

Metastable behavior of vortex matter in the electronic transport processes of homogenous superconductors

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We study numerically effect of vortex pinning on hysteresis voltage-temperature ($V - T$) loop of vortex matter. It is found that different types of the $V - T$ loops result from different densities of vortex pinning center. An anticlockwise $V - T$ loop is observed for the vortex system with dense pinning centers, whereas a clockwise $V - T$ loop is brought about for vortices with dilute pinning centers. It is shown that the size of the $V - T$ loop becomes smaller for lower experimental speed, higher magnetic field or weak pinning strength. Our numerical observation is in good agreement with experiments.

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Continuous efforts have been made to understand the metastable phenomena in vortex matter for low- T_c and high- T_c superconductors¹⁻¹⁰. Generally, metastable state is accessible when experimental speed is higher than the response speed of the vortices to external parameter such as temperature and magnetic field. For a vortex assemble, as the response speed of vortex system changing from its initial state to final one is too small to reach its final equilibrium state, vast metastable states that the system passes through are observed. As a result, relaxation which is the time evolution of the metastable state occurs². Then one can observe different states at different time windows. On the other hand, if experimental speed is smaller by far than the response speed of a vortex system, one can observe in fact its final equilibrium state and thus no metastable state can be probed. This is so-called the "quasi-static state" process. Therefore, one can study the metastability of vortex matter which is characterized by current-voltage $I - V$ or $R - T$ loops in transport measurements by adjusting either experimental speed (different scanning speed of temperature dT/dt , field dH/dt or current dI/dt) or system response speed by changing T , H , I or vortex pinning in experiments for different systems, such as the untwinned⁴⁻⁶ or the twinned $YBa_2Cu_3O_7$ ⁸, the pure⁹ or the doped $2H - NbSe_2$ ¹⁰, as well as in simulation^{11,12}. Despite a large number of research work which is mainly associated with the effects of external parameters on metastability, the role of vortex pinning is still not clear.

In this work, we use a standard two-dimensional Langevin simulation to study the effect of vortex pinning, temperature scanning speed (TSS), as well as magnetic field on the hysteresis $V - T$ loop. We demonstrate how the competition between experimental speed and the response speed of vortex systems with different densities of pinning centers affects the direction and size of the $R - T$ loop.

We model a transverse two dimension slice (in the x -

y plane) of an infinite superconducting slab containing rigid vortices that are parallel to the sample edges ($\mathbf{B} = B\hat{\mathbf{z}}$). The overdamped Langevin equation of motion for a vortex in position \mathbf{r}_i is

$$\mathbf{F}_i = \sum_{j \neq i}^{N_v} \mathbf{F}^{vv}(\mathbf{r}_i - \mathbf{r}_j) + \sum_k^{N_p} \mathbf{F}^{vp}(\mathbf{r}_i - \mathbf{r}_k^p) + \mathbf{F}^L + \mathbf{F}_i^T = \eta \frac{d\mathbf{r}_i}{dt}$$

