

1 **Title page**

2 **Semiconductor core fibres: a scalable platform for nonlinear**
3 **photonics**

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Abstract

Semiconductor core, glass cladding fibres that can be produced with scalable dimensions and unique waveguide designs are offering new opportunities for nonlinear photonics. This paper reviews developments in the fabrication and post-processing of such semiconductor core fibres and their enabling of low loss and high efficiency nonlinear components across wavelengths spanning the near- to mid-infrared. Through adaption and expansion of the production processes, routes to new core materials are being opened that could extend the application space, whilst all-fibre integration methods will result in more robust and practical semiconductor systems. Through continued improvement in the core materials, fibre designs and transmission losses, semiconductor fibres are poised to bring unique functionality to both the fibre and semiconductor research fields and their practical application into a myriad of optoelectronic devices.

Main body

1 Introduction

Nonlinear semiconductor photonics is a rapidly growing field of research and application. The combination of the tight mode confinement and the high nonlinear refractive indices available in these materials allows for very compact and efficient device development¹⁻³. Further, a wide range of semiconductor materials have been utilized for nonlinear signal processing, giving access to different optical transmission regions⁴ and nonlinear susceptibilities, to enable either second⁵ or third order processes⁶. Regardless of the choice of semiconductor, the majority of work in this field has made use of integrated on-chip platforms, primarily because of the optimized materials production methods (growth and deposition) for the planar formats⁷. Access to high quality materials is particularly important for the realization of low loss waveguides that are essential for nonlinear processing, which has led to a focus on Group IV semiconductors. Capitalizing on the excellent quality and the broad wavelength coverage offered by the Group IV materials, numerous nonlinear processes and demonstrations of device functionality have been reported, including optical signal processing⁸ and source generation in silicon^{9,10}, all-optical modulation in germanium¹¹, frequency comb generation in silicon nitride¹², and supercontinuum generation in silicon-germanium¹³, for example.

However, there are challenges to fabricating and working with the planar waveguides, particularly those with nanoscale dimensions, that still require some attention. Firstly, as the waveguides are formed by lithography and etching, the surface quality of the side walls can be difficult to control, and the nanoscale roughness values that are typical of most etch processes can result in sizable scattering losses at the high index contrast core/cladding boundaries¹⁴. Secondly, owing to the size-mismatch between nanoscale waveguides and micron-sized optical fibres, it is difficult to achieve efficient and robust coupling between the two

40 platforms. This is particularly true for use in nonlinear applications where the coupling must support relatively high pump powers
41 and broad operating bandwidths, so that end-fire coupling is really the only viable option^{15,16}. Although inverse tapers have been
42 developed for low loss coupling within the telecom band¹⁷, they are not suitable for mid-infrared systems when silica cladding
43 glasses are used, and also require very precise fabrication and alignment.

44 Over the past decade, semiconductor core optical fibres have emerged as an interesting alternative for the development of
45 nonlinear semiconductor devices. Semiconductor fibres retain many of the advantageous properties of the fibre platform such as
46 durability, polarization insensitivity, pristine core/cladding interfaces, and potential to access long waveguide lengths¹⁸. However,
47 compared with traditional silica fibres, the semiconductor cores bring benefits of the on-chip systems such as the tight light
48 confinement, so that they can be scaled to sub-micrometre dimensions, and the high nonlinear refractive indices reduce the power
49 requirements of the nonlinear systems. Moreover, as semiconductors usually have much broader transmission windows than
50 silica, semiconductor core fibres can extend the operation wavelength of fibre-based nonlinear systems from the near-infrared
51 into the mid-infrared region¹. Further, and notably, semiconductor fibres are now principally fabricated using conventional fibre
52 drawing methods, with the high drawing speeds allowing for the rapid production of many hundreds of metres of fibre to ensure
53 low costs and high yields of the in-fibre devices^{19,20}. Through careful consideration of the draw parameters and the starting
54 material, fibres can be produced with a wide range of unary and compound semiconductor cores²¹. Yet, the production of these
55 waveguides is not without its own challenges, and it is generally harder to control the crystalline alignment of the semiconductor
56 core phase during the draw process, with the as-drawn fibres being polycrystalline²². It can also be difficult to achieve the con-
57 tinuous nanoscale core dimensions desired for many nonlinear processes due to capillary instabilities in the draw²³. However,
58 taking advantage of the robust fibre cladding, well-established fibre post-processing techniques can be exploited to optimize the
59 core materials and achieve dimensions down to the nanoscale to not only enhance the nonlinear efficiency, but also improve the
60 coupling to other fibre and on-chip systems.

61 In this paper, we review recent developments in the fabrication and post-processing techniques that have enabled the produc-
62 tion of semiconductor core fibres that support efficient nonlinear processing. In particular, we describe the variety of materials
63 that are compatible with the thermal drawing process and the advantages that the fibre systems offer over planar production
64 methods. Following a review of the nonlinear processes that have been observed in the semiconductor fibres extending from the
65 telecom band up into the mid-infrared, some perspective will be provided on the possibility to expand the application potential
66 via access to new fibre materials and fully integrated systems. Through continued advancement of the fibre materials and pro-
67 cessing techniques, we anticipate that these fibres will open up new avenues of exploration for semiconductor nonlinear photon-
68 ics.

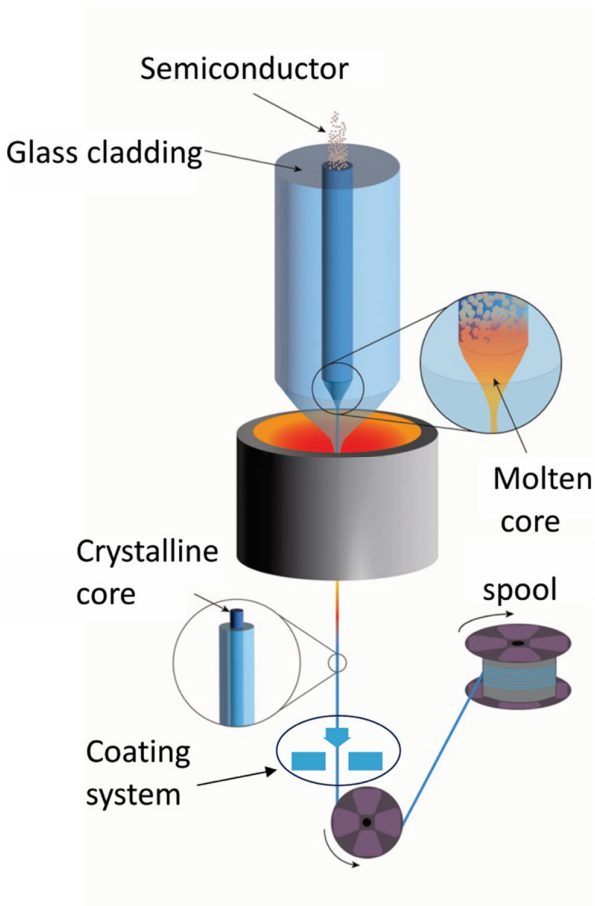
69 **2 Fabrication Procedures**

70 **2.1 Molten core drawing method**

71 To date, there have been several approaches employed to fabricate semiconductor optical fibres. Chronologically, the first was
72 the micro-pull-down method, which realized the first silicon fibres in 1996²⁴. The second, starting in 2006, employed high-
73 pressure chemical vapour deposition (HP-CVD) inside of a prefabricated glass capillary that subsequently acted as the fibre
74 cladding²⁵. The third, in 2008, was a variant of the core suction method but using a pressured (instead of evacuated) melt to flow
75 into a capillary²⁶. Two weeks later in 2008, the first use of the molten core method (MCM, also sometimes referred to as the
76 “melt-in-tube” method) to semiconductor core fibres was published¹⁹. Lastly, a “reactive” approach to semiconductor phase
77 formation was reported (first in 2011) whereby other phases are employed as precursors that react during fibre formation to yield
78 the desired core phase²⁷⁻²⁹. Of these, only optical transmission has been reported for fibres fabricated via HP-CVD, pressurized
79 filling, and the MCM. A series of review articles over the course of the past decade provides evolving details on each approach
80 and their relative advantages and disadvantages^{18,21,30}.

81 For the purposes of this article, the MCM will be the focus. This is because, of the noted fabrication approaches, only the
82 MCM relies on the thermal drawing of glass. This inherently makes the MCM a more scalable approach, permitting continuous
83 fibre fabrication at higher speeds and over long (> 100s of metre and longer) lengths³¹. As a result, it has become the principal
84 method of semiconductor and multi-material fibre fabrication³²⁻³⁶.

