

## Design of Optical Fibre Amplifiers J.E.Townsend and R.I.Laming

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### *Introduction*

During the 1980's there has been increasing interest in rare-earth doped, single mode fibres for devices and sensors, of which the most exciting results have been achieved with erbium doped fibre amplifiers (EDFA). Great progress has been made since the first EDFA was reported by Mears in 1987 [1]. A gain of 26dB was achieved whereas 46.5dB has now been demonstrated [2]. High speed, direct detection in fibre based transmission systems has allowed bit rates and transmission distances to be markedly increased. For example, use of twelve 1.48 $\mu\text{m}$  laser diode pumped EDFAs as power amplifiers has allowed transmission of 1.2 Gbit/s over more than 900km of fibre [3]. The power penalty arising from dispersion and the presence of the EDFAs was only 0.6dB at  $10^{-9}$  BER, with no observed error floor due to accumulated noise. A number of similar systems have also been demonstrated with alternative pump sources and device configurations [4,5]. Thus the compatibility of EDFAs with long-haul communications systems has clearly been demonstrated. However, simultaneous improvement in manufacture and design of doped fibres has been essential. Here, the techniques for manufacture of such fibres are described and design criteria for optimal performance outlined. First, however, the advantages of EDFAs are summarised.

All-optical amplifiers are attractive in fibre based communications networks for low maintenance and faster operating rates than conventional electronic repeaters. If amplification is achieved 'on line' by incorporating the amplifier into standard communications grade fibre then interface losses are reduced to negligible levels. Further, the geometry of optical fibres, with small waveguiding area and long interaction lengths allows efficient interaction of light with the dopant so that low power semiconductor lasers may be used to excite the active ions. A dopant is therefore required which can be incorporated into the fibre and is able to amplify signals at the chosen system wavelength, either 1.3 $\mu\text{m}$  or near 1.55 $\mu\text{m}$ .  $\text{Er}^{3+}$  ions fluoresce around 1.55 $\mu\text{m}$  when excited by optical pumping at one of a number wavelengths, as shown in figure 1 [6]. With reference to the figure it may be seen that emission occurs from the  ${}^4\text{I}_{13/2}$  level to the ground state ie. 3-level operation. Thus, a population inversion is essential to bleach out the optical loss at the emission frequency before gain can be achieved - a problem in bulk materials but easily overcome in the fibre geometry.

### *Fabrication methods*

Conventional manufacturing processes have been modified to incorporate erbium ions into standard fibres. Optical fibres are drawn, or pulled from preforms whose characteristics determine those of the fibre. Two processes are widely employed to prepare single mode or graded index preforms for telecommunications applications. The Modified Chemical Vapour Deposition process (MCVD) was proposed in the USA [7] and England [8] in the early 1970's and relies on gas phase oxidation of vapour phase reagents to form the glass. The Vapour Axial Deposition method relies on

hydrolysis of the same materials and was developed approximately 10 years ago by the Japanese [9]. The MCVD process is described here, with the modifications required to incorporate rare-earth dopants such as  $\text{Er}^{3+}$ .

#### MCVD

MCVD is based on the high temperature oxidation of reactant gases inside a rotating tube which is heated by an external source. A schematic of the process is shown in figure 2 [10]. Controlled quantities of the dopants are transported to the reaction zone by passing gases (e.g.  $\text{O}_2$ ) through the liquid precursors, usually halides, thereby avoiding contamination from transition metal ion impurities. Silicon tetrachloride ( $\text{SiCl}_4$ ) is used to form silica, the main component of optical fibres.  $\text{SiO}_2$  is attractive due to the low losses and low dispersion in the region 1300nm to 1550nm as well as its mechanical and chemical strength.  $\text{GeO}_2$  and  $\text{P}_2\text{O}_5$  are commonly added as fibre core material to raise the refractive index of silica. Conversely,  $\text{B}_2\text{O}_3$  or fluorine reduce the refractive index of the deposited glass.

The reactant gases flow, with  $\text{O}_2$ ,  $\text{N}_2$  and He, inside a rotating silica substrate tube mounted on a glass working lathe. The tube is heated by an external burner traversing in the same direction as the gas flow. Gas phase oxidation in the burner's hot zone (typically  $1700^\circ\text{C}$ ) results in the formation of oxide soot particles, which is deposited further downstream. As the burner passes over the deposited layer, it is fused to a glass by viscous sintering. Deposition temperatures depend on glass composition and are sufficiently high to sinter the deposited material but not so high as to cause distortion of the tube (usually between  $1300^\circ\text{C}$  and  $1650^\circ\text{C}$ ). Many layers of defined index are built up in this manner. After deposition of the required amount of material, the tube is collapsed to a solid rod by heating to approximately  $2000^\circ\text{C}$ , sufficient to soften the glass and allow collapse under surface tension.

