

Unravelling the photoactivity of metal-loaded TiO₂ for hydrogen production: Insights from a combined experimental and computational analysis

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ABSTRACT

Despite being the most employed material for photocatalytic hydrogen generation, TiO₂ suffers limitations such as a high rate of electron-hole recombination and poor light absorption in the visible spectrum. Among the various strategies developed to overcome these drawbacks, combining TiO₂ with a metal co-catalyst emerged as one of the most promising. In this study, we integrated experimental findings, advanced characterization techniques, and computational methods to shed light on how different noble metals influence the enhancement of the photocatalytic activity of TiO₂. Among the tested noble metal co-catalysts, the hydrogen production rate under UV and visible light irradiation followed the trend Pt > Au ≈ Pd > Ag > bare TiO₂, with Pt-decorated TiO₂ exhibiting a hydrogen production rate of 28 mmol/h g. The noble metals were found to significantly suppress the electron-hole recombination rate compared to bare TiO₂. Upon photodeposition, Pd and Pt formed the smallest nanoparticles with average sizes of 13.4 nm and 4.1 nm, respectively. Computational analyses were conducted to rationalize the difference in nanoparticle sizes by analyzing the binding and cohesive energies of the metal clusters on the TiO₂ surface. Additionally, calculations demonstrated the strong interaction of Pt, Au, and Pd nanoclusters with adsorbed hydrogen, with Pt achieving the closest-to-zero Gibbs free energy of hydrogen adsorption and displaying the most polar interaction with hydrogen. These findings align closely with the observed hydrogen production rates, where UV/Vis-driven hydrogen production is governed by the coupling of hydrogen radicals on the co-catalyst surface, while visible-light-driven production is limited by charge carrier lifetimes.

1. Introduction

Due to the imminent depletion of conventional energy sources, photocatalytic production of hydrogen is a critical challenge in the pursuit of clean and renewable energy [1–3]. Currently, the application of solar hydrogen production from renewable resources is limited, with the primary method being water electrolysis [4,5]. However, among the

techniques for hydrogen production, sacrificial photoreforming stands out as one of the most promising, attracting significant global research interests [6,7].

Titanium (IV) oxide (TiO₂) is a semiconductor recognized for its excellent photostability and highest sustained photocatalytic activity in these processes. However, one of the main limitations of TiO₂ is its high rate of photogenerated electron-hole pair recombination, which

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significantly reduces its photocatalytic efficiency. Indeed, high recombination rates, with a charge carrier lifespan of around 30 ns, result in insufficient time for further chemical reactions to occur. Another disadvantage of TiO_2 is its wide band gap (3.0–3.2 eV), which limits its light absorption primarily to the UV range [7]. Since UV irradiation accounts for only 4 % of the solar spectrum, a major challenge for TiO_2 -based photocatalytic processes is the limited utilization of solar energy. To address these drawbacks, previous studies have aimed to enhance the semiconductor activity by reducing the electron-hole recombination rate and extending the light adsorption to the visible range [8,9], typically through the combination of the semiconductor with a co-catalyst [10,11]. Among the various strategies employed, loading TiO_2 with noble or non-noble metals is one of the most promising approaches [12–14]. Several reviews have been published that discuss the behaviour of metals on the catalyst surface, offering possible explanations for the different observed activity [15,16]. For instance, Kumaravel et al. [17] reported that metal dopants can enhance the charge carrier separation and the spectral response of TiO_2 through various mechanisms, including the surface plasmon resonance (SPR) effect, the electron trap mechanism, the generation of gap states from interaction with TiO_2 valence band (VB) states, and the introduction of new energy levels. Defect engineering is a powerful technique for tuning the properties of a material to optimize its performance for specific functions [18]. It has been suggested that doping a photocatalyst with a metal of higher work function can enhance charge separation, thereby improving the photocatalyst's efficiency [19]. The role of metal loading is to primarily enhance charge carrier separation in the photocatalyst. This is achieved by introducing metal traps, which extend the lifetimes of these catalytically active charge carriers, resulting in more successful redox reactions on the surface. The formation of the Schottky barrier upon metal deposition on the semiconductor surface further promotes charge separation, by facilitating electron transfer from the bulk semiconductor to the metal sites.

From a theoretical perspective, DFT studies have proven to be powerful tools to select and design photocatalysts for efficient hydrogen production [20], as well as interpreting experimental results. However, these calculations require certain assumptions. The process of interest is typically studied on a small fragment of the photocatalyst's surface, assuming it has a periodic structure. For example, Assadi et al. conducted a theoretical study to explain the enhancement of anatase photoactivity through copper doping [21]. According to their computational results, enhancement was attributed to electronic interactions, facilitated by charge transfer and the formation of inter-bandgap states. Wen et al. utilized DFT to explore the photocatalytic behaviour of iron-doped photocatalysts [22]. That study revealed that the Fe^{3+} ions on the TiO_2 surface can provide an intermediate interfacial transfer pathway for the photogenerated charges, thus increasing the photocatalytic activity of TiO_2 . However, purely computational studies should be treated with extreme caution [23] and their combination with appropriate characterization techniques is crucial to fully understand how metals modify the photocatalytic properties of TiO_2 [7,24]. For this purpose, some authors have reported the possibility of determining the correlations between the material properties and the photocatalytic performance based on theoretical calculations and/or systematic characterizations [25]. However, despite the extensive body of research on TiO_2 -based photocatalysts, the primary factors influencing the variation of photoactivity depending on the metal co-catalyst remain unclear [12,17,26]. In this study, we propose an innovative synergistic approach that integrates state-of-the-art characterization techniques with advanced computational methodologies. This approach systematically investigates the role of metal co-catalysts in enhancing the photoactivity of TiO_2 for hydrogen production. By integrating experimental and theoretical insights, this framework offers a comprehensive model for designing photocatalytic systems, leveraging the interplay between semiconductors and transition metals.

2. Materials and methods

2.1. Materials

TiO_2 (anatase) nanopowder (<25 nm particle size, 99.7 % trace metal basis), palladium (II) chloride (PdCl_2 , ≥ 99.9 %), silver nitrate (AgNO_3 , ACS reagent, ≥ 99.0 %), gold (III) chloride trihydrate ($\text{AuCl}_3 \cdot 3\text{H}_2\text{O}$, ≥ 99.9 % trace metals basis), chloroplatinic acid hexahydrate ($\text{H}_2\text{PtCl}_6 \cdot 6\text{H}_2\text{O}$, ACS reagent, ≥ 37.50 % Pt basis) and methanol (99 %) were purchased from Sigma Aldrich. All reagents were used as received. Doubly glass-distilled water was used throughout this study.

2.2. Photocatalysis experiments

Photocatalytic experiments were conducted batchwise in an annular quartz batch reactor, placed in a quartz lamp jacket and fitted with high-pressure mercury vapour lamps (Nominal power $P = 400$ W) and a daylight lamp (Nominal power $P = 150$ W) both by Photochemical Reactors Ltd, as described in previous reports [27]. The effective irradiances of the medium pressure mercury vapour lamp were $1.17 \cdot 10^{-6} \text{ E s}^{-1}$ (302 nm), $3.19 \cdot 10^{-6} \text{ E s}^{-1}$ (313 nm), $3.41 \cdot 10^{-7} \text{ E s}^{-1}$ (334 nm), $6.97 \cdot 10^{-6} \text{ E s}^{-1}$ (366 nm) in the UV range, and $3.04 \cdot 10^{-5} \text{ E s}^{-1}$ (406 nm), and $6.62 \cdot 10^{-5} \text{ E s}^{-1}$ (436 nm) in the visible range. The effective irradiances of the daylight lamp were $2.84 \cdot 10^{-5} \text{ E s}^{-1}$ (420 nm), $5.63 \cdot 10^{-5} \text{ E s}^{-1}$ (535 nm), $8.14 \cdot 10^{-5} \text{ E s}^{-1}$ (585 nm), and $4.28 \cdot 10^{-5} \text{ E s}^{-1}$ (670 nm), all in the visible range. In a typical experimental run, 150 mg of TiO_2 nano powder was suspended in a 2.5 M methanol aqueous solution (total volume = 0.30 L) at natural pH. To prevent dissolved oxygen from reacting with photogenerated electrons, nitrogen gas was bubbled through the solution for 40 min before the photocatalytic experiments at a flow rate of 0.30 L min^{-1} . The temperature of the system, controlled with a thermostat, was fixed to 25°C . The composite materials were prepared through in-situ photodeposition of the metals onto the anatase surface, by adding selected concentrations of metal ions ($\text{Metal}/\text{TiO}_2 = 7.5 \text{ % mol}$) to the mixture in the form of salts. Specifically, depending on the metal, a selected amount of metallic salts was employed. To assess the activity under visible light irradiation, the visible lamp was used, after photodeposition of the metals in the presence of the UV lamp. In some cases, the photocatalyst was recovered at the end of the run, as previously described [28].