85 The MCM was originally developed to fabricate fibres with cores that were not sufficiently stable to first form into a glass
86 and draw^{37,38}. The approach employs a glass cladding tube that is closed at one end and then filled with a precursor phase that
87 melts at the draw temperature of the glass cladding. A canonical case applicable here is silicon (Si), with a melting temperature
88 of about 1414 °C, sleeved and drawn in a pure silica glass tube at a temperature of about 1950 °C¹⁹. As discussed in greater detail
89 in ref. ³⁷, the generalized MCM has been applied to a wide range of glassy and crystalline core phases, in addition to the crystalline
90 nonlinear semiconductor cores discussed herein. A schematic of the MCM is provided in Figure 1. Similar to traditional fibre
91 drawing, it includes a coating process to improve the mechanical strength³⁹. A key benefit of this approach to semiconductor
92 waveguide fabrication is that the core/cladding surfaces are defined by the atomically smooth silica glass tubing⁴⁰, which helps
93 to minimize scattering losses at the high index contrast boundary, of particular significance for waveguides with nanoscale
94 dimensions.



95

96 **Figure 1 A schematic of the molten core method for thermally drawing glass-clad semiconductor core fibres.** The crystalline core is
 97 melted and encapsulated by the viscous glass cladding during the drawing process. A coating system is used to improve the mechanical strength
 98 of the produced as-drawn semiconductor core fibres.

99 Though conceptually straightforward, as with any natural process, thermodynamics and kinetics always create limitations.
 100 For example, while amorphous cores are quite common when the MCM is applied to oxide systems (including with crystalline
 101 precursors³⁷), as-drawn semiconducting, i.e., semi-metal cores are always fully crystalline. This contrasts with HP-CVD, which
 102 can yield either crystalline or amorphous semiconductor core fibres⁴¹. A second important natural consequence of processing
 103 materials above their melt temperature is the propensity for the melt to corrode the crucible in which it is contained. In the case
 104 of the MCM, the core melt dissolves in some components of the softened glass cladding at temperatures approaching the draw
 105 temperature. Indeed, this is what partly contributes to the formation of glassy cores in many molten core oxide systems. For the
 106 semiconductors, where crystallization kinetics are many times faster than the fibre draw speed and associated solidification rates,
 107 the dissolution of silica into silicon, for example, leads to scattering inclusions and thus increased losses. Dissolution and the
 108 accompanying diffusion are both thermally driven processes that seek to reduce a compositional gradient (i.e., silicon core, silica
 109 cladding). Accordingly, drawing at higher temperatures to fibres of smaller core sizes always promotes greater concentrations

110 of cladding species in the core phase⁴². However, for nonlinear optics, small core sizes are highly desirable, if not wholly
111 required, making thermodynamics the enemy of optical performance. Unsurprisingly then, amongst the first optimization efforts
112 that followed the materials discovery phase of molten core semiconductor fibre development were approaches to mitigate
113 cladding dissolution into the core. These activities focused on reactive chemistries, first “after-the-fact” by reacting away
114 cladding oxides in the core melt⁴³ and then by proactively limiting the ingress from the initiating²³. The interfacial modifier
115 method by Gibson, et al.,²³ is arguably amongst the most important material and process advances in the practicality of
116 semiconductor optical fibres fabricated via conventional draw towers. However, even more recently, it has been shown that
117 swapping the usual thermal furnace with a laser heat source permits the use of faster drawing speeds, allowing for the production
118 of small core (1.3 μm diameter) silicon fibres with minimal cladding ingress³⁹. With or without the interfacial modifier to mitigate
119 cladding dissolution into the core, the cladding glass chemistry is also very important as it contributes significantly to the draw
120 temperature, thermal expansion mismatch (hence residual stresses), and which species are present for contaminating the core.
121 Cladding glasses employed for the molten core fabrication of semiconductor optical fibres include to-date silica^{19,23},
122 borosilicate^{44,45}, and phosphate^{46,47} glasses as well as proposals for property-matched heavy metal oxides and chalcogenides⁴⁸.

123 The MCM has enabled the realization of a vast variety of crystalline semiconductor fibre core phases. Of these materials, the
124 Group IV core phases are by far the most well-studied and developed^{19,23,29,43,44,48-50}. The principal reasons for such a focus are
125 several-fold and include (i) relative simplicity of the core phase (e.g., unary cubics), (ii) relative ease of fabrication with well-
126 behaved core melting and solidification, and generally stable and commercially available (in most cases) cladding glasses, and
127 (iii) great familiarity with Si and Ge, and their SiGe alloys, given their commercial ubiquity. For the most developed system, Si
128 core in a silica glass cladding, a kilometre of fibre (125 μm outer diameter and 8-10 μm core size) is straightforward to fabricate
129 at draw rates on the order of metres per minute. Moreover, even the first molten core Ge fibre was reported at about 250 metres,
130 despite the fact that a less durable borosilicate glass cladding was employed and that a two-draw process was required (preform
131 to cane, followed by cane to fibre) to limit surface crystallization of the borosilicate⁴⁴. Thus, the propensity of the MCM for such
132 large-scale waveguide production is highly beneficial for reducing costs and increasing component yields.

133 Beyond the Group IV core phases, several II-VI compound semiconductors were amongst the earliest to be fabricated and
134 studied, specifically InSb, in 2010⁴⁶, and GaSb, in 2013^{51,52}. As with the Group IV phases, these antimonides are well-behaved in
135 that they melt congruently at relatively low temperatures and with negligible vapour pressure. This latter point, negligible vapour
136 pressure, has been a historical limitation of the MCM since the build-up of vapour in the core during preform heat-up or fibre
137 draw can lead to the explosion of the preform. It had, until recently (2022), precluded the fabrication of semiconductor cores of
138 desirous nonlinear (and direct bandgap) phases such as GaAs⁵³ and ZnSe⁵⁴. In cases such as these, where the core phase exhibits
139 sufficient volatility upon heat-up to the (cladding glass) draw temperature and/or where the core incongruently melts, the solution

140 has been to include a low melting flux phase (i.e., “flux molten core”). The flux melts at a low temperature, dissolving the core
141 phase into a homogeneous and low volatility liquid, which then is drawn similarly to “conventional” molten core fibre
142 fabrication. The flux phase can be segregated away from the semiconductor phase using laser post-processing (see Section 3).
143 Thus, the MCM is opening a new route to the fabrication and optimization of semiconductor waveguides from materials that are
144 not readily available in an integrated on-chip format, and which could eventually be fully incorporated within all-fibre optical
145 and optoelectronic systems⁵⁵.

146 As previously mentioned, the as-drawn molten core semiconductor fibres are polycrystalline. Put another way, the
147 crystallographic orientation of the core changes discretely along the fibre. Single crystalline regions (i.e., grains) in the as-drawn
148 fibres of several millimetres, up to approximately a centimetre, are typical. The optical isotropy of the cubic Group IV
149 semiconductors permits light transmission even through polycrystalline cores, though scattering from impurities and defects at
150 grain boundaries might add to loss. Not long after the noted activities to manage cladding dissolution, efforts to induce single
151 crystallinity of the fibre cores over device-relevant lengths became active topics of study and development. The earliest
152 approaches to promoting single crystallinity involved thermal annealing⁵⁶⁻⁶⁰. However, this was later superseded, firstly by laser
153 heat treatments, then secondly by fibre tapering. Compared with the MCM, a key benefit of post-processing is that it is possible
154 to apply a more precise and stable heat treatment to a selected fibre region, allowing for the formation of longer semiconductor
155 crystal grains within the glass cladding. Thus, it is hoped that through combination of the recent advancements in the fibre
156 drawing methods^{39,50} and the post-processing treatments described below, there will be a wider availability of low loss
157 semiconductor fibres for distribution and application amongst the broader nonlinear research community.

