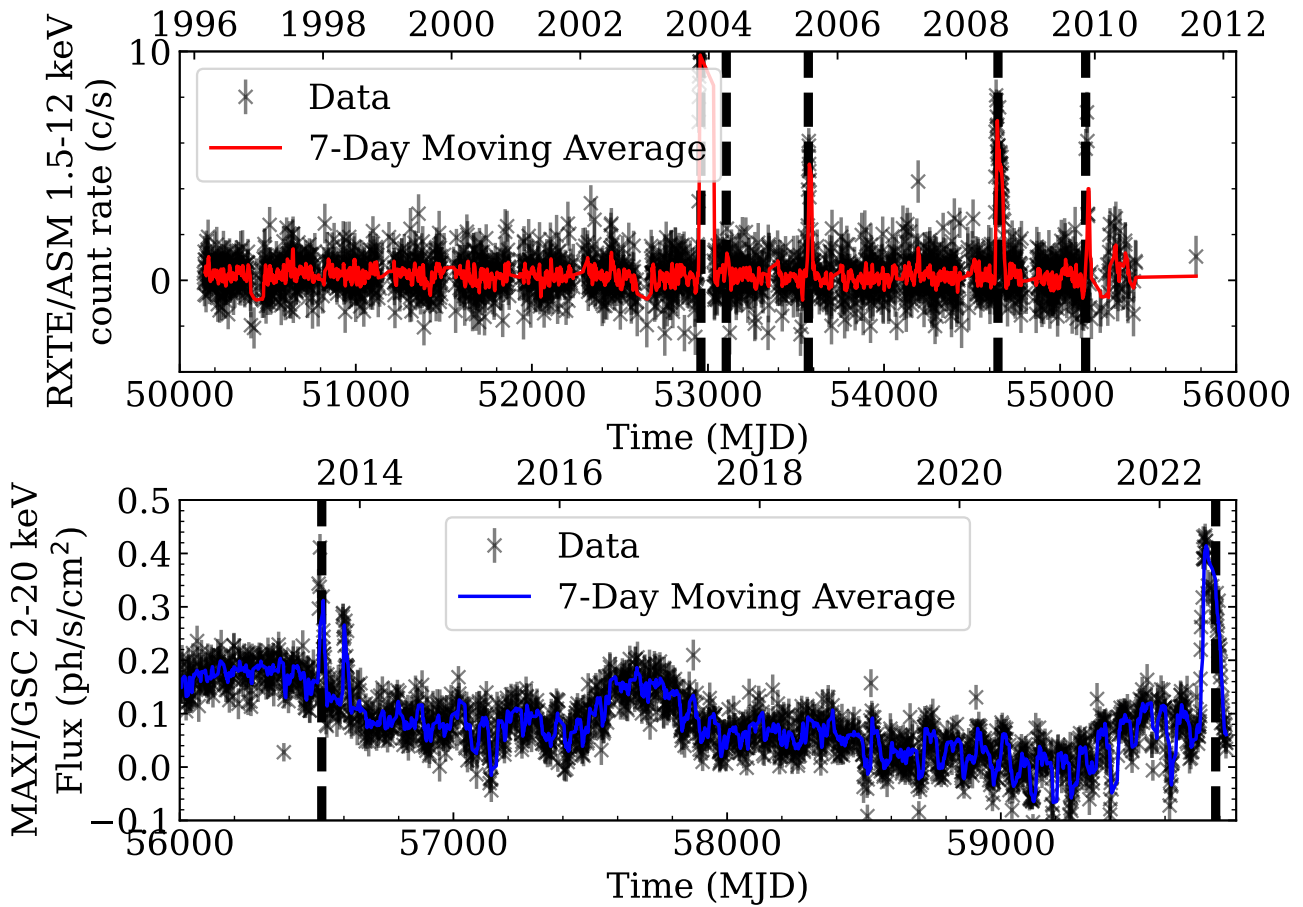


Figure 2. Long-term monitoring of 1A 1744–361 with RXTE/ASM (top; 1.5–12 keV count rate) and MAXI/GSC (bottom; 2–20 keV flux); black crosses and solid lines show the data points and the 7-day moving average, respectively. The secondary top x-axes for each panel displays the date in calendar years, with each tick corresponding to January 1 of the labeled year. It is worth noting that there is also another X-ray source (4U 1746–37) in the same MAXI/GSC field of view ($0^{\circ}9$ away) as 1A 1744–361. The black dashed vertical lines indicate rough dates of the outbursts of 1A 1744–361 that have been reported in the literature - they correspond to MJDs 52960 (2003-11-17), 53104 (2004-09-04), 53570 (2005-07-07), 54647 (2008-06-30), 55145 (2009-11-10), and 56507 (2013-08-03).

had the ratio of background-subtracted count rates in the 5.2–6.8 keV and 3.5–5.2 keV bands. In order to increase the signal-to-noise (S/N) in each data point, we increased the bin size to 256 s (from 128 s) after the broadband (0.3–12.0 keV) count rate fell below 50 c/s (around MJD 59811). The background-subtracted light curves were generated with `nicer13-1c` using the ‘spaceweather’ (sw) background model. Figure 3 clearly shows the source to be in the atoll-state throughout the outburst (Hasinger & van der Klis 1989; van der Klis 2004).

In the ObsID 5202800106, we found a ~ 100 s absorption dip, starting at MJD 59738.19431 (TT), in the light curve. The corresponding light curve and spectral color evolution from the absorption dip and the persis-

tent emission immediately after are shown in Figure 4, where the absorption dip is accompanied by an increase in the color, presumably due to the scattering of softer photons by an intervening gas cloud. The time intervals of the persistent emission post-absorption dip are indicated by the red crosses in the HID (Figure 3).

We also note that to increase the S/N of the observations at the tail end of the outburst, we attempted to combine several observations. The ObsID 5202800160 (see Table A2) was generated by combining (using the `ftmerge` tool) all housekeeping, auxiliary, and event files from ObsIDs 5202800149 and 5202800150. We were not able to combine subsequent observations as either the combined data exhibited a very large spread of colors (e.g., because too many observations were combined),

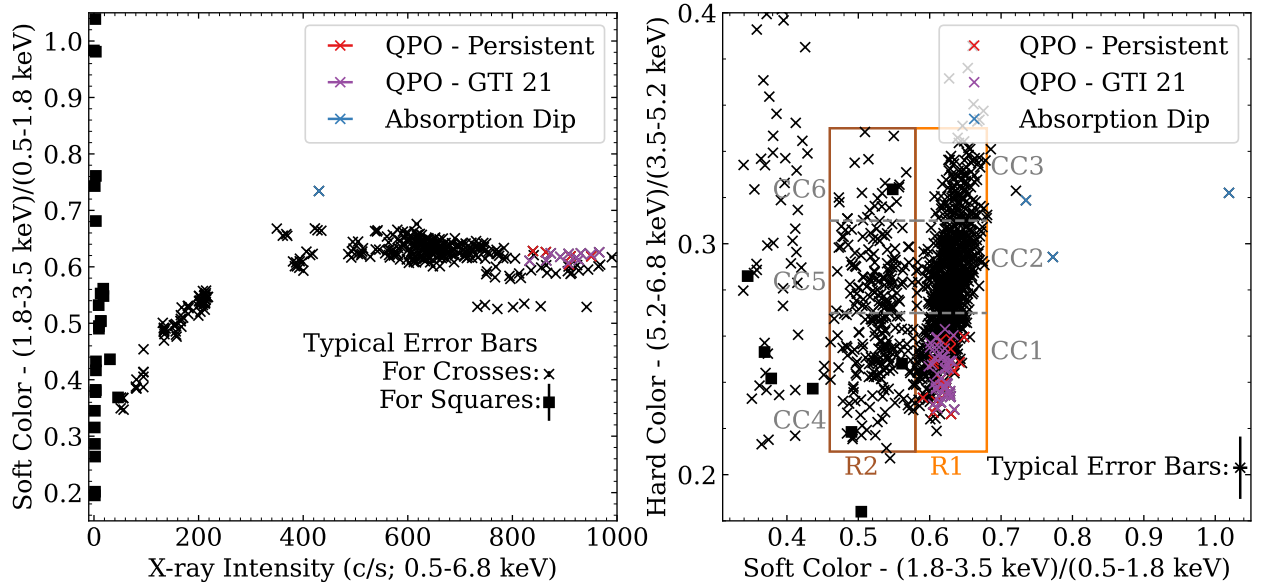


Figure 3. Left: Hardness-intensity diagram of 1A 1744–361 throughout the outburst, which is typical of the atoll-state. Crosses correspond to 128 s bins for the data, and squares correspond to 256 s bins (for bins with count rates ≤ 50 c/s, after MJD ~ 59811). The red crosses correspond to the time interval containing the persistent emission post-absorption dip; the blue cross corresponds to the absorption dip; purple crosses correspond to the GTI that showed the second instance of the QPO (see text for details). Right: Color-color diagram of 1A 1744–361 throughout the outburst. We display a typical error bar for the 32 s bins; we note the typical error bar for the 256 s bins was 6 times larger on average (not shown). Crosses correspond to 32 s bins for the data, and squares correspond to 256 s bins (for bins with count rates ≤ 50 c/s, after MJD ~ 59811). We also show two regions (denoted R1 and R2) in the CCD for which we conducted pulsation searches (see §3.1.2); we also broke up the two spectral regions into six smaller regions (CC1 to CC6) for further power spectral analysis. Background subtraction has been taken into account for all data points.

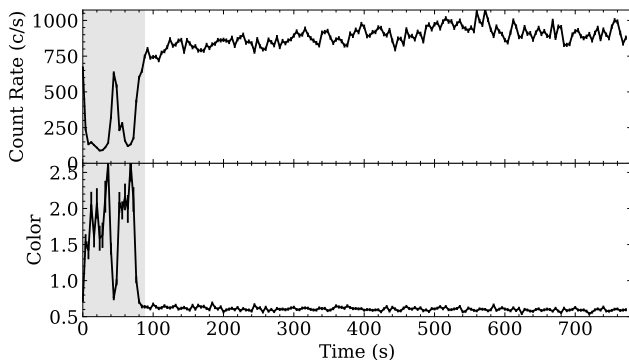


Figure 4. Top: 0.3–10.0 keV 4 s-binned NICER light curve corresponding to the absorption dip ($t = 0$ corresponds to MJD 59738.19431 (TT)). Bottom: The color is defined as the ratio of 2–10 keV and 0.3–2 keV background-subtracted count rates. The gray shaded region corresponds to the absorption dip; the rest of the GTI is referred to as the persistent emission post-absorption dip. The absorption dip corresponds to an increased spectral color/hardness.

or the background dominated the source spectrum well below 10 keV. Thus, we ended up not using data from

ObsIDs 5202800151 to 5202800157 (2022 August 25 to 2022 August 31) in our final analysis.

Finally, to track the outburst rise, we calculated the temporal evolution of the X-ray color from MAXI/GSC data (first three data points in panel b of Figure 1), and corroborated the evolution with the Swift/XRT spectrum of the first observation on 2022 June 1. We used the Windowed Timing observation (since 1A 1744–361 is bright) which had 1.4 ks of exposure. The 0.5–10.0 keV spectrum was well described by an absorbed cutoff power law with $n_H = 0.44 \times 10^{22} \text{ cm}^{-2}$ (see §3.2.1), power law index $\Gamma_{\text{co}} = 0.61^{+0.04}_{-0.04}$, e-folding energy $E_f = 3.21^{+0.18}_{-0.16} \text{ keV}$, and normalization $0.479^{+0.007}_{-0.007} \text{ photons s}^{-1} \text{ cm}^{-2} \text{ keV}^{-1}$. The 0.6–10.0 keV absorbed flux was approximately $F \approx 2.6 \times 10^{-9} \text{ erg s}^{-1} \text{ cm}^{-2}$ (uncertainties were reported at 90% confidence for this spectral fit).

3.1. X-ray Timing

3.1.1. Aperiodic Analysis

In order to look for QPOs, we first created an averaged Leahy-normalized power spectrum for each individual GTI longer than 50 s, with 32 s segments. We made

use of the `AveragedPowerspectrum` class in `Stingray` to perform the calculation (Huppenkothen et al. 2019a,b), where we also logarithmically binned the averaged power spectra with $f = 1.05$. Almost all of the power spectra were well-fit with a power law plus a constant; we show the results in Table A2 in the Appendix. However, the power spectra corresponding to GTIs 12 and 21 (see Table A2), observations that contained the absorption dip (see Figure 4) and an observation one day after, respectively, exhibited a residual feature at around $\nu \simeq 8$ Hz.

For GTI 12 (in ObsID 5202800106), we investigated the origins of the residual feature by isolating the dip emission and the persistent emission immediately after the dip (within the GTI). We found that the residual feature was seen only in the persistent emission and not in the absorption dip, though with a 2σ upper limit of 6.8%, so we cannot rule out the QPO being absent in the dip emission. Thus in what follows, we focused the analysis on the persistent emission immediately following the dip. We proceeded to model the residual by adding a Lorentzian on top of the initial continuum model (power law plus constant) and found a significant improvement. To quantify the improvement of the additional model component, we considered the difference in the Akaike information criterion (AIC; Akaike 1974), with $\Delta\text{AIC}_{12} = \text{AIC}_1 - \text{AIC}_2$, where

$$\text{AIC} = -2\ln\mathcal{L} + 2d, \quad (1)$$

\mathcal{L} is the maximum likelihood of any given model ($\mathcal{L} \propto e^{-\chi^2/2}$), χ^2 is the chi-squared of the fit, and d is the number of degrees of freedom (Liddle 2007; Tan & Biswas 2012; Arcodia et al. 2018). To evaluate the two models, we can define a rejection probability at a confidence level defined by

$$P_1/P_2 \approx e^{-\Delta\text{AIC}_{12}/2}. \quad (2)$$

In Figure 5, we show the fractional rms-squared-normalized power spectra, the residuals when fitting with an initial power law plus constant continuum, and the residuals when including an additional Lorentzian component. We found that the χ^2 values decreased from $\chi^2 = 304.6$ to $\chi^2 = 199.2$ (from 163 to 160 d.o.f.), thus $\Delta\text{AIC}_{12} = 111.4$, and so $P_1/P_2 \approx 6.5 \times 10^{-25}$. We found the Lorentzian (QPO) parameters to be $\nu_0 = 7.64 \pm 0.19$ Hz with $\text{FWHM} = 3.0 \pm 0.6$ Hz, which implied a quality factor, $Q = \nu_0/\text{FWHM} = 2.5 \pm 0.5$. The fractional rms amplitude of the QPO (from the normalization of the Lorentzian) was found to be $f_{\text{rms},1} = 6.0 \pm 0.4\%$ over 0.3–12.0 keV.