To evaluate the hydrogen production rate, gaseous samples were periodically collected using Tedlar gas sampling bags (1 L) at different reaction times. These samples were analysed using a gas chromatograph (Trace 1310), equipped with a HS-Q column (2 m, 60/80 mesh) and a TCD detector using argon as carried gas. The pH of the solution was monitored through an Orion 420 p pH-meter (Thermo).

2.3. Characterization techniques

A variety of advanced characterization tools were utilized in order to examine the chemical and physical properties of the photocatalyst synthesized in this study, including X-ray Diffraction (XRD), X-ray photoelectron spectroscopy (XPS), Scanning Transmission Electron Microscopy coupled with Energy Dispersive Spectroscopy (STEM-EDX), Brunauer-Emmett-Teller (BET) analysis, UV-Vis Diffuse Reflectance Spectroscopy (UV-DRS), Electrochemical Impedance Spectroscopy (EIS), and Photoluminescence (PL).

Structural properties of the photocatalysts were analysed using Empyrean Panalytical X-ray Diffractometer in 2θ range 20° – 80° angle with $\text{Cu K}\alpha$ radiation ($1.5406 \text{ \AA}$) in 0.02 step size using the program reflection transmission spinner (RTS). X-ray tension was 20 kV and current was 15 mA.

A PHI VersaProbe 5000 Scanning X-ray Photoelectron Spectrometer equipped with a $\text{MgK}\alpha$ X-ray source (1100 eV) was used to conduct the XPS analysis. The power of the X-ray beam was 50.17 W. The analysis was conducted using a E-neutralizer (1 V) and a I-neutralizer (0.11 kV

Ar⁺ ion). The samples were etched in-situ with a 3 kV Ar beam for 1 min to remove impurities from the top 20 nm layer prior to the XPS measurements. The analysis is done using Aventure software and all peaks are referenced to C 1s.

Transmission electron microscopy (TEM) and STEM-EDX analyses were carried out at 200 kV utilizing Tecnai G2 transmission electron microscopes. First, the powder was suspended in water and subjected to ultrasonic treatment. Three microliters of the suspension were placed twice in a row on a Tedpella-provided 400-mesh Cu grid. After that, water was evaporated at ambient temperature. To compute the particle size distribution, ImageJ was utilized.

Using N₂ as the adsorbate, the Quantachrome NOVA 2000e surface area and pore size analyzer and the Anton Paar Nova 800 physisorption analyzer were used to assess the adsorption/desorption isotherms of the catalysts. The samples were degassed under vacuum for 6 h at 250 °C before analysis. The multipoint BET technique was applied in the low-pressure regime (desorption points $P/P_0 = 0-0.35$) to obtain the specific surface area. The pore size distribution was calculated by fitting the adsorption curve over the whole pressure range to the Barrett-Joyner-Halenda (BJH) model.

The optical properties of the investigated photocatalyst samples were assessed using Shimadzu Europe-UV-2600-room temperature photoluminescence (PL) spectra, which were obtained with a PerkinElmer LS-55 at $\lambda_{ex} = 400$ nm, 4.5 nm emission and excitation slits, and a 100 nm/min scan speed.

The UV–vis diffuse reflectance spectra (DRS) were obtained using a UV–vis spectrophotometer (Shimadzu UV-2600) equipped with an integrating sphere attachment and BaSO₄ as the reflectance standard, in addition to the photoluminescence. The UV–vis DRS were recorded in the 250–750 nm range.

A PGSTAT302 N Potentiostat connected to a three-electrode system was used for measurements of electrochemical impedance spectroscopy (EIS) of the as-obtained samples. The EIS was performed in dark and under UV (370 nm) illumination in the frequency range of 0.1 Hz–10 kHz with an applied potential of 0.9 V vs NHE. The photoelectrochemical measurement was carried out using a Metrohm Autolab with the sample deposited on a fluorine doped tin oxide (FTO) as the working electrode, platinum electrode as the counter electrode, and a saturated silver chloride (Ag/AgCl) electrode as the reference electrode. A 0.5 M Na₂SO₄ electrolyte solution (pH 5.6) was used for the analysis. The working electrode was prepared by suspending 4 mg of the sample in 4 mL of 2-propanol. Followed by 15 min of ultrasonic treatment. The FTO glass surface (1 cm × 2.5 cm part) was covered with 10 μ L of each suspension using a micropipette, allowing it to dry completely before applying additional layers. This deposition was repeated 15 times, resulting in a total catalyst loading of 150 μ g on the FTO surface.

2.4. Computational methods

DFT calculations were performed using the 5.4.4 version of the Vienna ab initio Simulation Package (VASP) [29,30]. The Generalised Gradient Approximation (GGA) with the Perdew-Burke-Ernzerhof functional (PBE) was employed to obtain exchange-correlation energies in all calculations [31]. The Projector Augmented Wave (PAW) method was used to describe the inner electrons [32,33]. The cutoff energy of the plane-wave basis set was set to 520 eV. Single atoms and isolated clusters were calculated in boxes 13×14×15 Å. van der Waals interactions were represented by including Grimme's DFT-D3 dispersion correction method, along with the Becke-Johnson BJ damping method for interatomic distances [34,35]. Structural optimization allowed for symmetry breaking to ensure a more accurate representation of the system. Interactions within the semiconductor material were incorporated to correct the electron self-interaction by applying the Coulomb interaction U and the electronic exchange J following Dudarev's approach of GGA + U [36]. A U value of 2 was applied to the Ti 3d orbitals as suggested by Metiu and Hu [37]. Fig. S1 displays the impact

of the value of U on the Gibbs free energy of hydrogen adsorption on metal clusters on TiO₂ and confirms a consistent trend for different U values.

The lattice parameters for the anatase structure were optimised to $b = 7.49$ Å and $c = 10.47$ Å. The catalyst was modelled by a 2×2×3-extended primitive cell (Ti₂₄O₄₈) to mimic the bulk lattice structure. The system was periodic in two dimensions, forming a (101)-oriented anatase slab. A vacuum length of 20 Å was set in the non-periodic direction to avoid spurious interactions between neighbouring slabs. A single Gamma k-point (1 × 1 × 1) was deemed sufficient for these calculations. For the convergence criteria in all calculations, 10^{−6} eV was set for the self-consistent field and 0.01 eV/Å for the atomic forces. A conjugate-gradient optimization method was used to optimize the structures. The Vaspkit package was used to apply corrections to estimate the Gibbs free energy [38].

The photocatalysts were modelled using metal clusters (Pt_x, Au_x, Pd_x, Ag_x, with x ranging from 3 to 6) adsorbed on the anatase (101) surface. The adsorbed cluster structures were informed by earlier literature studies [39–42]. Each metal has several variations of optimised cluster configurations, with some clusters lying flat on the TiO₂ surface and others standing upright. The most stable optimised shapes of the adsorbed trimer clusters are reported in Fig. S2. These configurations prefer having at least two of the three atoms bonded to O_{2c} atoms on the semiconductor surface, indicating that this bonding may play a part in its stability.

The adsorption energy of a metal cluster on the TiO₂ surface (ΔE_{TM}) was defined by Equation (1), where E_{slab} is the energy of the relaxed pristine TiO₂, E_{TM_x} is the energy of a transition metal cluster (TM_x, where TM is Pt, Au, Pd, or Ag) in its most stable configuration, and E_{TM_x/TiO_2} is the energy of the TM_x/TiO₂ structure. The more negative the adsorption energy, the stronger the binding.

$$\Delta E_{TM} = E_{TM_x/TiO_2} - (E_{slab} + E_{TM_x}) \quad (1)$$

Cohesive energy (E_{coh}) was defined as the energy released per transition metal atom in forming the adsorbed structure TM_x/TiO₂ according to the formula:

$$E_{coh} = \frac{E_{TM_x/TiO_2} - (E_{slab} + x \cdot E_{TM_1})}{x} \quad (2)$$

where E_{TM_1} is the energy of a single isolated metal atom.

