

Monday
May 11, 1992

MORNING
CMG

CONVENTION CENTER ROOM B2

10:30 am Guided-Wave Photorefractivity

Charles L. Woods, U.S. Air Force, President

10:30 am Invited

CMG1 Photorefractive thin films and waveguides

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The fabrication of photorefractive waveguides for use in holographic processors, integrated-optic applications, and optical memory devices is a particularly desirable goal. To date, small diameter ($\leq 100 \mu\text{m}$) crystal fibers of materials such as SBN have successfully been grown, but their bulk and surface optical quality can be poor, and clad fibers are not readily available. Planar guides of Ti indiffused photorefractive LiNbO_3 are more easily fabricated, but the materials of most interest to the photorefractive community, such as BaTiO_3 , SBN or BSO, have not been produced as low-loss waveguide structures, apart from a few reports on highly multimode thin cleaved crystal wafers.¹

The use of guided wave structures in photorefractive media has the obvious advantage of optical confinement and its consequent reduction in material response times. For planar structures of typical guiding dimension 1–5 μm , this reduction is expected to be two or three orders of magnitude. Thin film optical materials have been produced previously using various evaporation or sputtering techniques, but the single crystal oriented nature of such films is often poor or irreproducible. We discuss here the three techniques of laser ablative growth, ion-beam implantation, and indiffusion for fabrication of planar guides in BaTiO_3 , SBN:Ce, and BSO/BGO.

Laser ablation is a relatively simple and convenient method for thin film optical growth. Provided attention is paid to such deposition parameters as substrate temperature, and incident laser flux, $\sim 1\text{-}\mu\text{m}$ thick single crystal epitaxial layers can be grown on suitable lattice matched substrates. So far, we have had success with BGO grown on ZrO_2 ,² for which single crystal growth of the correct phase $\text{Bi}_2\text{GeO}_{20}$ was achieved. The growth chamber used is illustrated in Fig. 1. BaTiO_3 on LiF has also been reported,³ but waveguide quality films were not produced as the substrate was cleaved rather than polished. There is also the problem of postpoling these ferroelectric films. We are currently growing $\text{BaTiO}_3/\text{LiF}$ thin films on polished substrates.

Ion-beam implantation via 2–3 MeV He^+ and H^+ ions has also been used to form guides in BaTiO_3 , SBN:65, SBN:52, and

BSO. Typical guide dimensions produced fall in the region of 2–20 μm , but so far the guide losses are rather high ($\sim 30 \text{ dB/cm}$), and the material's photorefractive response is minimal. Postannealing typically reduces these losses, but oxidation/reduction treatments may be required to restore the material's photorefractive response. Finally we are also starting work on indiffused (Ti and Cr) guides in BSO and BGO.

We discuss our progress so far in waveguide growth techniques, loss measurements, photorefractive characteristics, and prospects for the future.

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2. K. E. Youden, R. W. Eason, M. C. Gower, N. A. Vainos, *Appl. Phys. Lett.* **59**, 1989 (1991).
3. G. M. Davis, M. C. Gower, *Appl. Phys. Lett.* **55**, 112 (1989).

11:00 am

CMG2 Temporal evolution of the photorefractive effect in annealed proton-exchanged LiNbO_3 waveguides

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Waveguides based on proton exchange (PE) in LiNbO_3 substrates are of significant interest for integrated photonic applications; in particular, they are useful for efficient second-harmonic generation,^{1,2} because of their strong nonlinearity, large optical confinement, and simple fabrication process.

Recently, restoration of the nonlinear coefficients of the PE LiNbO_3 waveguides by thermal annealing was reported.^{3,4} Although it is known that the PE LiNbO_3 waveguides have higher resistance to the photorefractive effect than Ti-diffused LiNbO_3 waveguides,^{5,6} quantitative measurements of the effect in annealed proton-exchanged (APE) waveguides as well as the PE waveguides have not been reported to date.

In this paper, we report on the quantitative measurements of the photorefractive effect at irradiation wavelengths (λ_{ir}) of 532 and 633 nm in APE LiNbO_3 waveguides by using a technique similar to that used for Ti-diffused waveguides.⁷ From measurements of the temporal behavior of the photo-induced index change caused by irradiation for several hours, the saturated-index change, buildup time constant, and photorefractive sensitivity in APE LiNbO_3 waveguides were determined and compared to Ti-diffused waveguides.

Figure 1 shows the experimental setup to measure the photorefractive effect. The APE waveguides were fabricated by proton exchange at 220°C for 20 min in benzoic acid, and subsequent annealing at 350°C for 6 h. They guide only the lowest-order mode at 1.3 μm . Irradiation light produces an optical phase retardation between the two interferometer arms

due to the photo-induced index change which results in a change in the modulated output intensity of the 1.3- μm probe.

By measuring the time dependence of the probe power, we obtain the photo-induced index change $\Delta n(t)$. No measurable index change was observed at $\lambda_{\text{ir}} = 633 \text{ nm}$ for intensity of up to 7.1 W/cm^2 . However, the photorefractive effect caused the index to change when exposed to the 532-nm light, as shown in Fig. 2. The buildup of the index change is exponential and is governed by

$$\Delta n(T) = \Delta n_s(1 - e^{-t/\tau}), \quad (1)$$

where Δn_s is the saturated index change, and τ is the time constant.⁷ Figure 2 shows the time dependence of the photo-induced index change for two irradiation intensities fitted to Eq.(1) for 220 min of exposure.

The results are summarized in Table 1 where we have also listed the photorefractive sensitivity S (cm^2/J) defined by $S = \Delta n_s / I_{\text{ir}} \tau(I_{\text{ir}}$ is the irradiation intensity). For comparison, we have also listed the results on Ti-diffused waveguides.⁷ The photorefractive sensitivity S for APE waveguides, while immeasurable at $\lambda_{\text{ir}} = 633 \text{ nm}$, is at least three orders of magnitude smaller at $\lambda_{\text{ir}} = 532 \text{ nm}$ than that of Ti-diffused waveguides at $\lambda_{\text{ir}} = 633 \text{ nm}$. A comparison of the time constants indicates that the conductivity of the APE LiNbO_3 waveguides is lower than that of Ti-diffused waveguides. However, based on Glass's model, the smaller saturated index change Δn_s in the APE waveguides indicates that the absorption coefficient, the electro-optic coefficient, and Glass's constant may be smaller to account for the reduction in the sensitivity S , because S is independent of the conductivity.

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5. R. A. Becker, *Appl. Phys. Lett.* **43**, 131 (1983).
6. J. L. Jackel, A. M. Glass, G. E. Peterson, C. E. Rice, D. H. Olson, and J. J. Veselka, *J. Appl. Phys.* **55**, 269 (1984).
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11:15 am

CMG3 Photorefractivity in doped nonlinear polymers

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Photorefractive polymer materials offer a number of potential advantages over