We have found an additional instance of the QPO in one GTI of an observation about one day later, in GTI

21 (in ObsID 5202800107). The QPO had Lorentzian parameters $\nu_0 = 8.75 \pm 0.25$ Hz with $\text{FWHM} = 3.4 \pm 0.7$ Hz, which implied $Q = 2.6 \pm 0.5$ ($\chi^2 = 247.7$ to $\chi^2 = 164.1$; $P_1/P_2 \approx 3.5 \times 10^{-20}$). The fractional rms amplitude of the QPO was $f_{\text{rms},2} = 4.6 \pm 0.4\%$ over 0.3–12.0 keV. We also show the corresponding power spectrum and residuals in the right panels of Figure 5.

When we visually inspected the power spectral fits for all GTIs, we noticed a weak ($\sim 2\sigma$) indication of the QPO in GTI 6 (in ObsID 5202800103). To investigate this further, we ran `nicerclean` on the unfiltered event file (ufa file) and generated the cleaned event file. We found another instance of the QPO in the GTI with the highest count rates (MJDs 59735.31189 to 59735.31792 (TDB)), where we needed a power law and an additional Lorentzian ($\chi^2 = 257.9$ to $\chi^2 = 148.2$; $P_1/P_2 \approx 7.5 \times 10^{-26}$) to describe the power spectrum. The Lorentzian (QPO) parameters were $\nu_0 = 7.34 \pm 0.14$ Hz, $\text{FWHM} = 2.3 \pm 0.4$ Hz, thus $Q = 3.1 \pm 0.6$. Finally, the fractional rms amplitude was $f_{\text{rms},3} = 6.2 \pm 0.4\%$ over 0.3–12.0 keV (note that we did not plot the power spectrum given the non-standard filtering).

We also looked for energy dependence of the QPO, where we fixed the centroid frequency and FWHM, and subdivided the data into energy bands of 0.3–1.0, 1.0–2.0, 2.0–3.0, 3.0–4.0, and 4.0–12.0 keV, tabulating the results in Table 2. We did not find evidence for the energy dependence of the fractional rms amplitude for the QPO, which is consistent with the behavior of 1 Hz QPOs from dipping/eclipsing LMXBs (Jonker et al. 1999; Homan 2012) but not with sources that show NBO-like oscillations (Wijnands et al. 1999; Wijnands & van der Klis 1999). However, we note that the energy coverage and sensitivities differ between RXTE and NICER (and uncertainties here were large).

We characterized the power spectra from the different spectral regions in the CCD (based on the soft color) as shown in Figure 3, summarizing these results in Table 1. In order to understand the variability within each spectral region, we further subdivided each region into three spectral subregions based on the hard color and repeated the power spectral characterization (see Table 1). The 0.03–50 Hz variability (characterized by the fractional rms amplitude) in R1 turned out to be primarily due to the spectral subregion occupied by the QPO in the lower region of the soft state; the variability decreased going up the soft state (R1) from $7.3 \pm 1.4\%$ to $2.0 \pm 1.5\%$. For R2, there were no discernible differences between the spectral subregions within uncertainties. It is worth noting that the source moved erratically through the CCD throughout the outburst, so we cannot ascertain any smooth temporal evolution of the power spectrum.

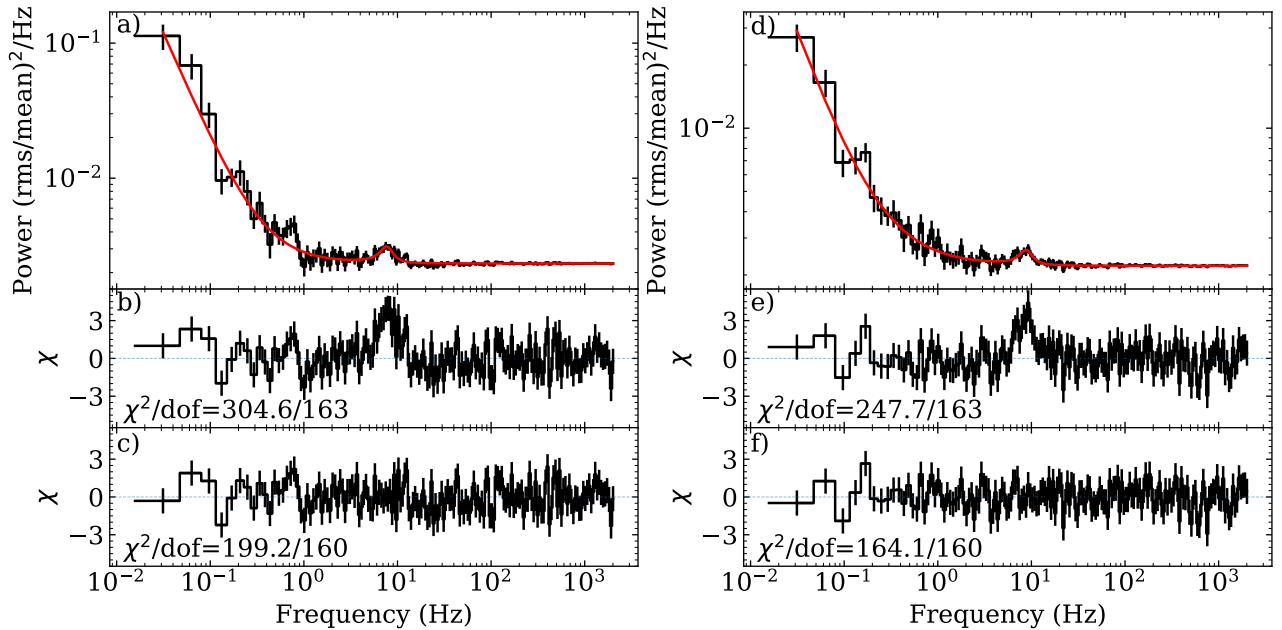


Figure 5. a) Power spectrum corresponding to the persistent emission post absorption-dip. The best-fit model, comprising a power law (and a constant for the Poisson level) and a Lorentzian for the QPO, is in red. b) Residuals [$\chi = (\text{data-model})/\text{error}$] when fitting the data to only a power law (plus constant) - the reduced χ^2 value is given by $\chi^2/\text{d.o.f.} = 304.6/163$. c) Residuals when the Lorentzian is included in the fit - the reduced χ^2 value is given by $\chi^2/\text{d.o.f.} = 199.2/160$. d), e), and f) are similar plots, but they correspond to GTI 21 (in ObsID 5202800107), the segment of data showing the second QPO instance.

We have also plotted the HID using data points from ObsID 5202800107, with an observation span of five hours, in Figure 6. The HID shows two separate tracks, and only Z-state sources are known to trace out multi-branched tracks on such short timescales (e.g., GX 13+1: Schnerr et al. 2003; Homan et al. 2004; Fridriksson et al. 2015; Cyg X-2: Fridriksson et al. 2015). However, given the limited data, we caution over-interpretation of these tracks as directly related to that of the branches in the Z-state.

3.1.2. Pulsation Searches

Given the discovery of coherent burst oscillations at ~ 530 Hz from 1A 1744-361 (Bhattacharyya et al. 2006a), the pulsation frequency must be around the same frequency (Chakrabarty et al. 2003; Watts 2012). Thus we conducted a coherent pulsation search with two independent techniques — stacked power spectra and acceleration searches (only for NICER data).

For the stacked power spectra, we first identified search regions in the CCD (see Figure 3), as the appearance of pulsations could be associated with different spectral states of the source (Vaughan et al. 1994). For each region (R1 and R2), the corresponding time windows of the 32-s bins were collated and joined, with

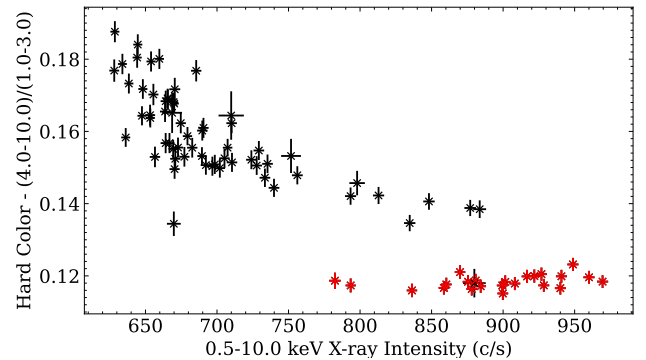


Figure 6. HID with data points from NICER ObsID 5202800107, where the third QPO instance appeared. The hardness is defined as the hard color, with the count rate ratio of 4–10 keV over 1–3 keV; the intensity is the NICER 0.5–10.0 keV count rate. The red and black points are data points corresponding to GTIs that do and do not exhibit the QPO, respectively. The HID shows two separate tracks, which only Z-state sources have shown before on short timescales, but given the limited data, we cannot confidently connect these branches to the Z-state tracks.

new GTI files generated for each region. We then used `ftselect` to extract the events given the filtered GTI file. The final averaged power spectra were calculated

Table 1. Best-fit parameters from fitting the power spectra with a power law plus constant continuum for each spectral region defined in Figure 3.

Spectral Region	$f_{\text{rms},0.03-50}$ (%)	m	No. photons	$\chi^2/\text{d.o.f.}$	PL Amplitude	PL α	PL Constant
R1	5.6 ± 1.4	411	17510405	474/360	$(3.56 \pm 0.11) \times 10^{-4}$	1.061 ± 0.012	$(3.0027 \pm 0.0006) \times 10^{-3}$
R2	10.5 ± 2.4	123	1995884	683/360	$(1.08 \pm 0.06) \times 10^{-3}$	1.13 ± 0.02	$(7.883 \pm 0.003) \times 10^{-3}$
CC1	7.3 ± 1.4	95	4919162	464/360	$(5.7 \pm 0.2) \times 10^{-4}$	1.106 ± 0.018	$(2.4706 \pm 0.0010) \times 10^{-3}$
CC2	4.5 ± 1.4	128	5126233	369/360	$(2.10 \pm 0.19) \times 10^{-4}$	1.02 ± 0.03	$(3.1937 \pm 0.0011) \times 10^{-3}$
CC3	2.0 ± 1.5	22	848125	425/360	$(2.6 \pm 2.1) \times 10^{-5}$	1.5 ± 0.2	$(3.316 \pm 0.003) \times 10^{-3}$
CC4	11.5 ± 2.1	34	847789	658/360	$(1.12 \pm 0.08) \times 10^{-3}$	1.12 ± 0.03	$(5.126 \pm 0.004) \times 10^{-3}$
CC5	10.1 ± 2.7	14	211120	551/360	$(1.2 \pm 0.7) \times 10^{-4}$	1.74 ± 0.19	$(8.492 \pm 0.009) \times 10^{-3}$
CC6	13.5 ± 9.5	3	3584	303/360	-0.021 ± 0.003	0.43 ± 0.06	0.1085 ± 0.0008

NOTE—We present the fractional rms amplitude (%) over 0.03–50 Hz, the number of averaged segments for the power spectrum (m), the number of photons, the χ^2 values, the power law amplitudes, spectral indices (α), and the constant value.**Table 2.** Energy dependence of the QPO.

Energy (keV)	$f_{\text{rms},1}$ (%)	$f_{\text{rms},2}$ (%)	$f_{\text{rms},3}$ (%)
0.3–1.0	...	3.9 ± 1.8	...
1.0–2.0	6.3 ± 0.5	4.4 ± 0.5	5.3 ± 0.5
2.0–3.0	5.8 ± 1.3	5.8 ± 0.9	6.8 ± 0.9
3.0–4.0	4.2 ± 3.6	< 9.8 (2σ)	...
4.0–12.0	< 45.9 (2σ)	< 11.9 (2σ)	7.7 ± 2.1
0.3–12.0	6.0 ± 0.4	4.6 ± 0.4	6.2 ± 0.4

NOTE—We explored the energy dependence of the QPO by fixing the centroid energy and FWHM for each respective QPO instance, and fitting for the normalization. We calculated 2σ upper limits if the uncertainty was larger than the best estimate. Energy bands for which no data are listed indicate non-convergent fits. The QPO instances occurred in GTIs 12, 21, and 6, respectively (see Table A2 for corresponding ObsID and MJDs).

using a segment size of 256 s and a bin size of $\Delta t = 2^{-11}$ s using the `AveragedPowerspectrum` routine in `Stingray` (Huppenkothen et al. 2019a,b).

The averaged power spectrum from each of the two search regions are shown in Figure 7. We did not find any significant coherent periodicity in any of the search regions, but we can place upper limits on the fractional rms amplitude. Since we did not know the precise spin frequency of the pulsar, and given the relatively low significance ($\sim 4\sigma$) of the burst oscillation (Bhattacharyya et al. 2006a), we calculated the 3σ upper limits on the fractional rms amplitude over 1–1000 Hz. We found that for R1 and R2, the upper limits were roughly 0.04% and 0.13%, respectively. We also found that for the 3–25 keV NuSTAR data (using FPMA), the 3σ upper limit

was 0.10%. We derived these upper limits by generating 1000 realizations of the event lists and GTIs corresponding to each of the regions (for NuSTAR, we used the entire FPMA event list). Across each frequency bin, we generate a cumulative distribution function (CDF) of the simulated powers, and took the power value where the CDF was equal to 0.9973 (3σ upper limit), and calculated the corresponding fractional rms amplitude (Vaughan et al. 1994; Munro et al. 2002a). Operationally, the simulations were carried out with `simulate_times` in `Stingray` (Huppenkothen et al. 2019a,b).