To evaluate the reactivity of each metal co-catalyst, the adsorption energy of a single hydrogen atom ΔE_H was calculated using Equation (3). $E_{H/TM_x/TiO_2}$ is the energy of the catalytic system with the hydrogen adsorbed on the TM_x cluster on anatase, and E_{H_2} is the energy of an isolated diatomic hydrogen molecule.

$$\Delta E_H = E_{H/TM_x/TiO_2} - E_{TM_x/TiO_2} - \frac{1}{2}E_{H_2} \quad (3)$$

The Gibbs free energy of hydrogen adsorption (ΔG_H) was calculated using Equation (4):

$$\Delta G_H = \Delta E_H + \Delta E_{ZPE} - T \cdot \Delta S \quad (4)$$

where ΔE_{ZPE} and ΔS are the difference in zero-point energy and the difference in entropy between adsorbed hydrogen and hydrogen molecule in the gas phase, respectively. T is the temperature of 298 K.

The charge density difference (CDD) $\Delta\rho$ was calculated based on Equation (5). $\rho_{H/TM_x/TiO_2}$ is the charge density of the entire system, ρ_H is the charge density of the adsorbed hydrogen, and ρ_{TM_x/TiO_2} is the charge density of the TM_x/TiO₂ system with no adsorbate present.

$$\Delta\rho = \rho_{H/TM_x/TiO_2} - \rho_H - \rho_{TM_x/TiO_2} \quad (5)$$

3. Results and discussion

3.1. Photocatalyst characterization

To investigate the crystalline characteristics of TiO₂ loaded with Au, Ag, Pd and Pt, X-ray diffraction (XRD) patterns were recorded as illustrated in Fig. 1. It is well reported that the anatase phase of TiO₂ exhibits major reflections at 25.5°, 38.2°, 48.2°, 54.6°, 55.2°, 63.3°, 69.4°, 70.8°, and 75.5° corresponding to (101), (004), (200), (105), (211), (204), (106), (220), and (215) planes, as referenced in JCPDS No. 21–1272 [43, 44]. Strong diffraction peaks at 25° and 48° confirm that TiO₂ used in this study is in the anatase phase. The observed XRD pattern of Ag-loaded TiO₂ sample shows unambiguously that a mixed phase of face-centered, cubic Ag and anatase has formed. Distinctive 2θ° values at 38.1°, 44.3°, 64.4° correspond to the metallic Ag in the face centered lattice at (113), (124), and (220) as well matched with JCPDS card No. 65–2871 [45]. The Pd peak observed at 40.4° for the Pd-loaded sample is assigned to the (111) plane. Because Pd nanoparticles are very small, the peaks appear to be considerably broad [46]. The face-centered cubic Au (JCPDS card No. 002–1095) and anatase TiO₂ (JCPDS card No. 21–1272) can be identified by reflections in the diffraction pattern of the Au-loaded sample. The diffraction peaks at 2θ° values of 44.7°, 62.9°, and 77.7° correspond to the Au nanoparticles' (200), (220), and (311) planes, respectively. This indicates the successful formation of metallic Au nanoparticles on the surface of TiO₂ [47]. Pt nanoparticles were detected at 2θ° angles of 39.9° and 48.3° which can be indexed as (111) and (210) reflections of metallic Pt (JCPDS No. 01–1194) [48]. Additionally, Pt nanoparticles exhibited the broadest peaks among the metals studied here, indicating a smaller particle size compared to the others [50].

The surface composition of the metal-loaded TiO₂ nanoparticles was examined using XPS. Fig. 2a illustrates the surface survey XPS spectra of bare TiO₂ and Pd-, Pt-, Ag- and Au-loaded TiO₂. Fig. 2b shows the Ti 2p spectrum of bare TiO₂ nanoparticles, exhibiting peaks at 458.1 eV and 463.6 eV, corresponding to Ti⁴⁺ in TiO₂ and a peak at 460.4 corresponding to Ti³⁺. These results confirm the full oxidation of titanium. The nearly symmetric peaks and the binding energy difference of approximately 6 eV are consistent with typical values reported for TiO₂ [49]. Furthermore, the high-resolution spectrum of O 1s of the samples was deconvoluted into three peaks at 530.6, 532.7 eV, and 534.1 eV as shown in Fig. 2c. The peaks indicate bridging oxygen, terminal hydroxyl groups and the last peak at 534.1 eV corresponds to physically adsorbed

water (H₂O) [50]. Fig. 2d displays the Pd 3d spectrum. Pd 3d_{5/2} and 3d_{3/2} states are located at 334.7 and 340.0 eV, and the peaks located at 335.7 and 340.8 eV correlate to the Pd²⁺ 3d_{5/2} and 3d_{3/2} states. Furthermore, the presence of Pd²⁺ peaks arise from the partial oxidation of the palladium surface. This mainly arises from the relatively small size of Pd nanoparticles [51]. In the case of Pt loaded TiO₂, it is observed in Fig. 2e the spectra for Pt 4f, revealing peaks corresponding to metallic Pt⁰ as well as a higher oxidation state, which can be attributed to Pt²⁺ and Pt⁴⁺.

The stoichiometry of the PtO_x species produced from the reaction between small Pt particles and oxygen is dependent on the size of the particles. In agreement with Wang et al. [51], the x value varies from 1 to 2 with the decrease in particle size, specifically below 2 nm. Given a range of Pt particle sizes, it can be expected that the smallest Pt particles will oxidize to form PtO₂, whereas the larger particles will produce PtO species. A PtO oxide layer may also form on the surface of the metallic Pt particles. Consequently, the detected Pt²⁺ oxidized species might be present either as a PtO coating on the Pt⁰ particles or as oxidized PtO particles. Zhang et al. [52] proposed that Pt⁰ sites can possibly be oxidized to Pt²⁺ species, which in turn substitutes for Ti⁴⁺ ions in the TiO₂ lattice [53]. Furthermore, The Pt⁴⁺ state suggests that there may be a strong interaction between platinum and TiO₂, possibly leading to the formation of Pt–O–Ti bonds, which stabilize the oxidized platinum state [54]. However, the oxidation of metallic Pd and Pt is reversible, as these can be reduced back to their metallic forms under the anaerobic conditions employed during hydrogen evolution reactions. This reversible behaviour has been documented in previous literature. For instance, Kumaravel et al. [55] reported that under anaerobic conditions, hydrogen can act as a reducing agent to restore oxidized noble metals to their metallic states. Therefore, Pt and Pd in this study got reduced from the oxidized states of Pd and Pt to their metallic states. The Ag 3d_{5/2} and Ag 3d_{3/2} peaks (Fig. 2f) are attributed to binding energies of 374.2 and 368.2 eV, respectively. The 6.0 eV separation between the peaks suggests the successful formation of metallic Ag, confirming that the in-situ reduction of silver has taken place [56]. Moreover, peaks of Ag⁺ appeared at binding energies of 367 and 373 eV, respectively. The Ag⁺ peak indicates the presence of silver in an oxidized state, likely Ag(I). This could be due to the interaction of silver with the TiO₂. The oxidation of silver can occur in the presence of oxygen or hydroxyl groups from TiO₂, leading to the formation of Ag⁺ (e.g., Ag₂O or AgTiO₃ species) [57]. Additionally, as illustrated in Fig. 2g, the XPS spectra of Au shows the occurrence of Au 4f_{5/2} and Au 4f_{7/2} doublet with binding energies of 93.3 and 86.9 eV respectively, which are characteristic for Au in the Au⁰ state [49]. Notably, Au–TiO₂ showed distinct peaks at 84.0 and 87.6 eV (Fig. 2g), corresponding to Au^{δ+} species, probably because Au^{δ+} can be stabilized on the surface of TiO₂ nanoparticles. It is speculated that such Au^{δ+} species may be incorporated into a part of the peripheral sites of the Au clusters on the TiO₂ surface, resulting in a strong interaction between the Au clusters and the TiO₂ NPs, thus avoiding sintering during the reaction [58].