158 **3 Post-processing Techniques**

159 **3.1 Laser processing**

160 A schematic of laser post-processing is shown in Figure 2a. The aim of this procedure is to heat, melt, and recrystallize the core
161 material to grow larger polycrystalline grain sizes. The laser is focused to a small spot size by a lens, and either the laser or the
162 fibre is scanned to control the heating and cooling dynamics of the semiconductor core, and to process longer lengths. To generate
163 sufficient heat to melt the semiconductor, the laser wavelengths are usually selected where the core materials have strong
164 absorption (e.g., 488 nm, 517 nm and 10.6 μm for Si^{61,62}, 514 nm, 632 nm and 10.6 μm for SiGe⁶³, 532 nm for Te³², 10.6 μm for
165 GaSb^{55,64}). By placing a camera perpendicular to the laser beam, it is possible to monitor the melting and recrystallizing
166 processes. Compared with previous thermal annealing methods (lamp-annealing⁵⁶, furnace-annealing⁵⁷⁻⁵⁹), the laser heating
167 process has many advantages. Firstly, the focused spot size is adjustable by changing the lens. Therefore, the size and position

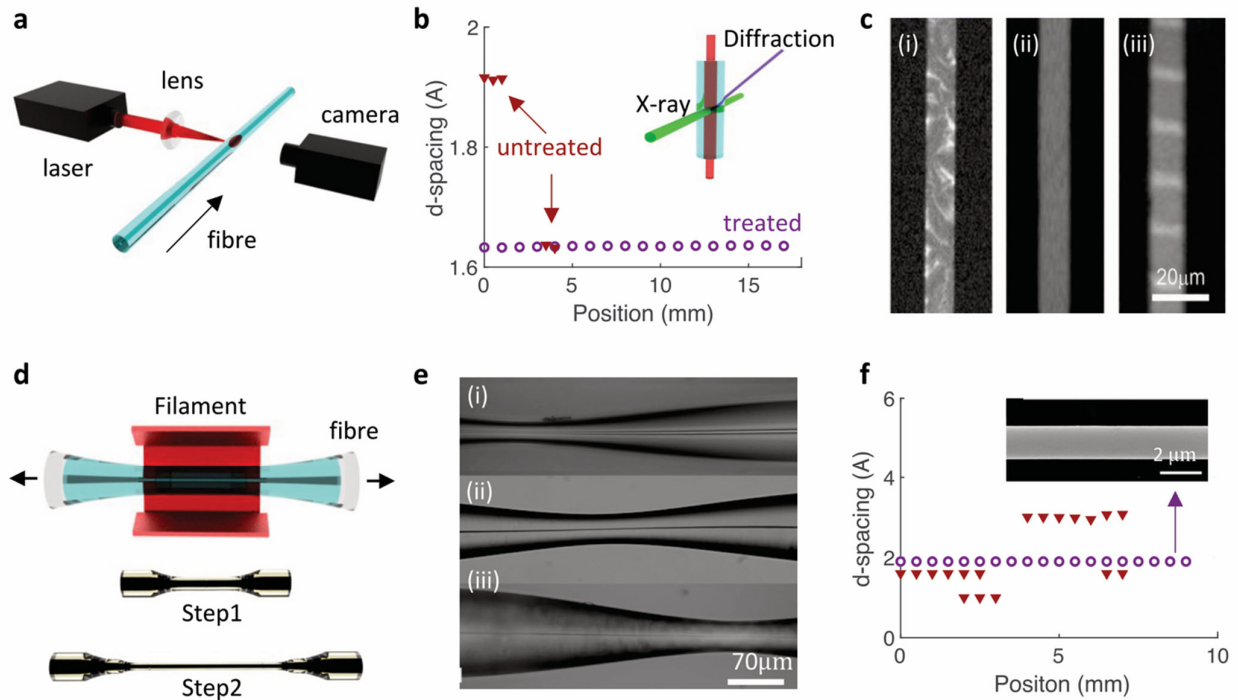
168 of the heating region can be more precisely controlled to be predominantly in the cladding or the core, and can even be used to
169 generate more complex structures within the fibre (e.g. Bragg gratings⁶⁵). Secondly, the laser processing has a gradient heating
170 distribution along the radial direction in a focused beam spot. As different semiconductor materials have different solid points,
171 this factor can be used to redistribute different components of a multi-material core fibre, making them either uniform⁶³ or
172 segregated^{52,64}.

173 The early work on laser processing was mainly focused on Si core fibres. In 2014, it was first used to engineer the bandgap of
174 Si from 1.11 eV down to 0.59 eV by introducing an anisotropic tensile stress within a HP-CVD fibre by crystallizing the
175 amorphous Si core to polycrystalline Si⁶¹. As the laser wavelength (488 nm) used in ref. ⁶¹ is transparent for the silica cladding,
176 the cylindrical fibre geometry offers a unique advantage as it allows the Si core to reach a molten state, whereas the cladding is
177 only modestly heated. As the crystalline Si core remains strongly adhered to the cladding but occupies a smaller volume than
178 the amorphous cladding material, this results in a large strain. The second, in 2016, used a CO₂ laser to recrystallize the poly-Si
179 cores in MCM fibres to single-crystalline like materials through control of the cooling dynamics, thus reducing the optical loss
180 from 14-20 dB/cm to only 2 dB/cm at 1.55 μm and 1 dB/cm at 2 μm wavelength⁶². The crystallography for the Si core after CO₂
181 laser writing was characterized by measuring the X-ray diffraction (XRD) at different positions along the fibre length, using the
182 configuration shown in the inset of Figure 2b. Figure 2b displays the recorded crystallographic d-spacing for both processed and
183 unprocessed fibres, showing that the laser treatment can produce single-crystal grain sizes with lengths of more than 15 nm,
184 highlighting the potential for the production of high quality semiconductor core fibres. Moreover, in 2017, a third body of work
185 made use of laser processing to fabricate resonators and Bragg gratings by using the heat treatment to manipulate the core
186 structure and glass cladding of the Si core fibres⁶⁵, highlighting the flexibility to modify both the material and shape of the fibres.

187 Subsequent interests in laser processing then shifted to compound materials, and predominantly SiGe core fibres, which aim
188 to extend the mid-infrared wavelength coverage over the pure Si cores⁶⁶. In 2016, a CO₂ laser was used to homogenize the uneven
189 segregation of Si and Ge within as-drawn SiGe core fibres, as illustrated by comparing the first two fibres in Figure 2c⁶³. The
190 results in ref. ⁶³ also showed that it was possible to manipulate the positioning of the Ge content to form gratings, shown by
191 bright sections in the third fibre in Figure 2c, or even a graded index Ge-rich core. By exploiting this latter feature, a fibre with
192 a Ge-rich central region (22 %) within a low Ge concentration (6 %) SiGe core fibre was formed in 2020, highlighting the
193 versatility of the post-processing technique⁶⁷.

194 Inspired by the successes of laser processing for semiconductor fibres with Group IV core materials, it was also used to
195 engineer the optical properties of fibres with other core materials, including GaSb^{52,64} and Te³². Moreover, as previously
196 mentioned, it is anticipated that laser processing could be used to segregate sections of pure GaAs and ZnSe core materials from
197 the flux phases to create fibres suitable for second order nonlinear optical processing. Beyond optimization of the optical

198 properties, laser processing has also been used to fabricate optoelectronic devices within semiconductor core fibres. In 2019, a
 199 CO₂ laser beam was used to segregate GaSb and Si within an as-drawn composite GaSb/Si fibre to enhance the
 200 photoluminescence⁶⁸.



201
 202 **Figure 2 Post-processing techniques for semiconductor core fibres.** **a** Schematic of laser processing. **b** Measured lattice spacing from X-
 203 ray diffraction as a function of position along as-fabricated and laser processed Si core fibres⁶². The inset shows the X-ray diffraction meas-
 204 urement setup. **c** Optical images of (i) an as-drawn SiGe core fibre, (ii) a laser recrystallized SiGe core and (iii) a Ge-rich grating formed
 205 within the SiGe core. **d** Schematic of the tapering procedure, showing how multiple taper steps can be employed. **e** Images of longitudinal
 206 taper profiles for starting fibre core diameters of (i) 5.6 μm, (ii) 2.7 μm, and (iii) 1.3 μm. **f** Measured X-ray diffraction along the as-fabricated
 207 and tapered Si core fibre. The inset shows an image of the smooth surface of an etched Si core fibre waist⁶⁹. **c** Reproduced from ref. ⁶³ with
 208 the permission from Springer Nature. **e** Reproduced from ref. ⁷⁰ with the permission from Optica Publishing Group.

209 3.2 Tapering

210 Another important post-processing method is fibre tapering. A schematic of the tapering process is shown in Figure 2d, where a
 211 filament is used to melt the semiconductor core and soften the glass cladding. The fibre is then stretched using controllable stages
 212 to reduce the core/cladding diameters along the length. As the tapering procedure essentially mimics a second MCM draw
 213 process, importantly it provides a means to control the longitudinal profile of the fibre as well as the crystallinity. However, as
 214 the material volume is now much smaller, a lower and more controllable level of heating powers can be used. To ensure the
 215 fibres have sufficient mechanical strength after being tapered to smaller sizes, the as-fabricated fibres can be sleeved inside

216 thicker capillaries before tapering. As illustrated in Figure 2d, two-step tapering can be used to tailor fibres with large core
217 diameters. Compared to single-step tapering, two-step tapering can use smaller tapering ratios and lower heating powers for the
218 final step, which is important for producing high-quality single-crystalline semiconductor cores with small diameters over long
219 fibre lengths⁶⁹.