We used acceleration searches to probe for coherent periodicity while accounting for the Doppler modulation of the pulsation frequency due to an orbit (Ransom et al. 2002). Acceleration searches are most sensitive with segments (with length T) where $T \lesssim P_{\text{orb}}/10$, such that the pulsar acceleration is roughly constant and so the frequency evolution is linear within the time segment (Ransom et al. 2002). We carried out the acceleration searches on the individual GTIs throughout the outburst of 1A 1744–361, with a minimum of 100 s in length (to have enough photons). A total of 108 GTIs were searched and the median segment length was 442 s, with a standard deviation of 299 s. We searched over the frequency range 525–535 Hz, energy ranges 0.3–2.0 keV and 2.0–10.0 keV. We found no significant coherent periodicity candidates from the acceleration searches.

3.1.3. NuSTAR

We calculated the averaged (temporal) cospectrum from the entire NuSTAR observation with 8 s segments and fit a simple power law to the observed cospectrum; the cospectrum and the corresponding residuals are shown in Figure 8. The cospectrum is well-described by a power law, $f(\nu) = A\nu^{-\alpha}$, with amplitude

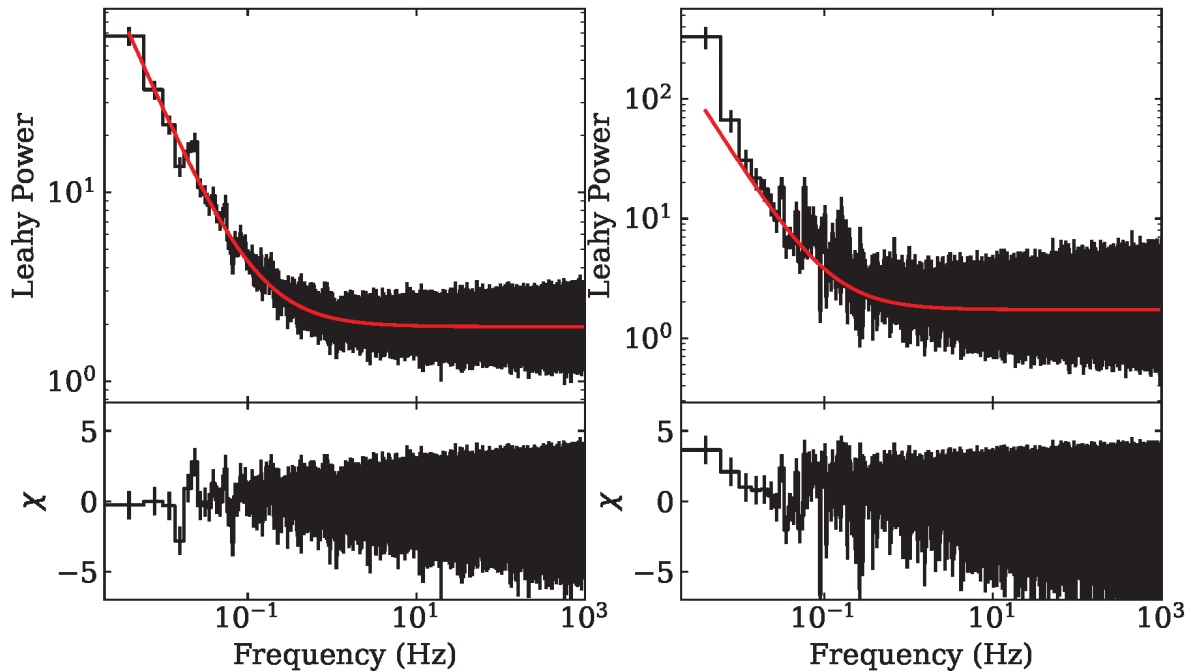


Figure 7. Averaged power spectra (256 s segments) using the events in R1 (left) and R2 (right), where $m = 76$ and $N_\gamma = 12897259$, and $m = 23$ and $N_\gamma = 1542304$, respectively. The red solid lines in the top panel are the best-fit power law (plus constant) model. The panels below show the residuals after subtracting the best-fit model from the data.

$A = (4.1 \pm 0.7) \times 10^{-4}$, spectral index $\alpha = 1.15 \pm 0.10$, with the reduced χ^2 being 0.74 (69 d.o.f.) The 0.1–50 Hz fractional rms amplitude was $2.7 \pm 2.4\%$. The cospectra were calculated using *Stingray* with data from the two independent FPMA and FPMB detectors, such that any periodic (or quasi-periodic) signals in phase between the two detectors can be detected and any unrelated variability is eliminated (e.g., dead time effects; Huppenkothen & Bachetti 2018; Huppenkothen et al. 2019a,b). The cospectrum is the real part of the cross power density spectrum (Bendat & Piersol 2011).

To further investigate the cospectral properties of 1A 1744–361, we first constructed the HID, which is shown in Figure 9. The HID (on the right) clearly shows three distinct spectral states which have been color-coded, and the corresponding points are shown in the light curve (on the left). In particular, the blue points corresponding to flaring activity in the light curve (Figure 9a) show up as a Z-state-like flaring branch in the HID (Figure 9b). For each of the three spectral regions, we calculated the averaged cospectrum (also 8 s segments) and we found that the power law amplitudes and spectral indices were roughly consistent with each other. There is weak evidence of a slightly shallower power law at lower count rates of the source (lower soft

colors). We did not see evidence of any QPOs in all of the cospectra. The best fit parameters are summarized in Table 3.

Table 3. Best-fit parameters from the power law fit to the NuSTAR cospectra defined from Figure 8.

Soft Color	PL Amplitude 10^{-4}	α	f_{rms} %	$\chi^2/\text{d.o.f.}$
All	4.1 ± 0.7	1.15 ± 0.10	2.7 ± 2.4	51/69
< 0.6	5.0 ± 0.9	0.87 ± 0.10	2.6 ± 1.9	94/69
[0.6, 0.78]	3.6 ± 1.4	1.26 ± 0.22	2.8 ± 2.6	62/69
> 0.78	2.1 ± 3.0	1.7 ± 0.7	3.1 ± 2.9	82/69

NOTE—The power law amplitudes and indices, the 0.1–50 Hz fractional RMS amplitudes (%), and χ^2 values are shown.

3.2. X-ray Spectroscopy

3.2.1. Evolution of the Continuum Components

The broadband 0.6–10.0 keV (see below for the choice of energy range) spectra were generally well-fit by several continuum models that were a function of the spectral states (R1, R2) as identified in Figure 3. We found

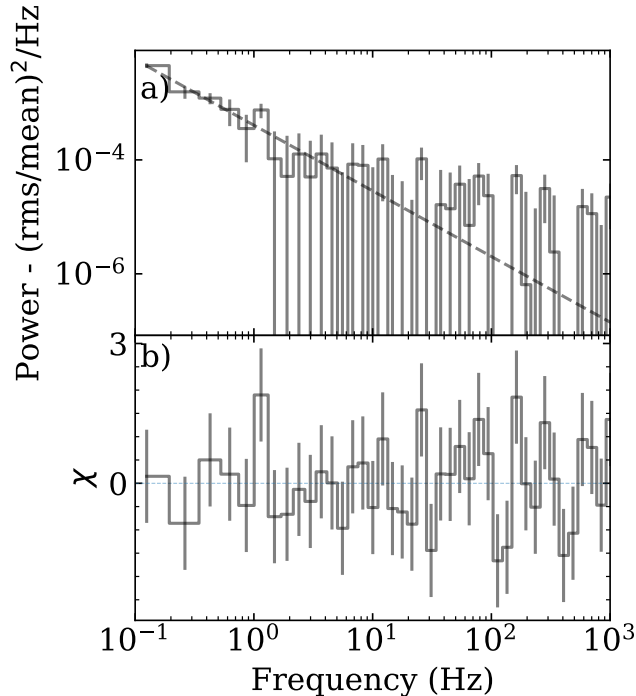


Figure 8. a) The 3–25 keV (temporal) cospectrum for the entire NuSTAR observation (constructed from 8 s segments); b) the residuals after fitting the observed cospectrum with a power law with index 1.15 ± 0.10 .

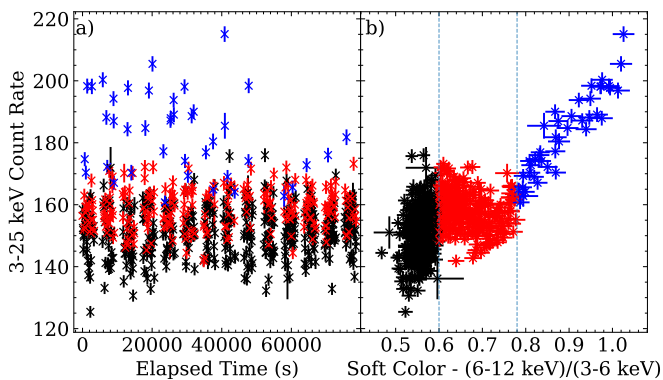


Figure 9. a) 3–25 keV NuSTAR light curve (combined FPMA and FPMB); b) 3–25 keV count rate vs. soft color, defined as the ratio of counts in the 6–12 keV and 3–6 keV bands. All data points are constructed from 64 s bins. The black, red, and blue data points corresponding to different spectral branches for which we investigated the cospectral properties - see Table 3 for the best fit parameters to the (temporal) cospectra.

that adopting a single spectral model for the entire outburst did not suffice. In particular, the source transitioned from R1 to R2 between MJDs 59797 and 59804 (NICER data gap). Thus we performed spectral fits within the following framework – for GTIs 0 to 93 (MJDs 59733 to 59797, inclusive), we adopted an underlying continuum consisting of a disk blackbody (for the accretion disk) and a thermal blackbody (possibly from the boundary layer at the NS surface). For subsequent observations, we performed spectral fitting on an individual ObsID basis to maximize the S/N. We tried a combination of a thermal blackbody (for the NS surface) and either a power law or a cutoff power law for the Comptonization component. We found that the cutoff power law worked best ($E_{\text{cut}} \approx 7 - 11$ keV) between ObsIDs 5202800140 and 5202800144 (MJDs 59804 to 59808), when the source was in R2, and the high energy cutoff after that was being pegged at 200 keV by the spectral fitting process, which indicated that a simple power law sufficed. Thus from ObsID 5202800145 onwards (MJD 59809), we employed a thermal blackbody and a power law. We note that we ultimately did not use Comptonization models (e.g., `thcomp`) as NICER did not have sufficient high energy coverage to provide meaningful constraints. We also tried including a partially ionizing absorber component `zxipcf` to the fit, but the data lacked constraining power. The fitting strategy described above was similarly adopted for NS LMXBs Aql X–1 and 4U 1608–52 (Lin et al. 2007).

The full results from the spectroscopic analysis are presented in Figure 10. For spectral region R1, we fit the spectra with an absorbed disk blackbody model and a thermal blackbody component. Some of the observations in R1 exhibited clear residuals around 1.0 keV (an emission feature) and around 7.0 keV (an absorption feature), so we added two Gaussian components (see §3.2.2). We restricted the energy range of the ~ 1 keV (~ 7 keV) component to be 0.85–1.15 keV (6.9–7.1 keV), and the line width to be between 0.001–0.2 keV. For the ~ 7 keV absorption line, we allowed the normalization of the Gaussian to be negative. In XSPEC parlance, the full model was expressed as `tbabs(diskbb+bbbodyrad+gauss+gauss)`. We tried fitting the absorption line with the multiplicative `gabs` model but we recovered very high (unphysical) values of the line depth.

In all spectral fits, the line of sight hydrogen column density was fixed at $n_H = 0.44 \times 10^{22} \text{ cm}^{-2}$, roughly consistent with the HI4PI value of $0.26 \times 10^{22} \text{ cm}^{-2}$ (HI4PI Collaboration et al. 2016), which was determined through the following. We initially performed the spectral fit (with the absorbed disk blackbody and

thermal blackbody model) to all observations on an individual GTI basis by letting all spectral parameters be free. The value of n_H typically ranged between $(0.39\text{--}0.45)\times 10^{22}\text{ cm}^{-2}$ with typical errors of $(0.01\text{--}0.03)\times 10^{22}\text{ cm}^{-2}$, so we calculated a weighted average of the n_H values, which resulted in the derived value above. Hereafter, we fixed n_H and repeated the spectral fitting. The spectral fitting was restricted to 0.6–10.0 keV as the column density absorption dominated below 0.6 keV, and the background dominated above 10 keV.

Returning to the results, we see that in R1, there is a strong anticorrelation between the blackbody temperature and normalization (**bbodyrad** in **XSPEC**), as well as between the temperature at the inner (accretion) disk radius and the disk blackbody normalization (**diskbb** in **XSPEC**). In this period, both temperatures were rising (and normalizations decreasing) while the overall 0.6–10.0 keV (absorbed) flux was gradually decreasing. In R2, where the absorbed blackbody and cutoff power law model was used, the flux was rapidly decreasing, while the blackbody temperature and associated normalization were roughly constant. There was a slight increase in the power law index and the normalization. Finally, after R2, fit with an absorbed blackbody and power law model, the flux was also decreasing but there was no monotonic evolution in the spectral parameters. We note that the thermal blackbody values changed dramatically across R1 to R2; however, the continua employed in both states were different so we caution any physical interpretation here.