Fibre is drawn from the preform by heating the rod and, in effect, stretching it to the required diameter (usually  $125\mu\text{m}$ ) under controlled conditions. During the process the fibre is coated with a plastic for protection.

#### *Fabrication of rare-earth doped fibres*

Several techniques for incorporating rare-earth ions have been developed, suitable for both VAD and MCVD. High vapour pressure organometallic reagents have been employed [11] in conventional MCVD, with the rare-earth ion being transported in the vapour phase to the reaction zone in the same manner as the glass forming material. The technique is simple to implement but the precursors are toxic, not widely available and very expensive.

An alternative, and widely used, vapour phase process was developed at Southampton University, to exploit the high purity achievable via the vapour phase, without needing novel reagents. Solid, low vapour pressure rare-earth halide precursors are used [12] and incorporated in the fabrication process as shown in figure 3. Rare-earth halide precursor is located in a dopant chamber, upstream of the substrate tube. During conventional deposition of the core glass the dopant chamber is heated to approximately  $1000^\circ\text{C}$  to achieve a significant vapour pressure of rare-earth halide. Accurate temperature control is essential for control of precursor vapour pressure and hence dopant concentration. The vapour is carried to the reaction zone with the glass forming reagents where it is oxidised and deposited as normal. Collapse and fibre drawing are conventional.

Fibres containing a few hundred parts per million (ppm) are readily prepared in this manner. Losses approach those of telecommunications fibres, as shown in figure 4, for a fibre doped with approximately 100ppm of  $\text{Er}^{3+}$ . Away from the dopant absorption bands losses of order 5dB/km are commonplace [13]. Further, backscatter measurements [14] show that the dopant is incorporated evenly in the glass. However, this method is not very versatile since it is difficult to add either high dopant concentrations or incorporate several types of ion simultaneously in predetermined ratios.

More recently, a novel method, whereby the ions are added from solution, has been developed at Southampton University [15]. The process, described in figure 5 overcomes the difficulties previously encountered. Although there is an increased risk of impurity contamination via the liquid phase, low loss fibres with high dopant concentrations are fabricated. Very simple modifications to the MCVD process allow the incorporation of dopant ions via solution. The conventional fabrication process is followed until the core layer is deposited. Here, a reduction in burner temperature of around 500°C allows the formation and deposition of the oxide soot without subsequent fusion of this porous layer. The tube is then removed from the lathe and the porous frit immersed in a dilute solution of the dopant ions. Additional glass modifying reagents can also be added at this stage, for example alumina can be included from  $\text{AlCl}_3$  in solution. Water and alcohols are commonly used as solvents since they are safe to handle and the halide precursors are soluble in both. After soaking, excess solution is removed, the tube replaced in the lathe and the porous layer dried. Dehydration is based on the VAD method [16].  $\text{Cl}_2$  is passed through the inside of the tube which is heated to approximately 1000°C, so evaporating residual solvent. As a result, hydroxyl ion contamination is insignificant, with residual concentrations of around 0.25ppm. Dopant ions are not volatile at this temperature and remained trapped in the frit. Fusion of the soot layer occurs at the temperature conventionally employed and collapse of the tube into the preform is also standard.

#### *Amplifier design parameters*

A number of parameters affect the design and performance of EDFAs, including host material, dopant concentration and pump wavelength. The importance of each of these on fibre design is discussed here.

#### *Host glass*

Telecommunications grade fibres typically comprise around 90% silica, with the remainder being compatible glass forming material of similar structure to control the refractive index. Although the spectral characteristics of the rare-earth ions are sensitive to surrounding ligand field (i.e. host), it might be anticipated that the spectral features are unlikely to be affected by such small changes in host composition. Nevertheless, the addition of a few percent of index modifying ions is found to be sufficient to alter the fluorescence spectra of dopant ions.

In addition to host effects, however, the gain spectrum is highly dependent on the relative populations of the metastable and ground states (hence pump power) [6], because  $\text{Er}^{3+}$  operates as a 3 level system. Gain spectra for highly inverted germanosilicate and aluminosilicate  $\text{Er}^{3+}$  doped fibre are shown in figure 6. Non-uniform emission spectra arise from Stark splitting of metastable and ground state