Furthermore, XPS analysis of Ti 2p and O 1s are represented in Figs. S5 and S6, respectively. In Fig. S5 (Ti 2p), two distinct peaks at 459.4 and 465.1 eV correspond to the Ti 2p_{3/2} and Ti 2p_{1/2} spin-orbit components, characteristic of the Ti⁴⁺ oxidation state [59]. However, a slight shift of 0.12 and 0.15 eV is observed from TiO₂ (458.94 eV) to Pd–TiO₂ (458.75 eV) and Pt–TiO₂ (458.78 eV), indicating electron transfer between the NPs and TiO₂. In contrast, Au–TiO₂ (458.69 eV) and Ag–TiO₂ (458.67 eV) exhibit a more pronounced Ti³⁺ signal at ~457.5 eV, suggesting an increase in oxygen vacancies and partial reduction of Ti⁴⁺. This electronic modification is known to enhance charge carrier mobility. The variations in satellite peak intensity further support these electronic changes [60].

Moreover, subtle shifts in the primary O1s peak (Fig. S6 (O1s)) corresponding to lattice oxygen (O–Ti–O) are observed in the spectra of all metal-loaded samples. In Pt–TiO₂ (530.76 eV) and Pd–TiO₂ (530.75 eV), the binding energy is slightly higher compared to bare TiO₂

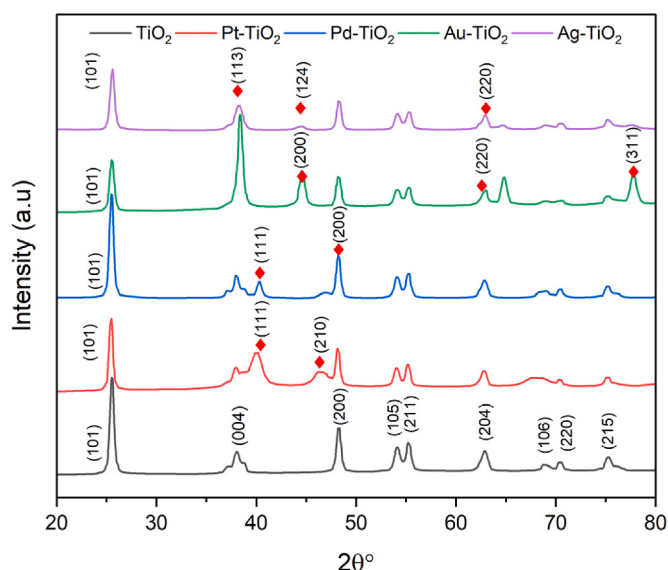


Fig. 1. X-ray diffraction peaks of metal-loaded TiO₂.

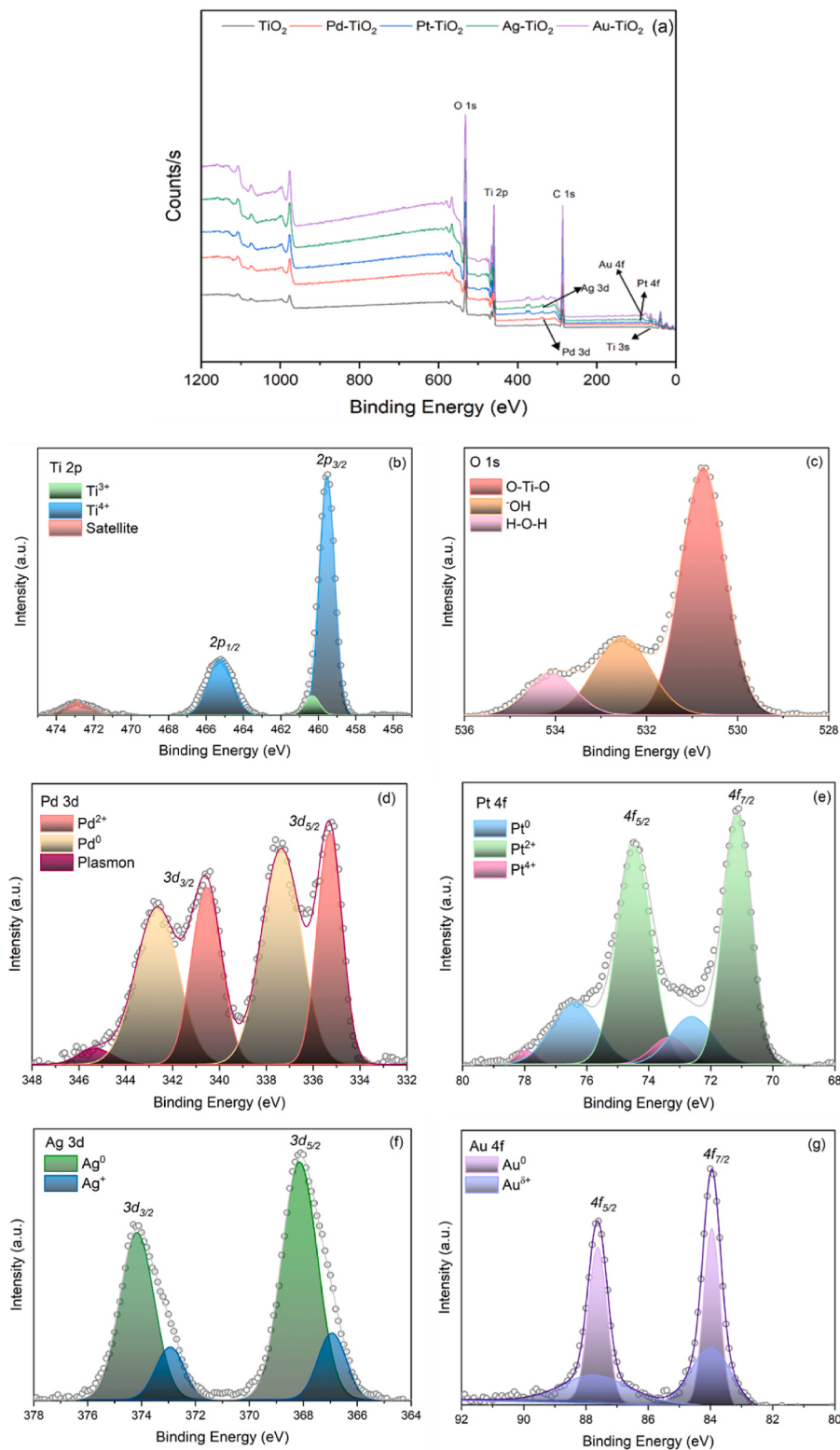


Fig. 2. XPS analysis of metal-loaded TiO₂. (a) XPS survey, (b) Ti 2p, (c) O 1s, (d) Pd 3d, (e) Pt 4f, (f) Au, 4f, (g) Ag 3d.

(530.63 eV), indicating electron withdrawal by these metals. This electron depletion reduces the electron density around oxygen, thereby stabilizing the Ti–O bond. In contrast, Au–TiO₂ (530.45 eV) and Ag–TiO₂ (530.63 eV) exhibit a lower or comparable binding energy, suggesting electron donation from these metals, leading to a slightly more reduced state [61].

The peaks at ~532–535 eV, corresponding to hydroxyl groups (-OH)

and molecularly adsorbed water (H₂O), also show intensity and position variations. Pt–TiO₂ (532.77 eV) and Pd–TiO₂ (532.55 eV) exhibit stronger hydroxyl signals, indicating enhanced surface hydroxylation, likely due to increased water dissociation facilitated by these metals. On the other hand, Au–TiO₂ (532.88 eV) and Ag–TiO₂ (532.20 eV) display lower hydroxyl intensities, implying reduced surface hydroxylation [62].

The TEM and HRTEM images presented in Fig. 3 provide insights into the morphology, dispersion, and crystalline structures of the metal-decorated TiO₂ nanocomposites, while the corresponding particle size distributions offer quantitative analysis of the nanoparticle sizes.

TEM images (Fig. 3a–d, g, j) confirm the successful deposition of metal nanoparticles on TiO₂. Ag–TiO₂ and Au–TiO₂ display larger, well-dispersed nanoparticles, with Ag exhibiting a more irregular morphology. In contrast, Pd–TiO₂ and Pt–TiO₂ show smaller, more

uniformly distributed particles, indicating enhanced dispersion.