where \mathbf{F}_i is the total force acting on vortex i due to vortex-vortex interactions (\mathbf{F}^{vv}), pinning centers (\mathbf{F}^{vp}), the driving current \mathbf{J} ($\mathbf{F}^L \propto \phi_0 \mathbf{J} \times \hat{\mathbf{z}}$) and thermal fluctuations (\mathbf{F}^T), η is the Bardeen-Stephen friction coefficient ($\eta \propto \phi_0 B_{c2} / \rho_n$), N_v the number of vortices, N_p the number of pinning centers and \mathbf{r}_k^p the position of the k th pinning center. We choice $\mathbf{F}^{vv}(\mathbf{r}_i - \mathbf{r}_j) = (\phi_0^2 s) (2\pi \mu_0 \lambda^2)^{-1} (\mathbf{r}_i - \mathbf{r}_j) (|\mathbf{r}_i - \mathbf{r}_j|)^{-2}$, where ϕ_0 is the flux quantum, s the length of the vortex, μ_0 the vacuum permeability, λ the temperature dependent penetration depth. We employ periodic boundary conditions and cut off the logarithmic vortex-vortex repulsion potential smoothly¹³. Random pinning center exerts elementary pinning force on the vortices: $\mathbf{F}^{vp}(\mathbf{r}_i - \mathbf{r}_k^p) = -f_{pv0} r_{ik} \exp(-(r_{ik}/r_p)^2) \hat{\mathbf{r}}_{ik}/r_p$, where f_{pv0} tunes the strength of this force at zero temperature, r_p which was assumed to equal to 0.2λ determines its range. The thermal force is implemented with a Box-Müller random number generator and has properties $\langle F_i^T \rangle = 0$ and $\langle F_i^T(t) F_j^T(t') \rangle = 2\eta k_B T \delta_{ij} \delta(t - t')$ at a given temperature T . To reflect the competing roles of \mathbf{F}^{vv} , \mathbf{F}^{vp} and \mathbf{F}^T as a function of temperature, we use $\lambda(T)/\lambda(0) = \xi(T)/\xi(0) = (1 - T/T_c)^{-1/2.1}$. The average velocities of vortices $\langle v_x \rangle = \frac{1}{N_v} \sum_i^{N_v} v_{ix}$ which is proportional to the mobility of the vortices and the resulting voltage (ie resistance R). We normalize lengths by $\lambda_0 = \lambda(0)$, the forces by $f_0 = (\phi_0^2 s) (2\pi \mu_0 \lambda_0^3)^{-1}$ and time by $t_0 = \lambda_0 \eta(0) / f_0$. All quantities shown in the following figures are expressed in the simulation unit. This simulation is always performed on cooling and followed heating cycle with a con-

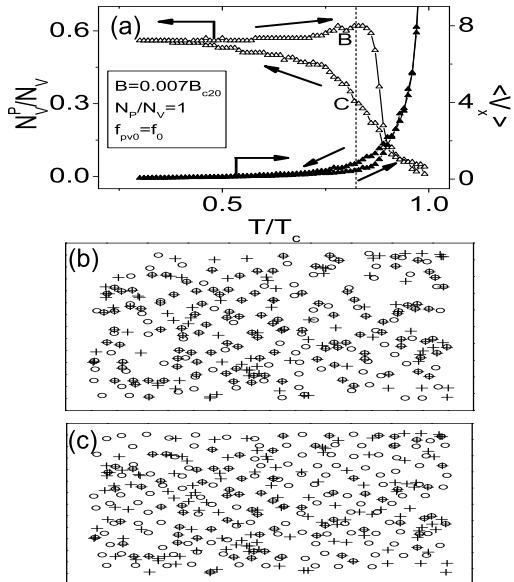


FIG. 1: (a) Typical anticlockwise $V-T$ loop (solid triangles) and the corresponding $N_p^v/N_v - T$ loop (open triangles). N_p^v is the number of pinned vortices. Temperature is varied in the direction indicated by the arrows. Solid lines are guide for eyes. (b) The vortices (open circles) and pinning centers (crosses) at position B in (a). (c) The vortices (open circles) and pinning centers (crosses) at position C in (a).

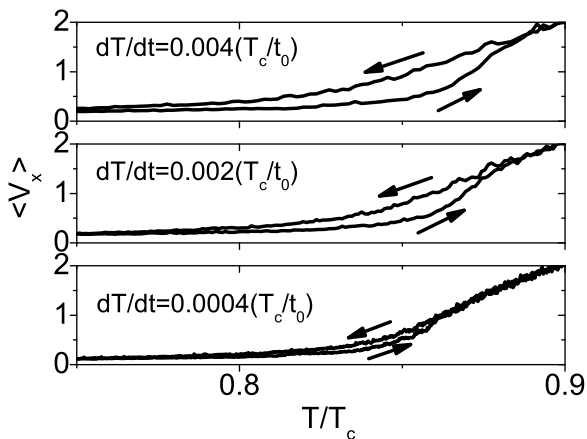


FIG. 2: Effect of temperature sweeping speed (TSS) on $V-T$ loop. $B = 0.007B_{c20}$, $N_p = N_v$, $f_{pv0} = 9f_0$. TSS affects only size of the $V-T$ loops but not their direction.

stant driving force ($F^L = 0.5f_0$) for the initial ordered vortex state. The random initial vortex arrangement has also been considered in this simulation, but the results obtained are qualitatively same as those for the ordered initial vortex state (see below). Besides, our results are insensitive to the choice of the number of vortices more than 400.