energy levels, with the selection of levels being host dependent. In  $\text{GeO}_2\text{-SiO}_2$ , the gain spectrum is highly wavelength dependent, peaking at  $1.535\mu\text{m}$ . Typically, the 3dB bandwidth is  $\sim 3\text{nm}$  for a 30dB amplifier. The peak gain cross section has been determined experimentally and is found to be higher in germanosilicate core EDFA's than those containing alumina. The peak gain is shifted to  $1.531\mu\text{m}$  in  $\text{Al}_2\text{O}_3\text{-SiO}_2$  core fibre and, for a 30dB amplifier, the 3dB bandwidth is  $7\text{nm}$ . However, with reference to the figure, it may be seen that there is a relatively constant gain region extending from approximately  $1.538\mu\text{m}$  to  $1.560\mu\text{m}$ , where gains in excess of 20dB have been demonstrated. Atkins [17] has exploited this region, under reduced inversion, to achieve a 25dB gain amplifier with 3dB bandwidth of  $35\text{nm}$ . More recently however, a novel method for generating a broad band, spectrally flat gain spectrum has been demonstrated by Tachibana [18]. An essentially flat gain of 27dB is obtained over  $33\text{nm}$  without loss of pump efficiency or saturation performance by incorporating an optical notch filter in the middle of the amplifier. The flattened gain spectrum is shown in figure 7. Thus, in general,  $\text{GeO}_2\text{-SiO}_2$  core EDFAs are preferred for preamplifier applications, where high efficiency and gain cross sections are required whilst the  $\text{Al}_2\text{O}_3\text{-SiO}_2$  host is preferred in broad band systems.

$\text{Er}^{3+}$  suffers parasitic, or excited state, absorption (ESA) from the metastable energy level at several wavelengths. A significant ESA cross section around potential pump energies or the signal band decreases the population inversion with increasing pump power and consequent reduction in performance. ESA is measured by comparing the spectral attenuation of an erbium doped fibre when pumped (i.e. under population inversion) and unpumped. Significant pump ESA is recorded at a number of potential pump wavelengths, as shown in figure 8 for an  $\text{Er}^{3+}$  doped  $\text{Al}_2\text{O}_3\text{-SiO}_2$  fibre. Bands around  $488\text{nm}$ ,  $514.5\text{nm}$ ,  $655\text{nm}$  and  $810\text{nm}$  are seriously affected. However, a strong host dependence is recorded, as shown in Table 1 [19], where the ratio of ESA and GSA cross sections is tabulated for a number of hosts.  $532\text{nm}$  may be employed to pump EDFAs, but  $980\text{nm}$  is a more efficient pump wavelength. In-band pumping around  $1485\text{nm}$  is also employed [20]. Signal ESA is measured in a similar manner and no significant ESA is recorded near  $1.55\mu\text{m}$ .

Host composition also determines emission decay times of the fluorescence. Differences in the lattice vibration frequency (phonon energy) of  $\text{GeO}_2\text{-SiO}_2$  and  $\text{Al}_2\text{O}_3\text{-SiO}_2$  hosts leads to lifetimes of around  $12\text{ms}$  and  $\sim 10\text{ms}$  for low  $\text{Er}^{3+}$  concentrations in the two materials, respectively. This parameter alone does not determine quantum efficiency of an amplifier thus the shorter emission decay time in  $\text{Al}_2\text{O}_3\text{-SiO}_2$  is not a drawback.

#### *Dopant concentration*

Both fluorescence lifetime and quantum efficiency (QE) are affected by the ion concentration in the fibre. QE data give the clearer indication of amplifier performance but are difficult to measure directly. Data presented in figure 9 show the dependence of pump energy conversion efficiency on dopant concentration [21] and give an indication of QE. A  $980\text{nm}$  pump source was employed to generate the  $1535\text{nm}$  emission so 64% conversion efficiency corresponds to quantum limited operation (100% photon conversion). A strong dependence on dopant concentration is recorded, particularly in  $\text{GeO}_2\text{-SiO}_2$  for  $\text{Er}^{3+}$  concentrations in excess of 200ppm, but the effect is alleviated by the addition of alumina to the host. The reduction in QE is attributed to a cooperative upconversion mechanism, described in figure 10 [22]. Energy transfer takes place between neighbouring ions, both excited to the

metastable state ( ${}^4I_{13/2}$ ). The acceptor ion is promoted to the  ${}^4I_{9/2}$  level whilst the donor returns to the ground state. For the process to occur the ions must be located close to each other and such an effect is not expected at low dopant concentrations. However, rare-earth ions are only slightly soluble in  $\text{GeO}_2\text{-SiO}_2$  glasses, tending to phase separate in the host, reducing average ion-ion separation. The slightly altered geometry of  $\text{Al}_2\text{O}_3\text{-SiO}_2$  glass allows the ions to redisperse. Nevertheless, a reduction in QE also seems to occur in  $\text{Al}_2\text{O}_3\text{-SiO}_2$  EDFAs, around 800ppm. In summary, the preferred dopant concentration is low, around 100ppm for efficient EDFAs based on  $\text{GeO}_2\text{-SiO}_2$  and a few hundred ppm in  $\text{Al}_2\text{O}_3\text{-SiO}_2$  core fibres.