3.2.2. Evolution of the Spectral Lines

As mentioned in §3.2.1, some of the observations in R1 exhibited clear residuals around 1.0 keV and 7.0 keV. To understand the spectral lines in more detail, we combined all of the observations that appeared in R1 and R2, as well as in CC1-CC6. As an example, we show the ratio of the data to the baseline model (absorbed disk blackbody and thermal blackbody) in Figure 11 for spectral region CC1 to illustrate the strong absorption features. For the combined spectra in R1 and R2, the best-fit parameters are given in Table 4, and the corresponding equivalent widths (and upper limits) of the lines are given in Table 5. We noticed that in the combined data in R1, there was a narrow 6.7 keV absorption feature as well as a potentially broad 8 keV absorption feature, which are likely due to the Fe XXV and Fe XXVI $K\beta$ absorption lines, respectively.

In order to look for any finer evolution of the 1 keV and 7 keV spectral lines in the spectral regions, we ran all the spectral fits without both Gaussians, with one

Table 4. 0.6–10.0 keV spectral fits for the spectral regions R1 and R2, as defined in Figure 3.

Model	Parameter	R1	R2
bbodyrad	kT (keV)	$1.675^{+0.032}_{-0.018}$	$0.468^{+0.016}_{-0.016}$
	norm _{bbodyrad}	$26.5^{+1.3}_{-1.3}$	178^{+24}_{-23}
diskbb	T_{in} (keV)	$0.913^{+0.019}_{-0.012}$...
	norm _{diskbb}	116^{+6}_{-3}	...
cutoffpl	Γ	...	$1.13^{+0.03}_{-0.03}$
	E_{cut} (keV)	...	$6.1^{+0.4}_{-0.3}$
	norm _{cutoffpl}	...	$0.1334^{+0.0021}_{-0.0013}$
gauss (Fe XXV)	E_{XXV} (keV)	$6.72^{+0.03}_{-0.04}$...
	σ_{XXV} (keV)	$0.001^{+0.064}_{-0.001}$...
	norm _{XXV} (10^{-4})	$-1.8^{+0.9}_{-0.9}$...
gauss ($K\alpha$)	$E_{K\alpha}$ (keV)	$6.985^{+0.009}_{-0.012}$	$6.99^{+0.07}_{-0.09}$
	$\sigma_{K\alpha}$ (keV)	$0.02^{+0.03}_{-0.02}$	Unconstrained
	norm _{$K\alpha$} (10^{-4})	$-7.1^{+0.9}_{-1.1}$	$-0.6^{+0.4}_{-0.4}$
gauss ($K\beta$)	$E_{K\beta}$ (keV)	$8.06^{+0.08}_{-0.09}$...
	$\sigma_{K\beta}$ (keV)	$0.5^{+0.2}_{-0.3}$...
	norm _{$K\beta$} (10^{-3})	$-1.4^{+0.8}_{-0.9}$...
...	Flux ^a	$2.949^{+0.006}_{-0.009}$	$0.7453^{+0.0017}_{-0.0056}$
...	$\chi^2/\text{d.o.f.}$	181/159	109/142

^a Absorbed flux over 0.6–10.0 keV in units of $10^{-9}\text{ erg s}^{-1}\text{ cm}^{-2}$

NOTE—The uncertainties reported here are 90% confidence limits. The spectral parameters are the blackbody temperature (kT , keV), blackbody normalization (scaled by R_{km}^2/D_{10}^2 , where R_{km} is the source radius in km and D_{10} is the source distance in units of 10 kpc), temperature at the inner disk radius (T_{in} , keV), and the disk blackbody normalization (scaled by $(R_{\text{in}}/D_{10})^2\cos(\theta)$; R_{in} is the apparent inner disk radius in km, and θ is the disk inclination angle (in degrees)). We also have the power law photon index (Γ), the cutoff energy (E_{cut} , keV), and the normalization (flux in photons/keV/cm²/s at 1 keV). $E_{K\alpha}$ and $E_{K\beta}$ are the Gaussian centroid energies for the Fe XXVI $K\alpha$ and $K\beta$ absorption lines, respectively; $\sigma_{K\alpha}$ and $\sigma_{K\beta}$ refer to the same lines but describe the Gaussian width; norm _{$K\alpha$} and norm _{$K\beta$} refer to the same lines are the line normalizations in photons/s/cm².

Gaussian (each line individually), and with both Gaussians. If the fit improved by adding a line, whether an individual line (1 or 7 keV) or adding a second line on top of the first fit, at the 99.73% confidence level as assessed by the difference in AIC values (see Equations 1 and 2), the spectral line was considered a detection. There was no obvious correlation between the presence (or absence) of the spectral lines with the position in R1. The equivalent widths of both lines also did not suggest an obvious correlation.

The 1.0 keV emission line feature could possibly be due to either incomplete modeling of the interstellar

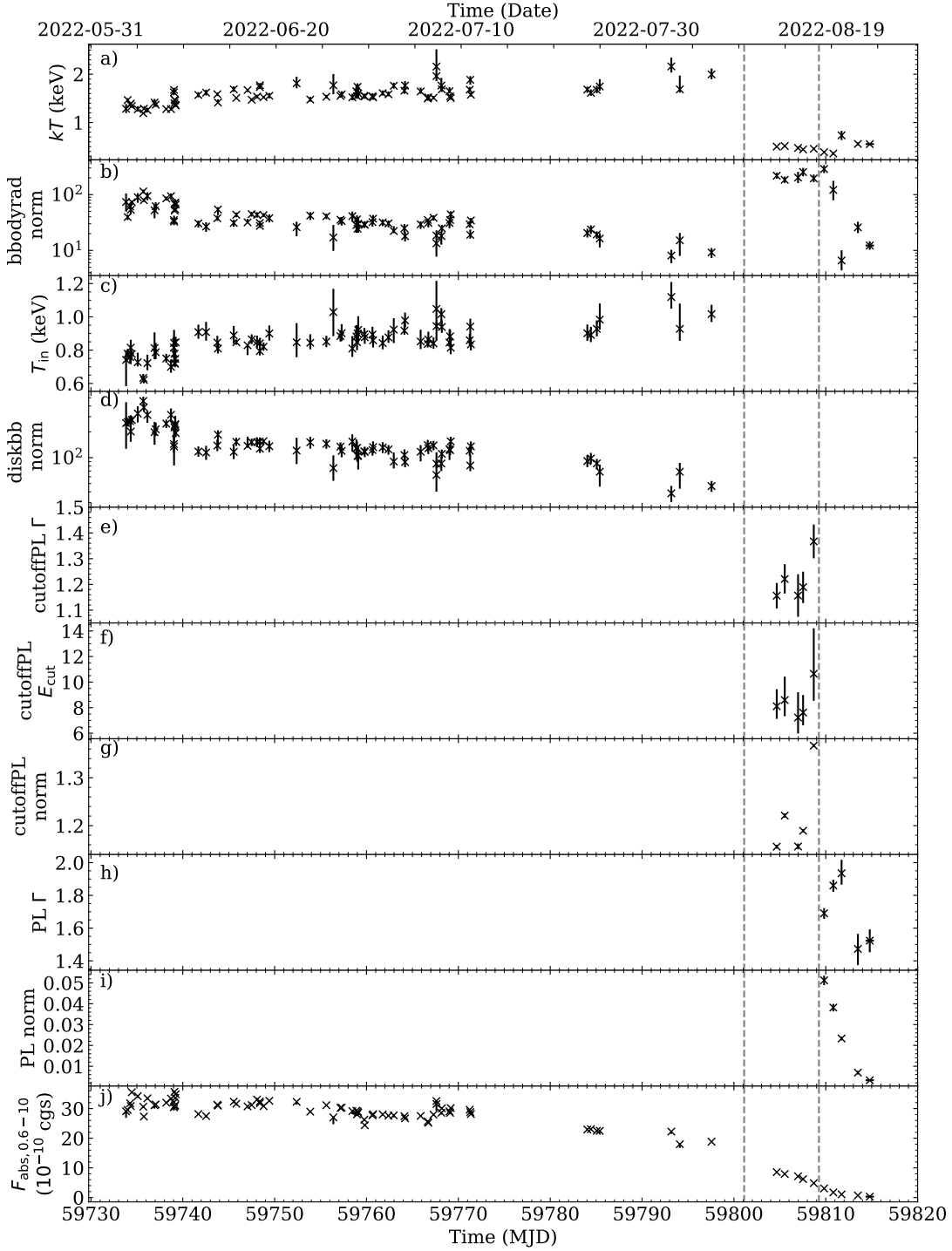


Figure 10. 0.6–10.0 keV spectral fit with XSPEC of individual NICER GTIs and observations (see text for details) throughout the outburst. The uncertainties reported here are 90% confidence limits. Panels a)–d) show results from R1 with the model $\text{tbabs}(\text{diskbb}+\text{bbodyrad}+\text{gauss}+\text{gauss})$; the results of the Gaussian lines are reported elsewhere in the text (see §3.2.2 and Table 5). Panels e)–g) show results from spectral fits in R2, with the model $\text{tbabs}(\text{bbodyrad}+\text{cutoffpl})$. The results corresponding to R2 are in between the two vertical dashed lines. Panels h) and i) show results from fitting the rest of the spectra (after the second vertical dashed line; after MJD 59809) with $\text{tbabs}(\text{bbodyrad}+\text{powerlaw})$. The power law photon index and

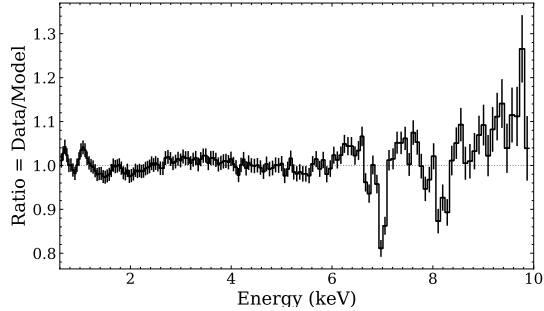


Figure 11. Ratio of the observed spectrum to the underlying model (`tbabs(diskbb+bbodyrad)`) as a function of energy for the spectral region CC1. Strong negative residuals, identified with Fe XXV, Fe K α , and Fe K β absorption features, are seen as unmodeled features in the ratio plot.

Table 5. Equivalent widths (in eV) for relevant lines.

Region	Exposure (s)	Line Energy			
		1 keV	6.7 keV	7 keV	8 keV
R1	34322	$4.8^{+2.7}_{-2.7}$	$-5.5^{+5.3}_{-4.3}$	-31^{+6}_{-4}	-46^{+14}_{-10}
R2	11451	< 2.7	> -20	-14^{+7}_{-9}	> -30
CC1	9710	$8.2^{+4.0}_{-3.0}$	-12^{+6}_{-8}	-37^{+6}_{-5}	-58^{+16}_{-16}
CC2	20590	$6.6^{+3.5}_{-4.3}$	> -7	-29^{+5}_{-4}	-36^{+14}_{-13}
CC3	4022	$2.6^{+2.2}_{-2.5}$	> -9	-22^{+6}_{-6}	> -40
CC4	3323	$4.9^{+2.4}_{-2.5}$	> -15	-29^{+14}_{-19}	> -55
CC5	8128	< 5	> -25	-12^{+9}_{-12}	> -55
CC6	1352	< 20

NOTE—The uncertainties reported here are 90% confidence limits. We report on the 1 keV emission line, the Fe XXV (6.7 keV), Fe XXVI K α (7 keV), and Fe XXVI K β (8 keV) absorption lines for the different spectral regions defined in Figure 3. All upper limits shown are 2σ upper limits from fixing the line energy and width to the values from R1. There were no upper limits on the equivalent widths of the three absorption lines for spectral region CC6 as the background dominated past 5 keV.

medium absorption², or a pseudo-continuum of weak narrow lines (e.g., IGR J17062-6143; van den Eijnden et al. 2018), or due to faint Fe L or Ne X emission (Degenaar et al. 2013; Bult et al. 2021; Ng et al. 2022).

3.2.3. Dip Spectrum vs Non-Dip Spectrum

We investigated the spectral variation of 1A 1744–361 during the absorption dip and in the persistent emission immediately following the dip. The uncertainties re-

ported in this subsection are 90% confidence limits. We performed a simultaneous fit of the two spectra with an absorbed disk blackbody, thermal blackbody, a neutral partial covering component, and a partial covering partially ionized absorber. We found that the column density for the partially ionized absorber, N_H , was unconstrained, so we set it to $N_H = 10^{23} \text{ cm}^{-2}$. The corresponding spectral parameters for both blackbody components (`diskbb` and `bbodyrad`) were tied across the dip and persistent emission spectra. The best-fit spectral parameters are shown in Table 6. The unfolded spectra corresponding to the absorption dip (in red pluses) and persistent emission (in black pluses) are shown in Figure 12, along with the disk blackbody (dashed lines) and thermal blackbody (dotted lines) model components.

We found that the best-fit values for both blackbody components are consistent with spectra around the GTI containing the absorption dip (see Figure 10). For the neutral partial covering component, we found that the column density decreased from $8.7^{+0.8}_{-0.9} \times 10^{22} \text{ cm}^{-2}$ to $4.8^{+1.3}_{-1.4} \times 10^{22} \text{ cm}^{-2}$ as the neutral absorber presumably moved away from the line of sight, though not highly significant. However, the partial covering fraction significantly dropped from $0.714^{+0.027}_{-0.015}$ to $0.16^{+0.06}_{-0.03}$. For the partial covering partially ionized absorber, the column density was $N_H = 100^{+39}_{-18} \times 10^{22} \text{ cm}^{-2}$ during the dip, and the ionization parameter decreased from $\log(\xi) = 4.37^{+0.30}_{-0.11}$ to $\log(\xi) = 3.20^{+0.12}_{-0.15}$ as the neutral absorber was presumably in the line of sight. There were no significant changes in the covering fraction.