HRTEM images (Fig. 3b–e, h, k) reveal clear lattice fringes in all samples, confirming high crystallinity. Distinct fringes corresponding to both the metal and TiO₂ phases further validate the successful nanoparticle incorporation. Particle size distributions (Fig. 3c–f, i, l) indicate that Ag–TiO₂ and Au–TiO₂ exhibit broader size ranges (20–110 nm and 10–90 nm, respectively), while Pd–TiO₂ and Pt–TiO₂ feature smaller, more uniform particles (~14 nm and ~5 nm, respectively). These results

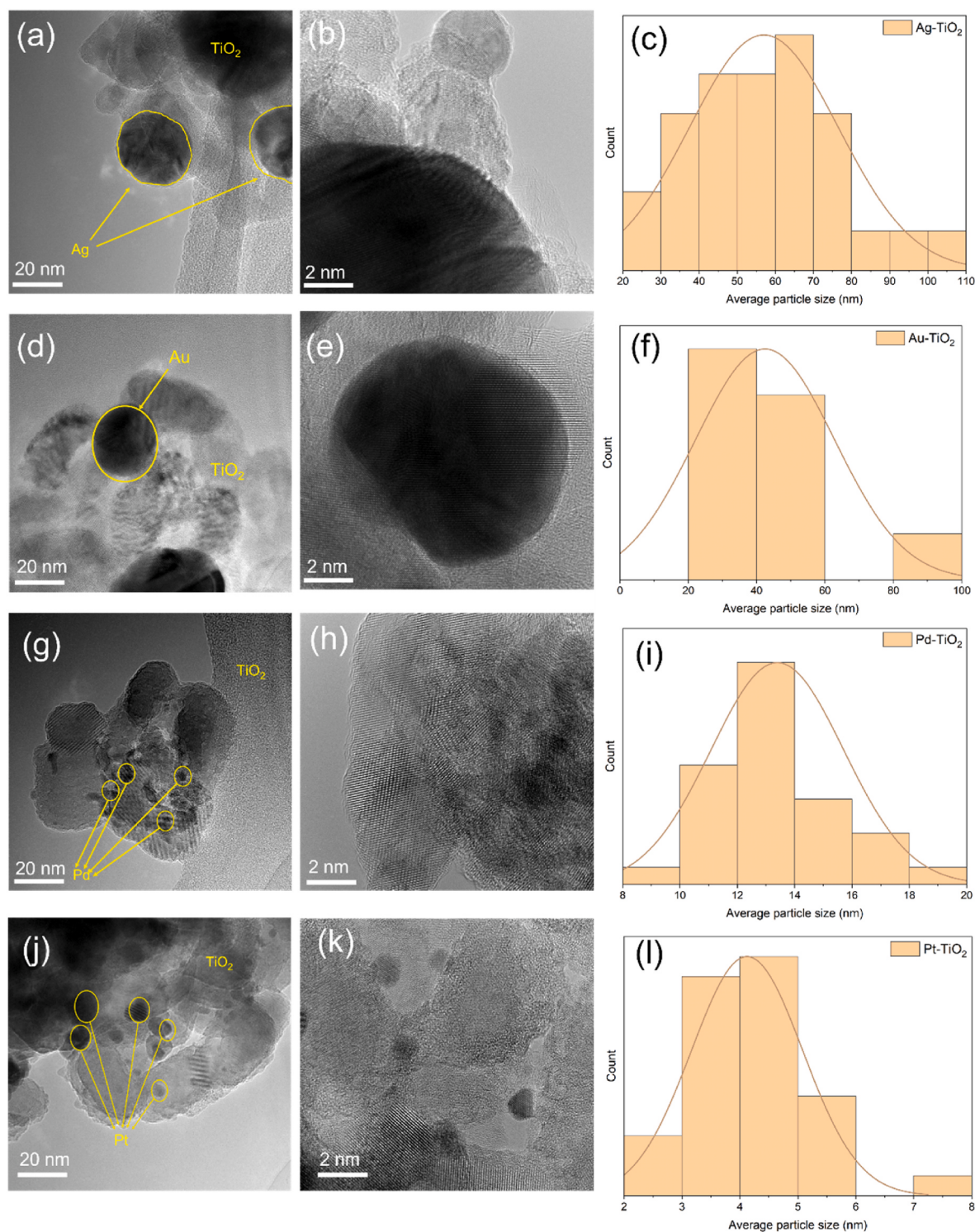


Fig. 3. Typical TEM and HRTEM images of Ag–TiO₂ (a, b), Au–TiO₂ (d, e), Pd–TiO₂ (g, h), and Pt–TiO₂ (j, k); particle size distributions of Ag–TiO₂ (c), Au–TiO₂ (f), Pd–TiO₂ (i), and Pt–TiO₂ (l).

align with the XRD analysis, which also indicates that Pt forms the smallest particles, followed by Pd, Au, and Ag.

The elemental distribution in TiO_2 samples doped with noble metals was evaluated using STEM/EDS mapping mode (Fig. 4). Pd and Pt appear to be evenly dispersed on the TiO_2 surface, while the Ag and Au nanoparticles distribution is less well distributed. The EDX analysis of Ag and Au samples were moreover confirmed. The Cu peaks observed in the EDX spectra are attributed to the use of copper tape for sample preparation. Furthermore, small amounts of gold and silver were detected in their volumes. The standard deviation for gold and silver

samples were very high (i.e., 19.56 and 20.37) against 2.31 and 0.96 for Pd and Pt nanoparticles, respectively. This indicates the agglomeration and non-uniform distribution of Ag and Au nanoparticles in the catalysts volume. This result is in agreement with the HR-TEM analysis, which shows few but large Au and Ag particles, compared to small and evenly distributed Pt and Pd particles.

We used DFT calculations to evaluate the stability of several model metal clusters adsorbed on the TiO_2 surface. The adsorption and cohesive energies calculated for the four investigated TM clusters is reported in Fig. 5. In all cases both the calculated adsorption and cohesive

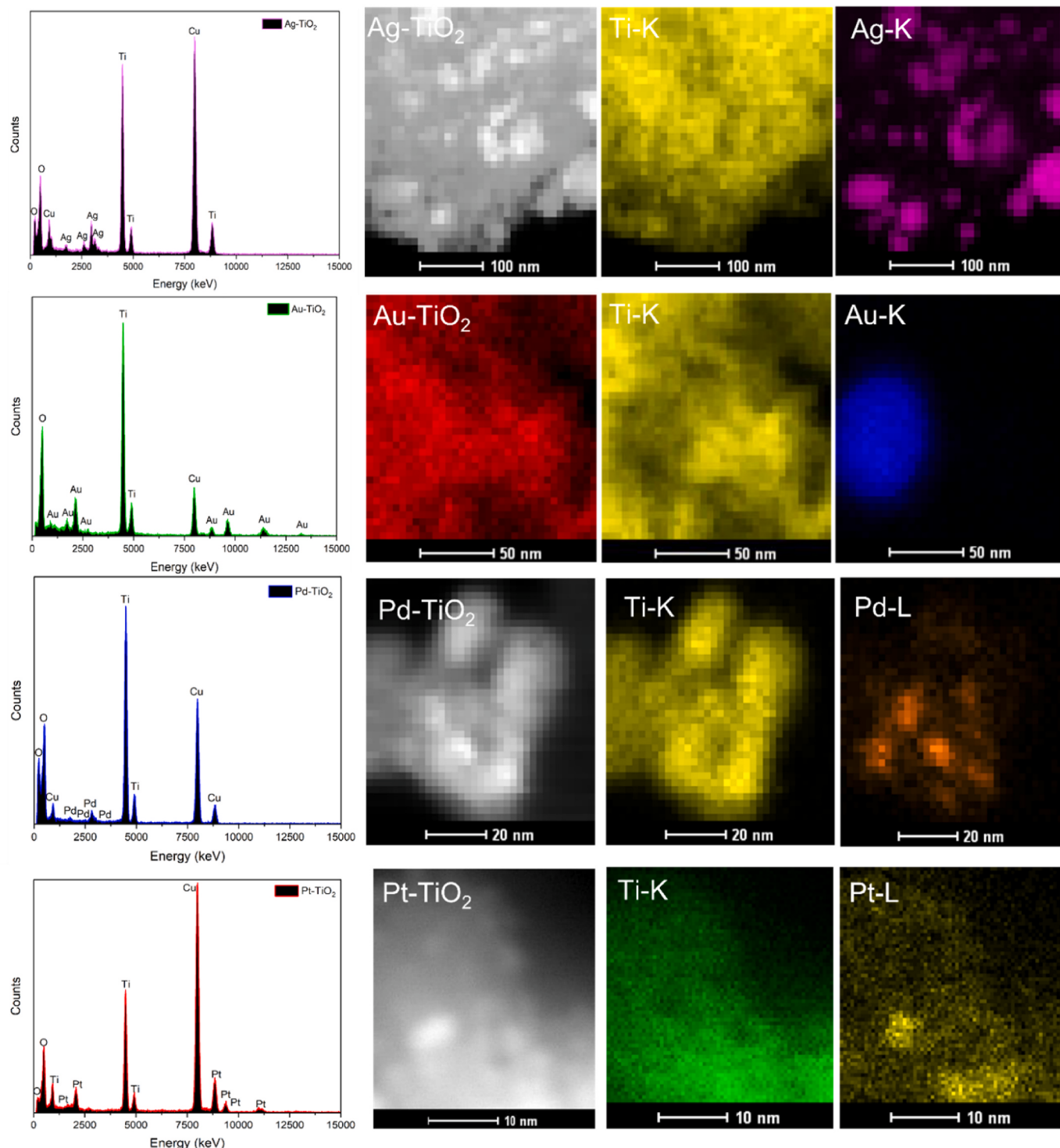


Fig. 4. HDAF-STEM EDX of (a) Ag-TiO₂, (b) Au-TiO₂, (c) Pd-TiO₂, (d) Pt-TiO₂.