220 The first tapering work was demonstrated in 2010, using a fusion splicer to adjust the core diameter of Si core fibres from
221 5.6 μm , 2.7 μm and 1.3 μm down to waist diameters of ~ 3 μm , ~ 2.7 μm and ~ 500 nm, respectively⁷⁰. The images of these three
222 tapered fibres are shown in Figure 2e. As can be seen, the tapering process produces a solid, continuous Si core with a smooth
223 transition from the untapered fibre to the taper waist, where at each point along the taper the core diameter changes proportionally
224 to the changing cladding diameter. In 2012, the tapering technique was extended to tailor the core diameter of a Ge core fibre
225 from ~ 150 μm to ~ 45 μm over a length of 16 mm, where it was shown that good control of the heating temperature (typically
226 ± 5 $^{\circ}\text{C}$) is important for increasing the single crystallinity²². However, it should be noted that no nonlinear effects were observed
227 in refs. ^{70,22} due to the relatively short waist lengths (~ 100 μm) and large diameters (~ 45 μm) for the Si and Ge core fibres,
228 respectively.

229 Low loss tapered Si core fibres suitable for the observation of nonlinear propagation, including nonlinear absorption and self-
230 phase modulation (SPM), were first produced in 2016. Specifically, a submicron-sized core diameter of 0.94 μm was achieved
231 over a length of 1 cm, with a loss of ~ 3 dB/cm⁷¹. XRD measurements along the tapered fibre length showed a similar level of
232 improvement in the crystal grain size to what was seen following laser processing, as illustrated in Figure 2f⁶⁹. Thus, tapering
233 also provides a route to producing single-crystalline like core materials. Importantly, the SEM image of an etched fibre core near
234 the tapered waist region (inset of Figure 2f) confirms that the surface of the core remains ultrasmooth (sub-nanometre roughness),
235 thanks to the pristine cladding glass crucible. Thus, this serves to highlight the advantage of the fibres to achieve nanoscale
236 waveguide dimensions with minimal loss due to surface scattering.

237 To date, tapering is the only post-processing method that can tailor down the MCM fibre core diameters to nanoscale
238 dimensions (< 1 μm). Although nanoscale-sized semiconductor cores can be produced directly via HP-CVD by using suitably
239 small silica capillary templates (~ 600 nm in ref. ²⁶), the deposition process is very time-consuming to achieve such small cores,
240 the dimensions are fixed by the template, and the fibre lengths are limited to ≤ 1 cm. In contrast, the tapering technique offers a
241 more flexible solution as the core size can be conveniently adjusted by changing the tapering ratios and the lengths are not limited
242 by the as-drawn fibre, but only the choice of tapering rig (with a tapered fibre length > 6 cm being recently achieved⁷²). Moreover,
243 interesting fibre profiles can be fabricated with this method to enhance the functionality of the fibres. For example, in 2017, a
244 nano-spike tip was formed by tapering the Si core to the point of collapse to form an inverse tapered coupler⁷³, analogous to what

245 has been used in on-chip formats⁷⁴, as will be described in Section 5. Then in 2019, an asymmetric tapered fibre profile was
246 designed to improve the output coupling of a supercontinuum generated in the mid-infrared, as detailed in Section 4⁷⁵. More
247 recently tapering approaches have been applied to SiGe fibres, showing that similar materials and structural improvements can
248 be made even in compound semiconductor systems⁷⁶.

249 **4 Optical Characterization**

250 Thanks to the materials advancements associated with the laser processing and tapering procedures, several semiconductor core
251 fibres have been produced with the low transmission losses required for nonlinear propagation. Table 1 summarizes the optical
252 losses that have been obtained in the telecom band and at selected mid-infrared wavelengths for semiconductor core fibres that
253 have been used for nonlinear applications. The most widely used platform has been the Si core fibres. By using laser post-
254 processing and tapering techniques, the optical loss of MCM Si core fibres (10 dB/cm) can be reduced to 2 dB/cm and 0.8 dB/cm
255 at a wavelength of 1.55 μm , respectively. Compared with laser post-processing, the lower losses associated with tapering can be
256 attributed to two main reasons. Firstly, the heating zone is larger and has a more uniform temperature distribution, thus making
257 it easier to generate a more uniform, continuous semiconductor core. Secondly, the heating zone of the filament is more stable
258 than a focused laser beam, resulting in fewer defects. To date, the most remarkable loss reduction achieved by laser processing
259 has been down to ~ 1 dB/cm at 1.55 μm wavelength in a HP-CVD amorphous Si core fibre, though the length was limited to
260 < 2 mm⁷⁷.

261 Although less well studied, SiGe core fibres have been attracting significant attention of late due to the material's potential
262 for higher nonlinear coefficients and extended mid-infrared wavelength coverage in comparison to unary Si⁶⁶. As the Si and Ge
263 tend to segregate unevenly during the MCM drawing (see Figure 2c), as-fabricated SiGe core fibres usually show very high
264 losses (> 20 dB/cm), dominated by scattering. To reduce the loss of SiGe core fibre, in 2016, laser processing was used to recrystallize
265 the SiGe core to achieve losses of 12 dB/cm and ~ 10 dB/cm at 1.55 μm and ~ 2 μm , respectively. However, such loss values
266 are still too high for practical use and the core diameter of the first-generation as-drawn fibres were also very large (~ 130 μm),
267 precluding their use for nonlinear applications⁶³. By employing tapering to reduce the core size and losses, these fibres have
268 recently been produced with loss values of ~ 3 dB/cm for core diameters of a few micrometers⁷⁶. It is worth mentioning that, in
269 addition to the tapering discussed previously, pure Ge core fibres have also been successfully laser processed, during which the
270 losses were reduced from ~ 7 dB/cm to ~ 2 dB/cm at wavelengths ~ 5 μm ⁸¹. However, the core diameters of the processed fibres
271 (> 22 μm) were too large to observe nonlinear propagation.

272 The only binary semiconductor fibres to be used for nonlinear demonstrations are the HP-CVD fabricated ZnSe core fibres.
 273 By optimizing the deposition recipe and temperature, these fibres could be produced with losses of only 1 dB/cm across the near-
 274 infrared and mid-infrared regimes, so that further post-processing was not required to observe nonlinear effects. However, the
 275 HP-CVD fibres are highly polycrystalline and thus it is hoped that post-processing the MCM ZnSe fibres will eventually produce
 276 fibres with higher quality core materials and controllable dimensions. Such work would be hugely significant as currently there
 277 are no practical fibres that can offer access to second order nonlinear processing.

278 **Table 1 Measured optical loss (unit: dB/cm) for semiconductor core fibres that have been used for nonlinear demonstrations. Top rows**
 279 **are for the telecom wavelength of 1.55 μm . Bottom rows for selected mid-infrared wavelengths of $\ast\sim 2.5 \mu\text{m}$ and $\dagger\sim 5 \mu\text{m}$.**

Platform	As-fabricated	Reference	Laser processed	Reference	Tapered	Reference
Si (MCM)	~ 10	23	2	62	0.8	78
a-Si (HP-CVD)	~ 50	79	1	77		
SiGe (MCM)	~ 20	63	12	63	2.2	76
ZnSe (HP-CVD)	1	80				
Si (MCM) \ast	~ 10	19	1	62	0.2	75
SiGe (MCM) \ast	~ 20	63	~ 10	63	4	76
ZnSe (HP-CVD) \ast	1	80				
Ge (MCM) \dagger	~ 7	81	~ 2	81		

280

281 5 Nonlinear Demonstrations

282 5.1 Description of nonlinear propagation

283 Following the realization of semiconductor fibres with losses of a few dB/cm and core diameters of a micrometre or less, it was
 284 not long before their potential for nonlinear application was demonstrated. Similar to the planar nanophotonic waveguides, pulse
 285 propagation in the semiconductor core fibres with a dominant χ^3 nonlinearity, as per the Group IV materials, can be described
 286 by the generalized nonlinear Schrödinger equation (GNLSE)⁸²:

$$287 \frac{\partial A}{\partial z} = -\frac{i\beta_2}{2} \frac{\partial^2 A}{\partial t^2} - \frac{1}{2} \left(\alpha + \sigma_f + i\alpha_1 \frac{\partial}{\partial t} \right) A + i \left(\gamma + i\gamma_1 \frac{\partial}{\partial t} \right) \left(A \int_0^\infty R(t) |A|^2 dt \right) |A|^2 A. \quad (1)$$

288 Here $A(z, t)$ is slowly varying pulse envelope, $\beta_2(z)$, is the group velocity dispersion (GVD), α is the linear loss and $\alpha_1 =$
 289 $d\alpha/d\omega$. σ_f is the free carrier contribution and $\sigma_f = \sigma(1 + i\mu)N_c$, where σ is the free-carrier absorption (FCA) coefficient and μ
 290 governs the free-carrier dispersion. The free-carrier density N_c is determined by the rate equation⁸³:

$$291 \quad \frac{\partial N_c(z, t)}{\partial t} = \frac{\beta_{TPA} |A(z, t)|^4}{2h\nu_0 A_{eff}^2} - \frac{N_c(z, t)}{\tau_c}, \quad (2)$$