The limit on dopant concentration determines the amplifier length, owing to the contribution to gain from each ion. The ratio of absorption and emission cross sections has been determined as 1.2 to 1 at the absorption peak, for both alumina and germania doped silica based hosts. Further, the absorption per unit ion is around  $0.05\text{dB}\cdot\text{m}^{-1}\cdot\text{ppm}^{-1}$  at 1535nm in a multimoded silica based fibre. Hence, to achieve 30dB gain from a finite pump power, 50dB absorption at 1535nm is typically employed, which, given the concentration limit of 100ppm, defines a minimum amplifier length of 10m.

#### *Dopant distribution*

The measured absorption in a length of  $\text{Er}^{3+}$  doped fibre is a function of the overlap between dopant and transmitted light as well as of concentration. First, the distribution of dopant in  $\text{GeO}_2\text{-SiO}_2$  core fibres does not follow a perfect step index and second, the modal overlap between dopant and pump or signal power is poor in a single mode fibre. The significance of these two features is discussed separately here.

The distribution of  $\text{Er}^{3+}$  in  $\text{GeO}_2\text{-SiO}_2$  fibre prepared by MCVD follows that of  $\text{GeO}_2$ , with a lower concentration of dopant in the centre of the fibre, as shown in figure 11. The distribution arises from evaporation of germania during the collapse stage of the MCVD process.  $\text{Er}^{3+}$  is assumed to be attached to the volatile component and is therefore, also evaporated. Conversely, in  $\text{Al}_2\text{O}_3\text{-SiO}_2$  fibres no loss of  $\text{Al}_2\text{O}_3$  or  $\text{Er}^{3+}$  occurs, as seen in figure 12. Thus, much of the light propagating along the germanosilicate core fibre will not interact with the dopant. The "effective" or measured dopant concentration is approximately half the true or local concentration, as determined by materials analysis.

Further, in single mode fibres, the overlap between dopant and pump mode is incomplete whilst as little as 50% of the signal power may propagate in the core of a fibre which is single mode around 950nm. The 3-level nature of  $\text{Er}^{3+}$  necessitates pumping half the ions from the ground to the metastable state to achieve gain. It is therefore desirable to locate the dopant where the pump intensity is highest, and to prevent there being any ions in regions of low pump power where a population inversion may not be obtained. Thus, the dopant should be confined to the region of highest power density for efficient operation of an EDFA, ie. the centre of the core [23]. The single mode signal beam also has its maximum intensity on the axis and therefore overlaps well with the confined excited ions. A reduced bleaching power (that required to just achieve an inversion) has been confirmed in confined EDFAs. Alumina is usually employed as a codopant in such a configuration to minimise loss of dopant from the centre and to allow higher dopant concentrations to be employed. A core occupancy of around 50% is effective in 980nm

pumped systems. An alternative approach, particularly suitable for in-band pumped systems, is to operate the EDFA well beyond cutoff, with only part on the power guided in the core region.

In both cases, the minimum amplifier length is increased by around 50%, giving a practical value of around 20m in GeO<sub>2</sub>-SiO<sub>2</sub> fibres. Baseline losses (e.g. waveguide imperfections) may then become significant. Typical values of  $5 \times 10^{-3}$  dB/m are recorded in doped GeO<sub>2</sub>-SiO<sub>2</sub> fibres, giving an excess loss of 0.1 dB per EDFA - an acceptable value. High performance amplifiers of 160m in length have been demonstrated. However, in Al<sub>2</sub>O<sub>3</sub>-SiO<sub>2</sub> fibres, losses typically of order 0.05 dB/m, due to the incorporation of impurities such as Fe<sup>2+</sup> in the doping process [13] although losses of 0.01 dB/m are obtained by the author. The loss scales with concentration of alumina. Hence, the excess loss can be significant in an Al<sub>2</sub>O<sub>3</sub>-SiO<sub>2</sub> EDFA more than 20m long, and elevated dopant concentrations should be used.

Given that it is preferable to operate the pump in a single mode configuration constraints are placed on the geometry of the core, particularly if a 980nm pump is employed. The modal parameters should be matched, at the signal wavelength, to dispersion matched telecommunications fibre. Also, the fibre should be single moded below 980nm. A high numerical aperture (around 0.2) is preferred for efficient pumping owing to the decreased core area and hence increased pump power density, although a core diameter in excess of 3µm is preferred for efficient launching. Further, the fibre waveguiding characteristics will be poor around 1550nm unless a high numerical aperture is used. It is found that all these criteria are fulfilled in a fibre of numerical aperture around 0.2 with second mode cutoff near 970nm. The geometric requirements are relaxed for in-band pumping since the pump wavelength (1485nm) is close to the signal wavelength (1535nm), but as stated above shorter cutoff wavelengths are preferred for effectively confining the dopant.