The presence of a partial covering partially ionized component is common among dipping NS LMXBs (Díaz Trigo et al. 2006; Raman et al. 2018; Gambino et al. 2019), but an additional neutral partial covering component is not usually required. To test the necessity of the additional component, we performed a similar fit to that presented in Table 6, but without the component (`tbpcf`). The column density of the partial covering partially ionized component remained fixed at $N_H = 10 \times 10^{22} \text{ cm}^{-2}$ in the persistent emission (unconstrained), but we found that $N_H = 14.0^{+1.3}_{-1.1} \times 10^{22} \text{ cm}^{-2}$ during the absorption dip. We also found that the ionization parameter decreased from $\log(\xi) = 4.39^{+0.56}_{-0.15}$ to $\log(\xi) = 1.55^{+0.18}_{-0.15}$ as the neutral absorber moved presumably in the line of sight. We did not see any statistically significant change in the covering fraction, with a lower limit of > 0.53 for the persistent emission and a value of $0.774^{+0.004}_{-0.004}$ during the absorption dip. Finally, for the blackbody components, we found that $T_{\text{in}} = 0.78^{+0.02}_{-0.02} \text{ keV}$, $\text{norm}_{\text{diskbb}} = 250^{+30}_{-20}$ (scaling defined in Table 4), $kT = 1.287^{+0.017}_{-0.021} \text{ keV}$, and $\text{norm}_{\text{bbodyrad}} = 83^{+7}_{-6}$. The fit statistic from this fit

² https://heasarc.gsfc.nasa.gov/docs/nicer/data_analysis/workshops/NICER-CalStatus-Markwardt-2021.pdf

is $\chi^2/\text{d.o.f.} = 281/236$; the fit presented in Table 6 is a statistically significant improvement with a rejection probability (see Equation 2) of 5.0×10^{-7} .

To investigate whether the varying absorption alone can account for the absorption dip, we performed an alternative fit to what was presented in Table 6 — we tied the absorption components across the dip and persistent spectra (N_H still constrained, so it was fixed at $10 \times 10^{22} \text{ cm}^{-2}$), but allowed the blackbody parameters (both `diskbb` and `bbodyrad`) to vary. For the absorption components, we found $n_{H,n} = 4.4_{-0.8}^{+1.4} \times 10^{22} \text{ cm}^{-2}$, a covering fraction for the neutral absorber of $0.21_{-0.06}^{+0.06}$, and an upper limit on the covering fraction for the ionized absorber of > 0.81 . We found that for the spectrum during the absorption dip (persistent emission), $T_{\text{in}} = 0.44_{-0.03}^{+0.04} \text{ keV}$ ($T_{\text{in}} = 0.78_{-0.03}^{+0.04} \text{ keV}$), $\text{norm}_{\text{diskbb}} = 430_{-110}^{+120}$ ($\text{norm}_{\text{diskbb}} = 320_{-50}^{+50}$; scaling defined in Table 4), $kT = 1.36_{-0.03}^{+0.03} \text{ keV}$ ($kT = 1.32_{-0.03}^{+0.03} \text{ keV}$), and $\text{norm}_{\text{bbodyrad}} = 51_{-4}^{+5}$ ($\text{norm}_{\text{bbodyrad}} = 72_{-6}^{+9}$). We see no evolution in the disk blackbody normalization and blackbody temperature (kT) within uncertainties, a small difference in the blackbody normalization, and a significant evolution in the inner disk temperature. We found the fit to describe the persistent and dip spectra well, with $\chi^2/\text{d.o.f.} = 252/233$. The rejection probability between the two models is $P_1/P_2 = 0.05$, so both fits are statistically consistent with the two spectra.

If we assume an upper limit of $D_{10} < 0.9$, and $i = 60 - 75^\circ$, we find that the upper limit of the apparent inner disk radius from the dip (persistent) spectrum is about $r_{\text{in}} < 22.8_{-1.9}^{+1.7} \text{ km}$ ($r_{\text{in}} < 26_{-4}^{+3} \text{ km}$). The realistic disk radius is related to the apparent inner disk radius, r_{in} , through $R_{\text{in}} = \xi \kappa^2 r_{\text{in}}$, where κ is the spectral hardening factor and $\xi = \sqrt{3/7}(6/7)^3$ (Kubota et al. 1998). In this case, for $L \approx 0.1 L_{\text{Edd}}$, we adopted $\kappa = 1.7$ (Shimura & Takahara 1995), thus for the dip (persistent) portion we obtain $R_{\text{in}} < 27.2_{-2.3}^{+2.0} \text{ km}$ ($R_{\text{in}} < 31_{-5}^{+4} \text{ km}$), or $R_{\text{in}} < 13.2_{-1.2}^{+0.9} R_G$ ($R_{\text{in}} < 15.0_{-2.4}^{+1.9} R_G$), where $R_G = GM/c^2$ is the gravitational radius with $M = 1.4 M_\odot$. These limits are consistent with those derived from reflection modeling of other NS LMXBs (Ludlam et al. 2022) as well as recent reflection modeling with NuSTAR observations Mondal et al. (2023). Additional detailed reflection modeling of 1A 1744–361 is outside of the scope of this work and will be reported in Pike et al. 2024, in prep. We will discuss the results in §4.4.

The 2008 outburst of 1A 1744–361 was about 10–20 times fainter than the latest outburst, so making fair comparisons is difficult, but Fe XXVI absorption lines were observed from the outburst with Chandra and Gavriil et al. (2012) found an upper limit of a redshift flow to be $v < 221 \text{ km s}^{-1}$. Therefore, a static disk atmo-

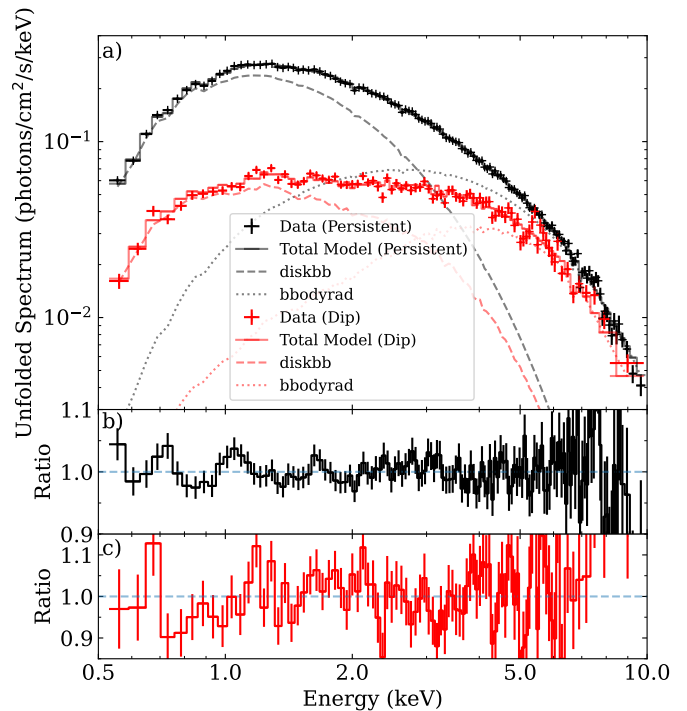


Figure 12. a) Unfolded spectra corresponding to the source emission during the absorption dip (red) and in the persistent emission immediately following the dip (black). The spectra were jointly fit with an absorbed disk blackbody (`diskbb`), a thermal blackbody (`bbodyrad`), a neutral partial covering component (`tbpcf`) and a partial covering partially ionized absorber component (`zxipcf`). In XSPEC parlance, the model is given by `tbabs(tbpcf(zxipcf(diskbb+bbodyrad)))`. The best-fit spectral parameters are given in Table 6; b) Ratio of the spectrum from the persistent emission to the model (specified above); c) Similar to b), but for the emission during the absorption dip. The spectra have been rebinned for visual purposes.

sphere cannot be ruled out. Unfortunately, the spectral resolution of the NICER observations is too low for us to ascertain any useful limits from the spectral fit of the absorption dip and the persistent emission post dip. Letting the redshift parameter freely vary yielded a best fit value of $z = 0.0003_{-0.0049}^{+0.0064}$, which is consistent with zero. High-resolution X-ray spectroscopic studies with the recently launched X-ray Imaging and Spectroscopy Mission (XRISM) of NS LMXBs will be able to ascertain the nature of the disk atmosphere (Tashiro et al. 2018; Trueba et al. 2020, 2022).

3.3. Correlated X-ray Behavior

We can combine the insights from the X-ray timing and spectroscopic analyses we have taken by looking at

Table 6. 0.6–10.0 keV spectral fits for the absorption dip interval and the persistent emission post-dip.

Model	Parameter	Dip	Persistent
tbpcf	$n_{H,n}$ (10^{22} cm $^{-2}$)	$8.7^{+0.8}_{-0.9}$	$4.8^{+1.3}_{-1.4}$
	pcf	$0.714^{+0.027}_{-0.015}$	$0.16^{+0.06}_{-0.03}$
	Redshift	0	0
zxipcf	N_H (10^{22} cm $^{-2}$)	100^{+39}_{-18}	10 (fixed)
	$\log(\xi)$	$3.20^{+0.12}_{-0.15}$	$4.37^{+0.30}_{-0.11}$
	f_c	$0.62^{+0.13}_{-0.13}$	> 0.74
	z	0	0
diskbb	T_{in} (keV)	Tied	$0.81^{+0.03}_{-0.04}$
	norm _{diskbb}	Tied	270^{+60}_{-30}
bbodyrad	kT (keV)	Tied	$1.33^{+0.03}_{-0.03}$
	norm _{bbodyrad}	Tied	69^{+10}_{-8}
...	$F_{0.6-10.0,unabs}$ (10^{-9} cgs)	$1.795^{+0.019}_{-0.020}$	$3.870^{+0.013}_{-0.013}$
$\chi^2/\text{d.o.f.}$...	260/232	

NOTE—The uncertainties reported here are 90% confidence limits. For interstellar absorption, we used the **tbabs** model (Wilms et al. 2000) and fixed $n_H = 0.44 \times 10^{22}$ cm $^{-2}$ (see §3.2.1 for more information). The variable $n_{H,n}$ is the column density of the neutral absorber, “pcf” is the partial covering fraction, N_H is the column density of the (partial covering) partially ionized absorber, $\log(\xi)$ is the ionization parameter, f_c is the covering fraction, and z is the redshift. The unabsorbed 0.6–10.0 keV flux was determined using **cflux**.

correlations between the different timing and spectroscopic variables. To do so, we calculated the Spearman correlation coefficient (SCC) and the corresponding p-values. There were some correlations that can already be seen from the overall spectroscopic evolution of the source (see Figure 10) — the blackbody temperature (kT) and normalization (from **bbodyrad** in **XSPEC**) had an SCC of -0.95 , which translated to a p-value of $\sim 10^{-35}$, which is unsurprising as they are intrinsically correlated for a given flux. For the correlation between the 0.6–10.0 keV absorbed flux and blackbody components, the SCC and p-values were -0.53 and 1.1×10^{-6} (T_{in}), 0.61 and 6.9×10^{-9} (disk blackbody normalization), -0.39 and 6.0×10^{-4} (blackbody temperature), and 0.61 and 1.1×10^{-8} (blackbody normalization).

The potentially more interesting correlations come from comparing quantities derived from spectroscopy and from timing for each data segment. For example, we found a strong correlation between the blackbody temperature (kT) and the fractional rms amplitude of the power spectrum over 0.03–50 Hz with an SCC of -0.64 and p-value of 4.0×10^{-9} . The corresponding SCC and p-values for the correlation between the fractional rms

amplitude and the blackbody normalization were 0.55 and 1.8×10^{-6} , respectively.

3.4. MeerKAT

We have shown the MeerKAT 1.3 GHz flux densities in the bottom panel of Figure 1 and Table A1 in the Appendix. The measured flux densities in each of the 15-minute observations are shown in blue, where we plot both $> 4\sigma$ detections (squares) and 3σ upper limits (triangles). During the first three observations, the source was radio-bright with a flux density $\gtrsim 500 \mu\text{Jy}$. As a result, we were able to measure the (intra-band) spectral index in the bright epochs by breaking the bandwidth in half (i.e., 856–1284 and 1284–1712 MHz) and measuring the flux density in each sub-band. We measured radio spectral indices ($S_\nu \propto \nu^\alpha$) of $\alpha = -0.2 \pm 0.1$, -0.3 ± 0.1 , and -0.6 ± 0.2 during the observations on MJD 59731, 59733, and 59742, respectively. Following the bright detections, the source flux density rapidly decayed over ~ 7 days, dropping below our single observation detection threshold of $\sim 80 \mu\text{Jy}$ for all but one of our remaining observations. We performed image stacking to increase the S/N, grouping the last eleven epochs into three stacked images consisting of three, four, and four independent observations. The stacked flux densities are shown in black, where we detected the source at $\gtrsim 4.5\sigma$ in the first two stacked images (diamonds). However, even with the improved S/N we still did not detect the source in the final stacked image, and thus represent it as a 3σ upper limit. We did not observe any correlated evolution (in time) when comparing the X-ray and radio light curves, as would be expected from hard-state jet emission.

To constrain the position of 1A 1744–361 on the L_R – L_X plane, we calculated the radio luminosity at 5 GHz ($L_{R,5}$) through

$$L_{R,5} = 4\pi S_\nu \nu d^2 \approx 3.9 \times 10^{26} \text{ erg s}^{-1} \left(\frac{S_\nu}{1 \mu\text{Jy}} \right) \left(\frac{d}{8 \text{ kpc}} \right)^2, \quad (3)$$

where $1 \text{ Jy} = 10^{-23} \text{ erg s}^{-1} \text{ cm}^{-2} \text{ Hz}^{-1}$, $\nu = 5 \text{ GHz}$, and adopting the standard assumption of a flat radio spectral index ($\alpha = 0$) in $S_\nu \propto \nu^\alpha$. We find an upper limit on $L_{R,5}$ for the last stacked data point to be $L_{R,5} < 1.2 \times 10^{28} \text{ erg s}^{-1}$, assuming the source is at a distance of 8 kpc.