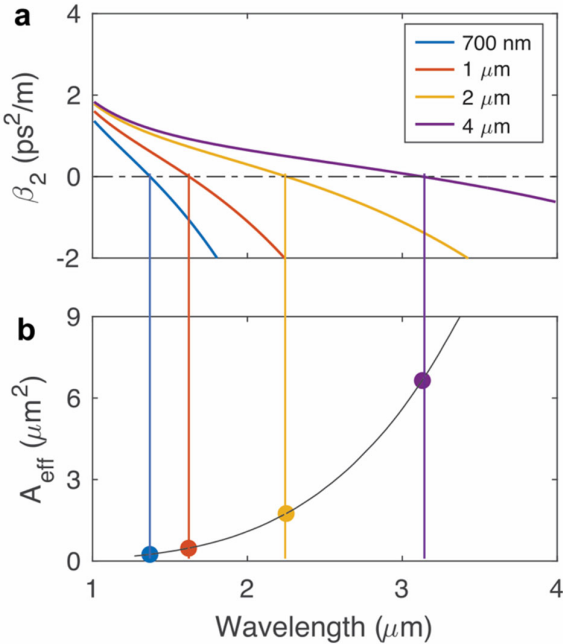
292 where τ_c is the estimated carrier lifetime. $\gamma(z)$ represents the nonlinear parameter and $\gamma_1 = d\gamma/d\omega \approx \gamma/\omega_0$. For semiconductors,
 293 it is necessary to include both the Kerr and two-photon absorption (TPA) contributions to the nonlinearity as: $\gamma = k_0 n_2 / A_{eff} +$
 294 $i\beta_{TPA}/2A_{eff}$, where β_{TPA} is the TPA coefficient, n_2 is the nonlinear refractive index and A_{eff} is the effective mode area. The
 295 GNLSE also includes Raman effects, and both the electronic and vibrational contributions are included in the response function
 296 as⁸⁴:

$$297 \quad R(t) = (1 - f_R)\delta(t) + f_R h_R(t), \quad (3)$$

298 where f_R represents the fractional contribution of the delayed Raman response given by the function $h_R(t)$.

299 The parameters for the GNLSE mainly depend on the waveguide material and pump wavelength. For example, the parameters
 300 for single-crystalline Si in the telecom band (e.g., $\lambda = 1550$ nm) are $n_2 = 5.6 \times 10^{-18}$ m²/W, $\beta_{TPA} = 5 \times 10^{-12}$ m/W, whilst in
 301 the mid-infrared (e.g., $\lambda = 3$ μ m), $n_2 = 3.98 \times 10^{-18}$ m²/W and β_{TPA} is negligible when operating beyond the TPA edge
 302 ($\lambda \sim 2.2$ μ m)¹. Moreover, the GVD and effective mode area, and thus $\gamma(z)$, depend strongly on the core diameter.

303 To illustrate this, Figure 3a shows the GVD as a function of wavelength for Si core fibres with different waist diameters. As
 304 can be seen, the zero-dispersion wavelength (ZDW) shifts quite dramatically as the core is adjusted from a few microns to sub-
 305 micrometre sizes, allowing for the dispersion to be tailored for different wavelength regions, extending across the near to mid-
 306 infrared spectral bands. For example, core diameters of <900 nm are required to access the anomalous dispersion regime in the
 307 telecom band, which is favoured for applications such as four-wave mixing (FWM) and soliton propagation. Figure 3b then plots
 308 the A_{eff} for the various ZDW crossings in Figure 3a, which corresponds to the different fibre sizes. It is clear that for the larger
 309 mid-infrared fibres, the larger mode area will result in a reduced Kerr effect, though this can be compensated somewhat by the
 310 lower nonlinear losses in this region. However, another advantage of having access to large core sizes is that they can support
 311 higher power handling, and thus the nonlinear semiconductor fibre systems can be designed to deliver more practical output
 312 powers when compared to on-chip platforms.



313

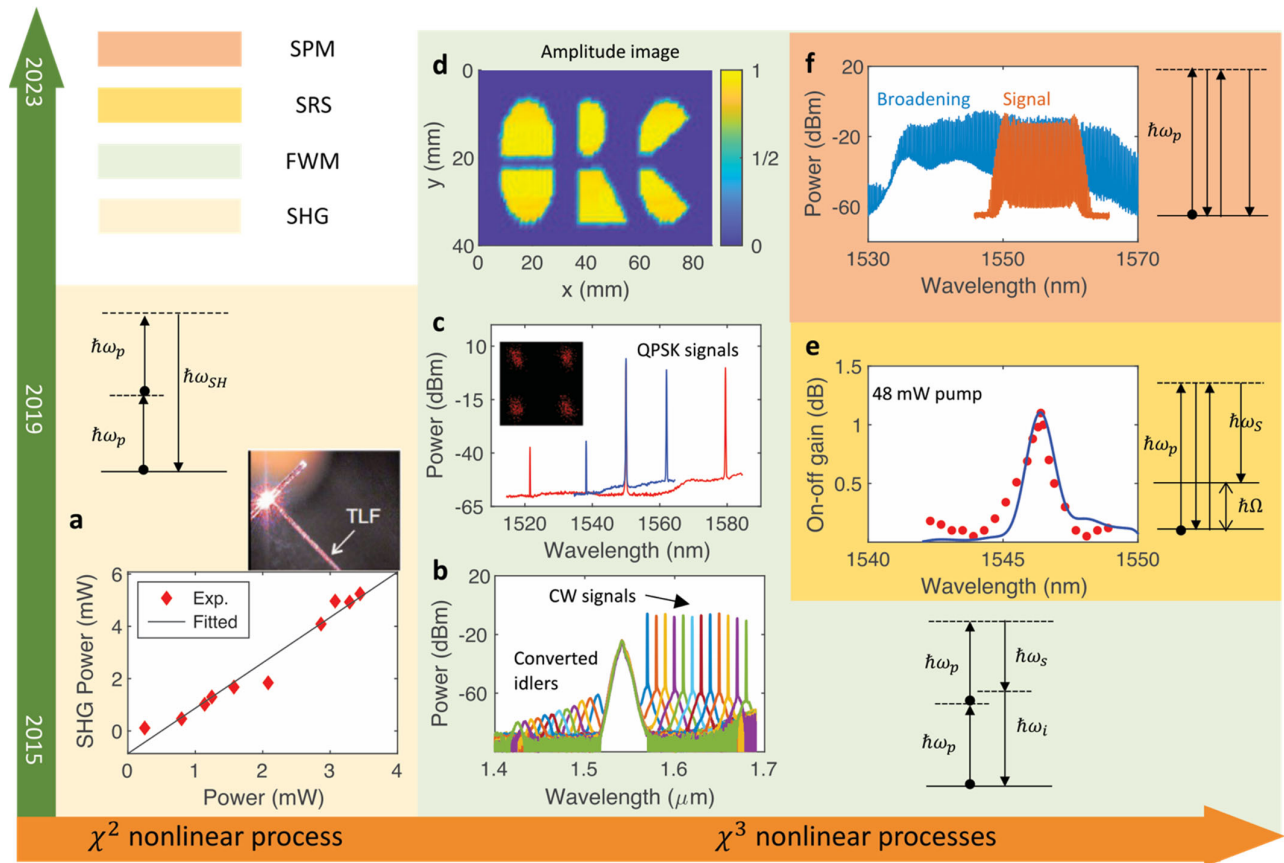
314 **Figure 3 Dispersion engineering for Si core fibres. a** Group velocity dispersion as a function of wavelength for different core diameters, as
 315 labelled in the legend. **b** Effective mode area variation for Si core fibres with different zero-dispersion wavelengths.

316 5.2 Telecom band applications

317 Owing to the ready availability of pump sources and diagnostics in the telecom band, most of the early work to characterize the
 318 nonlinear performance of the semiconductor core fibres was conducted in this wavelength region. Initially, the focus was very
 319 much on the HP-CVD fibres, as these could be produced directly with losses that were sufficiently low to observe nonlinear
 320 propagation⁸⁵. However, most of this work was dedicated to characterizing the nonlinear parameters of the deposited
 321 materials^{85,86} and, due to their relatively large core sizes, it was not possible to access the anomalous dispersion region with
 322 telecom pumps⁸⁷.

323 However, one of the first functional demonstrations of nonlinear processing was in fact made using a HP-CVD ZnSe core
 324 fibre in 2015, as shown in the timeline of Figure 4 (χ^2 process). In this first case, the fibre was used in a resonator geometry
 325 (15 μm diameter core) to observe second-harmonic generation (SHG), converting the telecom pump into the visible spectrum, as
 326 illustrated in Figure 4a⁸⁸. Here the resonator helped to enhance the light-matter interactions and obtain some phase matching
 327 between the pump and the SH. Whilst this work remains as the only demonstration of second order processes in the
 328 semiconductor core fibres, it clearly motivates the significance of obtaining high quality compound semiconductor core fibres
 329 for future applications in areas such as quantum information and imaging.