### Summary

In summary, for an optimum EDFA, low dopant concentrations, of order 100ppm, should be employed in GeO<sub>2</sub>-SiO<sub>2</sub> core fibres. Long lengths may be used owing to the low attenuation, similar to that of conventional fibres. Higher concentrations (a few hundred ppm) may be employed in Al<sub>2</sub>O<sub>3</sub>-SiO<sub>2</sub> core fibres to reduce the amplifier length, owing to the excess loss introduced by Al<sub>2</sub>O<sub>3</sub>. A host dependent gain spectrum obtained, with higher cross sections (more gain per unit pump power) being recorded in GeO<sub>2</sub>-SiO<sub>2</sub> fibres whilst the gain bandwidth is broader in Al<sub>2</sub>O<sub>3</sub>-SiO<sub>2</sub> EDFAs. Two pump wavelengths are preferred, 980nm or 1485nm, neither of which suffers ESA in either host. Fibres which are single mode at the pump wavelength are employed, with a high numerical aperture being essential in 980nm pumped systems. For efficient pumping the dopant is either confined to the centre of the core or the EDFA operated well beyond cutoff.

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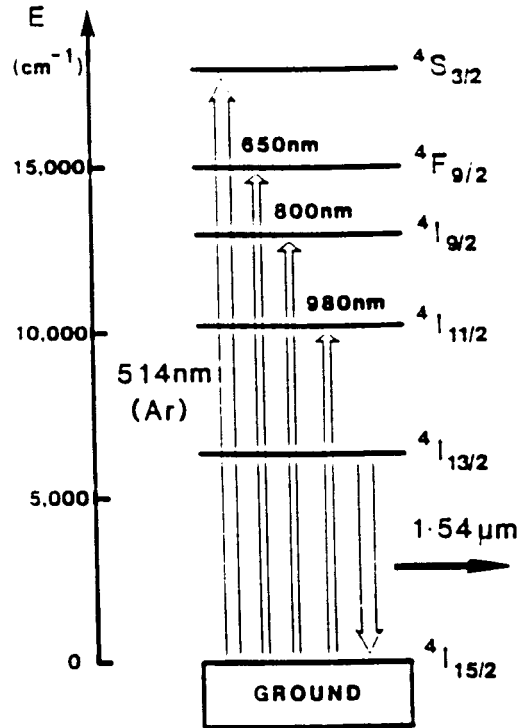


Figure 1. Energy level diagram of  $\text{Er}^{3+}$  showing possible pump wavelengths

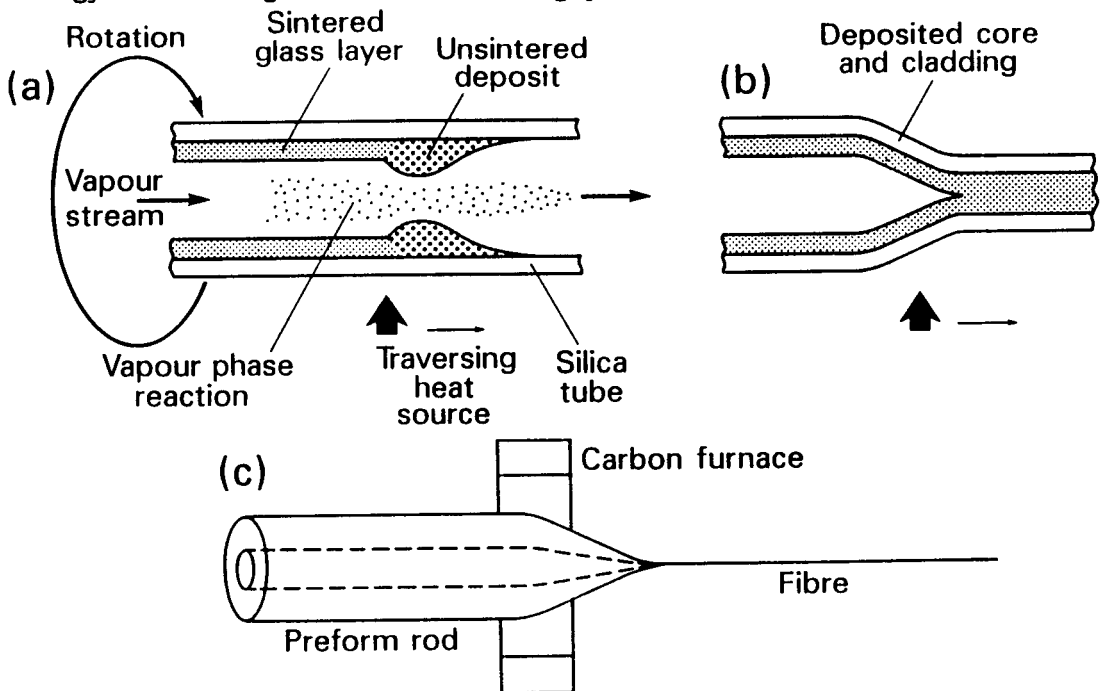


Figure 2. Schematic showing the MCVD method for the preparation of optical fibres; a) deposition, b) collapse to produce a preform and c) fibre drawing