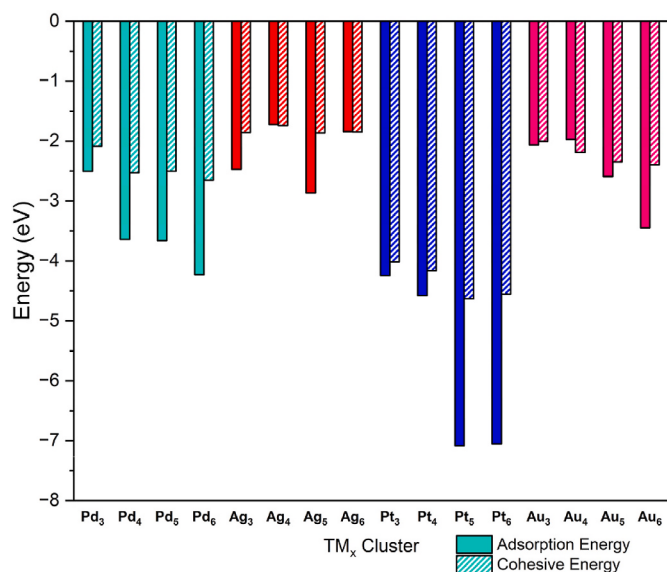


Fig. 5. Adsorption (solid bars) and cohesive (dashed bars) energies of different metal clusters (TM_x) on anatase (101) surface.

energies are negative, indicating the stability of the metal clusters on the host substrate [63]. The metal adsorption energies for the trimer clusters range from -2.0 eV to -4.2 eV, with Pt exhibiting the most favourable adsorption value. These values are within the range of other calculated adsorption energies reported in the literature [64–66]. It is worth noting that in the case of Pt and Pd clusters the adsorption energies are consistently and significantly more negative than the corresponding cohesive energies. This indicates a preference for these metals to form smaller and more dispersed nanoparticles. In contrast, Au and Ag exhibit adsorption and cohesive energies of similar magnitude, suggesting a more pronounced tendency for these metals to agglomerate. These trends align well with the experimentally observed nanoparticle sizes shown in Fig. 3.

In order to determine the specific surface area and assess the porosity of the photocatalysts, BET N_2 adsorption/desorption measurements were carried out (Fig. S8, Table S4). Based on the IUPAC classification, the isotherms of the synthesized catalysts can be classified as type IV isotherms [67]. This characteristic is typical of mesoporous adsorbents and indicates monolayer adsorption at low pressures and multilayer adsorption at higher pressures. The separation of the adsorption-desorption branches at low pressure ($P/P_0 = 0.1$) suggests

that capillary condensation, represented by type H3 hysteresis loops, begins at lower pressures in these samples. The metal photodeposition reduces the specific surface area and the total pore volume compared to pristine TiO_2 . This can be associated to alterations in the morphology, crystalline structure, and size [68,69]. The addition of Ag resulted in a strong decrease in the peak from 146 Å for bare TiO_2 to 86 Å. This decrease can be attributed to metal particles blocking the pores, thereby reducing the overall surface area available for nitrogen adsorption in the BET experiment. In contrast, the addition of Pd slightly increased the peak to 153 Å. This increase is likely due to the highly dispersed Pd nanoparticles on the TiO_2 surface, which not only preserve the pore structure but may also create new nanopores with a radius of approximately 15.3 nm, thus decreasing the overall surface area.

The optical characteristics of TiO_2 particles, both pure and those loaded with noble metals, were examined using UV–vis DRS, as depicted Fig. 6a. The DRS spectra span from 200 to 800 nm. Unlike bare TiO_2 , the noble metal-loaded TiO_2 particles exhibit visible light absorption. This broad absorption feature is due to the surface plasmon resonance of the noble metal nanoparticles. The introduction of Pd, Pt, Au, and Ag into TiO_2 lattice increases the likelihood of electron transitions from the valence band to the conduction band [70]. Moreover, the localized surface plasmonic resonance (LSPR) effect of Ag, Au, Pt, and Pd nanoparticles further improves the solar energy conversion efficiency. This enhancement is achieved by broadening light absorption to longer wavelengths, light scattering increase, and promotion of the photo-generated carriers in the semiconductor through the transfer of plasmonic energy from the noble metal nanoparticle to the semiconductor. However, it is worth noting that, in some cases, surface-located metal nanoparticles may serve as recombination sites [71].

According to Sellappan et al. [72], and Ghosh et al. [73], the LSPR effect in Ag- TiO_2 and Au- TiO_2 photocatalysts is influenced by the number of nanoparticles on the TiO_2 surface. Increasing Ag and Au ion concentrations during the catalyst preparation results in a red shift in the LSPR band of their absorption spectra. Besides, the photoconversion quantum yield may be improved by modification of the photocatalyst using noble metal nanoparticles (e.g., Pt) [74]. This finding might be explained considering two distinct phenomena: (i) the existence of Ti^{3+} , as a result of platinum being deposited on the TiO_2 surface; and (ii) the establishment of the Schottky barrier at the interface between Pt and TiO_2 , which prevents the recombination of the photogenerated charge carriers [70].

To better understand the behavior of the electron-hole pairs in the semiconductor particles, the effectiveness of charge migration and carrier trapping has been investigated using PL spectroscopic technology. Fig. 6b shows the PL spectra of bare TiO_2 and representative metal-

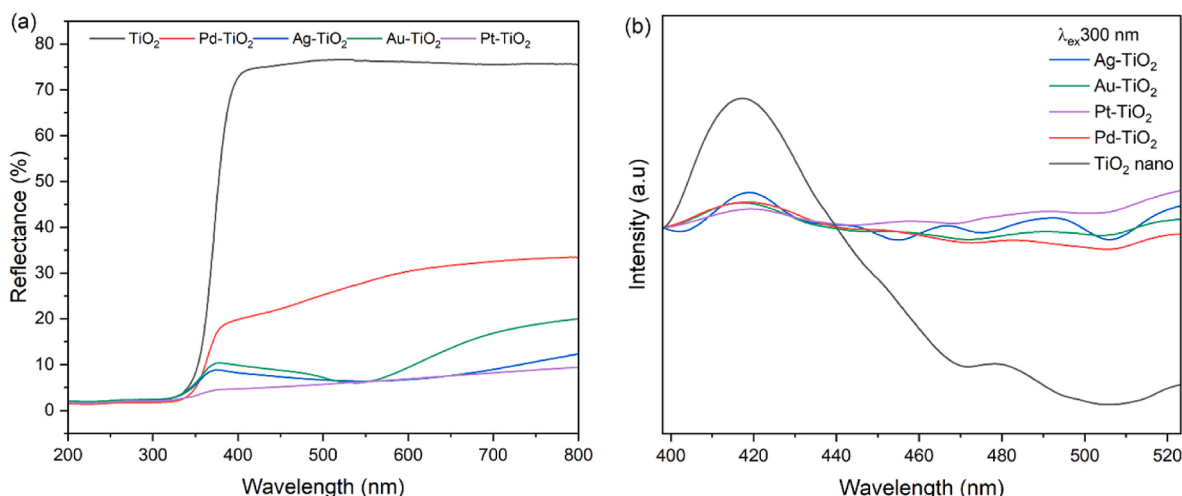


Fig. 6. (a) UV–vis diffuse reflectance spectra of bare and metal-loaded TiO_2 , (b) Photoluminescence spectra of bare and metal-loaded TiO_2 .

loaded TiO₂, after background subtraction.