330 The remaining demonstrations summarized in Figure 4 (χ^3 processes), have been obtained using tapered MCM Si core fibres.
331 Here the tapering has been critical not only for reducing the losses, but also engineering the dispersion. As a result, these fibres
332 have allowed for the first examples of FWM processes that are used extensively in signal processing applications. In particular,
333 in 2019, a fibre with a tapered waist of ~ 915 nm (waist length 5 mm), and a loss of 2.8 dB/cm, was used as an optical parametric
334 amplifier to generate 9 dB of gain with an average pump power of only 0.63 mW, as shown in Figure 4b⁸⁹. Not only was the
335 gain higher than what had been achieved in a similar planar silicon waveguides (5.2 dB in ref. ⁸), but the optimized coupling that
336 was facilitated by the taper profile meant that this was the first example of a net parametric gain in a silicon system. Subsequently,
337 in 2022, a similar sub-micrometre Si core fibre was used for FWM-based wavelength conversion of QPSK signals, as shown in
338 Figure 4c. Consistent conversion efficiencies and constellation diagrams were recorded for both C- and L-band signal
339 wavelengths, highlighting the benefits of precise dispersion engineering for use in broadband or multi-wavelength systems⁹⁰.
340 Another noteworthy example of a FWM-based application was reported in 2023, when two identical tapered Si core fibres were
341 fabricated (waist diameters of ~ 915 nm, lengths of 1.5 cm, losses of 0.8 dB/cm) for use in an undetected-photon imaging
342 system⁷⁸. In such systems, the signal wave is used to probe the object whilst the generated idler beams, which have no interaction
343 with the object, are used for the detection. As all the components used to construct the imaging system were fiberized, it was
344 remarkably stable, allowing for a high degree of spatial and phase correlation, as illustrated by the amplitude image shown in
345 Figure 4d. Moreover, the use of FWM allows for great flexibility in the positioning of the signal and idler waves, which could
346 potentially span both near and mid-infrared regions, presenting a unique advantage over the more traditional bulk imaging
347 systems employing χ^2 materials^{91,92}.



348

349 **Figure 4 Summary of nonlinear semiconductor core fibre devices in the telecom band.** a Second-harmonic generator. b Optical parametric
 350 amplifier based on FWM. c FWM wavelength convertor, with inset showing a retrieved constellation diagram at the 1521 nm wavelength. d
 351 FWM-based undetected-photon imaging system. e Raman amplifier. f Parametric mixer based on SPM as a broadband comb source. d Repro-
 352 duced from ref. ⁷⁸ with the permission from Chinese Laser Press. f Reproduced from ref. ⁹³ with the permission from Springer Nature.

353 Beyond FWM, Raman scattering is another important χ^3 nonlinear process that can also be used for optical amplification or
 354 new wavelength generation. In fact, Raman scattering was the first nonlinear process to be demonstrated in planar silicon
 355 waveguide systems as, unlike FWM, it does not require phase matching⁹⁴. However, owing to the relatively short lengths and
 356 high losses (~ 2 dB/cm) of the early tapered Si core fibres, Raman scattering was only recently observed in 2021. The
 357 demonstration was achieved by producing a Si core fibre with a core diameter of only 750 nm over a length of ~ 2 cm, and a
 358 reduced loss of ~ 1 dB/cm⁹⁵. A maximum Raman gain of 1.1 dB was observed with a CW pump power of only 48 mW, as shown
 359 in Figure 4e, comparable to planar systems with similar waveguide dimensions and pumps. Further simulations conducted using
 360 the GNLSE (Eq. (1)) have indicated that the Raman gain could be increased substantially (up to 6 dB) by simply increasing the
 361 fibre length to ~ 10 cm, highlighting the importance of obtaining longer tapered fibres in future work.

362 As a final demonstration, Figure 4f shows self-phase modulation (SPM) broadening of a comb source, of potential use in
 363 wavelength division multiplexing systems. This work was reported in 2022, using a Si core fibre that was fully integrated with

364 standard single mode fibres (SMF), as will be discussed in Section 5. Significantly, the simple use of SPM in this work helped
365 to preserve the important comb features such as spectral flatness, low noise levels, narrow tone bandwidths and high tone powers,
366 that are important for many communication applications⁹³. Thus, this serves to highlight the value of having access to the full
367 range of nonlinear processes to ensure maximal impact of the fibre systems.

368 **5.3 Applications into the mid-infrared**

369 Inspired by the successes in the telecom band, recent attention has shifted to explorations in the mid-infrared spectral region
370 where there are important applications in gas sensing⁹⁶, environmental monitoring⁹⁷, and medical diagnostics⁹⁸. In this regard,
371 the semiconductor core fibres offer a unique advantage over their planar waveguide counterparts in that, as indicated in Figure 3a,
372 they can be readily produced with larger, few micron-sized cores, that are well-suited to mid-infrared operation. Moreover, when
373 using the Si core fibres, the TPA parameter becomes negligible when operating at wavelengths $> 2.2 \mu\text{m}$, so that short-pulsed
374 lasers with high peak powers can be used to boost the nonlinear performance⁹⁹. Although chalcogenides have long been a popular
375 material of choice for mid-infrared nonlinear fibre optics owing to their broad transmission windows and low transmission losses
376 ($<5 \text{ dB/m}$)¹⁰⁰, they typically have lower nonlinear refractive indices than crystalline semiconductors (by around an order of
377 magnitude)¹ and they can also suffer from stability and photosensitivity issues¹⁰⁰. Thus, the silica-clad semiconductor core fibres
378 offer an advantage over chalcogenide fibres in terms of their high nonlinearity, strong durability, and mechanical robustness,
379 particularly for systems requiring short fibre lengths and high operating powers.

380 By further adapting the fibre tapering, it has been possible to produce long lengths of micron-sized Si core fibres with losses
381 $<1 \text{ dB/cm}$. Access to such low losses becomes even more important for the longer wavelength pumps due to the lower intensities
382 associated with the larger A_{eff} , also seen in Figure 3b. Figure 5a shows results of simulations conducted using the GNLSE (Eq.
383 (1)) indicating how the FWM conversion bandwidth depends on the core diameter for a pump positioned at $\sim 2 \mu\text{m}$, a typical
384 wavelength for a thulium doped fibre laser system¹⁰¹. The maximum bandwidth exceeds 1200 nm for a waist diameter of 1.6 μm .
385 By fabricating a tapered Si core fibre with the target core diameter over a length of 4 cm, and a loss of 0.5 dB/cm, an experimental
386 conversion bandwidth of 690 nm was obtained, as shown in Figure 5b. It should be noted that the measured bandwidth was only
387 limited by the tuneability of the available seed signal and the low pump power (6 dBm average power, corresponding a peak
388 power of $\sim 2 \text{ W}$).

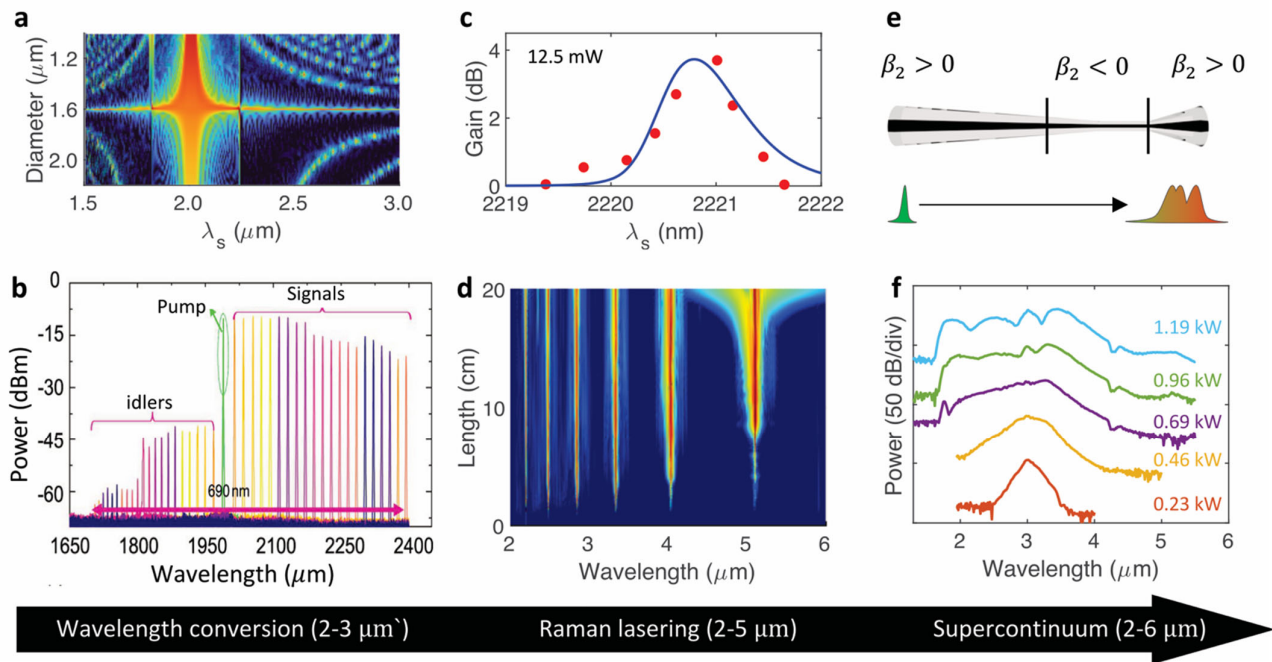
389 By further extending the Si core fibre length to 6 cm and reducing the transmission loss to 0.2 dB/cm, Raman scattering and
390 amplification was observed with the same 2 μm pump source⁷². A maximum time-averaged Raman gain of 3.7 dB was obtained
391 for a pump power of 12.4 mW, as shown in Figure 5c, corresponding to a peak on-off Raman gain of $\sim 30 \text{ dB}$ when considering