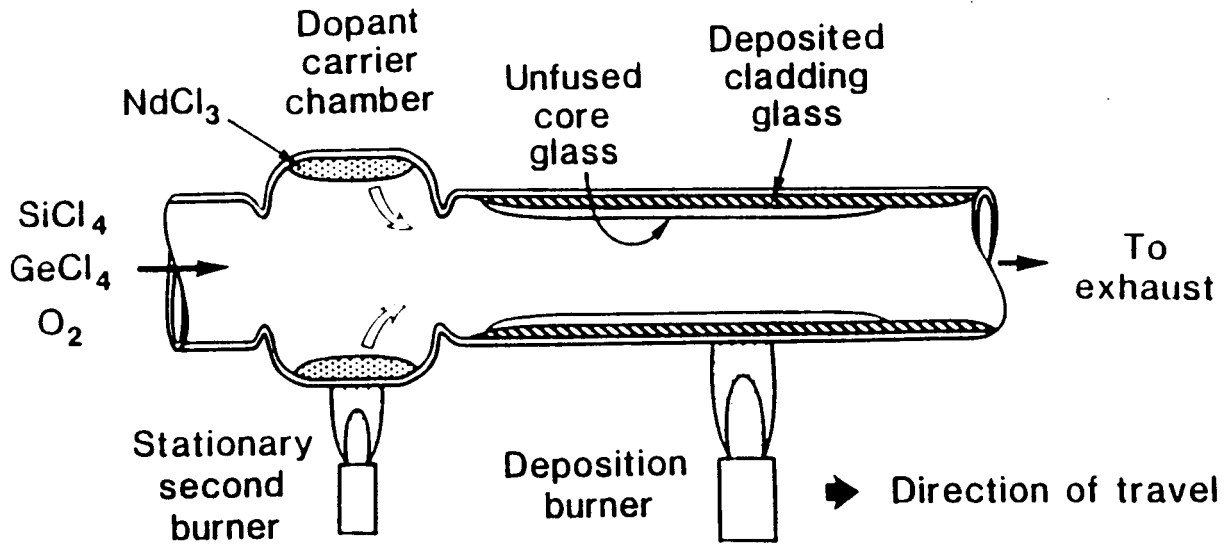


Figure 3. Vapour phase process for fabrication of rare-earth doped optical fibres

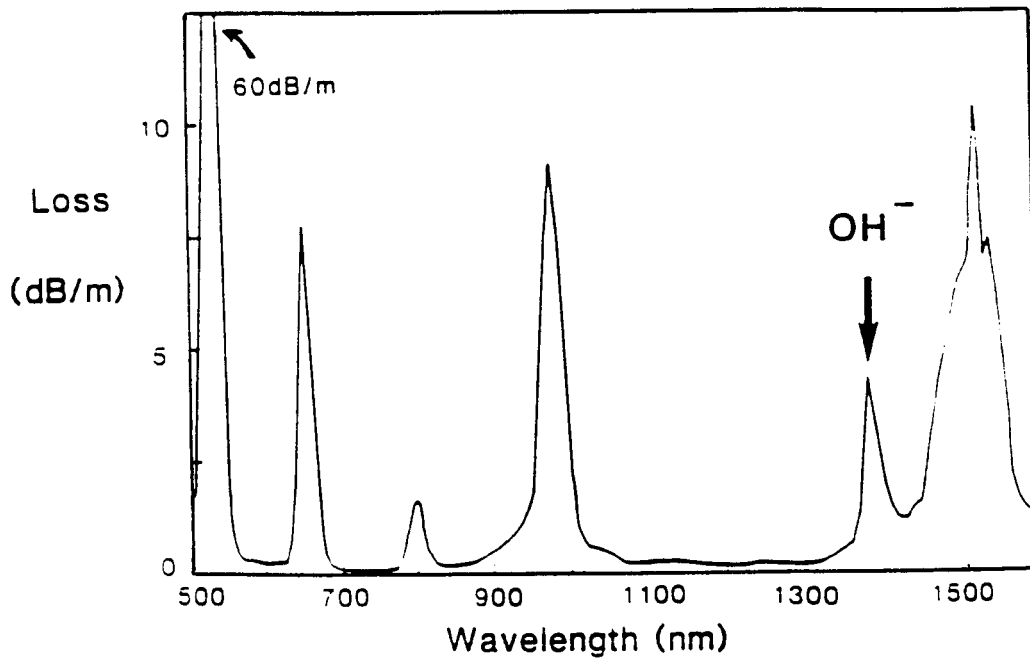


Figure 4. Absorption spectrum of optical fibre doped with  $\text{Er}^{3+}$  ions