The primary emission peak observed around 420–423 nm can be ascribed to the band-to-band transition taking place owing to the migration of electrons from the semiconductor's conduction band back to the valence band [43]. A distinguished reduction in the main emission peak is noted with the doping of Ag, Pd, Pt and Au into TiO₂ surface. Thus, although noble metals do not induce new PL phenomena, they reduce the excitonic PL intensity. This was attributed to the capture of photoexcited electrons or holes by the metal ions. Consequently, in noble metal-loaded systems, the lower excitonic PL intensity is indicative of a higher rate of photo-induced charge carrier separation [75,76]. It can be seen in Fig. 6b that the extent of PL quenching varies depending on the metal type. Specifically, the addition of Pt, Au, and Pd results in a significant reduction in PL intensity compared to Ag deposition and bare TiO₂, while the addition of Ag results in a moderate reduction in PL intensity. Therefore, surface modification with Au, Pt, Pd and Ag can effectively suppress the recombination process of the photogenerated carriers in TiO₂, in the following order of effectiveness: Pt > Pd ≈ Au > Ag > bare TiO₂.

To further validate the impact of noble metal incorporation on the charge transfer dynamics, an electrochemical impedance spectroscopy (EIS) test was conducted as depicted in Fig. S7 and Table S3. Pristine TiO₂ exhibits the highest R_{ct} value of 152 kΩ, confirming its poor electrical conductivity and inefficient charge transfer. Doping with noble metals significantly reduces R_{ct} , with Ag–TiO₂ (0.15 kΩ) and Pd–TiO₂ (0.71 kΩ) showing the most pronounced improvements, suggesting enhanced charge transport facilitated by the formation of Schottky junctions [77]. Pt–TiO₂ also demonstrates a substantially lower R_{ct} (7.52 kΩ) compared to bare TiO₂, whereas Au–TiO₂ exhibits a moderate reduction to 87.7 kΩ, indicating relatively limited conductivity enhancement [78]. The trends in the constant phase element (CPE) values, presented in Table S3, further support these findings. Pristine TiO₂ exhibits the highest capacitance (50.7 μF), indicative of greater charge accumulation and recombination, while noble metal-doped samples such as Pt–TiO₂ (0.131 μF) and Pd–TiO₂ (0.163 μF) display significantly lower CPE values, reflecting improved charge separation [79].

3.2. Photocatalyst activity

Fig. 7a shows the photocatalytic hydrogen production rate in the presence of different metals (i.e., Ag, Pd, Au, Pt) photodeposited on TiO₂, using methanol as the sacrificial agent, under UV and visible light irradiation. As clearly indicated in the figure, the loading of Pt resulted in the best performance in terms of hydrogen evolution, reaching a hydrogen production rate of 28 mmol/h g, which is about 35-fold higher than bare TiO₂. In contrast, Ag shows the lowest hydrogen productivity compared to the other metals, resulting in a hydrogen production rate of approximately 2 mmol/h g. Overall, the H₂ evolution under UV/visible irradiation follows the trend: Pt > Au ≈ Pd > Ag > bare TiO₂.

The particle size of metal nanoparticles plays a critical role in the photocatalytic hydrogen production efficiency of metal-decorated TiO₂ nanocomposites [80]. Smaller nanoparticles typically exhibit a higher surface area-to-volume ratio, which increases the number of active sites available for the hydrogen evolution reaction (HER) and enhances charge separation by acting as electron sinks [81]. This often leads to improved hydrogen production rates, with small nanoparticles demonstrating superior catalytic activity due to their high work function and near-zero overpotential for HER [82]. The enhanced performance of Pt and the relatively weaker performance of Ag under UV/visible light observed in Fig. 7a align with the observed nanoparticle sizes (see Fig. 3). However, the performance trend cannot be solely attributed to nanoparticle sizes, indicating that other factors should be taken into account. The H₂ evolution trend under UV/visible light is consistent with the suppression of the photogenerated charge recombination, as evidenced by the PL measurements in Fig. 6b. A slightly different trend

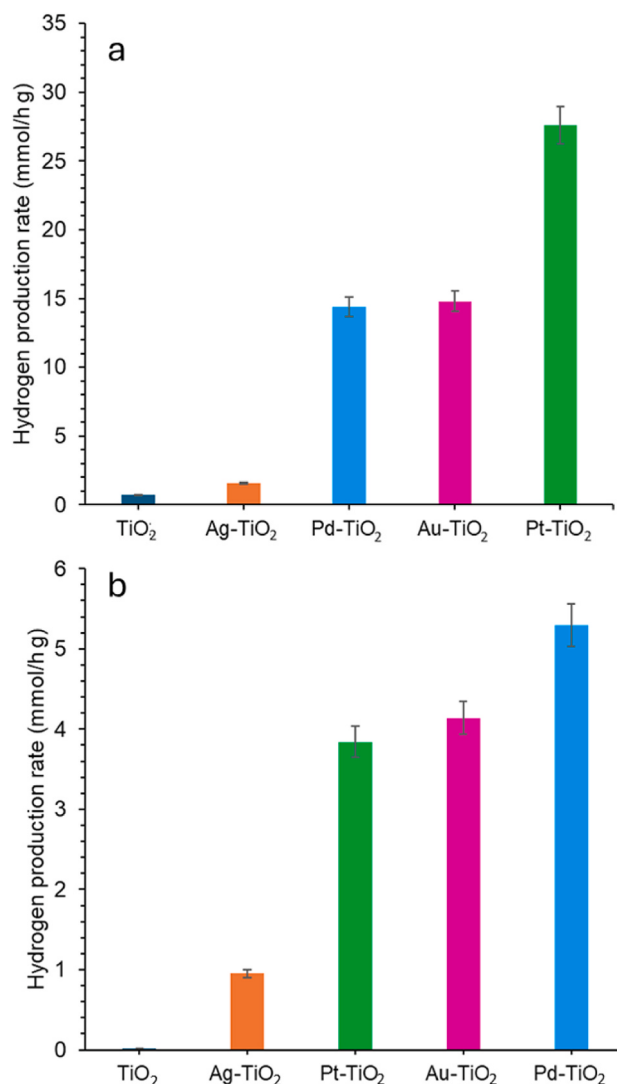


Fig. 7. Hydrogen production rate in the presence of bare and metal-loaded photocatalysts under (a) UV/visible and (b) visible light irradiation. Experimental conditions: [Catalyst] = 500 ppm; [Methanol] = 2.5 M; neutral pH; T = 25 °C.

was observed under visible light irradiation (Fig. 7b), with Pd exhibiting the highest activity amongst the investigated metals. This sample is the one exhibiting the least extent of PL low energy emission (430–500 nm), likely correlated to a lower density of intraband gap states working as detrimental recombination sites [83]. The lower density of these gap states implies that under visible light, the concentration of catalytically active photoholes is rate limiting, with Pd displaying the longest lifetimes out of all metals, which promotes steps 2–3, 4–5, and 7–8 in Fig. S9, resulting in higher hydrogen production. All the metal-loaded samples demonstrated a measurable H₂ production rate under visible light irradiation, although this was substantially lower than under UV/visible light, where hydrogen radical coupling is the rate limiting factor.

Reactant species, such as methanol or water, may initially undergo thermal dissociation, generating surface anions and protons. Under suitable photocatalytic conditions, methoxy and hydroxyl anions can capture holes, while protons migrate to co-catalysts to capture electrons, leading to the formation of surface radicals. Hydrogen radicals can subsequently combine to form molecular hydrogen, which desorbs from the surface, while methoxy and hydroxyl radicals undergo further oxidation, producing additional protons during steps 3–4 and 6–7

(Fig. S9). The reaction mechanism remains consistent under both UV/visible and visible light irradiation; however, the lower activity observed under visible light is attributed to the reduced concentration and shorter lifetimes of charge carriers, reducing the likelihood of steps 2–3, 4–5, and 7–8 occurring.