392 the 125 ps pulse duration. Interestingly, because of the long pulse duration, nonlinear absorption at the 2 μm pump wavelength
393 is still significant, so that increasing the pump power will not result in a significant increase in the Raman gain. However, if the
394 pump could be shifted to a longer wavelength of 2.2 μm , where TPA is negligible, a significant increase in the gain can be
395 achieved. In this case, the first Stokes wave can grow sufficiently strong to act as a seed for higher order cascaded processes, as
396 displayed via the GNLSE simulations in Figure 5d. Thus, this worked illustrated the potential for Si core fibres to form the basis
397 of tuneable wavelength sources across the 2-5 μm range.

398 Beyond FWM wavelength conversion and Raman amplification, tapered Si core fibres have also been used for
399 supercontinuum generation - a process that draws on all the full suite of nonlinear effects. Figure 5e shows a tapered fibre profile
400 that was designed especially for the observation of supercontinuum generation at mid-infrared wavelengths. Particularly, the
401 fibre had an asymmetric profile such that the long-tapered input facilitated optimum coupling into the waist, whilst the short
402 output taper ensured minimal interaction of the newly generated long wavelength light with the lossy silica cladding⁷⁵. The
403 fabricated fibre was designed for a pump wavelength of 3 μm and had a waist diameter of 2.8 μm over a length of only 1.7 mm
404 to keep the cladding absorption to a minimum, with input and output core diameters of 10 μm . The resulting output spectra
405 obtained with input pump powers up to ~ 10 mW is shown in Figure 5e. The broadest spectrum spans almost two octaves (~ 1.7),
406 covering wavelengths from 1.62 μm to 5.34 μm , with the red edge of the spectrum extending well beyond what had been
407 previously achieved in any planar silicon-on-insulator (SOI) waveguide by around 2 μm .

408 Despite these achievements, mid-infrared transmission of the Si fibres will always be limited by the core transmission to
409 wavelengths up to ~ 8 μm . However, Ge materials can offer extended transmission up to ~ 14 μm as well as higher nonlinear
410 coefficients, so that Ge core fibres present as an interesting alternative for long wavelength source generation. As previously
411 mentioned, although Ge core fibres have been successfully tapered and laser processed, so far their nonlinear performance has
412 yet to be fully characterized. However, Ge core fibres have been used to demonstrate a nonlinear response from the core, firstly
413 through Raman emission at 6.8 μm when pumped with a QCL source at 5.6 μm ¹⁰², and secondly via the observation of detuning
414 oscillations in the frequency-resolved response of pump probe experiments conducted at 4.6 μm ¹⁰³. Thus, these works highlight
415 the potential of these fibres for mid-infrared source generation provided the losses can be reduced from the current levels, to
416 values closer to what has been achieved in silicon (see Table 1). Alternatively, the SiGe core fibres also offer potential to extend
417 the operational window beyond the pure Si cores, but still retain access to the more advanced production methods that have
418 enabled low losses. Thus, further work to optimize the core dimensions and losses of the SiGe core fibres for wavelengths > 2 μm
419 could help open new routes to achieving high quality mid-infrared semiconductor fibres. We note that as the operating

420 wavelengths get longer, reducing the losses due to the silica cladding may not be possible by simply altering the fibre design as
 421 in ref. ⁷⁵, and alternative cladding materials may need to be considered that offer lower losses in the mid-infrared⁴⁸.



423 **Figure 5 Summary for nonlinear semiconductor core fibre demonstrations in the mid-infrared regime.** **a** Simulation results of FWM
 424 efficiency as a function of signal wavelength and Si core diameters when pumped at 2 μm. **b** Measured mid-infrared FWM output spectra for
 425 a Si core fibre with a core diameter of 1.6 μm over a length of 4 cm. **c** Measured average on-off Raman gain in a Si core fibre with a waist
 426 diameter of 1.59 μm and a length of 6 cm. **d** Simulation results of cascaded Raman scattering when pumped at 2.2 μm wavelength with a peak
 427 power of 20 W. **e** Si core fibre profile designed for mid-infrared supercontinuum generation. **f** Measured output supercontinuum generation
 428 spectra obtained in an asymmetric tapered Si core fibre for different pump peak powers, as labelled. **a** & **b** Reproduced from ref. ¹⁰¹ with the
 429 permission from AIP Publishing. **c**, **d** & **f** Reproduced from ref. ^{72,75} with the permission from Springer Nature.

430 6 Perspective and Outlook

431 6.1 New materials – second order nonlinear systems

432 As described in Section 2, the MCM has become the focus of semiconductor fibre fabrication, owing to its scalability and
 433 alignment with commercial fibre production methods. Through adapting the post-processing methods, there is a clear route for
 434 the optimization and nonlinear application of the SiGe and Ge fibres in the near future, which will be important for applications
 435 within the longer mid-infrared wavelength region. It is worth noting that although there is currently considerable interest within
 436 the chip-based community in other third order nonlinear materials such as silicon nitride and titanium dioxide¹⁰⁴, principally as

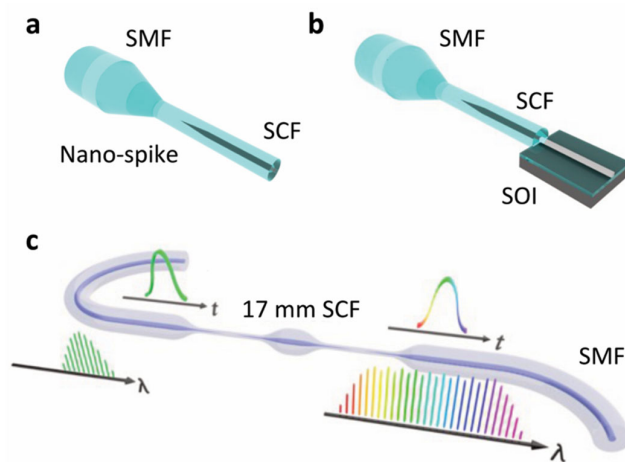
437 they offer broad transmission windows from the visible to the mid-infrared, these are not compatible with MCM due to their
438 high melting temperatures. Thus, in terms of expanding the material systems, focusing on the development of second order core
439 materials, e.g., GaAs⁵³ and ZnSe⁵⁴, which also offer very broad transparency windows, presents the most fruitful approach for
440 the fibre work. However, so far, the MCM fibre quality has precluded direct measurement of important optical properties, such
441 as loss or nonlinearities, so that their maturity lags significantly behind the Si core fibres. It is, therefore, perhaps best here to
442 muse instead on what is possible, what seemingly is not possible, and why.

443 From the perspective of scalable fabrication, the requirement of a (typically metal) flux to permit thermal drawing in a glass
444 cladding mandates subsequent laser post-processing to segregate the nonlinear phase from metal. However, the more phase-pure
445 the nonlinear phase, the more likely the issues of volatility and incongruency that restricted their “conventional” MCM
446 fabrication. Although the high nonlinearity and tight mode confinement of the small core semiconductor fibres means that only
447 short phase-pure sections are needed for device functionality, it may be that lateral segregation (across the core cross-section),
448 rather than longitudinal, ends up as the more practical approach. This would be somewhat analogous to a planar waveguide
449 inside a glass-clad, fibre core. However, one advantage of this approach is that laser drilling holes through the cladding could
450 then enable electrical contacts to the metal (flux) phase, thus marrying χ^2 (or χ^3) nonlinear functionality, with direct bandgap
451 light emission all within a fibre platform. Whatever the final geometry, it is clear that having access to the broadest range of in-
452 fibre semiconductor materials is important for maximizing the application potential of this platform.