Fig. 1 illustrates a typical anticlockwise $V-T$ loop and the corresponding number of pinned vortices (N_p^v) normalized by N_v vs temperature curve ($N_p^v/N_v - T$) in one

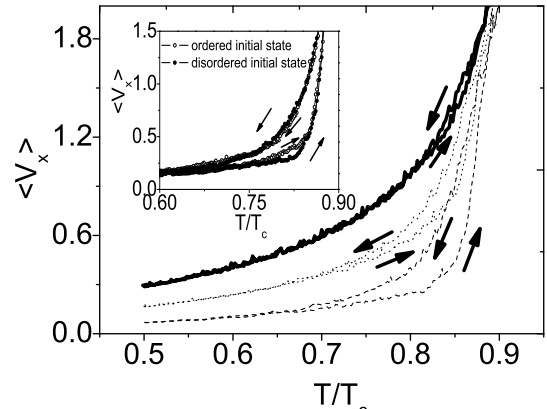


FIG. 3: Effect of field on $V-T$ loop. $dT/dt = 0.04T_c/t_0$, $f_{pv0} = 9f_0$, $B = 0.007B_{c20}$ (dash lines), $0.028B_{c20}$ (dot lines) and $0.056B_{c20}$ (solid lines) corresponding to $N_p = N_v$, $0.25N_v$ and $0.125N_v$. Inset: The $V-T$ loops for different initial vortex states. $dT/dt = 0.04T_c/t_0$, $f_{pv0} = 9f_0$, $N_p = N_v$ and $B = 0.007B_{c20}$. The field affects only the size of $V-T$ loops but not their direction.

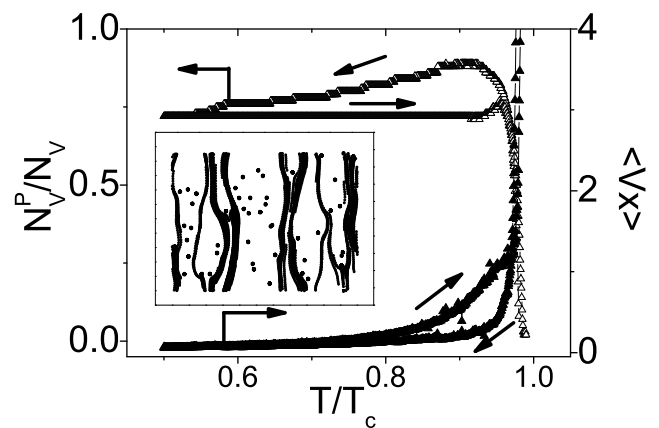


FIG. 4: Typical clockwise $V-T$ loop (solid triangles) and the corresponding $N_p^v/N_v - T$ curve (open triangles) for a systems with dilute pinning centers. $B = 0.007B_{c20}$, $N_p = 0.1N_v$, $dT/dt = 0.001T_c/t_0$, and $f_{pv0} = 90f_0$. Inset: Steady channel flow forms in the warming branch from $0.5T_c$ to $0.91T_c$. The circles and lines represent the random pinning centers and vortex trajectories, respectively.

cooling/heating circle for the vortex system with relative dense and weak pinning centers ($N_p = N_v$, $f_{pv0} = f_0$). With decreasing temperature, the number of pinned vortices increases at first because of the enhanced pinning forces for lower temperatures. However, the speed of vortex moving into the random weak pinning potentials is very slow due to the small f_{pv0} , which results in the slow but continuous increment of the number of pinned vortices N_p^v till the temperature reaching T_B on heating. At still higher temperatures, the pinned vortices decrease