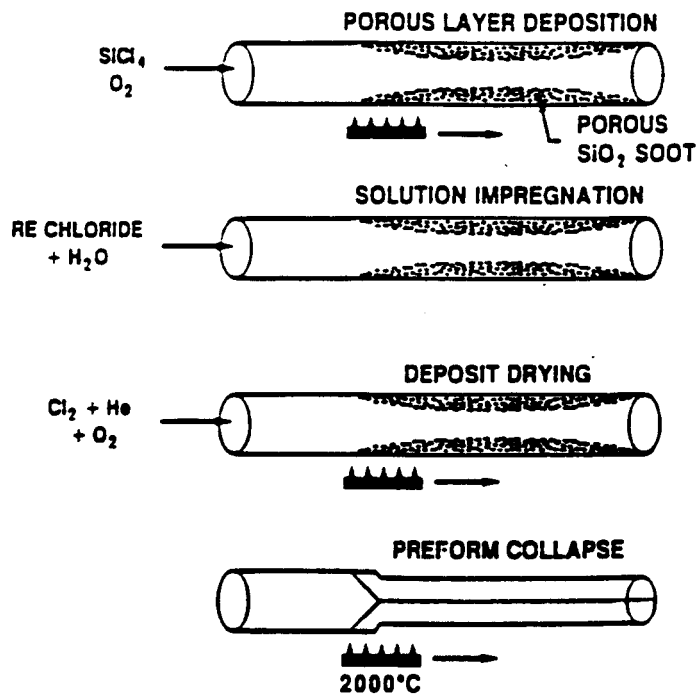


Figure 5. Schematic of the solution doping process for manufacture of rare-earth doped fibres via MCVD

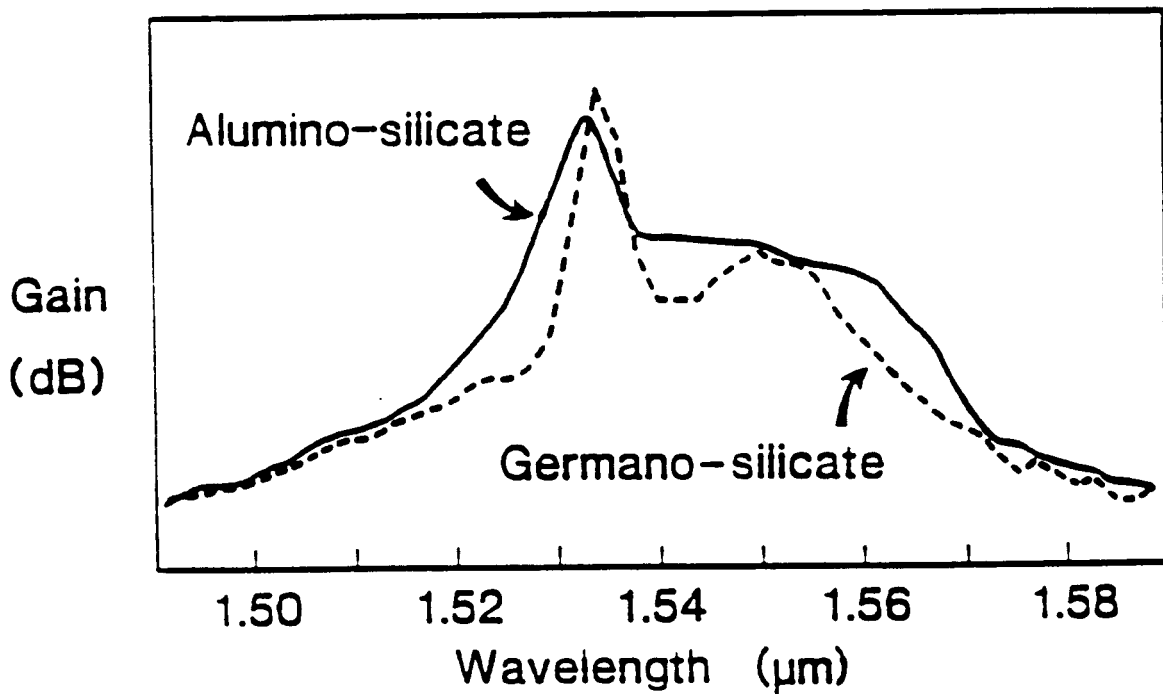


Figure 6. Gain spectrum of  $\text{Er}^{3+}$  ions in two optical fibre host glasses prepared by MCVD