453 **6.2 Fibre connections – all-fibre integrated systems**

454 One of the main motivations for incorporating the semiconductor materials inside fibre geometries is the potential for seamless
455 integration with more conventional fibre networks. Despite this promise, most of the nonlinear demonstrations described in
456 Section 4 have been achieved using free space coupling systems, either using lenses or tapered lens fibres. Although the lensed
457 fibre approach allows for high coupling efficiencies (~ 3 dB/facet) and ensures the systems are essentially all-fiberized, the
458 connection points are not robust and still require precise alignment, limiting their portability. To achieve robust and user-friendly
459 connections between SMF and the semiconductor core fibres, in 2019, a nano-spike coupler approach was developed to facilitate
460 the connection, as shown in Figure 6a. Here, the spike helps to transition the light from the low index glass core to the high index
461 semiconductor through a graded mode conversion. The benefits of this coupler are: (i) improved mode matching between the
462 fibres, (ii) reduced reflection losses, and (iii) robust splicing between the silica cladding structures. Although the first fabricated
463 nano-spike coupler had a loss of ~ 3 dB⁷³, simulations of the mode coupling predict that losses of ~ 1 dB are achievable with the
464 correct fibre designs.

465 Moreover, these semiconductor fibre couplers could also serve as an interface between SMF and planar semiconductor
466 waveguides, provided the fibre and planar core dimensions were well matched, as illustrated in Figure 6b. Currently, the main
467 challenge to the fabrication of these couplers is that the outer cladding diameter of the nano-spike region is very small ($\sim 30\ \mu\text{m}$),
468 so the mechanical strength of coupler is not very strong⁷³. Thus, there is work on-going to fabricate fibres with more suitable
469 core/cladding ratios to better match both the fibre-to-fibre and fibre-to-chip interfaces. However, it is also possible to make use
470 of capillary sleeves to support the fibres in the coupling region, similar to the role of a conventional splice protector. Using such
471 a scheme, in 2022 a Si core fibre was pig-tailed on both ends with SMF fibres to produce an all-fibre parametric mixer, as shown
472 schematically in Figure 6c. The fabricated Si fibre had a core diameter of $1.1\ \mu\text{m}$ over a length of $1.7\ \text{cm}$, which was used to
473 broaden the bandwidth of a frequency comb source through SPM, as shown in Figure 4f. As previously mentioned, thanks to the
474 relatively short length of the Si core fibre section, the key performance metrics of the comb source required for
475 telecommunications applications was preserved. However, the high nonlinearity of the silicon section resulted in a tripling of
476 the comb bandwidth with only modest telecom pump powers. Thus, this work demonstrates an important step towards practical
477 all-fibre integrated nonlinear semiconductor photonic systems.



478
479 **Figure 6 Fibre integrated nonlinear semiconductor photonic systems.** **a** Schematic of a nano-spike coupler to enhance the coupling between
480 SMF and a semiconductor core fibre (SCF). **b** Concept of using a Si core fibre with a nano-spike coupler to transition light from SMF to a Si-
481 on-insulator (SOI) waveguide on-chip. **c** Schematic of an all-fibre integrated Si core fibre parametric mixer for comb source broadening in the
482 telecom band⁹³.

483 7 Conclusion

484 This paper has reviewed developments in the fabrication and post-processing of semiconductor core fibres that are emerging as
485 robust and scalable platforms for nonlinear applications. These fibres are now primarily fabricated by the molten core drawing

486 method, owing to its capacity for high speed and high-volume production, as well as its suitability for adaption to both unary
487 and compound semiconductor core materials. When combined with post-processing methods such as laser heating and tapering,
488 low loss fibres can be produced over lengths of several centimetres, with core dimensions that can be tailored from a few microns
489 down to hundreds of nanometres. Through precise control of the longitudinal profile to engineer the dispersion and enhance the
490 coupling, it has been possible to observe a wide range of nonlinear effects, starting in the telecom band and extending up into
491 the mid-infrared regime. Moreover, with access to new materials and integration schemes starting to emerge, it is expected that
492 these fibres will continue to expand the avenues of exploration for advanced fibre systems in terms of the available nonlinear
493 processes and transmission windows. Although the semiconductor core fibres will not replace their chip-based waveguide
494 counterparts, they will help to extend the rich landscape of nonlinear semiconductor photonics.

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500 **Competing interests**

501 The authors declare no conflicts of interest.

502 **Author contributions**

503 M.H., J.B, and A.C.P. prepared the manuscript text together. A.C.P and M.H were responsible for the figure preparation and final editing. All
504 authors reviewed the manuscript.

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670 **Figure and Table legends**

671 Figure 1

672 Title: **A schematic of the molten core method for thermally drawing glass-clad semiconductor core fibres.**

673 Legend: The crystalline core is melted and encapsulated by the viscous glass cladding during the drawing process. A coating system is used to
674 improve the mechanical strength of the produced as-drawn semiconductor core fibres.

675 Figure 2

676 Title: **Post-processing techniques for semiconductor core fibres.**

677 Legend: **a** Schematic of laser processing. **b** Measured lattice spacing from X-ray diffraction as a function of position along as-fabricated and
678 laser processed Si core fibres⁶². The inset shows the X-ray diffraction measurement setup. **c** Optical images of (i) an as-drawn SiGe core fibre,
679 (ii) a laser recrystallized SiGe core and (iii) a Ge-rich grating formed within the SiGe core. **d** Schematic of the tapering procedure, showing
680 how multiple taper steps can be employed. **e** Images of longitudinal taper profiles for starting fibre core diameters of (i) 5.6 μm , (ii) 2.7 μm ,
681 and (iii) 1.3 μm . **f** Measured X-ray diffraction along the as-fabricated and tapered Si core fibre. The insert shows an image of the smooth
682 surface of an etched Si core fibre waist⁶⁹. **c** Reproduced from ref. ⁶³ with the permission from Springer Nature. **e** Reproduced from ref. ⁷⁰ with
683 the permission from Optica Publishing Group.

684 Figure 3

685 Title: **Dispersion engineering for Si core fibres.**

686 Legend: **a** Group velocity dispersion as a function of wavelength for different core diameters, as labelled in the legend. **b** Effective mode area
687 variation for Si core fibres with different zero-dispersion wavelengths.

688 Figure 4

689 Title: **Summary of nonlinear semiconductor core fibre devices in the telecom band.**

690 Legend: **a** Second-harmonic generator. **b** Optical parametric amplifier based on FWM. **c** FWM wavelength convertor, with inset showing a
691 retrieved constellation diagram at the 1521 nm wavelength. **d** FWM-based undetected-photon imaging system. **e** Raman amplifier. **f** Parametric
692 mixer based on SPM as a broadband comb source. **d** Reproduced from ref. ⁷⁸ with the permission from Chinese Laser Press. **f** Reproduced
693 from ref. ⁹³ with the permission from Springer Nature.

694 Figure 5

695 Title: **Summary for nonlinear semiconductor core fibre demonstrations in the mid-infrared regime.**

696 Legend: **a** Simulation results of FWM efficiency as a function of signal wavelength and Si core diameters when pumped at 2 μm . **b** Measured
697 mid-infrared FWM output spectra for a Si core fibre with a core diameter of 1.6 μm over a length of 4 cm. **c** Measured average on-off Raman
698 gain in a Si core fibre with a waist diameter of 1.59 μm and a length of 6 cm. **d** Simulation results of cascaded Raman scattering when pumped
699 at 2.2 μm wavelength with a peak power of 20 W. **e** Si core fibre profile designed for mid-infrared supercontinuum generation. **f** Measured
700 output supercontinuum generation spectra obtained in an asymmetric tapered Si core fibre for different pump peak powers, as labelled. **a** & **b**
701 Reproduced from ref. ¹⁰¹ with the permission from AIP Publishing. **c**, **d** & **f** Reproduced from ref. ^{72,75} with the permission from Springer
702 Nature.

703 Figure 6

704 Title: **Fibre integrated nonlinear semiconductor photonic systems.**

705 Legend: **a** Schematic of a nano-spike coupler to enhance the coupling between SMF and a semiconductor core fibre (SCF). **b** Concept of using
706 a Si core fibre with a nano-spike coupler to transition light from SMF to a Si-on-insulator (SOI) waveguide on-chip. **c** Schematic of an all-
707 fibre integrated Si core fibre parametric mixer for comb source broadening in the telecom band⁹³.

708 Table 1

709 Title: **Measured optical loss (unit: dB/cm) for semiconductor core fibres that have been used for nonlinear demonstrations. Top rows**
710 **are for the telecom wavelength of 1.55 μm . Bottom rows for selected mid-infrared wavelengths of $\ast\sim 2.5 \mu\text{m}$ and $\dagger\sim 5 \mu\text{m}$.**

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712