4. DISCUSSION

We have reported on the X-ray timing and spectroscopic analysis with NICER as well as radio monitoring observations with MeerKAT, throughout the nearly three-month long outburst of the dipping NS LMXB

1A 1744–361. We also presented supplementary NuSTAR observations (timing only), and monitoring data with MAXI and Swift. With NICER we have discovered, for the first time, several instances of a $\nu_0 \approx 8$ Hz QPO in the high intensity, banana spectral state of the source. The QPO had fractional rms amplitudes of $\sim 5\%$. We identify this QPO as an NBO-like QPO, and argue that through 1A 1744–361, we are probing the lower accretion rate boundary at which NBOs (can) occur. Additionally, we argue that since: a) the source exhibited two separate tracks in a very short timescale, only seen in the Z-state of other sources, in a NICER observation (Figure 6); b) we observed a flaring branch in the 3–25 keV NuSTAR HID (Figure 9); c) we saw NBOs at the peak of the outburst, the source is exhibiting Z-state-behavior at the outburst peak. Finally, MeerKAT observations confirmed that at the peak of the outburst, 1A 1744–361 was at its radio-brightest. We discuss below that the results are consistent with the picture of a transient ejection of material accompanying the island-to-banana spectral state transition of the NS LMXB in the atoll-state. The observations provided the strongest observational evidence for radio flaring (and jet ejection) during spectral state transitions in the atoll-state. The spectroscopic results are possibly consistent with the presence of a disk atmosphere and a contracting boundary layer of the NS during the banana state, where the source spent most of the outburst in.

4.1. X-ray and Radio Evolution

The X-ray spectral and timing properties of 1A 1744–361 showed clear signatures of a spectral state transition. When we began our NICER monitoring (MJD 59734), the X-ray spectra were best-fit with thermal components (i.e., a blackbody and multi-color disk blackbody), characteristic of a banana state (Lin et al. 2007). At later times, for GTI segments after MJD 59800, the X-ray spectra favored power law models, which are often used to describe X-ray spectra dominated by high energy Comptonized photons, and thus an island state. The change in the best-fit X-ray spectral model coincided with the evolution of the timing parameters, namely the transition from the R1 to the R2 spectral states (highlighted in Figure 3). Moreover, an increase in the average (continuum) fractional rms (from 3% to 10%) is typical of atoll-state sources transitioning from banana to island spectral states (Hasinger & van der Klis 1989; van der Klis 2004). The changes in the X-ray spectral continua, position on the CCD, and average fractional rms strongly suggest that a physical change in the accretion flow occurred around MJD 59800.

The radio evolution consisted of a bright initial detection on our first day of observing (MJD 59731), followed by a rapid, monotonic decrease. Initially, the radio detections were thought to be associated with a steady compact jet (Hughes et al. 2022). However, the best-fit model for the X-ray spectra (i.e., the source was in a banana state) and the temporal evolution of the radio is more consistent with emission originating from a transient jet ejection (otherwise known as an optically-thick flare; Fender & Bright 2019). Moreover, the steepening of the radio spectral index ($\alpha = 0.2 \pm 0.1$ to -0.6 ± 0.2 ; optically thick-to-thin synchrotron emission) is also suggestive of a transient jet ejection. There is an alternative explanation for optically thin radio emission in NS LMXBs; Russell et al. (2021) observed optically thin radio emission in the “high X-ray mode” (equivalent to the banana state) for the ultracompact NS LMXB 4U 1820–30. In this system, the emission was attributed to a heavily quenched compact jet with a break frequency that had moved into the radio regime. However, it is difficult to apply a similar explanation to our observations of 1A 1744–361. The optically thin radio emission in 4U 1820–30 is persistent and modestly variable (a factor of ~ 2), whereas 1A 1744–361 showed a rapid decay corresponding to, at a minimum, a factor of ~ 100 decrease in the radio flux. These decay properties, combined with the temporal coincidence between the radio flare and the X-ray state transition, lead us to favor the jet ejection scenario.

We looked for evidence of ballistic motion as ejecta have been spatially resolved in several BH (e.g., Hjellming & Rupen 1995; Bright et al. 2020) and NS LMXBs (e.g., Fomalont et al. 2001; Miller-Jones et al. 2012; Motta & Fender 2019), although for NSs the ejections have only been resolved for Z-state sources. We saw no evidence of movement of the radio emission in 1A 1744–361. However, this was expected as the MeerKAT beam size ($\sim 6''$) and decay timescale (undetectable after ~ 40 days) would make proper motion undetectable for distances > 1 kpc. We use the maximum flux density (1.2 mJy) and observing frequency (1.3 GHz) to derive a distance-scaling relation for the minimum energy (E_{\min}) of the jet ejection (following the method described in Fender & Bright 2019),

$$E_{\min} \approx 2.2 \times 10^{37} \text{ erg} \left(\frac{d}{8 \text{ kpc}} \right)^{40/17} \quad (4)$$

A minimum energy of $\sim 10^{37}$ erg is consistent with past observations of flaring NS LMXBs (e.g., Swift J1858.6-0814; Rhodes et al. 2022). Measuring an accurate distance (e.g., through the photospheric radial expansion of a type I X-ray burst; Kuulkers et al. 2003)

will remove the distance ambiguities and confirm (or refute) the apparent consistency between the minimum energy calculations.

Motta & Fender (2019) investigated simultaneous X-ray and radio observations of the NS LMXB Z-state source Sco X-1 and found that ultrarelativistic outflows were connected to the simultaneous manifestation of an NBO/HBO pair. For 1A 1744-361, we did not observe any HBOs, and the NBO-like QPOs were not continuously observed — the three GTIs for which NBO-like QPOs were observed had multiple GTIs in between where the QPOs were not observed. Additionally, simultaneous X-ray and radio observations of another Z-state source, GX 17+2, associated the occurrences of NBO with jet formation (Penninx et al. 1988; Migliari et al. 2007). Unlike the simultaneous X-ray and radio observations that Motta & Fender (2019) analyzed, we only observed the QPO instances days after the ejecta would have been launched. Our X-ray and radio monitoring, although not simultaneous, is consistent with the picture of NBO-like QPOs and transient ejections occurring around the outburst peak. In order to confirm (or refute) this suggestive picture, more dense monitoring of future outbursts of similar sources in the radio and X-rays at the onset of the outburst are required, which could be achieved by coordinated all-sky X-ray monitoring and rapid follow-up observations³.

4.2. Observational Evidence for Radio Flaring During Spectral State Transitions in Atoll-State Sources

Using the behavior of BH LMXBs as a reference, ejection events in NS LMXBs have been predicted to occur during a transition from a low to high mass accretion rate (i.e., ‘island-to-banana’; Migliari & Fender 2006; Muñoz-Darias et al. 2014). Given that we did not observe the rise of the radio flare and that our early-time NICER X-ray observations were already consistent with banana state emission, it is likely that we missed the initial island-to-banana state transition. Looking at the MAXI/GSC light curves (see Figure 1), the onset of the outburst began on MJD 59728, six days before our first NICER observations. During this time, the MAXI hardness ratio showed a drop from ~ 0.6 on MJD 59729 (the highest hardness ratio of the entire outburst) to ~ 0.3 on MJD 59731 (the same day as the initial radio detection). It is plausible that the observed hardness ratio evolution was a signature of a state transition that coincided with the launching of a jet ejection, and the source rapidly transitioned ($\lesssim 3$ days) during the rise of the X-ray flux.

This is corroborated by the Swift/XRT observation on 2022 June 1 (see §3), where the 0.5–10.0 keV spectrum was well described with an absorbed cutoff power law with a photon index of $\Gamma_{\text{co}} \approx 0.6$ and e-folding energy $E_f \approx 3.2$ keV, suggesting that 1A 1744-361 was still in the island state. The behavior of 1A 1744-361 constitutes some of the strongest observational evidence for radio flaring (and jet ejecta) during state transitions in atoll-state LMXBs.

We did not detect any radio re-brightening after the source transitioned back to the island state (towards the end of the outburst). Following the transition, assuming 1A 1744-361 did, in fact, re-form the compact jet, we used the final stacked image upper limit ($\sim 30 \mu\text{Jy}$) to place constraints on the position of the source on the L_R - L_X diagram as this observation coincided with the hard state X-ray emission. We present these results for three assumed distances (4, 8, and 12 kpc) in Figure 13. The luminosities at all three distances put 1A 1744-361 among the lower-end of radio luminosities for NS LMXBs at our measured X-ray luminosity. Higher sensitivity radio monitoring of future outbursts will allow us to directly detect (or better constrain) the radio flux density of the compact jet, thereby determining whether 1A 1744-361 is anomalously radio-quiet for an NS LMXB.

4.3. QPO Phenomenology

In the most recent high-intensity outburst in 2022, we discovered instances of a $\nu \approx 8$ Hz QPO over 1.0–3.0 keV (see Table 2) around the peak of the outburst. This is similar to the NBO-like QPOs observed in Aql X-1 (Reig et al. 2004), XTE J1806-246 (Wijnands & van der Klis 1999), and 4U 1820-30 (Wijnands et al. 1999). While the QPOs (including from 1A 1744-361) manifested near the peak luminosities of the outburst, the luminosities were below $0.5 L_{\text{Edd}}$ (except XTE J1806-246), unlike the NBOs in Z-state sources ($0.5 - 1.0 L_{\text{Edd}}$; van Paradijs et al. 1988; Wijnands et al. 1996; Piraino et al. 2002). For 4U 1820-30, the peak luminosity was approximately $0.2 - 0.4 L_{\text{Edd}}$, and for Aql X-1 it was approximately $0.05 L_{\text{Edd}}$. In the other high-intensity outbursts of 1A 1744-361 in 2005 and 2008 (see Figure 2) with RXTE, no QPOs were observed (except for a low significance 3σ indication at 29 Hz in 2008) though the bolometric luminosity was on the order of $0.01 L_{\text{Edd}}$. Additionally, the high-intensity outburst in 2003 did not show any QPOs in the data.

The QPOs from this outburst were found in the banana state in the CCD. The QPOs found in XTE J1806-246 and 4U 1820-30 exhibited a clear energy dependence below 12 keV, whereas we found no

³ https://heasarc.gsfc.nasa.gov/docs/nicer/science_nuggets/20201203.html

such relation in the NICER data. However, this could just be due to the differences in sensitivities in the different energy bands. Detections of NBOs/NBO-like QPOs below 3 keV are rare. At the higher accretion rates exhibited in Z-state sources, NBOs have been observed in the soft energy range with Z-state sources: over 2–3.1 keV from Cyg X–2 (Wijnands & van der Klis 2001) with RXTE and down to ~ 0.4 keV from Cyg X–2 and Sco X–1 with NICER (Jia et al. 2023). We have seen in various atoll-state and Z-state sources that the appearance of the different QPOs is a function of the spectral state of the source (Motta et al. 2017). The fact that we have observed NBO-like QPOs from 1A 1744–361 implies that at the peak of the outburst, the source was a Z-state source, and the QPOs were undetectable by the time the source was below a luminosity threshold. Thus we suggest that the ‘NBO-like QPOs’ we have observed from other atoll-state sources (Wijnands et al. 1999; Wijnands & van der Klis 1999) and NBOs observed from Z-state sources are simply the same phenomena (NBOs) observed at different mass accretion rates. With 1A 1744–361, we are probing the lower accretion rate boundary at which NBOs are observed. Observing atoll-state sources at higher luminosities will confirm this picture.

We also note from the low-intensity 2004 outburst from 1A 1744–361 that a ~ 3.5 Hz QPO was discovered (Bhattacharyya et al. 2006a). Seven sources in total have also exhibited ~ 1 Hz QPOs which are empirically observed in the island spectral states and is thought to originate from a geometrically modulated precessing inner accretion flow (Homan et al. 1999; Jonker et al. 2000; Homan 2012; Homan et al. 2015).

4.4. Absorption Dip

The source light curve exhibited an absorption dip during the outburst, as shown in Figure 4, a known feature of the source (Bhattacharyya et al. 2006a; Gavriil et al. 2012). We note that during the roughly 90 s duration of the absorption dip, the source flux temporarily increased to near-persistent levels (approximately 16–20 s; 18–22% of the dip duration), and the color correspondingly fell.

We also investigated the spectra during the absorption dip and the persistent emission immediately following the dip (Table 6). The spectral results, as well as the lone absorption dip observed, are consistent with the idea of an intervening orbital partial covering absorber along the line of sight appearing for about 100 s (Jimenez-Garate et al. 2002; Díaz Trigo et al. 2009). In particular, we required a partial covering partially ionized absorber and a neutral partial cover-

ing absorber (see §3.2.3) whose parameters varied across the dip and persistent emission spectra. We also tested for the possibility that the different observed spectra were due to changes in the blackbody parameters. We found a significant increase in the inner disk temperature from $T_{\text{in}} = 0.44_{-0.03}^{+0.04}$ keV to $T_{\text{in}} = 0.78_{-0.03}^{+0.04}$ keV and a small increase in the blackbody normalization ($\text{norm}_{\text{bodyrad}} = 51_{-4}^{+5}$ to 72_{-6}^{+9} ; scaling in Table 4) as the intervening absorber left the line of sight. From the light curve shown in Figure 4, these spectroscopic results imply that the inner disk temperature increased by almost a factor of two in a matter of seconds, which seems implausible. It is more likely that the spectral changes during the dip can be explained by the presence of a neutral partial covering absorber. In fact, the partial covering fraction of $0.16_{-0.03}^{+0.06}$ is similar to the fraction of time that the source flux temporarily increased during the absorption dip, which is consistent with the idea of the covering fraction being the fraction of time that the source was fully obscured by the neutral absorber.