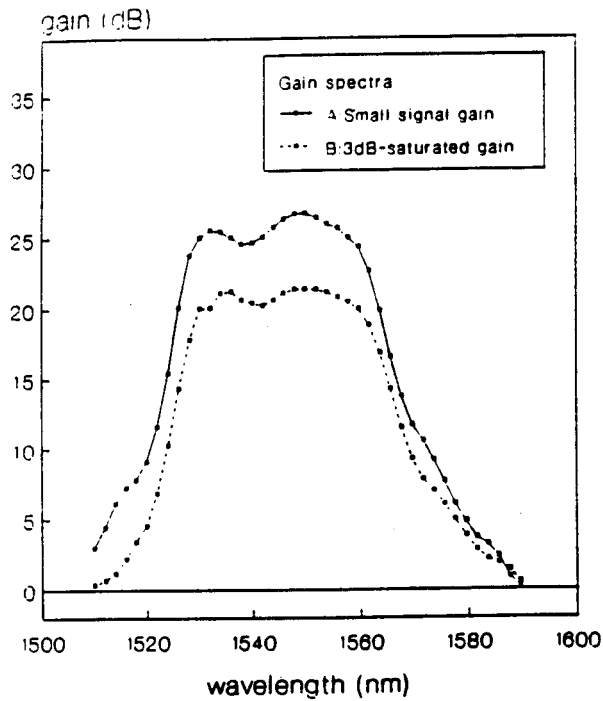


Figure 7. Flattened gain spectra for small signal and 3dB saturated operation of EDFA, including optical notch filter (after Tachibana et al)

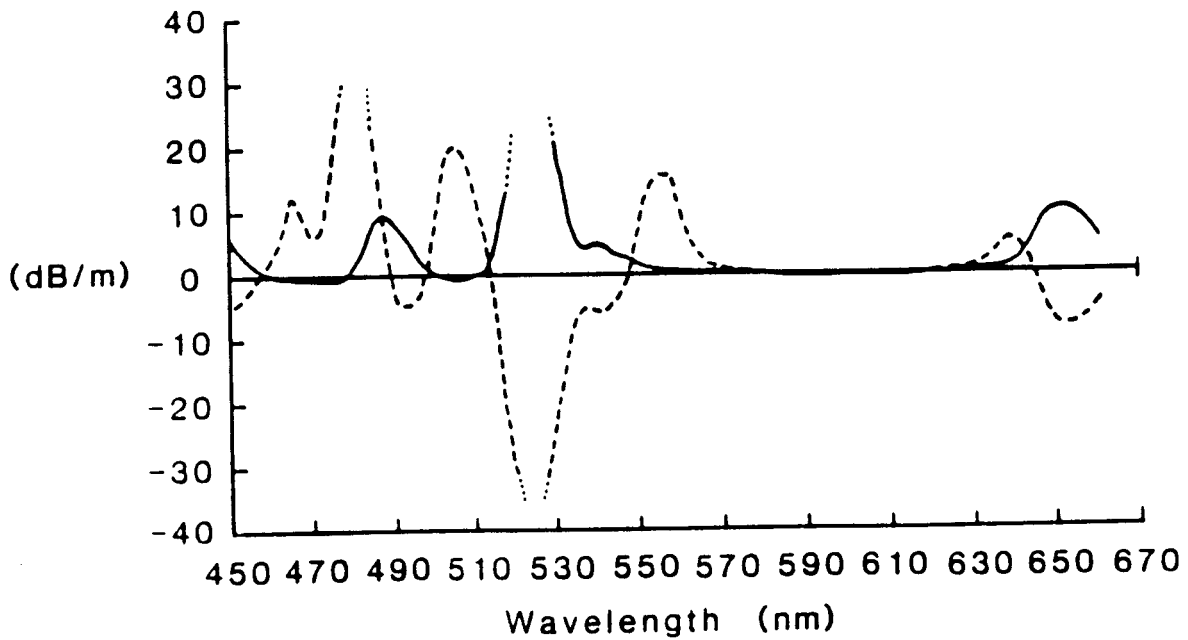


Figure 8. ESA and GSA spectra of Er<sup>3+</sup> doped aluminosilicate core fibre

| Fibre Type<br>Wavelength (nm) | GeO <sub>2</sub> | GeO <sub>2</sub> /B <sub>2</sub> O <sub>3</sub> | GeO <sub>2</sub> /P <sub>2</sub> O <sub>5</sub> | Al <sub>2</sub> O <sub>3</sub> |
|-------------------------------|------------------|---|---|--------------------------------|
| 488                           | 2.9              |   | 1.86  | 1.74                           |
| 514.5                         | 0.95             |   | 0.55  | 0.5                            |
| 655                           | 0.28             | 0.25  | 0.13  | 0.14                           |
| 810                           | 2.0              | 2.0   | 1.0   | 1.0                            |
| 980                           |                  |   |   | 0                              |

Table 1. Ratio of ESA and GSA absorption cross sections,  $\sigma_{\text{ESA}}/\sigma_{\text{GSA}}$  as a function of wavelength (after Laming et al)