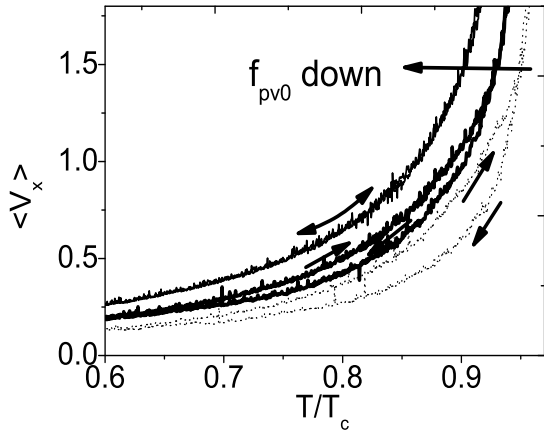


FIG. 5: Effect of pinning force strength f_{pv0} on the $V-T$ loops. $B = 0.007B_{c20}$, $N_p = 0.1N_v$, $dT/dt = 0.001T_c/t_0$, $f_{pv0} = (\text{right}) 60f_0, 45f_0, 30f_0$ (left). The size of $V-T$ loop becomes smaller with the decrement of f_{pv0} , but its direction does not change.

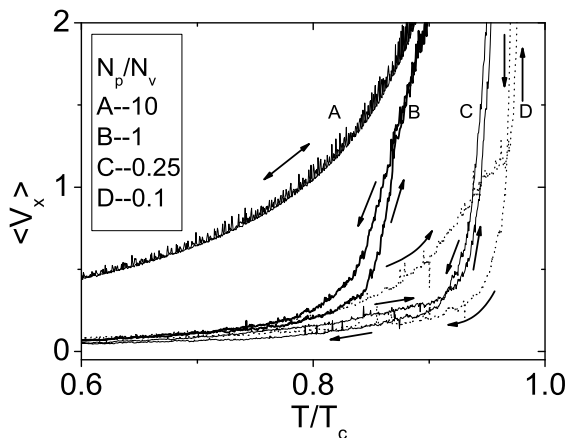


FIG. 6: Effects of vortex pinning density on the $V-T$ loops. $B = 0.007B_{c20}$, $dT/dt = 0.001T_c/t_0$, $N_p f_{pv0}/N_v \equiv 9f_0$. The direction of $V-T$ loops depends on vortex pinning density.

quickly because of the stronger thermal fluctuation. That is to say, the number of pinned vortices N_p^p is increasing not only on cooling but also in most temperature regime on heating (see Fig. 1(b) and Fig. 1(c), respectively). Therefore, the $N_p^p/N_v - T$ curve manifests itself as the clockwise loop and the corresponding anticlockwise $V-T$ loop⁴. Because the response speed of vortices increases with increasing temperature, there is no loop in the region of temperature near T_c where thermal fluctuation is important. In contrast, at very low temperatures where both elasticity and pinning interactions are strong and thermal fluctuation is weak, the loop is too small to be observed.

Fig. 2 shows the influence of TSS on $V-T$ loop. It is seen that the size of the loops becomes smaller with decrement of TSS . The reason is that the vortex sys-

tem has enough time to relax into more stable pinned states in a slower transport measurement. Therefore, for a very clean sample in which vortex pinning is very weak and thus response speed is very large, only has a fast experimental speed such as fast temperature or current scanning rates been used, the hysteresis loop can be observed⁹.

We have also investigated the effects of magnetic field on the hysteretic behaviors. Because the interaction between vortices (the elasticity of vortex matter) increases with increasing field, the response speed of vortex system will also increase correspondingly. Hence it is expected that the $V-T$ loop becomes smaller with increasing field. Shown in Fig. 3 is the influence of applied field on the $V-T$ loop, which is just as our anticipation. The inset in Fig. 3 shows typical $V-T$ loops for different initial vortex states, indicating that the choice of initial vortex states does not influence qualitatively the results in this simulation. Therefore, we conclude that either TSS or field strength affects only the size and shape but not the direction of the loops.