The best-fit value for the ionization parameter, $\log(\xi) \approx 4.4$, is consistent with the lower limit of $\log(\xi) > 3.6$, derived from the absorption dip in the July 2008 outburst (Gavriil et al. 2012). Absorption dips observed in other dipping LMXBs also revealed ionization parameters in the persistent emission $\log(\xi) > 4.0$ (Iaria et al. 2007a,b; Marino et al. 2022). However, we do not rule out ionized absorption during the dip.

4.5. Accretion Disk Atmosphere

Looking at the combined spectra from spectral regions R1 and R2 (see Table 4), the blackbody normalization from R1 of 26.5 (scaled by $(R_{\text{km}}/D_{10})^2$) implied an emitting region of size $R \approx 4.6$ km. In fact, for R1, the normalization took on values over 10 – 100 ($R \approx 2.8 - 9$ km), though note these are all upper limits as we adopt $d < 9$ kpc. The decreasing blackbody normalizations could be a result of a contracting boundary layer.

The presence of significant absorption lines from Fe XXV $K\alpha$, Fe XXVI $K\alpha$, and Fe XXVI $K\beta$ in R1 suggested the presence of a highly ionized disk atmosphere (plasma with an inferred equatorial geometry), which have been observed in several NS LMXBs (Parmar et al. 2002; Iaria et al. 2007a,b; Hyodo et al. 2009; Marino et al. 2022). They are also a common feature in NS LMXBs exhibiting absorption dips (Díaz Trigo et al. 2006). On the other hand, in R2, we saw significantly weaker Fe XXVI $K\alpha$ and a non-detection of the Fe XXVI $K\beta$ line (albeit with a high upper limit), which possibly implies a cooler disk atmosphere. This could also be due to the photoionization of the disk atmosphere from

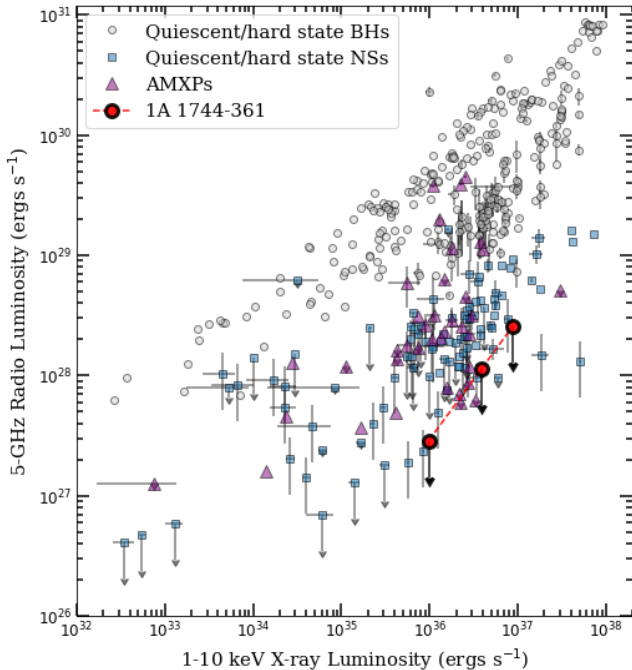


Figure 13. 1–10 keV X-ray and 5 GHz radio luminosity plane (i.e., L_R – L_X plane). The plot includes archival hard state BH LMXBs (gray circles), hard state NS LMXBs (blue squares), and accreting millisecond X-ray pulsars (AMXPs; purple triangles) as seen in Figure 4 of van den Eijnden et al. 2022, which contains a large number of sources from Bahramian & Rushton (2022). As 1A 1744–361 was never detected in the hard state, its position on the L_R – L_X plane is limited to upper limits (red circles). Furthermore, due to the unknown distance of 1A 1744–361, the plot includes three data points at distances of 4, 8, and 12 kpc. The upper limits reside within the lower end of radio luminosities for known hard-state NS LMXBs. Future, higher sensitivity observations (and an accurate distance measurement) are critical for determining the source’s exact position on the L_R – L_X plane, and whether it has an anomalously radio-quiet compact jet.

the corona (Ko & Kallman 1991; Neilsen & Lee 2009). We tried to investigate individual segments within the spectral regions (R1 and R2), but within uncertainties and due to the high upper limits, we could not ascertain any differences between the spectral subregions. In addition, the spectrum during the persistent emission after the absorption dip (Table 6) exhibited a high ionization parameter ($\log(\xi) \approx 4.4$) and a large covering fraction (> 0.7), suggesting the presence of an ionized medium possibly in the form of an accretion disk atmosphere. Thus, with the expected $60 - 75^\circ$ orbital inclination of the system (observed dip, but no eclipses), we are likely probing into the highly ionized accretion disk atmosphere, with the absorption lines originating in the disk atmosphere.

We also note that this outburst was the first time that the Fe XXV $K\alpha$ and Fe XXVI $K\beta$ lines were detected from this source (Tobrej et al. 2023, Mondal et al. 2023, Pike et al. 2024, in prep.). Higher spectral resolution observations, such as that of XRISM, will be crucial in revealing the dynamic (or static) nature of the accretion disk atmosphere or wind from NS LMXBs (Trueba et al. 2020, 2022; Gandhi et al. 2022)

5. CONCLUSIONS

In this work, we have reported on the X-ray timing and spectroscopic analysis with NICER (timing only for NuSTAR; and monitoring with MAXI and Swift) of the dipping atoll-state NS LMXB 1A 1744–361, as well as radio monitoring observations with MeerKAT, throughout its nearly three-month long outburst in 2022. The overall outburst proceeded as follows:

1. On \sim MJD 59728, 1A 1744–361 entered into an outbursting state, its X-ray flux rapidly increased, and it likely transitioned from the island to the banana spectral states within ~ 3 days of the onset of the outburst, launching transient jet ejecta.
2. The source remained at a high accretion rate (i.e., in the banana state) from \sim MJD 59728 to 59800, where any radio detection was the result of the fading jet ejecta.
3. At around \sim MJD 59800, the source transitioned back to the island state. We detected no radio emission, and, therefore, if a compact jet formed, its radio emission was below our luminosity detection threshold of $L_{R,5} < 1.2 \times 10^{28} (d/8 \text{ kpc})^2 \text{ erg s}^{-1}$
4. For the remainder of our monitoring, the X-ray luminosity decreased, dipping below NICER’s sensitivity limit on \sim MJD 59820, as the source returned to quiescence.

The results presented in this work support the notion that atoll sources and Z sources are not distinct source classes, but instead are two states on a continuum defined by the mass accretion rate, as previously suggested in detailed spectral and timing studies of XTE J1701–462 (Lin et al. 2009; Homan et al. 2010) and other NS LMXBs (Fridriksson et al. 2015). Thus, it is more appropriate to refer to atoll-state sources and Z-state sources. We have discovered several instances of a $\nu_0 \approx 8$ Hz QPO in the high intensity, banana state outburst of the source with fractional rms amplitudes of $\sim 5\%$ for the first time. The HID of a NICER observation around the outburst peak (Figure 6) suggests that the source exhibited two distinct tracks on a very short time scale ($\lesssim 2$ hr, which is typically seen in Z-

state sources). The HID for the NuSTAR observation, taken around the peak of the outburst (Figure 9), shows a flaring branch. While the HID/CCD corresponding to the outburst with the NICER data suggested an atoll-state source, the aperiodic timing features (along with the HIDs) observed by NICER and NuSTAR of 1A 1744–361 resemble that of other atoll-state sources that exhibit Z-state-like behavior around the peak of the outbursts, such as XTE J1806–246 (Wijnands et al. 1999), 4U 1820–30 (Wijnands & van der Klis 1999), and Aql X–1 (Reig et al. 2004). In fact, with the growing understanding that the spectral evolution of atoll-state sources and Z-state sources differ only in the mass accretion rate (Lin et al. 2009; Homan et al. 2010; Fridriksson et al. 2015), we posit that the ‘NBO-like QPOs’ observed in several atoll-state sources (Wijnands et al. 1999; Wijnands & van der Klis 1999; Reig et al. 2004) are in fact NBOs manifesting at lower accretion rates. The NBO production mechanism is likely activated at a luminosity threshold below that of the Eddington luminosity, as suggested from observations of 4U 1820–30 and XTE 1806–246 (Wijnands et al. 1999; Wijnands & van der Klis 1999). To probe this further, dense multiwavelength monitoring and rapid follow-up of future outbursts of similar sources across different accretion states are required.

Around the peak of the outburst of 1A 1744–361, we detected NBOs and inferred the transient ejection of material accompanying the island-to-banana state transition with quasi-simultaneous X-ray and radio monitoring. This observed scenario is possibly consistent with the simultaneous X-ray and radio observations of Sco X–1 which found that ultrarelativistic outflows were connected to the appearance of an NBO/HBO pair (Motta & Fender 2019). However, the robust confirmation of such a scenario will require dense monitoring of future outbursts of such sources in the X-rays and radio wavelengths with sensitive instruments.

We also observed an absorption dip in the light curve of one observation throughout the outburst. The X-ray spectroscopy results are consistent with the idea of a non-orbiting neutral partial covering absorber with a non-trivial interior structure along the line of sight. The large covering fraction ($> 70\%$) of the ionized absorber during the persistent emission suggests the presence of a highly ionized medium surrounding the NS. We have also investigated the evolution of the 6.7 keV, 6.96 keV, and 8.0 keV Fe XXV, Fe XXVI $K\alpha$ and $K\beta$ absorption lines, respectively, where the presence of these absorption lines also suggest the presence of a highly ionized medium, with origins likely in the disk atmosphere of the source.

Facilities: NICER, MAXI, NuSTAR, Swift, MeerKAT

Software: Astropy (Astropy Collaboration et al. 2013, 2018), NumPy and SciPy (Virtanen et al. 2020), Matplotlib (Hunter 2007), IPython (Perez & Granger 2007), tqdm (da Costa-Luis et al. 2022), HEASoft 6.31.1⁴ (NASA High Energy Astrophysics Science Archive Research Center (HEASARC) 2014)

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AKH and GRS respectfully acknowledge that they perform the majority of their research from Treaty 6 territory, a traditional gathering place for diverse Indigenous peoples, including the Cree, Blackfoot, Métis, Nakota Sioux, Iroquois, Dene, Ojib-

⁴ <http://heasarc.gsfc.nasa.gov/ftools>

way/Saulteaux/Anishinaabe, Inuit, and many others

whose histories, languages, and cultures continue to influence our vibrant community.

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APPENDIX

A. TABLES

Table A1. Measured MeerKAT flux densities (μJy) as a function of time for the stacked data points and per epoch (see description in §3.4). Upper limits are quoted to 3σ .

Type	Date (MJD)	Flux (μJy)
Stacked	59730.9693	1248 ± 23
	59733.9204	1066 ± 23
	59742.9526	563 ± 20
	59757.2662	72 ± 11
	59779.8157	42 ± 9
	59808.2435	< 30
Epoch	59730.9693	1248 ± 23
	59733.9204	1066 ± 23
	59742.9526	563 ± 20
	59747.9011	< 53
	59759.9651	< 59
	59763.9325	79 ± 17
	59769.8848	< 56
	59775.8213	< 54
	59782.8344	< 52
	59790.7224	< 56
	59797.9283	< 53
	59804.6530	< 71
	59811.7461	< 55
	59818.6467	< 54

Table A2. Properties of the NICER observations throughout the outburst of 1A 1744–361. We presented the GTI numbers, corresponding ObsID, start time of the interval in MJD (TT units), the amount of filtered exposure (in s), the averaged total (source plus background) count rate (c/s), the background count rate (c/s) as modeled with the 3C50 model (Remillard et al. 2022), the number of averaged segments in the power spectra (m), the fractional RMS amplitude (in %) over 0.03–50 Hz, the power law amplitude, the power law index, and the reduced χ^2 between the model (power law plus constant) and the observed power spectrum. Entries with ellipses denote unconstrained values.

GTI	ObsID	MJD (TT)	Exposure	r_{tot}	r_{bg}	m	f_{rms}	A	PL Index	χ^2
			s	c/s	c/s		% (0.03–50 Hz)			
0	5202800101	59733.826	54.0	780.3 ± 3.8	1.5 ± 0.17	1	5.5 ± 2.6	-0.66 ± 0.42	0.15 ± 0.22	0.53
1	5202800102	59733.993	862.0	811.5 ± 1.0	2.47 ± 0.01	26	9.1 ± 2.0	0.72 ± 0.05	1.11 ± 0.04	0.91
2	5202800102	59734.267	277.0	828.3 ± 1.8	2.47 ± 0.02	8	8.2 ± 2.2	0.5 ± 0.1	1.3 ± 0.1	0.7

Table A2 continued

Table A2 (*continued*)