To elucidate the mechanism of charge transfer upon hydrogen adsorption on the TM, we evaluated the change in electron density upon adsorption of hydrogen on a model TM_3/TiO_2 system. Fig. 8 shows the charge density distribution plots (CDDs) for the investigated TMs. When hydrogen adsorbs on Pt, Au, and Pd, migration of electrons from the adsorbed hydrogen to the TM cluster is observed. In these cases, electron depletion regions appear on top of the adsorbed hydrogen atom. In contrast, when hydrogen adsorbs on Ag, there is significant electron transfer from TiO_2 and the TM cluster to the adsorbed hydrogen. Fig. 9 shows the Gibbs free energies of hydrogen adsorption on the TM_3/TiO_2 clusters (ΔG_H) calculated using Equation (4). Qualitative correlations can be observed between charge transfer and hydrogen adsorption Gibbs free energies: Pt, Au, and Pd, which gain some electron density from hydrogen, have zero or small negative hydrogen adsorption energies, while Ag, which lose electron density to hydrogen, have large positive ΔG_H values. According to the Sabatier principle, a photocatalyst active site with a ΔG_H close to zero presents the most promising activity towards the hydrogen evolution reaction rate [84]. In this context, a single hydrogen atom must not adsorb too weakly ($\Delta G_H \gg 0$) nor too strongly ($\Delta G_H \ll 0$) to avoid a significant energy barrier to reaction.

As observed in Fig. 9, Pt exhibits an absolute ΔG_H value close to 0. Consistent with the Sabatier principle, Pt provides the highest hydrogen production rate amongst the investigated metals. Both Au and Pd show similar small negative ΔG_H values. These metals also achieve relatively high hydrogen production rates. In contrast, a large positive value of ΔG_H for Ag (0.92 eV) aligns with the lowest hydrogen production rate observed experimentally. By comparing Fig. 7a and b, we observe that

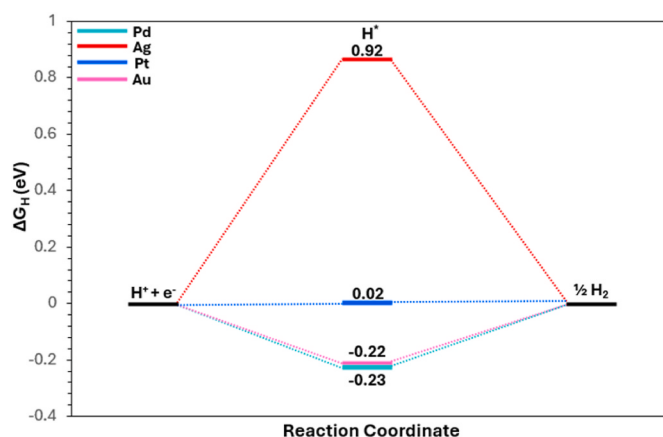


Fig. 9. Reaction coordinate diagram showing the Gibbs free energy for hydrogen evolution reaction on different TM_3/TiO_2 systems. H^* stands for hydrogen adsorbed on TM_3/TiO_2 .

hydrogen radical coupling on Ag remains the rate limiting step, discounting any benefit in additional absorbance of visible light when compared to other metals. This analysis suggests that ΔG_H is a critical descriptor for evaluating the activity of a photocatalytic material in hydrogen production.

4. Conclusion

This study provides a comprehensive analysis of the photocatalytic performance of anatase TiO_2 loaded with various noble metals, such as Pt, Au, Pd, and Ag. Through a combination of experimental and computational approaches, we have elucidated the main factors through

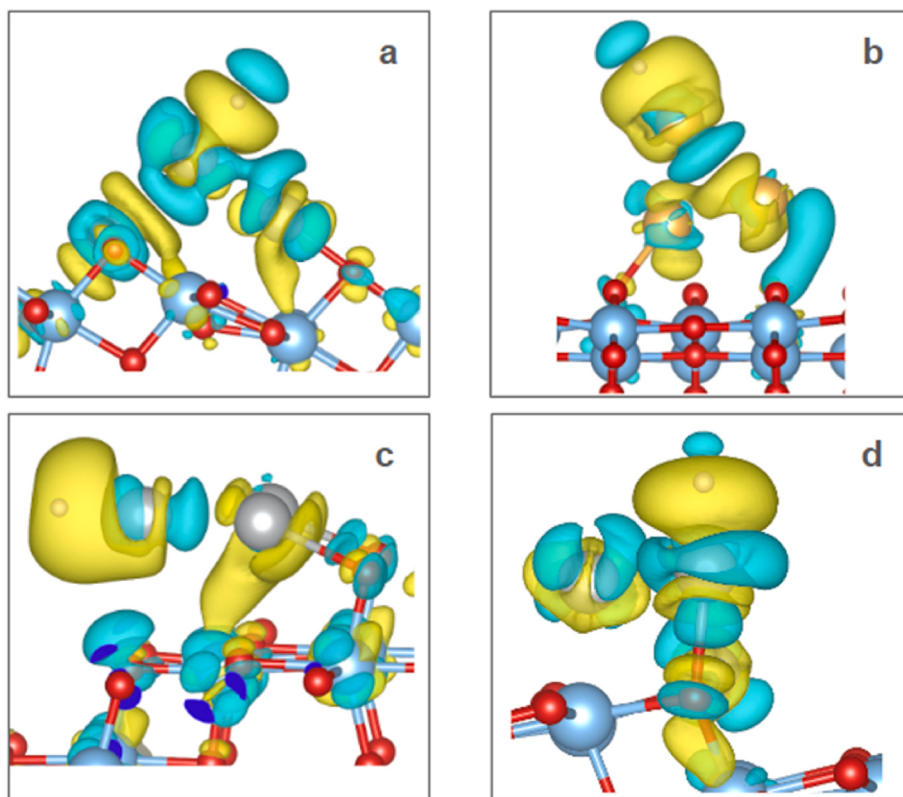


Fig. 8. Charge density distribution plots for hydrogen adsorbed on Pt (a), Au (b), Ag (c), and Pd (d) model TM_3/TiO_2 structures. The yellow region represents electron accumulation, and the blue region represents electron depletion. (Isosurface value = 0.0011 bohr^{-3}). (For interpretation of the references to colour in this figure legend, the reader is referred to the Web version of this article.)

which these co-catalysts influence the charge carrier dynamics, light absorption properties, and overall hydrogen production efficiencies of TiO₂ photocatalyst. The hydrogen production rate under simultaneous UV and visible light irradiation followed the trend Pt > Au ≈ Pd > Ag > bare TiO₂. Among the co-catalysts, Pt and Pd were well distributed on the TiO₂ surface and formed the smallest nanoparticles, while Ag and Au exhibited less uniform distribution. All noble metal-loaded TiO₂ samples showed enhanced absorption in the visible region, primarily due to the surface plasmon resonance of the metal nanoparticles. Pt, Au, and Pd were particularly effective in suppressing the electron-hole recombination process, as evidenced by the significant reduction in photoluminescence intensity compared to Ag deposition and bare TiO₂. The effectiveness in suppressing the electron-hole recombination followed the order Pt > Pd ≈ Au > Ag > bare TiO₂.

Additionally, our DFT calculations showed that hydrogen adsorption on Pt, Au, and Pd co-catalyst clusters results in the electron transfer from the adsorbed hydrogen to the metal cluster. The analysis of Gibbs free energies of hydrogen adsorption revealed a correlation between adsorption Gibbs energies and hydrogen production rates, confirming the order of effectiveness as Pt > Pd ≈ Au > Ag. In conclusion, this study underscores the critical role of metal co-catalysts in optimizing the photocatalytic performance of TiO₂ for hydrogen production, with Pt emerging as the most effective due to its superior charge carrier management, stability, and interaction with hydrogen. This research provides a framework based on the combination of experimental and computational methods for the design and development of novel photocatalysts based on the combination of semiconductors with transition metals for sustainable applications.

CRediT authorship contribution statement

Sarah Hamdan: Writing – original draft, Investigation, Data curation. **Matthew J. Wigglesworth:** Writing – original draft, Investigation, Formal analysis. **Marica Muscetta:** Writing – review & editing, Methodology, Data curation, Conceptualization. **Ruiman Ma:** Writing – original draft, Investigation, Data curation. **Mohamed I. Helal:** Investigation, Data curation. **Natalia Martsinovich:** Writing – review & editing, Supervision, Project administration, Conceptualization. **Giovanni Palmisano:** Writing – review & editing, Supervision, Project administration, Conceptualization. **Sergio Vernuccio:** Writing – review & editing, Supervision, Funding acquisition, Conceptualization.

Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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Appendix A. Supplementary data

Supplementary data to this article can be found online at <https://doi.org/10.1016/j.ijhydene.2025.03.184>.

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