Shown in Fig. 4 is a typical clockwise $V-T$ loop of the vortex system with relative dilute and strong pinning force ($N_p = 0.1N_v$, $f_{pv0} = 90f_0$). On cooling from the temperature near T_c , most of the vortices enter quickly into the random pinning potentials (N_p^p quickly increases) at the temperatures near T_c because of the strong elementary pinning force and thermal fluctuation, which results in the occurrence of disordered pinned vortex states. With further decreasing temperature, vortex depinning occurs due to the faster increment of F^{vv} ($\sim \lambda^{-2}$) than F^{vp} ($\sim \lambda^{-1}$), as indicating by the corresponding decrement of the N_p^p/N_v in Fig. 4. The physics underlying the depinning is that a softer vortex is more efficient in vortex pinning than a stiffer one. Practically the measured voltage V is determined by the global pinning force $F(\sim N_p^p F^{vp})$. That is, V is determined not only by the elementary pinning force F^{vp} but also by the number of pinned vortex N_p^p . If the contribution of N_p^p is more important than the one of F^{vp} , $F(\sim J_c H)$ decreases and thus V increases. On the contrary, if the contribution of F^{vp} is dominant, F increases and hence V decreases continuously, just as seen in Fig. 4. Furthermore, in the low temperature regime far from T_c , the depinned vortices form stable channels of vortex flow with the help of F^L as well as F^T , which is shown in the inset of Fig. 4^{12,14-19}. The moving vortices in the channels are more order than the pinned ones. This moving vortex structure maintains from $0.5T_c$ to $0.91T_c$ in the following heating process. At still higher temperatures on the heating branch, the moving structure is destroyed by strong thermal fluctuation. It is noted that at a given temperature N_p^p of vortex system with channel structure on the heating branch is smaller than the one on the cooling one. Therefore, voltages are higher in the heating process than in the cooling one, i.e., the $V-T$ loop is clockwise. This suggests that the clockwise $V-T$ or $R-T$ hysteresis curves could be observed in those su-

perconductors with few artificial strong pinning centers, just as observed in doped single crystals $2H - NbSe_2$ ¹⁰ and the twinned $YBa_2Cu_3O_7$ ⁸.

In order to check whether the strength or the density of pinning centers can govern the direction of the $V - T$ loop, we have to study the effect of the pinning strength f_{pv0} and density on the direction of the $V - T$ loop. Displayed in Fig. 5 are effects of pinning strength f_{pv0} on the $V - T$ loop. It is clear that the size of the $V - T$ loops becomes smaller with decreasing f_{pv0} , but the clockwise direction does not change. That is to say, whether the loop is clockwise or anticlockwise is not determined by vortex pinning strength (also by TSS or field).

Fig. 6 demonstrates the effects of density of pinning center on the $V - T$ loop. One can see a crossover in the $V - T$ curves, showing directly the change of the direction as well as the size of the loops for different density N_p/N_v . For a larger N_p/N_v , such as $N_p/N_v = 10$ (see curve A in Fig. 6), no hysteresis can be seen, suggesting that it is difficult to probe the metastable states by the $V - T$ loop measurement for a sample with the

denser and weaker pinning centers. Whereas with decreasing N_p/N_v , it is noted that an anticlockwise $V - T$ loop (curve B in Fig. 6), a crossover one (curve C in Fig. 6), and a clockwise one (curve D in Fig. 6, the anticlockwise part disappears gradually) occur successively. This confirms the dominant rule of pinning center density in determining the direction of $V - T$ loop. Besides, such a crossover of the $V - T$ loop (curve C in Fig. 6) hasn't been reported so far as we know.