GTI	ObsID	MJD (TT)	Exposure	r_{tot}	r_{bg}	m	f_{rms}	A	PL Index	χ^2
			s	c/s	c/s		% (0.03–50 Hz)			
3	5202800102	59734.332	547.0	805.0 ± 1.2	2.47 ± 0.01	17	9.1 ± 2.1	0.72 ± 0.07	1.17 ± 0.05	0.99
4	5202800102	59734.445	1053.0	934.4 ± 0.9	2.48 ± 0.01	32	7.5 ± 1.9	0.62 ± 0.05	1.24 ± 0.04	0.72
5	5202800103	59735.094	83.0	915.5 ± 3.3	2.3 ± 0.01	2	6.8 ± 2.3	-0.23 ± 0.13	0.52 ± 0.36	0.7
6	5202800103	59735.696	306.0	809.9 ± 1.6	0.774 ± 0.004	9	5.6 ± 2.0	0.29 ± 0.08	1.28 ± 0.13	0.8
7	5202800103	59735.759	406.0	656.8 ± 1.3	0.782 ± 0.003	12	5.7 ± 2.1	0.17 ± 0.06	1.38 ± 0.14	0.9
8	5202800104	59736.15	121.0	866.3 ± 2.7	1.04 ± 0.03	3	7.8 ± 2.1	0.0 ± 0.01	2.86 ± 1.61	0.79
9	5202800104	59736.922	38.0	783.3 ± 4.5	1.07 ± 0.01	1	6.5 ± 2.6	0.01 ± 0.08	1.94 ± 1.98	0.57
10	5202800104	59736.924	140.0	773.7 ± 2.4	1.08 ± 0.03	4	6.7 ± 2.2	0.11 ± 0.09	1.6 ± 0.34	0.82
11	5202800105	59737.05	165.0	790.5 ± 2.2	1.23 ± 0.01	5	6.7 ± 2.1	0.24 ± 0.11	1.28 ± 0.2	0.81
12	5202800106	59738.194	779.0	795.3 ± 1.0	0.98 ± 0.01	24	12.9 ± 4.5	0.44 ± 0.06	1.62 ± 0.08	1.86
13	5202800106	59738.215	41.0	993.2 ± 4.9	1.0 ± 0.03	1	3.4 ± 2.3	-0.82 ± 0.13	0.38 ± 0.09	0.48
14	5202800106	59738.713	176.0	838.2 ± 2.2	2.57 ± 0.08	5	3.5 ± 1.9	0.08 ± 0.09	1.15 ± 0.37	0.93
15	5202800107	59739.027	131.0	640.7 ± 2.2	0.88 ± 0.01	4	4.1 ± 2.2	-0.11 ± 0.08	0.33 ± 0.43	0.71
16	5202800107	59739.041	1138.0	666.6 ± 0.8	0.877 ± 0.004	35	4.1 ± 2.0	0.12 ± 0.03	1.19 ± 0.11	0.91
17	5202800107	59739.092	277.0	710.2 ± 1.7	0.74 ± 0.01	8	4.2 ± 2.0	0.01 ± 0.03	1.73 ± 0.83	0.81
18	5202800107	59739.108	995.0	572.1 ± 0.8	0.741 ± 0.004	31	5.4 ± 2.1	0.15 ± 0.04	1.08 ± 0.1	0.96
19	5202800107	59739.156	1030.0	580.8 ± 0.8	0.741 ± 0.004	32	5.2 ± 2.1	0.15 ± 0.04	1.06 ± 0.1	0.78
20	5202800107	59739.182	158.0	678.0 ± 2.1	0.74 ± 0.01	4	4.2 ± 2.2	-0.15 ± 0.08	0.52 ± 0.28	0.78
21	5202800107	59739.221	1470.0	901.5 ± 0.8	0.879 ± 0.003	45	7.3 ± 1.8	0.46 ± 0.04	1.07 ± 0.04	1.48
22	5202800108	59741.693	579.0	656.4 ± 1.1	1.06 ± 0.01	18	4.3 ± 2.0	0.1 ± 0.04	1.22 ± 0.16	0.91
23	5202800109	59742.533	517.0	636.1 ± 1.1	1.06 ± 0.01	16	4.9 ± 2.1	0.05 ± 0.03	1.52 ± 0.24	1.1
24	5202800110	59743.756	775.0	680.9 ± 0.9	1.25 ± 0.01	24	2.9 ± 2.0	0.09 ± 0.04	1.17 ± 0.16	0.73
25	5202800110	59743.821	773.0	757.4 ± 1.0	1.18 ± 0.01	24	5.0 ± 1.9	0.21 ± 0.04	1.04 ± 0.09	0.92
26	5406620101	59745.544	599.0	664.6 ± 1.1	1.95 ± 0.01	18	2.9 ± 2.0	0.02 ± 0.02	1.51 ± 0.43	0.94
27	5202800111	59745.819	900.0	732.3 ± 0.9	1.171 ± 0.004	28	3.7 ± 1.9	0.19 ± 0.04	0.96 ± 0.08	0.95
28	5202800112	59747.048	676.0	632.7 ± 1.0	1.322 ± 0.005	21	3.8 ± 2.0	0.02 ± 0.03	1.37 ± 0.38	0.93
29	5202800112	59747.492	1267.0	756.3 ± 0.8	1.322 ± 0.004	39	4.0 ± 1.9	0.16 ± 0.03	1.12 ± 0.08	1.05
30	5202800113	59748.082	492.0	743.6 ± 1.2	1.34 ± 0.01	15	5.3 ± 1.9	0.08 ± 0.05	1.0 ± 0.21	0.88
31	5202800113	59748.383	497.0	626.9 ± 1.1	1.08 ± 0.01	15	3.3 ± 2.1	-0.0 ± 0.04	0.23 ± 5.71	0.93
32	5202800113	59748.392	565.0	619.7 ± 1.0	1.08 ± 0.01	17	3.6 ± 2.1	0.0 ± 0.01	1.91 ± 0.91	0.84
33	5202800113	59748.785	1085.0	687.3 ± 0.8	1.08 ± 0.01	33	3.2 ± 1.9	0.08 ± 0.03	1.21 ± 0.15	0.98
34	5202800114	59749.415	377.0	769.3 ± 1.4	1.075 ± 0.004	11	2.0 ± 1.9	0.07 ± 0.05	1.48 ± 0.29	0.79
35	5202800115	59752.385	67.0	622.8 ± 3.0	1.2 ± 0.01	1	...	-0.72 ± 0.14	0.33 ± 0.1	0.52
36	5202800115	59752.393	114.0	613.4 ± 2.3	1.2 ± 0.01	3	...	-0.36 ± 0.07	0.46 ± 0.1	0.7
39	5202800116	59753.353	46.0	765.8 ± 4.1	1.19 ± 0.01	1	...	-0.76 ± 0.14	0.4 ± 0.12	0.48
40	5202800116	59753.873	346.0	699.5 ± 1.4	1.179 ± 0.004	10	5.3 ± 2.0	0.04 ± 0.04	1.43 ± 0.38	0.97
41	5202800117	59755.628	561.0	707.9 ± 1.1	1.03 ± 0.01	17	5.1 ± 2.0	0.11 ± 0.05	1.0 ± 0.19	0.81
44	5202800118	59756.392	115.0	630.1 ± 2.3	1.39 ± 0.01	3	4.5 ± 2.3	-0.13 ± 0.09	0.29 ± 0.45	0.65
46	5202800119	59757.158	805.0	715.1 ± 0.9	1.86 ± 0.01	24	5.2 ± 1.9	0.1 ± 0.04	1.17 ± 0.14	0.93
47	5202800119	59757.293	343.0	692.9 ± 1.4	1.062 ± 0.004	10	3.0 ± 2.0	...	0.0 ± 0.59	0.96

Table A2 *continued*

Table A2 (*continued*)

GTI	ObsID	MJD (TT)	Exposure	r_{tot}	r_{bg}	m	f_{rms}	A	PL Index	χ^2
			s	c/s	c/s		% (0.03–50 Hz)			
49	5202800120	59758.457	87.0	651.7 ± 2.7	0.845 ± 0.003	2	7.1 ± 2.4	-0.37 ± 0.1	0.45 ± 0.16	0.69
51	5202800120	59758.9	90.0	667.7 ± 2.7	1.128 ± 0.004	2	6.7 ± 2.4	-0.34 ± 0.11	0.7 ± 0.14	0.67
52	5202800120	59758.964	206.0	631.9 ± 1.8	1.128 ± 0.004	6	5.6 ± 2.2	0.0 ± 0.01	2.04 ± 2.88	0.85
53	5202800121	59759.029	327.0	607.2 ± 1.4	1.05 ± 0.01	10	...	-0.15 ± 0.04	0.41 ± 0.16	1.03
54	5202800121	59759.095	143.0	653.2 ± 2.1	1.0 ± 0.01	4	5.0 ± 2.2	-0.15 ± 0.07	0.58 ± 0.23	0.78
55	5202800121	59759.759	697.0	636.8 ± 1.0	2.86 ± 0.02	21	4.4 ± 2.0	0.07 ± 0.04	1.17 ± 0.2	0.64
56	5202800121	59759.808	423.0	592.4 ± 1.2	2.86 ± 0.02	13	3.9 ± 2.2	0.03 ± 0.03	1.6 ± 0.34	0.72
57	5202800122	59760.648	406.0	666.4 ± 1.3	3.21 ± 0.01	12	2.4 ± 2.0	-0.03 ± 0.04	0.3 ± 0.93	0.94
58	5202800122	59760.712	307.0	658.9 ± 1.5	3.22 ± 0.01	9	3.2 ± 2.1	0.01 ± 0.02	1.91 ± 0.74	0.68
59	5202800123	59761.762	273.0	613.9 ± 1.5	0.88 ± 0.01	8	3.1 ± 2.1	...	0.01 ± 0.49	0.73
61	5202800124	59762.39	439.0	642.2 ± 1.2	1.24 ± 0.01	13	3.7 ± 2.1	...	0.0 ± 1.33	0.92
62	5202800124	59762.967	313.0	578.9 ± 1.4	1.24 ± 0.01	9	1.0 ± 2.2	0.03 ± 0.05	1.22 ± 0.68	0.75
63	5202800125	59764.126	393.0	621.0 ± 1.3	1.36 ± 0.01	11	1.9 ± 2.1	0.02 ± 0.03	1.64 ± 0.42	0.88
64	5202800125	59764.193	627.0	629.2 ± 1.0	1.38 ± 0.01	19	3.4 ± 2.1	0.07 ± 0.04	1.22 ± 0.21	0.75
65	5202800126	59765.892	349.0	582.1 ± 1.3	2.7 ± 0.03	10	4.5 ± 2.2	0.03 ± 0.05	1.24 ± 0.53	0.71
66	5202800126	59765.933	43.0	618.1 ± 3.8	2.73 ± 0.03	1	...	-1.08 ± 0.11	0.41 ± 0.05	0.59
67	5202800126	59765.934	140.0	588.6 ± 2.1	2.73 ± 0.03	3	5.0 ± 2.4	-0.31 ± 0.08	0.38 ± 0.13	0.69
69	5202800127	59766.666	294.0	616.5 ± 1.4	0.93 ± 0.01	9	5.9 ± 2.2	...	0.0 ± 1.21	0.96
70	5202800127	59766.731	314.0	615.2 ± 1.4	0.93 ± 0.01	9	5.2 ± 2.2	0.04 ± 0.05	1.32 ± 0.43	0.83
71	5202800128	59767.291	613.0	634.9 ± 1.0	...	19	...	0.04 ± 0.04	1.26 ± 0.28	0.89
73	5202800128	59767.612	107.0	602.4 ± 2.4	2.11 ± 0.01	3	...	-0.3 ± 0.08	0.46 ± 0.08	0.67
74	5202800128	59767.614	864.0	585.9 ± 0.8	...	26	3.3 ± 2.1	0.0 ± 0.01	1.85 ± 0.75	0.81
75	5202800129	59768.149	328.0	661.1 ± 1.4	2.1 ± 0.02	9	1.9 ± 2.1	0.1 ± 0.06	1.18 ± 0.22	0.79
76	5202800129	59768.195	634.0	683.5 ± 1.0	2.1 ± 0.02	19	3.1 ± 2.0	0.05 ± 0.04	1.23 ± 0.23	0.88
77	5202800130	59769.035	334.0	599.7 ± 1.3	0.871 ± 0.004	10	3.9 ± 2.2	-0.06 ± 0.07	0.18 ± 0.6	1.08
78	5202800130	59769.099	548.0	648.0 ± 1.1	0.862 ± 0.003	15	5.7 ± 2.1	0.08 ± 0.05	0.98 ± 0.22	0.87
79	5202800130	59769.168	199.0	659.6 ± 1.8	0.862 ± 0.005	6	1.3 ± 2.1	-0.06 ± 0.31	0.11 ± 1.11	0.83
80	5202800131	59771.233	372.0	603.7 ± 1.3	2.22 ± 0.01	11	0.8 ± 2.1	...	0.0 ± 0.95	0.79
81	5202800131	59771.297	375.0	548.7 ± 1.2	2.22 ± 0.01	11	0.01 ± 0.6	0.93
82	5202800131	59771.362	373.0	608.2 ± 1.3	0.776 ± 0.003	11	3.9 ± 2.1	-0.03 ± 0.04	0.49 ± 0.77	0.95
83	5202800132	59783.974	311.0	491.9 ± 1.3	1.04 ± 0.02	9	7.6 ± 2.4	0.03 ± 0.05	1.15 ± 0.68	0.75
84	5202800133	59784.039	298.0	503.7 ± 1.3	1.04 ± 0.02	9	...	-0.12 ± 0.06	0.2 ± 0.28	0.92
85	5202800133	59784.426	359.0	522.3 ± 1.2	0.94 ± 0.01	11	4.9 ± 2.3	-0.03 ± 0.08	0.16 ± 1.07	0.99
86	5202800134	59785.071	394.0	510.1 ± 1.1	1.03 ± 0.01	12	3.7 ± 2.3	-0.01 ± 0.04	0.26 ± 1.86	0.84
87	5202800134	59785.393	318.0	499.4 ± 1.3	2.66 ± 0.01	9	4.2 ± 2.4	-0.03 ± 0.05	0.35 ± 0.77	0.63
89	5202800135	59793.183	550.0	410.0 ± 0.9	1.0 ± 0.02	17	2.8 ± 2.5	...	0.0 ± 0.44	1.21
90	5202800136	59794.108	83.0	413.9 ± 2.2	1.02 ± 0.02	2	3.5 ± 3.0	-0.63 ± 0.09	0.38 ± 0.08	0.88
91	5202800137	59795.721	43.0	409.9 ± 3.1	1.08 ± 0.02	1	0.3 ± 3.5	-1.3 ± 0.11	0.35 ± 0.05	0.8
92	5202800138	59796.79	1159.0	392.9 ± 0.6	2.41 ± 0.01	36	...	0.01 ± 0.02	1.24 ± 0.78	0.85
93	5202800139	59797.563	443.0	368.4 ± 0.9	2.43 ± 0.01	13	...	-0.01 ± 0.02	1.39 ± 0.99	0.64

Table A2 *continued*