In summary, we have performed simulations to study the effects of vortex pinning on metastable phenomenon in transport process. It is shown that the vortex systems with very dense pinning centers will manifest their metastable states as an anticlockwise $V - T$ loop. And the density of pinning center plays a crucial role in determining the direction of $V - T$ loop. Meanwhile pinning strength (also magnetic field and TSS) can only influence the size of $V - T$ loop. The simulation results agree well with the reported experiments and predict the existence of crossover in $V - T$ loops.

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- ¹ G. Blatter, M. V. Feigel'man, V. B. Geshkenbein, A. I. Larkin, and V. M. Vinokur, *Rev. Mod. Phys.* **66**, 1125(1994); E. H. Brandt, *Rep. Prog. Phys.* **58**, 1465(1995).
- ² Y. Yeshurun, A. P. Malozemoff, A. Shaulov, *Rev. Mod. Phys.* **68**, 911(1996).
- ³ T. Giamarchi and S. Bhattacharya, *cond-mat/0111052*.
- ⁴ J. A. Fendrich, U. Welp, W. K. Kwok, A. E. Koshelev, G. W. Crabtree, and B. W. Veal, *Phys. Rev. Lett.* **77**, 2073(1996).
- ⁵ H. Safar, P. L. Gammel, D. A. Huse, and D. J. Bishop, J. P. Rice, and D. M. Ginsberg, *Phys. Rev. Lett.* **69**, 824(1992).
- ⁶ W. Jiang, N. -C. Yeh, D. S. Reed, U. Kriplani, and F. Holtzberg, *Phys. Rev. Lett.* **74**, 1438(1995).
- ⁷ M. Charalambous, J. Chaussy, P. Lejay, and V. Vinokur, *Phys. Rev. Lett.* **71**, 436(1993).
- ⁸ S. O. Valenzuela, B. Maierov, E. Osquiguil, and V. Bekeris, *Phys. Rev. B* **65**, 60504(2002).
- ⁹ Z. L. Xiao, E. Y. Andrei, and M. J. Higgins, *Phys. Rev. Lett.* **83**, 1664(1999); Z. L. Xiao, E. Y. Andrei, P. Shuk, and M. Greenblatt, *Phys. Rev. Lett.* **85**, 3265(2000); *Phys. Rev. Lett.* **86**, 2431(2001), and references therein.
- ¹⁰ W. Henderson, E. Y. Andrei, M. J. Higgins, and S. Bhat-

- tacharya, *Phys. Rev. Lett.* **77**, 2077(1996).
- ¹¹ P. Moretti, M. -C. Miguel, and S. Zapperi, *Phys. Rev. B* **72**, 014505(2005).
- ¹² C. J. Olson, C. Reichhardt, R. T. Scalettar, G. T. Zimányi, and Niels Grønbech-Jensen, *Phys. Rev. B* **67**, 184523(2003); C. J. Olson, C. Reichhardt, and Franco Nori, *Phys. Rev. Lett.* **81**, 3757(1998).
- ¹³ H. Fangohr, A. Price, S. Cox, P. A. J. de Groot, G. J. Daniell, and K. S. Thomas, *J. Comput. Phys.* **162**, 372(2000).
- ¹⁴ Seungoh Ryu, M. Hellerqvist, S. Doniach, A. Kapitulnik, and D. Stroud, *Phys. Rev. Lett.* **77**, 5114(1996).
- ¹⁵ C. Reichhardt, G. T. Zimányi, and Niels Grønbech-Jensen, *Phys. Rev. B* **64**, 014501(2001).
- ¹⁶ H. Fangohr, Simon. J. Cox, and P. A. J. de Groot, *Phys. Rev. B* **64**, 64505(2001).
- ¹⁷ A. B. Kolton, D. Domínguez, and Niels Grønbech-Jensen, *Phys. Rev. Lett.* **83**, 3061(1999).
- ¹⁸ F. Pardo, F. de la Cruz, P. L. Gammel, E. Bucher, and D. J. Bishop, *Nature (London)* **396**, 348(1998).
- ¹⁹ K. Moon, R. T. Scalettar, G. T. Zimányi, *Phys. Rev. Lett.* **77**, 2778(1996).